

Range expansion switches clonal interference patterns in evolving populations

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ABSTRACT

Predicting g evolutionary processes is a sought after goal in a range of fields ranging from medicine, biology, engineering, and industry. However, our ability to predict evolution in real populations is limited in part by our understanding of how spatial structure modulates evolutionary dynamics. In this study, we examine beneficial mutant fixation dynamics along the front of an asexual population expanding its range. Using an agent-based simulation of range expansion, we observe that multiple distinct evolutionary regimes of beneficial mutant origin-fixation dynamics are maintained at characteristic length scales along the front of the expanding population. We find that the length scale at the tip of the front within which 'local' mutant fixation occurs in a successive mode decreases with increasing mutation rate, as well as population size in a manner predicted by our derived analytic expressions. Additionally, we discuss the relevance of our findings to real cellular populations.

1 Introduction

Beneficial mutants and their substitution within a population encapsulate the very crux of natural selection, especially in its classic conceptualization. Using his now ubiquitous plots of genotypic frequency over time, H.J. Muller demonstrated how beneficial mutants can fix long before subsequent beneficial mutants arise, thus accumulating in a successive manner over time.^{1,2} Using a theoretical population genetics approach, J.H. Gillespie influentially termed these dynamics strong selection weak mutation (SSWM), after the relative balance between population size, mutation frequency, and selective advantage he derived for evolution under this model.³⁻⁵

Gillespie further demonstrated that evolution proceeding in this regime can be conveniently approximated using the robust mathematical construct of a Markov chain, in which evolution proceeds as an adaptive walk across a mutational landscape. Gillespie, and many others thereafter, have exploited the mathematical convenience of this regime to explore adaptation within molecular evolutionary data, devise algorithms to control the speed and direction of evolution, and examine the computational limits of evolution.⁶⁻⁹ More generally, this framework has proven a useful qualitative description for beneficial mutation and selection, as these ideas have been applied to experimental biology since Gillespie's theoretical work.¹⁰ It is useful to be able recognize evolution under the SSWM regime given this Markov-chain representation of evolution, and the powerful analyses it allows (**Box 1**).

However, making such assumptions once spatial complexity is introduced presents a challenge. The spatio-temporal dynamics of mutants introduced in a migrating population can be surprising, complex, and resistant to analytical solutions. After spread into a new environment, the individuals at the very tip of the expanding population are preferentially represented in the resulting population's genetic makeup, as observed across scales and domains of life.¹¹⁻¹³ Relatedly, mutants at the front of a population expanding its range have an increased probability of survival and further proliferation at the wave front, a phenomenon known as gene surfing.¹⁴ Such results illustrate a fascinating positional dependence of evolutionary dynamics in spreading populations, an observation that has sparked many lines of investigation both experimental and theoretical.¹⁵⁻¹⁸ Great theoretical progress on population expansions has been made analyzing modelling frameworks based on stochastic reaction-diffusion systems.¹⁹ Recently this has included work on the properties of range expansions including range expansion with neutral mutations,²⁰ range expansion in the strong noise limit,²¹ selection in the absence of mutations,²² long-range dispersals,²³ and cooperation during range expansion.^{24,25}

The interference of beneficial clones as they fix through a spatially extended population undergoing mutation and selection, however, represents an analytical challenge. The dynamics of clonal interference, even within well-mixed

populations, are complex, and an active area of investigation in which recent advances rely heavily on simulation-based studies.^{10,26} In the case of spatially extended populations, Korolev et al. have shown (in a one-dimensional stepping stone model with two alleles with unequal fitness benefits and a non-zero mutation rate between them) that a closed-form analytic solution to the spatiotemporal dynamics of the allele frequencies cannot be achieved without approximations that ultimately yield largely inaccurate results.²⁰

Despite these limitations, Martens and Hallatschek have made progress in modeling the interference of beneficial mutants in a spatially structured population by considering low mutation rates and relatively large spatial scales.²⁷ The authors show that during migrations across large habitats, waves of beneficial mutants arising on distinct genetic backgrounds are likely to interfere beyond a certain characteristic habitat length, determined by the frequency and strength of the mutations. Within this habitat length, one beneficial mutant wave fixes locally on average, while interfering with previous waves further from the edge of the mutant beyond this habitat length. We note that in this work Martens and Hallatschek refer to this as ‘clonal interference’, while this term, as originally coined by Gerrish and Lenski, refers to interference between two beneficial mutant arising in the same genetic background.²⁸ In the language of Gerrish and Lenski, Desai and Fisher, and others, Martens and Hallatscheks’ work on the inference of slowly advancing waves of beneficial mutants with distinct genetic backgrounds during long-range migration is the ‘multiple mutation’ phenomenon, and hereafter we will refer to these separate phenomena using this convention.^{5,29}

In this study we will build upon previous work showing dynamical patterns that emerge during the process of mutants arising and fixing within their local environment are positionally dependent with respect to origin along the population front of expanding populations. We employ agent-based simulations of a modified one-dimensional stepping stone model and analysis based off of the estimated time and length scales of mutant establishment and fixation. We demonstrate that within a co-moving frame, a stable sub-population at the tip of a spreading population front ‘locally’ fixes beneficial mutants analogous to a well-mixed population within the successive-mutant regime, while simultaneously, proximal regions of the wave front experience clonal interference. In addition, we discuss the biological relevance of our results as a whole.

Model

Our model begins with a one-dimensional lattice configuration representing a linear habitat. Along the habitat, at each lattice site are colonization sites with a fixed capacity of K particles, called demes. Within each deme, these particles can either represent a vacancy or an individual, with some number of mutations. For each simulation step the following three steps are carried out:

1. A migration step in which two neighboring demes are chosen at random and a random particle from each deme is swapped with one another.
2. A duplication step, in which two particles are chosen from a randomly chosen deme, and the second particle is replaced by a duplicate of the first with some probability.
3. A mutation step, in which a random particle is granted a mutation with some probability.

These steps are schematized in **Fig. 1**. The probability of duplication during the duplication step is 1, except for the case of replacing an individual with a duplicate of a vacancy (i.e. death), which occurs with probability $1-r_m$, where r is a constant that varies with m , the number of mutations the particle has acquired. The transition probability of a non-vacancy particle acquiring one mutation during a time step is U_b and 0 for all other possible mutational events (i.e. no backward or multiple mutations are allowed within a time step). The above model is akin to a modified stepping stone model used in previous literature and generates stochastic Fisher-Kolomogrov-Petrovsky-Piscounov (sFKPP) waves undergoing selection and mutation.³⁰ The speed of the Fisher wave established from any individual is dependent upon the parameter r_m , which can shown to be analogous to the growth rate of the corresponding sFKPP system when written in the continuous limit with appropriate assumptions (See **SI, sec. ??**).

Simulations are initialized using a standing wave solution to the deterministic Fisher equation and we use an initial wild type wave of growth rate $r_0 = 0.1$. As the cells spread throughout the linear habitat they undergo rare mutations that confer the organism a growth rate of $r_m = r_0\alpha^m$, where α quantifies the average strength of the fitness benefit conferred by the mutation, and $\alpha = 1 + s$ for $s > 0$. As we are interested only in the advancing wave front, during the simulation full demes furthest behind the wave tip are periodically dropped and empty demes are added past the wave tip.

In this way, the simulation field represents the wave front, traveling at the average velocity of the spreading population. As mutants accumulate over time, their clonal sub-populations grow and establish into travelling waves

81 with some probability, eventually causing the wild-type population to fall behind the mutant wave-front. Simulations
 82 stop at the extinction of a specified number of mutants within the wave front, with the number of distinct clonal
 83 populations in each deme recorded. All simulations were performed in the python programming language with the
 84 Numba package to optimize computational speed.³¹

85 Results

86 Fate of a beneficial mutant

87 At short spatial and temporal scales and sufficiently large deme sizes, we estimate that the dynamics of mutant
 88 establishment are roughly as in a well-mixed population. To emphasize this analogy we define $s \equiv \alpha - 1$. At
 89 sufficiently low mutational supply and strong selective mutant strength, a mutant establishes locally with probability
 90 $\sim s$, analogous to the ultimate fixation probability of a beneficial mutant in a well-mixed population of constant
 91 size (See Box 1).⁵ Accordingly, the reciprocal of this probability is the approximately average timescale of local
 92 establishment, $\tau_{est} \sim 1/s$. As shown by Lehe et al., the probability that a single mutant wave emerges from the
 93 wild-type wave after local establishment and goes on to outpace and take-over the wild-type wave (or surf on the
 94 wave front), in the absence of clonal interference, saturates at $r_0(1+s)$ (conditioned on establishment) as the position
 95 at which the mutant is introduced goes further into the sparsely populated tip of the wave front (See Box 2).^{14,22}
 96 Before a mutant establishes, at a time $t < \tau_{est}$, it diffuses along the axis, and, if it fails to successfully establish, is
 97 left behind by the traveling population front. With sufficient mutational supply and strength, multiple mutants
 98 may arise along the wave front and interfere if their spatial scales of diffusion overlap or if one or more subclones
 99 establishes into a traveling wave and migrate to neighboring demes (See Fig. 3).

100 Length scales of evolutionary regimes

101 In our initial simulations, the genetic diversity (number of distinct subclones) in each deme throughout the wave
 102 front was recorded at the time of wild-type extinction. From our simulation data, with sufficient mutation rate,
 103 deme size, and mutation strength, we observe a length scale extending from the very wave tip to some position
 104 within the bulk of the wave front in which the population remains genetically pure on average (Fig. 5). Following
 105 the scaling arguments of Martens and Hallateschek, we can quantify the time scales of the previously described
 106 scenario of mutant inference to illustrate the quantitative relationships between these parameters and this length
 107 scale, which we denote as L_{SSWM} .²⁷ A single beneficial mutant that has established into a traveling wave will fix in
 108 time τ_{fix} by travelling across a length L such that

$$\tau_{fix} \propto L/v_m, \quad (1)$$

109 where v_m is the relative velocity of the mutant wave. Within the length scale L , the average time for mutants to
 110 establish (τ_{est}) locally is estimated by the following relation:

$$\tau_{est} \propto 1/(LU_b s), \quad (2)$$

111 somewhat analogously to a well-mixed population (See Box 1). These times are comparable at a particular length
 112 scale at which mutants begin to interfere locally as they fix. We find L_{SSWM} by setting $\tau_{fix} \sim \tau_{est}$:

$$L_{SSWM} \propto \sqrt{\frac{v}{U_b s}}. \quad (3)$$

113 Within this length scale, we expect local fixation of an established mutant to occur before another mutant establishes
 114 locally, on average, analogous to a well-mixed population within the SSWM regime. Thus, from Eq. (3) we expect
 115 L_{SSWM} , the length scale at the tip of the wave in which we expect mutants to appear successively to go as $U_b^{-1/2}$
 116 which agrees well with our stochastic simulations, see (Fig. 5A).

117 As expected, this behavior varies strongly with deme size. As the total length of the wave front, L_{front} , defined as
 118 the furthest deme with greater than one individual, varies linearly with total population size, within the wave front
 119 we expect a linear relationship with $\log K$. This can be demonstrated by noting that the initial wild-type population
 120 density at the tip of the front in the continuous limit $b_w(x)$ (As in Box 2 Fig. L), can be approximated as an

121 exponential decay function,

$$b_w(x) \approx e^{-v_w x/2}, \quad (4)$$

122 for $x > 0$ and $b_w(x) \equiv 1$ everywhere else, with position x and velocity v_w . As x for which $K * b_w(x) = 1$ determines
123 the length of the wave front, L_{front} , solving Eq. (4), we obtain

$$L_{front} \approx 2 \log K / v_w. \quad (5)$$

124 Equation (5) implies that L_{SSWM} expressed as a proportion of the total front, L_{SSWM}/L_{front} , hereafter \tilde{L}_{SSWM}
125 for simplicity, is inversely proportional to $\log K$. This relationship is observed within our simulation data as well (**Fig.**
126 **5B**).

127 We additionally observe that with relatively weak mutational supply or mutation strength, this heterogeneity in
128 evolutionary regime is not observed, and instead the entire wave front accumulates beneficial mutants successively.
129 We estimate this to occur when the mutational supply relative to its selective advantage is sufficiently small:
130 $U_b \ll \sqrt{s}/K$. When this constraint holds, it is expected that $\tilde{L}_{SSWM} = 1$ on average throughout the wavefront
131 (See **SI, Section 2.3**).

132 To further examine the relationship between s and L_{SSWM}/L , we simulated evolution in our system through
133 multiple extinction events and observe that if the evolutionary regime switches at some position in the wave tip, this
134 regime switching is sustained and roughly invariant over multiple extinction events and across evolutionary scales
135 (**Fig. 5C**). As each mutation event on average occurs at a particular time in the ‘cycle’ of mutant establishment
136 and fixation (which parameter specific), $\tilde{L}_{SSWM}L$ is seen to vary stochastically over time, though a predictable
137 relationship between \tilde{L}_{SSWM} with s is observed over an average of many extinction events. Our simulations of
138 multiple extinction events were limited by the maximum fitness at the end of simulations, which we ensured obeyed
139 our analytic assumptions, i.e. $r_m \ll 1$ by performing these simulations at low values of $s \in 0.01, 0.03, 0.05$. Multiple
140 fixation events at higher values of s or K would allow for mutants that violate the assumption of strong migration
141 ($r_m \gtrsim 1$), under which our analysis holds.

142 The relationship between \tilde{L}_{SSWM} and the selective advantage at each extinction event in our simulation is more
143 nuanced as s determines mutant establishment probability, velocity, and the shape of the front profile. While the
144 velocity v for a classic Fisher wave $v \sim 2\sqrt{Dr}$ with growth-rate r , there are two important deviations from this
145 classic behavior.

146 First, as waves interfere with each other, as is the case within the parameter regime we simulated, Martens and
147 Hallatschek demonstrated that their velocity slow with increasing interference, though this velocity does saturate
148 at a critical value. Secondly with decreased s relative to the strength of the noise, the velocity deviates from the
149 deterministic velocity $2\sqrt{s}$, and becomes roughly proportional to the product of s and wave density, as opposed to
150 \sqrt{s} in the deterministic limit²¹. These two phenomenon underscore the complexity in the relation between L_{SSWM}
151 and v and s . At strong noise relative to s as in the parameter regime simulated, for example, we expect $v \propto s$ and
152 thusly \tilde{L}_{SSWM} to be invariant with s , from equation (3). Simulation results from $s = 0.01, 0.03, .05$, $K = 100, 200, 500$
153 and $\mu = 10^{-3}$ show \tilde{L}_{SSWM} to be roughly invariant with s in this regime (**Fig 5D**).

154 Discussion

155 An enduring goal of population genetics is to understand the temporal dynamics of genotypic frequency and the
156 selective forces that shape them. When mathematical analyses or simulation models are used for these means, it is
157 important, then, to consider the extent to which our model does and does not reflect known biology to be able to
158 evaluate the validity of the results it provides.

159 To this end, we first consider the range of experimentally observed beneficial mutation rates. Beneficial mutation
160 rates have historically been posited to be very low, and early experiments supported this with estimates on the order
161 of $10^{-9} - 10^{-8}$ per genome per generation.^{28,32} However, subsequent experiments that have sought to eliminate
162 the bias introduced by non-optimized growth conditions, and unaccounted-for clonal interference, variously suggest
163 beneficial mutation rates may be on the order of 10^{-4} . Accurately measuring and even identifying mutations as
164 beneficial in an empirical context is still an active area of investigation.^{28,32} Our model examines parameters an
165 order of magnitude above below this upper limit.

166 We secondly consider the strength of beneficial mutations observed previously. Recent studies have been in closer
167 agreement with measurements often from .01 – 0.05. In our study, we examine this range of selective advantages
168 within and up to an order of magnitude above the empirical values.

169 Thirdly, we address deme size in our model and its relation to real populations. We note that in a one dimensional
170 stepping stone model, deme size is effectively a linear density term, however its relation to real population is not
171 straight-forward. In this study, as in others we explore a relatively wide range of deme size to capture effects specific
172 to any given regime. However we can attempt to gain a rough idea of the relationship of deme size to population
173 density of real populations by stating deme size in terms of maximum individuals per diffusion length, or average
174 distance diffused within a generation. Given $D = 1$ and $r_0 = 0.1$, in our model $\sqrt{D/r_0}$ describes a diffusion length,
175 and $K\sqrt{D/r_0}$ the maximum number of individuals within a diffusion length. For values of K ranging from 100 to
176 1000 as examined in our simulations this corresponds to $\sim 300 - 3000$ individuals per diffusion length.

177 To compare this estimate to real cellular populations, we can take the example of bacterial colonies, which are
178 observed to exhibit diffusive motility within a large range of densities.³³ To get a sense for this range we refer to
179 Butler et al. and Perfeito et al. who report report a migrating cell surface density of $\sim 3 \times 10^6 \text{ cells/mm}^2$ on agar
180 plates. We further use an average generation time (i.e. experimental r_0 of $.7 \text{ hr}^{-1}$) and a range of cell diffusivities
181 from 0.7 – 5.7 mm/hr, observed with varying agar concentration for *E. Coli*.³⁴ We can further suppose a minimum
182 possible surface density of close packed *E. Coli* without deformation as 1 cell per μm^2 . Converting the experimental
183 surface densities to a linear density ρ , we find individual per diffusion length as $\propto \rho * \sqrt{D/r}$.

184 This yields a range of 1000 – 5000 cells per diffusion length and overlaps significantly with the estimated linear
185 density examined in our model, suggesting the model parameters examined and the corresponding behavior observed
186 are biologically plausible. Furthermore, at higher densities than those examined in our model, shifting clonal
187 interference patterns within a wavefront can be observed at even lower mutation rates (See **SI ??**).

188 The aforementioned discussion suggests a biological relevance for interference with relatively high beneficial
189 mutation rate and short spacial scales as examined in our model. We have shown that in this regime within a
190 relatively wide range of biologically relevant parameters, a region at the tip of the wave front fixes mutations ‘locally’
191 in a successive manner, analogous to the SSWM regime in well-mixed populations. Such a property is potentially
192 useful for applying analyses such as a mutational landscape model, or considering theoretical work built on models
193 assuming SSWM regime for the crucial tip of a wave expanding its population. In a variety of contexts, as the
194 population ‘pioneers’, the individuals at the tip of the population and their dynamics can be considered exclusively
195 as they disproportionately influence the genetic diversity of the population following migration. In addition, our
196 results demonstrate how the effective evolutionary parameters of a population shift through range expansion. The
197 implications of the results for inferring evolutionary dynamics, for antibiotic resistance assays in clinical contexts, for
198 instance, is that the dynamics observed at a specific time and place in the population maybe transient and highly
199 unrepresentative of the entire population. While the spatiotemporal dynamics of a population expanding its range
200 and acquiring mutations subject to selection can be highly complex, evading exact analytic treatment, we believe the
201 our work demonstrates that analyses from well-mixed populations can be applied to the tip of the wave as long as it
202 is maintained within the SSWM regime. To this end, we have estimated, in terms of ‘real parameters’ when this
203 phenomenon is observed and the dependence of this scale length of local successive mutation fixation on relevant
204 population genetic parameters. In summary, we shed light on the phenomenon of shifting clonal interference patterns
205 during range expansion with relevance and application across evolutionary and scientific contexts.

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276 **Author contributions statement**

277 NK performed the mathematical analysis, wrote the code, performed the simulations, analyzed the data and wrote
278 the manuscript. DF and JGS analyzed the data and wrote the manuscript.

279 **Code availability**

280 The code to perform the numerical simulations is available via github at [https://github.com/nkrishnan94/](https://github.com/nkrishnan94/Range-Expansion-Eco-Evo-Regime)
281 [Range-Expansion-Eco-Evo-Regime](https://github.com/nkrishnan94/Range-Expansion-Eco-Evo-Regime).

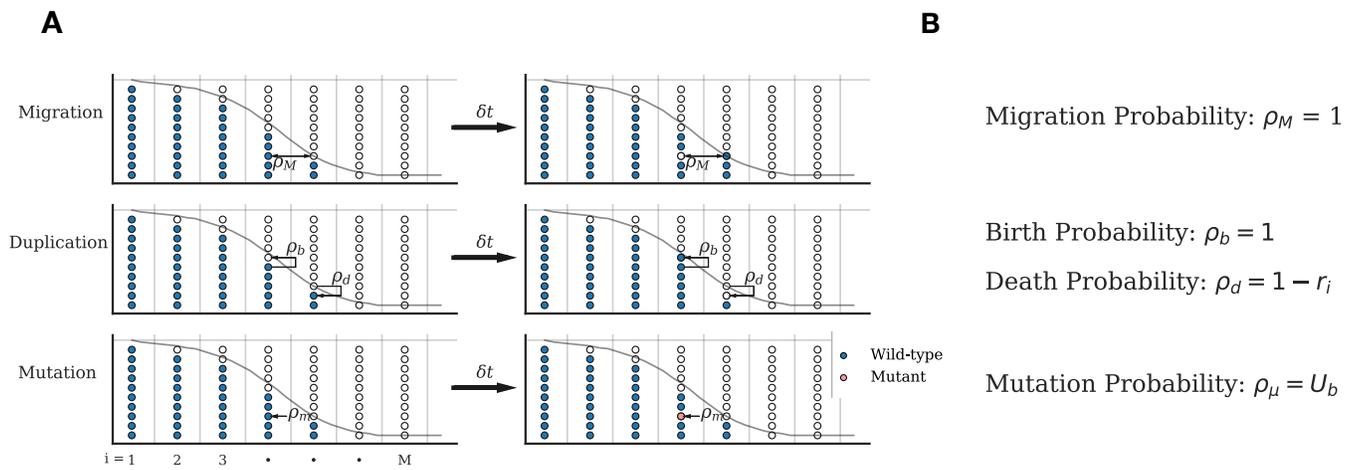


Figure 1. Simulation algorithm: one-dimensional stepping stone model with selection and mutation. A) The simulation field is comprised up to M demes distributed along a line, shown here along the horizontal axis. Each deme can each contain up to K particles, shown here as distributed along the y axis within each deme. Each particle is either a vacancy (white) or individual with m mutations, with, in this diagram, $m = 0$ (blue) or 1 (pink). The simulation is comprised of a Migration, Duplication, and Mutation step, illustrated from top to bottom respectively, which each occur once per time step δt . Each of the 10 possible events is diagrammed with an arrow, illustrating the change in the system state per time step $\Delta \vec{x}$, and its associated probability p_E . The simulation is equivalent to a sFKPP wave in the continuous limit, arbitrarily re-scaled, and superimposed on the simulation field as a light gray line. B) Each possible event and the corresponding probabilities with which they occur during a time step for migration, duplication, and mutation.

Figure 2

Beneficial mutations in a well-mixed population

Beneficial mutant fixation and establishment times :

For a large asexual population of fixed size N , with a large, finite number of beneficial mutations conferring fitness advantage s , acquired at rate U_b , the mutant lineage population size, n , grows as $\langle n | \text{not extinct} \rangle \approx \frac{1}{s} e^{st}$ at long times, t . Thus we have the approximate average time for a mutant to establish, $\langle \tau_{est} \rangle$, and average time for a mutant to fix $\langle \tau_{fix} \rangle$:

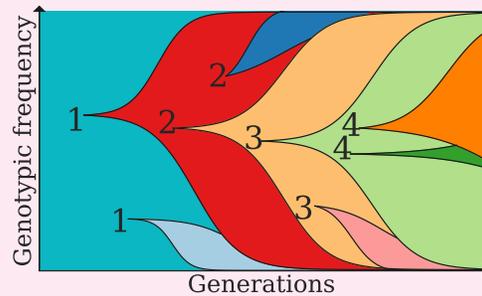
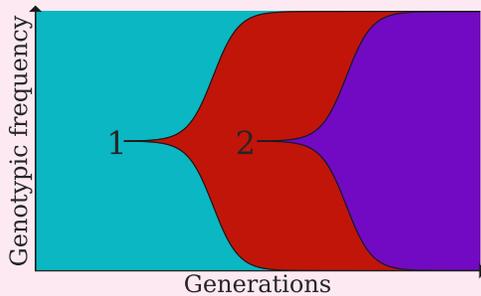
$$\langle \tau_{fix} \rangle \approx \frac{1}{s} \log(Ns).$$

$$\langle \tau_{est} \rangle \approx 1/NU_b s.$$

Dynamical regimes of beneficial mutations:

When $\langle \tau_{est} \rangle \gg \langle \tau_{fix} \rangle$, or $NU_b \log(Ns) \ll 1$, evolution is mutation limited, this is termed strong selection weak mutation (SSWM).

When $\langle \tau_{fix} \rangle \gtrsim \langle \tau_{est} \rangle$, $NU_b \log(Ns) \gtrsim 1$, and clones are expected to interfere as they fix, this is termed strong selection strong mutation (SSSM).⁵



With even higher mutation frequency, the population is likely to have many coexisting mutations, with successful selection of any given mutant limiting the rate of adaptation. This regime can be characterized as weak selection strong mutation (WSSM).¹⁰

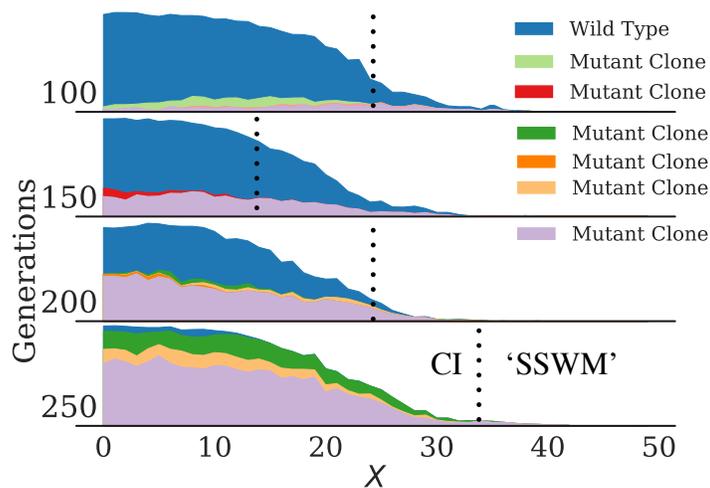


Figure 3. Mutant clones interfere with each other in various stages of diffusion, moving as an establishing traveling wave, and successfully surfing on the wave front. The wave profile in the co-moving frame of a population that can acquire beneficial mutations during range expansion is shown at a particular time point on each row, with increasing time from top to bottom. The y - axis represents the total population density while the x- axis represents the position x in the co-moving frame. Wild-type individuals are indicated in blue, and each other color a distinct mutant clonal population, each with the same selective advantage relative to its genetic background of origin. Mutant clonal populations are variously observed as traveling waves beginning to surf at the front of the population (violet), established mutants that have failed to surf at the wave front (red, generation 150), and spreading diffusively along the x-axis without apparently establishing as a traveling wave (salmon, generation 100). Dotted lines are placed at the position at which the population transitions from being polyclonal (labeled 'CI' for clonal interference) to monoclonal (labeled 'SSWM' for strong selection weak mutation) within local spatial scales. Simulations were performed with mutation rate $U_b = 10^{-3}$, selective advantage $s = .1$, and diffusion constant $K = 500$.

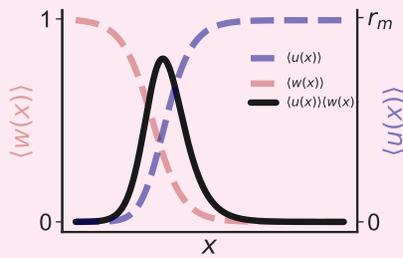
Figure 4

Mutant surfing on the wave of an expanding population

Beneficial mutant surfing probability: Mutant surfing probability with respect to position, $u(x)$ is well approximated as a stochastic branching process with a phenomenological correction (u^2 term):

$$0 = \partial_x^2 u + ur_m(1 - \langle b_w \rangle) - v_w \partial_x u - u^2,$$

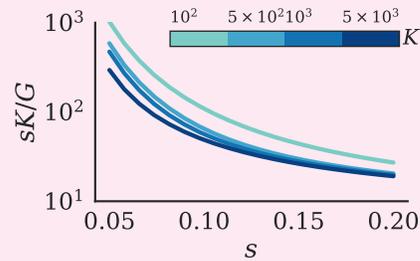
The balance between increased surfing probability at the wave tip, and larger mutational supply toward the wave bulk can result in the surfing mutant supply being maximized somewhere in between (See **SI Section 1**)



Beneficial mutant substitution rate: The supply of successful surfing mutants can approximately be quantified as:

$$G = K \int \langle b_w(x)u(x) \rangle dx,$$

and the substitution rate for the population front is given by GU_b . In a well-mixed population the analogous expression to G is represented as Ks , for total population size K and selective advantage s . Compared to a well-mixed population, where the substitution rate increases weakly with s , especially at low values, demonstrates that selection is rendered inefficient by range expansion.²²



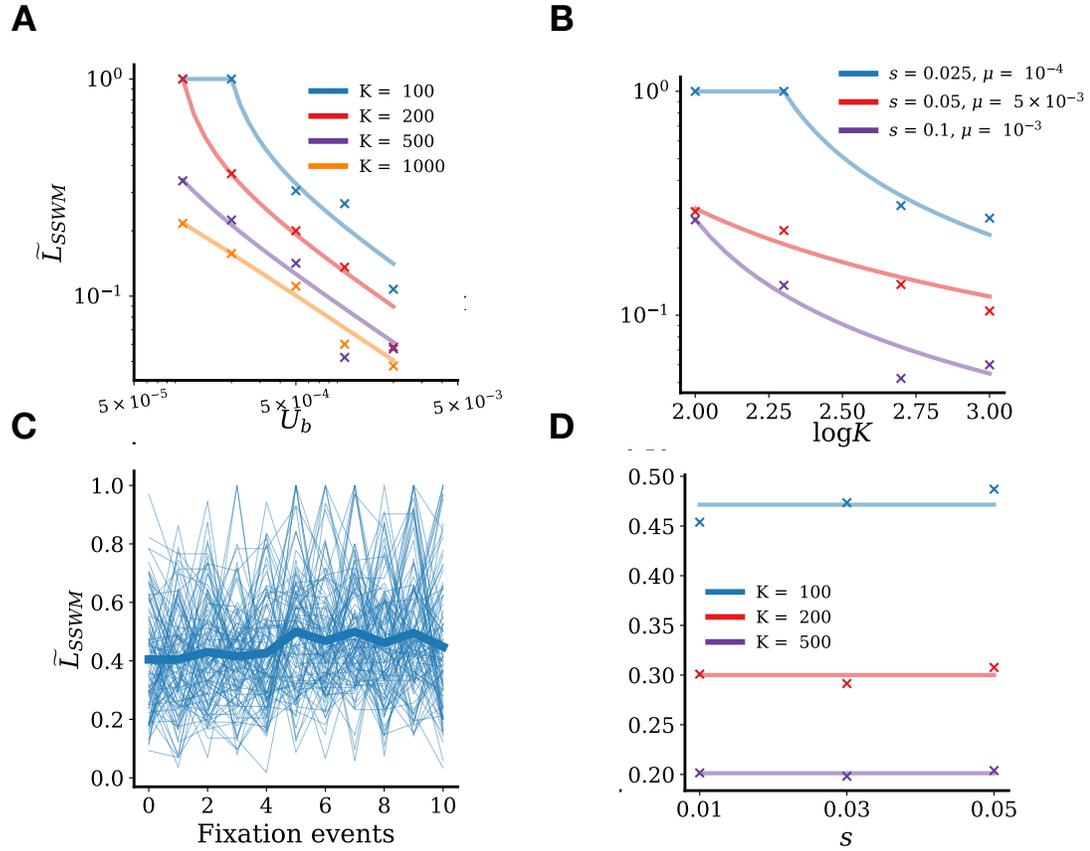


Figure 5. Proportion of wave front in ‘local’ successive mutation regime goes as $U_b^{-1/2}$ and $\log^{-1} K$, and saturates at 1 with sufficiently weak mutational supply and strength. Over an ensemble of simulations (~ 40 for panels A, B; ~ 150 for panels C, D) \tilde{L}_{SSWM} is measured as the distance in demes from the wave tip to the position at which no more than two clonal populations exist, divided by the total number of demes occupied by the wave front in the co-moving frame. **A)** Over a range of U_b and K with $s = 0.1$, the median \tilde{L}_{SSWM} from simulation results and standard deviation (shown with X’s and error bars respectively) are plotted showing approximately approximate agreement with an inverse square root fit (colored lines) well, as predicted by Eq. (3). $\tilde{L}_{SSWM} = 1$ corresponds to successive mutant fixation throughout the wave front. **B)** \tilde{L}_{SSWM} is again measured from simulations and plotted against $\log K$ at three particular values of U_b and s , showing approximate agreement with an inverse fit, as predicted by Eq. (3). **C)** \tilde{L}_{SSWM} shown for each of 10 extinction events over 100 instantiations for $s = 0.01$, $K = 100$, $\mu = 10^{-3}$ show strong fluctuations, but remain roughly constant on average (thick blue line). **D)** Median \tilde{L}_{SSWM} over 10 extinction events at $K = 100, 200, 500$ and $\mu = 10^{-3}$, \tilde{L}_{SSWM} does not vary over $s = 0.01, 0.03, 0.05$. Simulation data is fit with line of slope 0 as predicted by Eq. (3) for relatively weak selection assuming a noisy Fisher wave velocity proportional to s .