

1 **Sensitivity of Simulated Fire-Generated Circulations to Fuel Characteristics During**
2 **Large Wildfires**

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8 **Key Points:**

- 9 • Coupled fire-atmosphere models struggle to simulate critical fire-generated winds and
10 plume rise during large wildland fires
- 11 • Deficient fire-generated winds are linked to inadequate fuel loads and burnout timescale
12 in the model
- 13 • Adjustment of the fuel characteristics results in more realistic simulated plumes and fire-
14 generated winds

15 **Abstract**

16 Coupled fire-atmosphere models often struggle to simulate important fire processes like fire
17 generated flows, deep flaming fronts, extreme updrafts, and stratospheric smoke injection during
18 large wildfires. This study uses the coupled fire-atmosphere model, WRF-Fire to examine the
19 sensitivities of some of these phenomena to the modeled surface fuel load. Specifically, the 2020
20 Bear Fire and 2021 Caldor Fire in California's Sierra Nevada are simulated using three fuel
21 loading scenarios (1x, 4x, and 8x LANDFIRE derived surface fuel), while controlling the fire
22 rate of spread, to isolate the fuel loading needed to produce fire-generated flows and plume rise
23 comparable to NEXRAD radar observations of these events. Increasing fuel loads and
24 corresponding fire residence time in WRF-Fire leads to deep plumes in excess of 10 km, strong
25 vertical velocities of 40-45 m s⁻¹, and combustion fronts several kilometers in width (in the along
26 wind direction). These results indicate that LANDFIRE-based surface fuel loads in WRF-Fire
27 likely under-represent fuel loading, having significant implications for simulating landscape-
28 scale wildfire processes, associated impacts on spread, and fire-atmosphere feedbacks.

29 **Plain Language Summary**

30 Coupled fire-atmosphere models poorly depict large-scale fire processes, such as fire generated
31 winds and deep smoke plumes. In this study, the 2020 Bear Fire and 2021 Caldor Fire in
32 California are simulated under various fuel scenarios. The simulations show that fuel
33 characteristics used in the fire-atmosphere model under-represent observed conditions and thus
34 produce inadequate fire-generated winds and plume characteristics. When the modeled fuels are
35 augmented to match observed fuel load and burnout time, simulated fire-atmosphere feedbacks
36 better resemble fire generated winds and deep convective plumes seen in radar observations. The
37 results of these simulations will help inform future improvements to coupled fire-atmosphere
38 models to better simulate large wildland fires.

39 **1 Introduction**

40 Fire size and intensity has been increasing in the western United States in recent decades
41 (Westerling et al., 2006, 2016; Williams, 2013; Holden et al., 2018; Parks and Abatzoglou 2020).
42 These larger, more intense fires are often characterized by 1000s of acres of simultaneous
43 combustion (i.e., mass fire, Finney and McAllister, 2011), deep convective columns,
44 pyrocumulonimbus (pyroCu/Cb) capable of injecting smoke into the stratosphere (Fromm et al.,
45 2006, 2010; Rodriguez et al., 2020; Peterson et al., 2021), and extreme fire-generated winds
46 including fire-generated tornadic vortices (FGTVs) (Fromm et al., 2006, 2010; Cunningham and
47 Reeder, 2009; Lareau et al., 2018, 2022a). Given the complex threats posed by landscape fires on
48 the social, ecological, and built environments and the expected increase in fire frequency and
49 intensity in a warming climate (Abatzoglou and Williams, 2016; Dowdy et al., 2019), accurate
50 simulation of fires and their impacts are necessary for improved societal resilience, pre-fire
51 planning, and active-fire situational awareness.

52 Uncertainties in combustion processes, fire spread, fuel representation, and atmospheric
53 feedbacks make simulations of large real-world fires challenging (Peace et al., 2020; Shamsaei et
54 al., 2023a). For example, current fire spread models used in fire-fighting operations such as
55 FARSITE (Finney, 1998) and ELMFire (Lautenberger, 2013, 2017) rely on the semi-empirical
56 Rothermel (1972) rate of spread model but are not coupled to the atmosphere. Thus, these
57 models cannot simulate turbulent flow fields or the feedbacks between fire and atmospheric

58 processes, such as fire-induced updrafts and associated inflow winds that alter the rate and
59 direction of fire spread. This is problematic in that these fire-induced winds can become the
60 dominant driver of large wildland fires (Coen et al., 2018). Rather, to simulate these feedbacks,
61 coupled fire-atmosphere modes are required, wherein an atmospheric model resolves the wind
62 field that drives fire spread. In turn, the fire's heat and moisture fluxes are released back into the
63 atmosphere, thereby perturbing the wind field, which are then passed back to the fire spread code
64 to represent coupling between the fire and atmosphere (Clark et al., 2004).

65 WRF-Fire, and the similar WRF-SFIRE, are examples of coupled fire-atmosphere
66 simulation platforms that link the Weather Research and Forecasting (WRF) atmospheric model
67 (Skamarock and Klemp, 2008; Skamarock et al., 2019) with the Rothermel rate of spread model
68 (Rothermel, 1972) to simulate fire spread along with atmospheric responses and feedbacks on the
69 fire (Clark et al., 2004; Mandel et al., 2011; Coen et al., 2013). While these coupled models show
70 promise in simulating perimeter changes in landscape scale fires (Kochanski et al., 2013;
71 Jimenez et al., 2018; DeCastro et al., 2022; Shamsaei et al. 2023a,b; Juliano et al., 2023),
72 thorough validation of the atmospheric response and feedbacks to the fire are lacking outside of
73 small-scale grass fire experiments (e.g., FIREFLUX, FIREFLUX II). For example, most studies
74 validate perimeter changes without providing validation of plume responses or flow
75 modifications, and thus it is possible that these models sometimes produce the right answer (e.g.,
76 a correct perimeter) for the wrong reason. This can be problematic in simulations of landscape-
77 scale fires, where atmospheric responses and feedbacks become more important in dictating fire
78 spread and its impacts.

79 To investigate these model deficiencies, we conduct a sequence of sensitivity tests
80 designed to isolate the role of fuel loading and consumption on simulated fire-generated
81 circulations, including the plume rise and inflow winds. We first motivate this work with an
82 example of the model deficiencies (Section 2), before moving on to our methods, results, and
83 implications (Sections 3-5).

84 **2 Problem Statement**

85 Shamsaei et al. (2023a, b) showed in two recent simulations of California's deadliest fire,
86 the Camp Fire in 2018, that burn area was relatively well depicted by WRF-Fire, however fire
87 and atmospheric feedbacks were deficient in terms of heat release, fire-generated flows, and
88 plume depth. With this in mind, a preliminary simulation of another landscape-scale fire (2020
89 Bear Fire in California's northern Sierra Nevada; Fig. 1) was conducted using a similar WRF-
90 Fire configuration to that of Shamsaei et al. (2023a, b) based on the operational Colorado Fire
91 Prediction System (CO-FPS; Jimenez et al., 2018). The details of this simulation, including the
92 namelist are contained in supplements S1 and S2. In this preliminary simulation, although WRF-
93 Fire depicts a similar fire perimeter (Fig. 1c) to the observed perimeter (Fig. 1a), comparison
94 with radar observed winds reveals that the simulation lacks both the pronounced region of fire-
95 generated flow reversal and inflow wind opposing the background flow to the west of the head
96 fire (note red shading in Fig. 1a, b) and the deep plume structure that lofts smoke and ash into the
97 mid-troposphere (Fig. 1a-d). Thus, while this operational WRF-Fire configuration produces
98 adequate fire spread, it does not produce the fire-generated winds and plume dynamics that are
99 critical drivers of the fire behavior. The preliminary simulations are further deficient in that they
100 inadequately represent the breadth of the combustion, measured in terms of the satellite observed
101 infrared footprint of the fire (Fig. 2). For example, the broad region of high heat release rates in

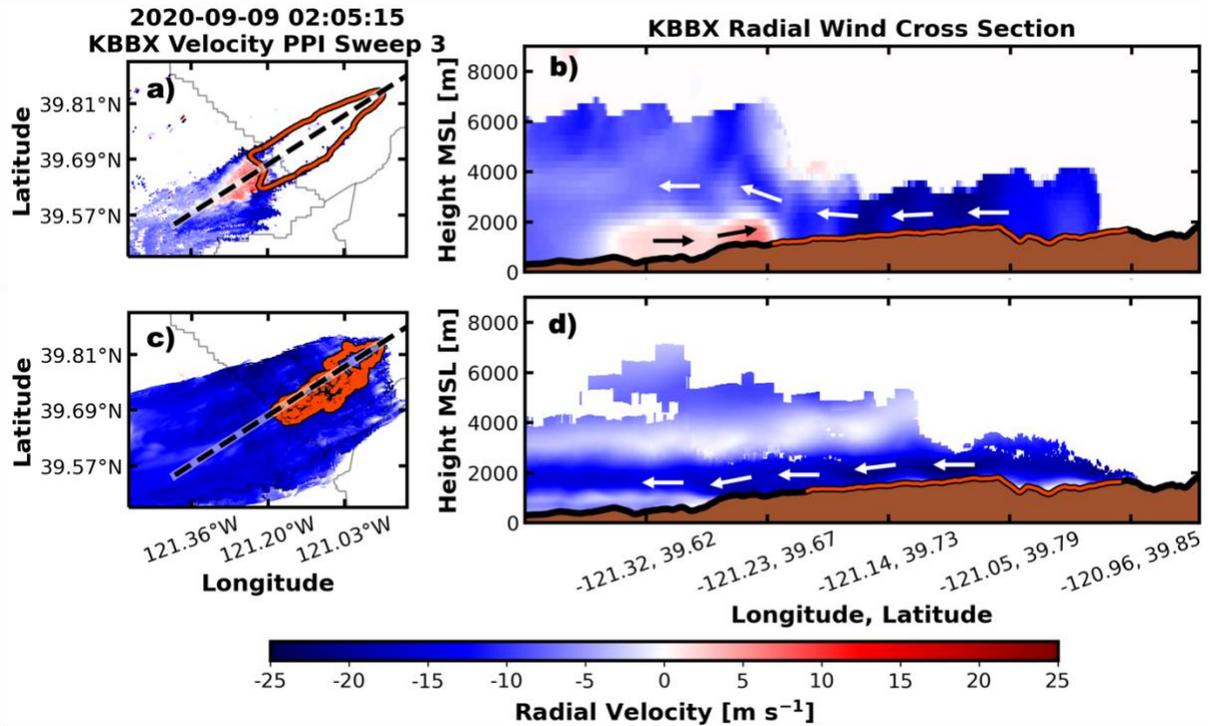


Figure 1. Comparison of observed and simulated fire properties. (a) Beale Air Force Base (KBBX) NEXRAD radial velocity PPI (shaded) and radar-estimated fire perimeter (red contour), (b) radial wind and radar-estimated fire perimeter cross section (red line) along black dashed line in (a), (c) WRF-Fire simulated in-plume radial velocity PPI and fire perimeter, and (d) simulated in-plume radial wind and fire perimeter cross section (red line) along black dashed line in (c) during a period of pronounced fire atmosphere coupling on the Bear Fire around 0200 UTC September 9, 2020. In-plane directional flow vectors annotated in b and d.

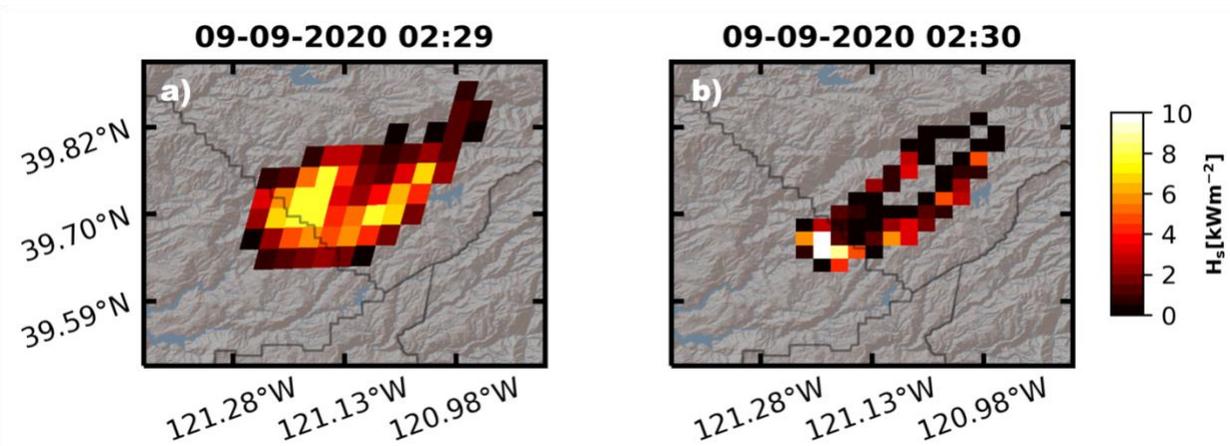


Figure 2. Comparison of (a) GOES-17 Fire-Radiative Power (FRP) converted to sensible heat flux (FRPx10; from Val Martin et al., 2012) with (b) preliminary WRF-Fire sensible heat flux down-sampled to a 2x2 km grid in the Bear Fire.

102 observations (Fig. 2a) is much larger than that of the preliminary WRF-Fire simulation, even
103 when we resample the WRF output to match the satellite's spatial resolutions (Fig. 2b). While
104 previous studies have noted deficiencies with fuel representations in fire models and their impact
105 on fire perimeter changes (DeCastro et al., 2022; Stephens et al., 2022), the goal of this work is
106 to isolate how fuel characteristics affect fire-generated winds and plume development using
107 observations of these processes as a validation metric.

108 We hypothesize that existing coupled fire-atmosphere models are deficient in producing
109 the observed fire-atmosphere coupling during landscape-scale fires because they have (1)
110 insufficient fuel loads and consumption and (2) inadequate representations of how fires move
111 through the landscape due to inherent limitations of the fire spread model (e.g., lack of mass fire
112 and spotting).

113 To test these hypotheses, we use WRF-Fire to simulate two landscape-scale wildfires
114 (details below) during periods of strong fire-atmosphere coupling and conduct a series of fuel
115 load sensitivity tests while prescribing the fire's rate of spread. This is accomplished by turning
116 off the model's fire spread code and using a "time-of-arrival" grid (similar to the process
117 described in Farguell et al., 2021) based on radar observations (methodology described in Lareau
118 et al., 2022b). We also modify the fire residence time (i.e., the time required for the fuel to burn
119 down to ~37% of its initial load) to generate broader combusting regions more consistent with
120 the observations. These permutations allow us to determine the threshold fuel loading for WRF-
121 Fire to generate reasonable fire-atmosphere coupling comparable to observations.

122 **3 Data and Methods**

123 3.1 The Fires

124 The Bear and Caldor Fires in California's Sierra Nevada (see Table 1, Fig. 3) provide
125 ideal test cases to examine WRF-Fire's ability to simulate fire-atmosphere coupling during high-
126 intensity landscape scale fires. Both fires developed deep convective plumes and strong fire-
127 induced winds in similar terrain and fuels, but under strong (i.e., 30 m/s) and light (i.e., 10 m/s)
128 wind scenarios, respectively. Details of the fires are as follows:

129 *The Bear Fire* was ignited by lightning on 17 August 2020 in Plumas National Forest in the
130 northern Sierra Nevada. On 8 September the fire was affected by a strong downslope wind event
131 with wind gusts up to 30 m s^{-1} which drove extreme rates of spread, deep pyroCb-topped plumes,
132 and FGTVs (Lareau et al., 2022a, b). The fire ultimately burned approximately 318,935 acres
133 (129,068 ha), destroyed 2,455 buildings, and resulted in 16 fatalities.

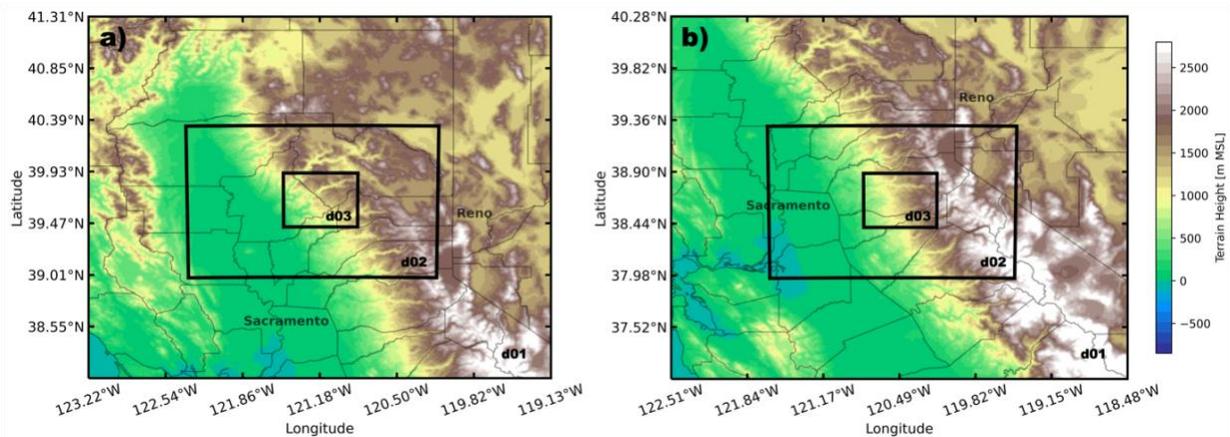
134 *The Caldor Fire* ignited on 14 August 2021 in Eldorado National Forest in the central Sierra
135 Nevada. On 17 August the fire experienced rapid fire spread and deep pyroCb plumes while
136 advancing eastward across the Sierra Nevada west slope. The fire ultimately burned
137 approximately 221,835 acres (89,773 ha) and destroyed 1,003 buildings (USDA Forest Service,
138 2021 and CalFire Incident Archive, 2021).

139 Both fires produced pronounced flow reversals downwind of the head fire (e.g., Fig. 1a,b)
140 and plume echo tops episodically exceeding 10km above mean sea level (MSL) in NEXRAD
141 radar imagery. These strong fire-generated circulations make these cases well suited for model
142 sensitivity tests.

143

Table 1. Two fire cases identified for sensitivity analysis.

Fire Name	Date of Ignition	Analysis Date(s)	Location	Acres (ha) Burned on Analysis Date(s)	Total Acreage (ha)	Dominant SB40 Fuel Type
Bear Fire	17 August 2020	1900 UTC 8 Sep – 0400 UTC 9 Sep 2020	Plumas National Forest	193,759 (78,411)	318,935 (129,068)	TU5 (69%)
Caldor Fire	14 August 2021	1500 UTC 17 Aug – 0000 UTC 18 Aug 2021	Eldorado National Forest	20,939 (8,474)	221,835 (89,773)	TU5 (73%)

**Figure 3.** Outer (d01), middle (d02), and inner (d03) domain configuration for the (a) Bear Fire and (b) Caldor Fire with WRF terrain (shaded).

144 3.2 WRF-Fire

145 Our simulations are conducted with WRF-Fire (Mandel et al., 2011; Coen et al., 2013).
 146 The configuration closely follows that of Jimenez et al. (2018) and Shamsaei et al. (2023a, b).
 147 The atmospheric model uses one-way nesting across three domains containing 41 vertical levels
 148 up to 50 hPa. The outermost domain has a horizontal grid spacing of 1 km with inner nests of
 149 333 m and 111 m on the atmospheric mesh, with the innermost domain resolving the fire on a
 150 further refined mesh with spacing of ~28 m, centered over the fire areas (Fig. 3). The terrain in
 151 the inner fire domain is derived from the 30-meter resolution NASA SRTM topographic dataset
 152 (van Zyl, 2001; Farr and Kobrick, 2000). The simulations use the 2011 National Land Coverage
 153 Database (NLCD2011) (Homer et al., 2015) with Noah land-surface (Chen and Dudhia, 2001)
 154 and Revised Monin-Obukhov surface layer (Jimenez et al., 2012) parameterization schemes. The
 155 Dudhia (1989) shortwave radiation, Rapid Radiative Transfer Model (RRTMG) longwave
 156 radiation (Iacono et al., 2008), and Hong and Lim (2006) WRF single-moment 6-class (WSM6)
 157 microphysics schemes are also used. The Mellor-Yamada-Nakanishi-Niino (MYNN; Nakanishi
 158 and Niino, 2006) PBL scheme is used on the two outer domains, with the innermost domain
 159 resolving turbulence using the subgrid-scale model of Lilly (1966a, b) and Deardorff (1980).
 160 Initial and boundary conditions are set using High Resolution Rapid Refresh (HRRR) analysis
 161 data (3 km spatial resolution) that update every hour through completion of the simulation.

162 3.3 Fire Spread and Perimeters

163 In its operational configuration, WRF-Fire uses the Rothermel ROS model (Rothermel,
164 1972) to propagate fire across a landscape. The Rothermel model uses a semi-empirical
165 relationship amongst the wind speed at flame height and terrain slope to produce fire spread. The
166 fire and atmosphere are coupled by fire-generated heat and moisture fluxes which then perturb
167 the lower atmospheric model layers via an exponential decay function with height (described in
168 Clark et al., 1996a,b and Coen et al., 2013). The Rothermel model has known limitations in that
169 it assumes a narrow, linear fire line and neglects key landscape-scale fire components such as
170 spotting and mass fire (Andrews, 2018). It is hypothesized that these limitations play a role in the
171 poorly developed plume structure seen in the control simulations (Shamsaei et al., 2023b). To
172 bypass this deficiency, the Rothermel model is replaced by continuously updated fire perimeters
173 derived from NEXRAD radar data. This technique is based on locating local maxima in the radar
174 reflectivity and associated active combustion, and then aggregating these points into an evolving
175 fire polygon (Lareau et al., 2022b). The process has been validated against infrared observations
176 for several fires, including the Bear and Caldor fires. These radar perimeters are converted to a
177 time-of-arrival grid that is passed into WRF-Fire, which controls the time at which a given cell in
178 the fire mesh ignites. This process is similar to the satellite-based time-of-arrival approach used
179 by Farguell et al. (2021). This “forced fire” approach maintains consistent fire rate and direction
180 of spread across all of the sensitivity tests, allowing us to isolate the impact of fuel load on the
181 heat fluxes and plume development without having to interpret changes in fire ROS, which itself
182 is a function of fuel load in the Rothermel model.

183 3.4 Fuel Depiction and Fire Residence Time

184 The WRF-Fire simulations use the Scott and Burgan 40 (SB40) fuel categories (Scott and
185 Burgan, 2005) derived from the LANDFIRE 2016 (Rollins, 2009) dataset to represent fuel type
186 and load in the model domain. The LANDFIRE dataset is widely used among the wildfire
187 modeling community because of its high resolution (30 x 30 m) coverage of fuel type, fuel load,
188 fuel bed depth, and surface area to volume ratio across the contiguous United States (DeCastro et
189 al., 2022). The dominant SB40 fuel category in the central Sierra Nevada is Timber-Understory 5
190 (TU5), comprising 69% and 73% of the simulated burn area in the Bear and Caldor Fires,
191 respectively (Fig. 4). The TU5 fuel type is a high-load conifer litter and shrub understory with a
192 combined 1-, 10-, and 100-hour fuel load of 2.47 kg m^{-2} (11 t ac^{-1}) and moderate flame length
193 and spread rate (Scott and Burgan, 2005).

194 This default fuel load of $\sim 2.5 \text{ kg m}^{-2}$, however, is a drastic underestimate of the fuels
195 available-for and consumed-in large fires, especially fuels consumed after the passage of the
196 initial fire front. For example, using pre- and post-fire fuel measurements in in the central Sierra
197 Nevada, Cansler et al. (2019) found an average fuel consumption of 15.1 kg m^{-2} (151 Mg ha^{-1})
198 during the 2013 Rim Fire in Yosemite National Park. Similarly, McCarley et al. (2020) showed
199 airborne laser scanning estimated fuel consumption in large wildfires exceeding 20 kg m^{-2} (200
200 Mg ha^{-1}) over large areal expanses. These observations suggest that, even in the best-case
201 simulations with WRF-Fire and SB40 fuels, fires may not yield total released heat comparable to
202 those in real fires, and thus cannot simulate the strong fire-generated circulations (e.g., updrafts

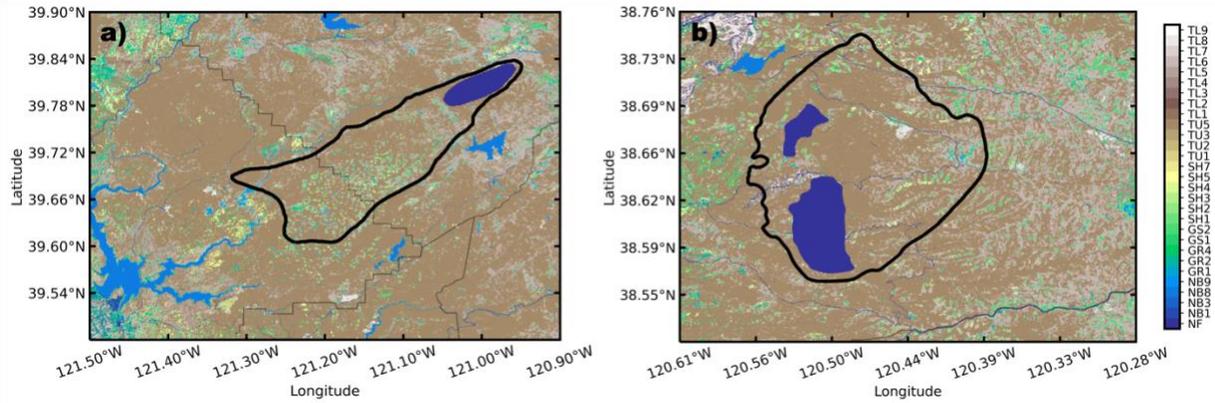


Figure 4. SB40 fuel category map for the (a) Bear Fire and (b) Caldor Fire. Dark blue no fuel (NF) region shows estimated initial perimeter used to initiate WRF-Fire simulation with final fire perimeter shown in black.

Table 2. Summary of case studies and variables.

Case Name	TU5 Fuel Load (kg m^{-2})	w	Fuel Moisture (%)	Fire Spread Method
BearControl	2.47	900	5	Rothermel
BearFuelx1	2.47	4080	5	NEXRAD
BearFuelx4	9.86	4080	5	NEXRAD
BearFuelx8	19.73	4080	5	NEXRAD
CaldorControl	2.47	900	5	Rothermel
CaldorFuelx1	2.47	3825	5	NEXRAD
CaldorFuelx4	9.86	3825	5	NEXRAD
CaldorFuelx8	19.73	3825	5	NEXRAD

203 and inflows) that feedback on fire processes. This deficiency is apparent in Figure 1 when
 204 comparing the preliminary simulation (Fig. 1c,d) to observed flow perturbations (Fig. 1a,b).

205 To examine the sensitivity of fire-generated circulations to fuel loads we devise three
 206 sensitivity tests all using the same observationally-based prescribed fire spread. Due to the
 207 dominance of TU5 fuels in the study area, and to eliminate further uncertainties in fuel types,
 208 only the TU5 fuel loads are adjusted in this study. We first use a control case with the default
 209 TU5 load of 2.47 kg m^{-2} (Fuelx1) and two augmented fuel loads of 9.86 kg m^{-2} (Fuelx4) and
 210 19.73 kg m^{-2} (Fuelx8) (Table 2). Note that the Fuelx8 cases are similar to observed loads and
 211 consumption of $15\text{-}20 \text{ kg m}^{-2}$ described above, and thus a priori we expect these simulations to
 212 best match observations.

213 In addition to the fuel load, in WRF-Fire each SB40 fuel category has a weighting
 214 parameter controlling the fire's residence time. This weighting factor is defined as

$$w = 0.8514 \times T_f, \quad (1)$$

215 where T_f is the time for the fuel to burn down to $e^{-1} \approx 0.3689$ of the initial fuel load (Mandel et
216 al., 2011). The default values for w are derived from approximations of mass-loss curves from
217 the Albin and Reinhardt (1995) BURNUP algorithm (Clark et al., 2004); however, Mandel et al.
218 (2011) noted there is significant uncertainty in the default values used in WRF-Fire. Due to the
219 relationship between fuel load and burnout time, w must proportionally change with fuel loads to
220 avoid unphysically large heat release rates (i.e. burning the fuel too quickly) under increased fuel
221 scenarios. The value of w also impacts the breadth of the combusting region: for a given fuel
222 load a larger w produces a broader combusting region when we force the perimeters to
223 observations, and thus constrain the rate of spread. We note that, when using the Rothermel
224 spread model after increasing the fuel load and the weighting parameter, the fire spread
225 unrealistically decreases, thus highlighting the need for forced fire perimeter approach in the
226 sensitivity analyses.

227 Since the values of w used in WRF-Fire are uncertain, we use GOES-17 Fire Radiative
228 Power (FRP) data to estimate representative values for T_f (burnout time) and thus w using the
229 relationship shown by Eq. 1 from Mandel et al. (2011). The FRP algorithm uses the $3.9 \mu\text{m}$ and
230 $11.2 \mu\text{m}$ brightness temperatures along with a number of temporal and spatial checks to
231 characterize fire temperature, size, and FRP and is thus a useful metric in identifying regions of
232 active fire (Schmidt et al., 2012) and how long fire resides within a given pixel ($2 \text{ km} \times 2 \text{ km}$).
233 We estimate this “pixel” residence time by evaluating each GOES-17 pixel during the simulation
234 timeframe (Table 1) to determine when the pixel reached maximum FRP (Fig. 5). Then, we
235 evaluated how long each pixel took to cool to e^{-1} of its normalized FRP maxima and defined the
236 value as T_f (interquartile ranges for all pixels depicted with red shading in Fig. 5). Individual
237 pixel T_f values are shown in the insets of Fig. 5 for both the Bear (Fig. 5a) and Caldor (Fig. 5b)
238 Fires. We note that the pixel residence time is not purely the physical burndown time of the fuels
239 since it includes information about both the rate of spread through the pixel and the consumption
240 of fuel. Nonetheless, it is a useful approach for grounding our simulations in an observational
241 framework. T_f values for all fire pixels during the timeframe are then averaged (maroon line in
242 Fig. 5) to produce a representative T_f and w value for each fire. The resulting analysis suggests
243 values of 4080 and 3825 seconds are appropriate for the Bear (Fig. 5a) and Caldor (Fig. 5b) fires,
244 respectively, which is about four times larger than the default value in WRF-Fire (900 s). These
245 w values are not intended to be physical or universal for improving freely evolving WRF-Fire
246 simulations, but rather an approach at producing realistic breadth of the combusting zone for the
247 given cases using the forced fire perimeters, described below.

248 3.5 Radar Observations of Plume Processes

249 In addition to providing estimated fire perimeters, NEXRAD radar data are also used to
250 compare simulated fire-generated circulations with observed plume injection heights and fire-
251 generated flows. Specifically, we use radar reflectivity and radial velocity cross sections
252 extracted from a cartesian gridded version of the NEXRAD observations (see Lareau et al.,
253 2022a) to document plume structure, plume injection height, and radial wind components due to
254 the ambient and fire-generated winds. These data provide a useful validation approach for
255 landscape scale fires where in-situ measurements are otherwise unavailable (e.g., Jones et al.
256 2022).

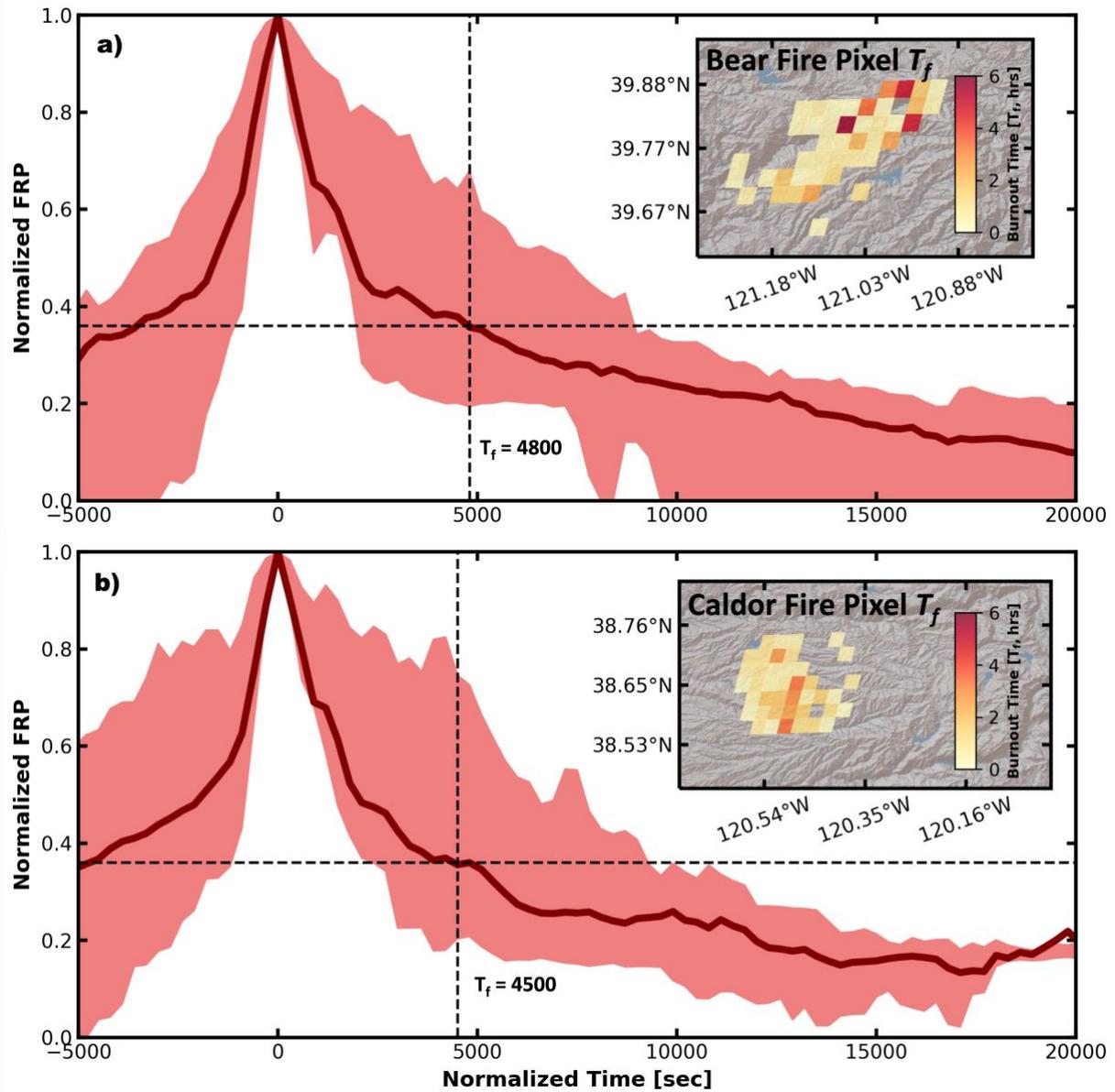


Figure 5. Fire-averaged FRP timeseries (maroon line) and interquartile ranges (light red shading) of individual pixel FRP normalized by maximum detected FRP in the scene. Pixel T_f (inset, shaded) for the (a) Bear Fire and (b) Caldor Fire. Horizontal black dashed line indicates e^{-1} of normalized FRP, and vertical dashed black line indicates T_f where average FRP crosses e^{-1} .

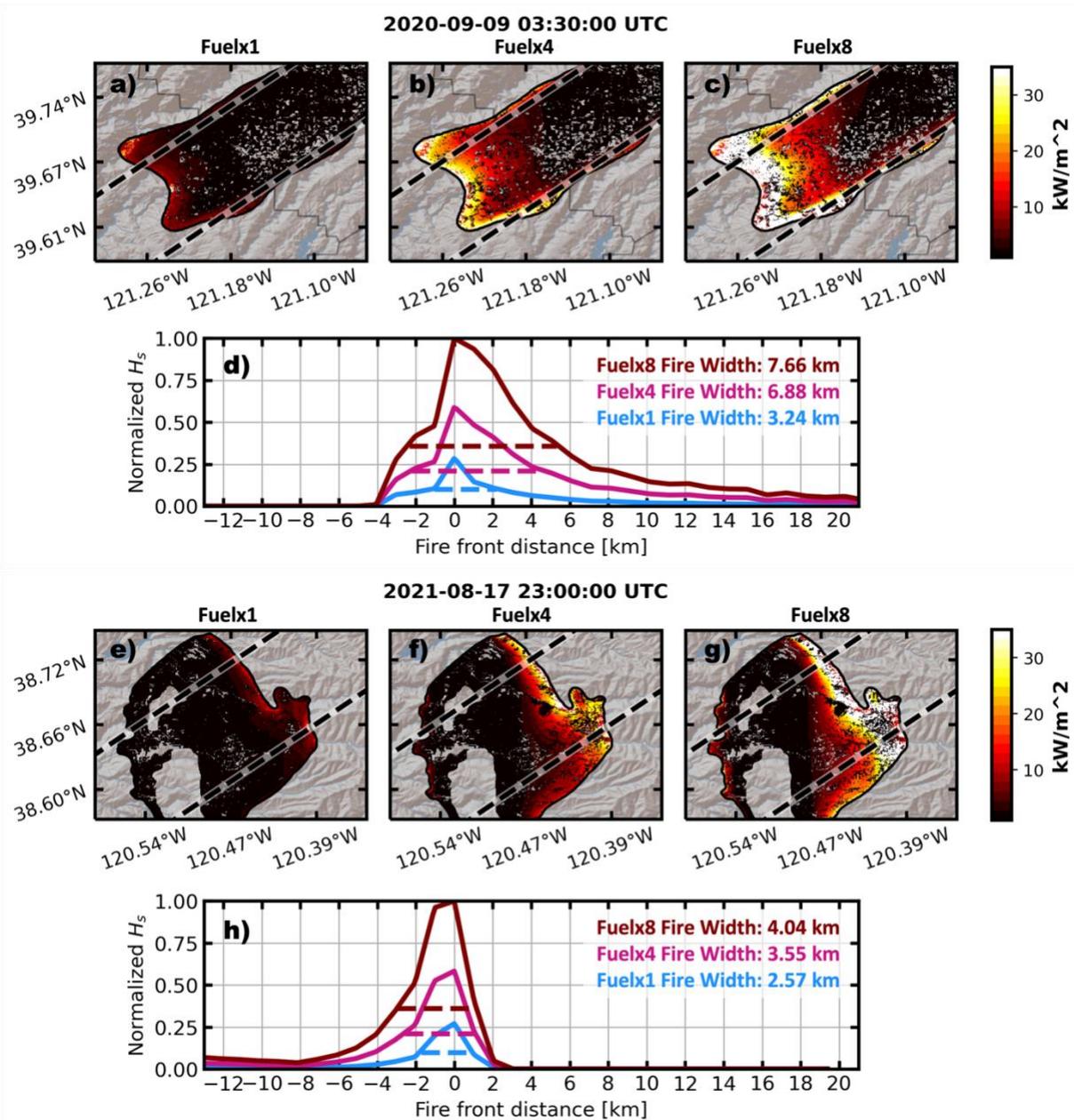


Figure 6. Fire generated heat flux for the (a-c) Bear Fire and (e-g) Caldor Fire. Normalized cross sections of fire heat flux with fire width indicated by dashed line at e^{-1} of the peak heat release for the Fuelx1 (light blue), Fuelx4 (magenta), and Fuelx8 (maroon) scenarios for the (d) Bear and (h) Caldor Fires.

259 4 Results

260 4.1 Fireline Width and Intensity

261 Our simulations show that increasing TU5 fuel loads and burnout timescale (T_f) while
262 forcing the fire spread increases the areal extent and width of intense ($>10 \text{ kW m}^{-2}$) fire-
263 generated sensible heat fluxes for both the Bear and Caldor fires (Fig. 6). To quantify these
264 changes, we use an e-folding scale (e.g., ~ 0.37) of the peak fire-generated sensible heat flux to
265 identify the width of the head fire (see dashed lines in Fig. 6d, h) in each simulation, where “fire
266 front distance” corresponds to the horizontal distance from the normalized maximum heat flux in
267 the head fire region.

268 The results show that at 330 UTC the Bear Fire Fuelx1 scenario generates a narrow (~ 3.2
269 km, Fig. 6d) head fire with a maximum sensible heat flux of $\sim 9 \text{ kW m}^{-2}$ (Fig. 6a) whereas the
270 Fuelx4 scenario has a comparatively wider (~ 6.9 km, Fig. 6d) head fire with a maximum of 34
271 kW m^{-2} (Fig. 6b). Finally, the Fuelx8 scenario produces the widest (~ 7.7 km, Fig. 6d) head fire
272 with a peak sensible heat flux of 68 kW m^{-2} (Fig. 6c). Heat fluxes of this magnitude are
273 consistent with those estimated in recent observational studies of plume rise (e.g., Lareau and
274 Clements, 2017).

275 Similar changes in the head fire width and heat fluxes are simulated for the Caldor Fire.
276 Specifically, at 2300 UTC, the Caldor Fuelx1 scenario produces a relatively narrow (~ 2.6 km,
277 Fig. 6h) head fire with a maximum heat flux of $\sim 9 \text{ kW m}^{-2}$ (Fig. 6e), whereas the Fuelx4 scenario
278 has a head fire width of ~ 3.6 km (Fig. 6h) with a peak heat flux of $\sim 36 \text{ kW m}^{-2}$ (Fig. 6f), and the
279 Fuelx8 scenario has the widest fire head (~ 4 km, Fig. 6h) and highest maximum fire heat flux
280 ($\sim 73 \text{ kW m}^{-2}$, Fig. 6g).

281 These results show that increasing the fuel load and weighting factor increases the
282 maximum fire-generated heat flux and the areal extent of the head fire, thus implying a wider
283 “flaming region” that better agrees with available infrared (e.g., GOES-17) observations (Fig. 7).
284 For example, Fig. 7 shows down-sampled versions of the WRF-Fire sensible heat fluxes to
285 mimic the GOES-17 FRP satellite footprint (2x2km). For this comparison the FRP data are
286 converted to sensible heat fluxes using the assumption that FRP is approximately one tenth the
287 sensible heat flux (Val Martin et al., 2012). We note that there is uncertainty in these
288 measurements due to sensor saturation and shading from pyroCb, likely resulting in artificially
289 low observed intensities. Nonetheless, these comparisons show that both the Bear and Caldor
290 Fire Fuelx8 (Fig. 7d,h) simulations compare favorably with the observations (Fig. 7a,e), whereas
291 the Fuelx1 (Fig. 7b,f) and Fuelx4 (Fig. 7c,g) cases insufficiently represent the breadth of intense
292 combustion. As we show in the next two sections, only the simulations with wider and higher
293 intensity combustion zones yield atmospheric response comparable to the observations.

294 4.2 Fire-Generated Horizontal Flow Perturbations

295 Commensurate with the changes in fire heat flux and head fire width, our simulations
296 show improved representation of the fire-generated horizontal flow perturbations with increasing
297 fuel loads. The horizontal component of the flow is evaluated by comparing “radial velocity”
298 observations from the NEXRAD radar with the flow component in the simulations that would be
299 observed with a hypothetical radar in the same location. This is accomplished by computing the
300 component of the simulated winds that projects onto radials originating from the radar base
301 locations (KBBX and KDAX for the Bear and Caldor Fires, respectively), and thus provide a

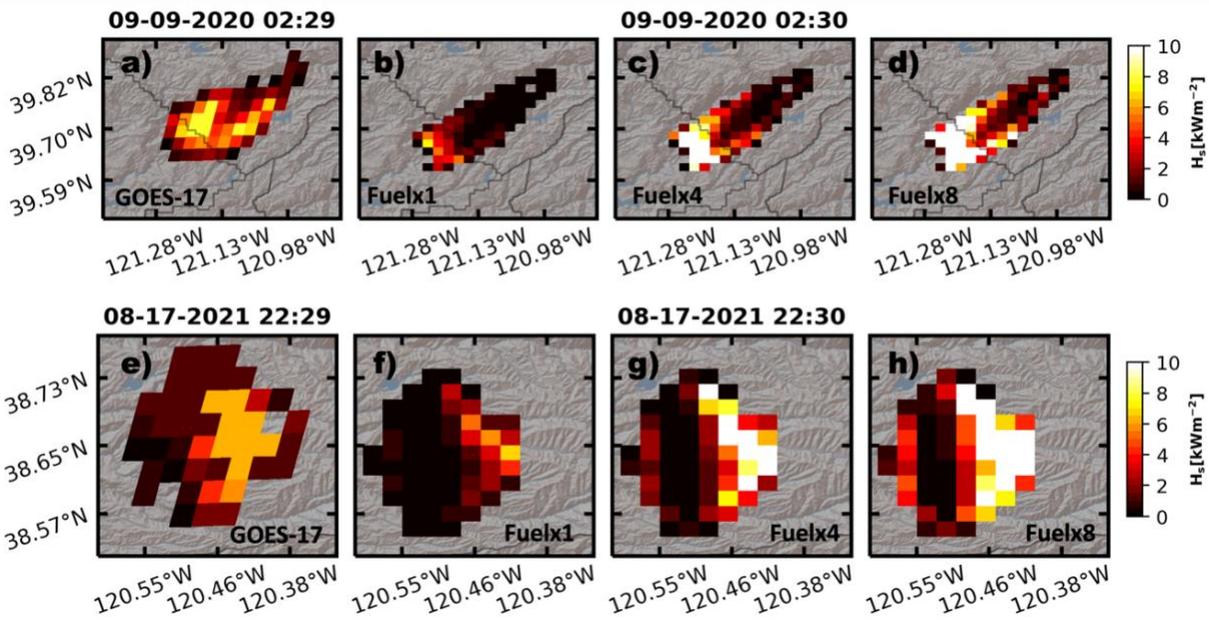


Figure 7. Observed GOES fire intensity using FRP converted to sensible heat flux for the (a) Bear Fire and (e) Caldor Fire. WRF-Fire sensible heat flux for the (b-d) Bear and (f-h) Caldor Fires. WRF-Fire sensible heat fluxes are down-sampled to a 2x2 km grid to emulate the GOES FRP data resolution.

302 direct comparison with the wind components observed by the radars. In this framework, all
 303 winds are either “inbound” (shown in blue) or “outbound” (shown in red) relative to the radar
 304 location.

305 For the Bear Fire, the observed radial velocities from the KBBX NEXRAD on 9
 306 September 2020 at 330 UTC indicate strong downslope winds of $\sim 25 \text{ m s}^{-1}$ towards the radar
 307 (blue shading and white arrows, Fig. 8a, b), with a pronounced region of flow reversal in the lee
 308 of the fire head indicated by outbound radial velocities of $5\text{-}10 \text{ m s}^{-1}$ (red shading and black
 309 arrows, Fig. 8a, b). The flow reversal is clear observational evidence for a mesoscale fire-
 310 generated wind that produces strong convergence at the fire front and feeds the vigorous fire-
 311 generated updrafts. While all three fuel scenarios depict strong, downslope winds and inbound
 312 radial velocities greater than 20 m s^{-1} (Fig. 8c-h), they differ in the magnitude and extent of the
 313 fire-generated flow reversal. The Fuelx1 scenario shows no flow reversal, with inbound radial
 314 velocities of $20\text{-}25 \text{ m s}^{-1}$ spanning the fire head (Fig. 8c,d) and no evidence of feedback from the
 315 fire (e.g., no flow weakening or reversal to the west of the fire head), which is clearly deficient.
 316 The Fuelx4 scenario has a small region of near-zero to slightly positive radial velocities west of
 317 the fire head (Fig. 8e,f). The Fuelx8 scenario has the greatest extent of outbound radial velocities
 318 in the lee of the head fire and covering a greater areal extent than the Fuelx4 scenario (Fig. 8g,
 319 h). This area of positive ($5\text{-}8 \text{ m s}^{-1}$) radial velocities is around 2 km MSL with small regions of
 320 stagnant flow extending up to 3 km MSL. Strong negative radial velocities upwind of fire front
 321 indicate a region of strong convergence with the fire-generated wind at the head fire. While the
 322 Fuelx8 scenario is in best agreement with the observations, it still underestimates the strength
 323 and spatial extent of the fire-generated winds apparent in the

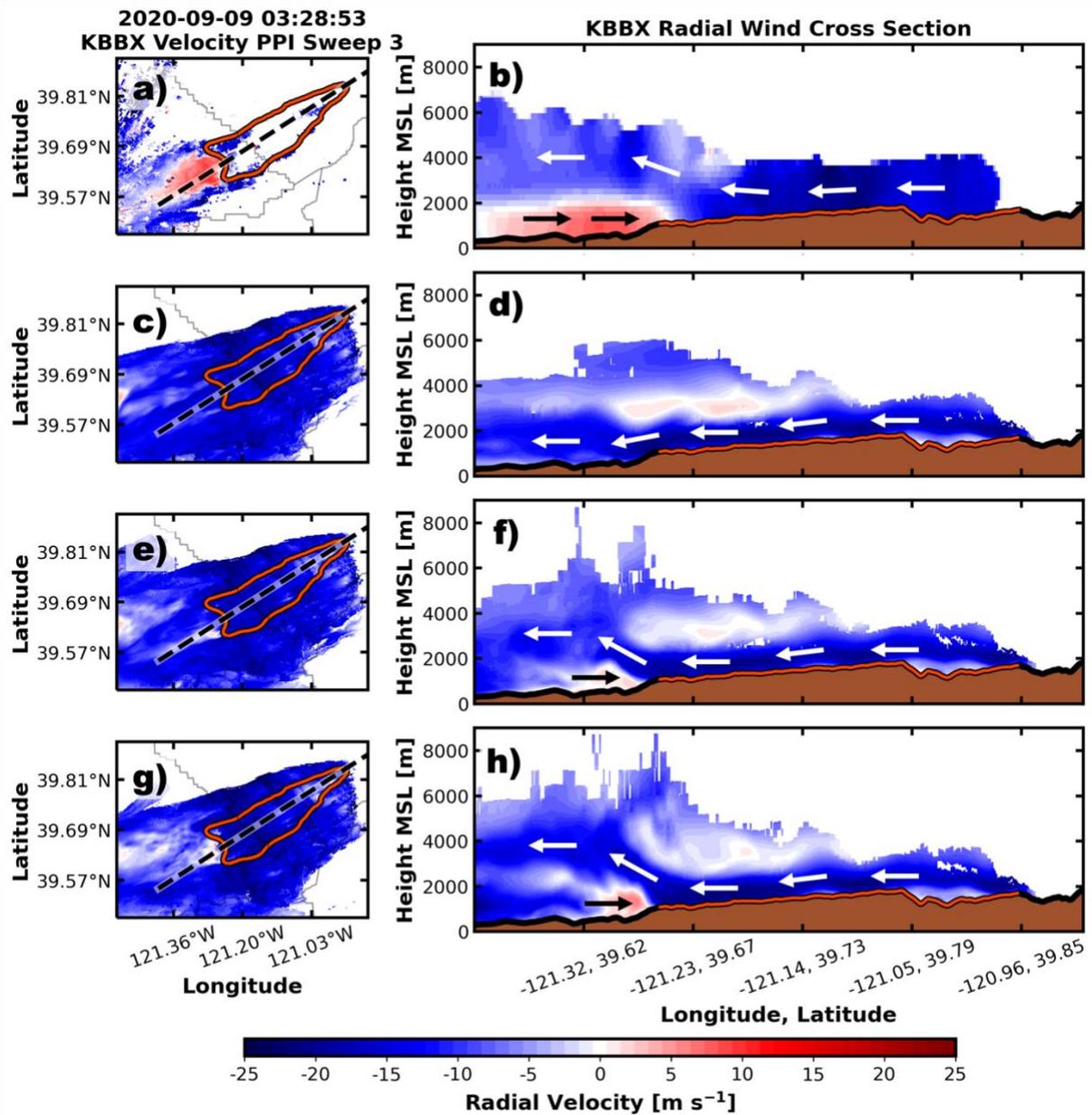


Figure 8. Comparison of observed and simulated fire flows during the Bear Fire. (a) Beale Air Force Base (KBBX) NEXRAD radial velocity PPI (shaded) and radar-estimated fire perimeter (red contour), (b) radial wind and radar-estimated fire perimeter (red line) cross section along black dashed line, WRF-Fire simulated PPI and cross section of in-plume radial velocity and fire perimeter for (c-d) Fuelx1, (e-f) Fuelx4, and (g-h) Fuelx8 scenarios, around 0330 UTC September 9, 2020. In-plane directional flow vectors annotated in b, d, f, h.

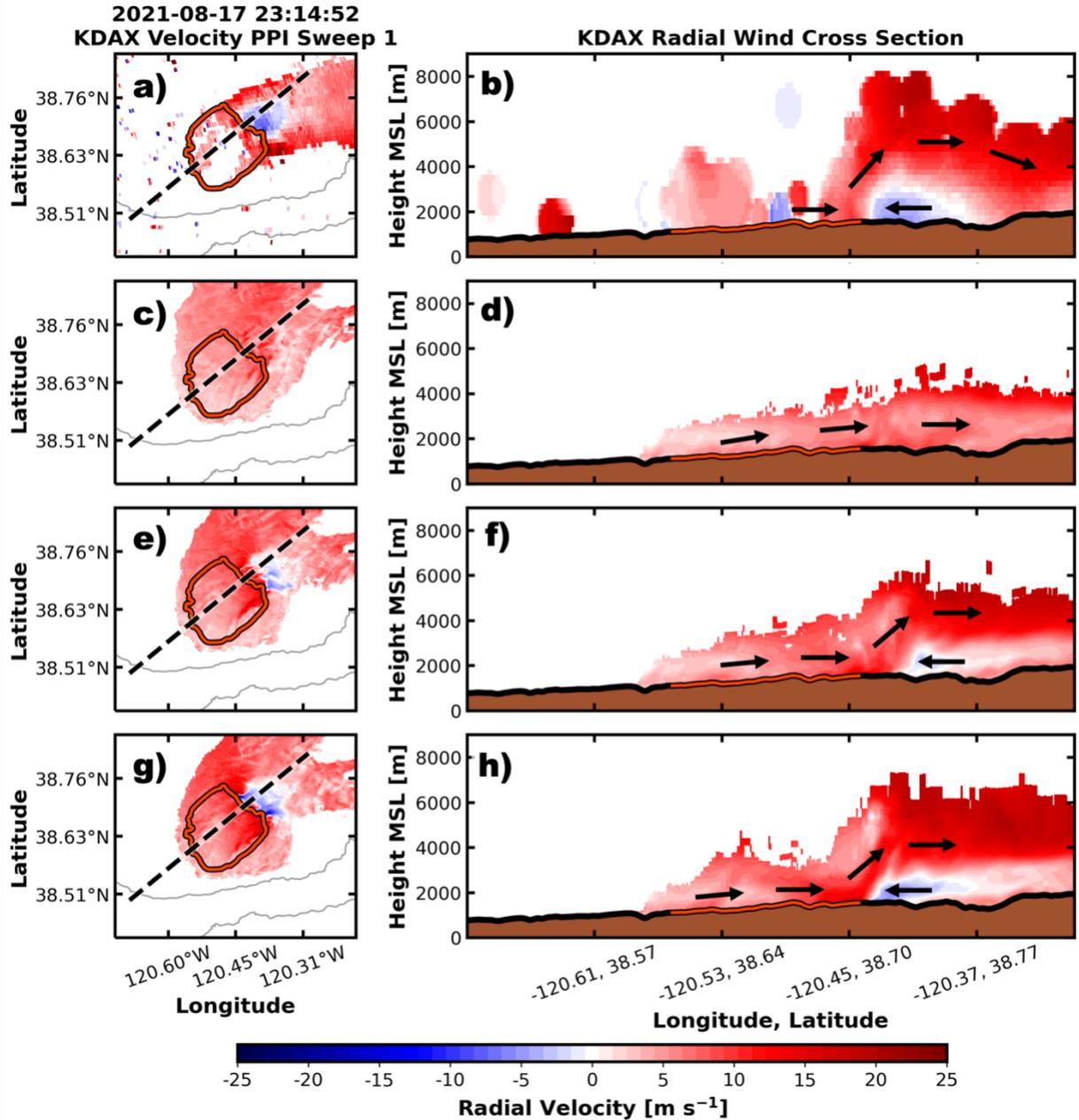


Figure 9. Comparison of observed and simulated fire flows during the Caldor Fire. (a) Sacramento (KDAX) NEXRAD radial velocity PPI (shaded) and radar-estimated fire perimeter (red contour), (b) radial wind and radar-estimated fire perimeter (red line) cross section along black dashed line, WRF-Fire simulated PPI and cross section of in-plume radial velocity and fire perimeter for (c-d) Fuelx1, (e-f) Fuelx4, and (g-h) Fuelx8 scenarios, around 2315 UTC August 17, 2021. In-plane directional flow vectors annotated in b, d, f, h.

324 observations, suggesting that the actual fuel consumption, or rate of consumption, during the
325 Bear Fire may exceed our simulated results.

326 We find similar sensitivity to fire-generated horizontal winds for the Caldor Fire when
327 we compare simulated radial velocities with those observed by the Sacramento NEXRAD
328 (KDAX) on 17 August 2021 around 2315 UTC. To be specific, the observations indicate upslope
329 flow with generally positive (outbound) background radial velocities of around 10 m s^{-1} (red
330 shading in Fig. 9a, b) with a pronounced region of fire induced flow reversal (inbound, blue
331 shading) radial velocities in the lee of the fire head that extends up to approximately 4 km MSL.
332 Similar to the Bear Fire, the Caldor Fuelx1 scenario shows no evidence of fire generated flow
333 reversal, with positive radial velocities of $5\text{-}10 \text{ m s}^{-1}$ extending across the fire front, and no
334 region of flow weakening or reversal near the fire head (Fig. 9c,d). The Fuelx4 scenario shows a
335 small region of stagnant to negative radial velocities ($0\text{-}5 \text{ m s}^{-1}$) on the northeast lobe of the fire
336 head (Fig. 9e). Fig. 9f shows this region of inbound radial velocities extends up to about 3 km
337 MSL and is slightly displaced downstream of the fire head. The Fuelx8 scenario again shows the
338 most pronounced flow perturbation by the fire with a larger region of inbound radial velocities in
339 the lee of the fire front with peak values around 10 m s^{-1} (Fig. 9g). Fig. 9h shows the maximum
340 of this flow reversal region is situated immediately downstream of the fire front, with flow
341 stagnation extending well downstream of the fire head as evidenced by the region of weakly
342 positive radial velocities.

343 A clear takeaway from these results is that when models produce too little heat flux, they
344 also produce deficient fire-generated horizontal winds and thus do not capture critical
345 components of the feedback between the fire and the atmosphere. In that these fire-generated
346 winds have been identified as contributors to the onset of extreme events, such as FGTVs
347 (Lareau et al., 2022a), this data deficiency urgently needs to be resolved.

348 4.3 Plume Depth and Updraft Strength

349 Consistent with the increase in horizontal flow perturbations, our simulations also show
350 increases in vertical velocity, plume verticality, pyroCu/Cb initiation, and smoke injection height
351 which are proportional to the increase in fuel load and thus heat flux (Fig. 10a-c).

352 To frame the simulation results, we first examine radar observations of the plume
353 structure. During the Bear Fire, representative radar cross sections indicate an upright plume core
354 (e.g., corridor of high reflectivity) rising from the head fire with plume tops near 9 km MSL (Fig.
355 10d). Further analysis of the plume evolution from 0000-0300 UTC (not shown, see also Lareau
356 et al., 2022a,b) indicates plume tops ranging from 8 to almost 12 km MSL including
357 considerable pyroCu/Cb development, with cloud bases near 6 km MSL. While we do not have
358 updraft observations, it is reasonable to conclude that these deep, nearly vertical plume cores
359 must possess very strong (e.g., $>30 \text{ m s}^{-1}$) updrafts that can compete with the strong cross flow
360 ($25\text{-}30 \text{ m s}^{-1}$) to produce an upright plume core.

361 Unsurprisingly, these upright plume structures and high vertical velocities are absent
362 from the Fuelx1 simulations, but present in the high fuel load cases (Fig. 10a-c). To be specific,
363 at 0200 UTC in the Bear Fire Fuelx1 simulation, maximum updraft velocities are less than 10 m s^{-1}
364 and do not penetrate above 3 km MSL (Fig. 10a). The maximum smoke plume depth in this
365 scenario is less than 4 km MSL. The Fuelx4 simulation has maximum updraft velocities of just
366 over 20 m s^{-1} with comparatively wider and deeper updraft cores of ~ 3 km wide and 5-6 km
367 MSL deep, respectively (Fig. 10b). The smoke plume depth reaches 6 km MSL in this scenario,

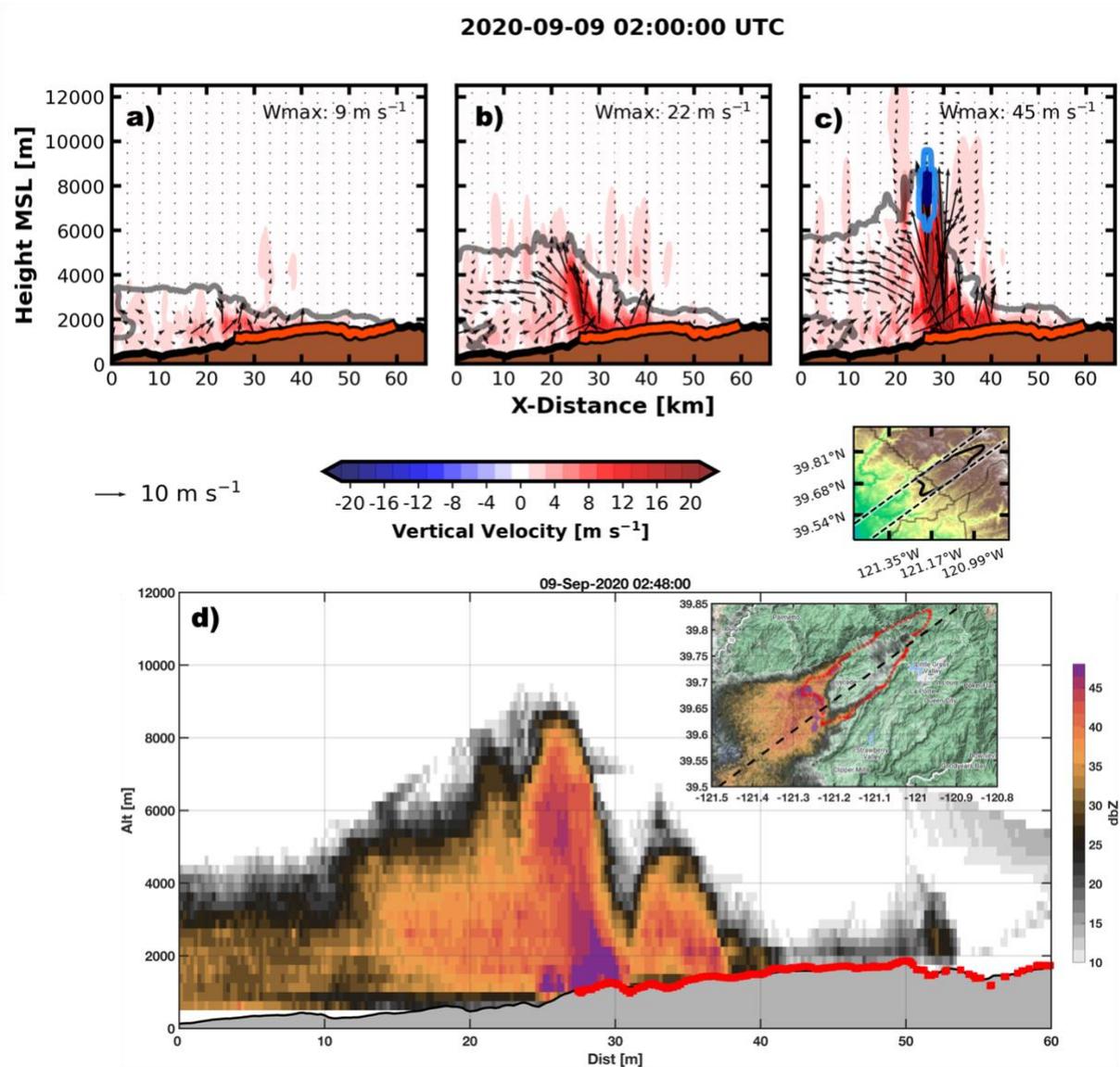


Figure 10. Bulk cross section normal to the Bear Fire head of fire-generated maximum vertical velocity (shaded), in-plane average fire-generated wind vectors, smoke plume extent (gray contour), cloud water (navy) and ice (light blue) contours in the (a) Fuelx1, (b) Fuelx4, and (c) Fuelx8 simulations, and (d) observed NEXRAD reflectivity (shaded) cross-section.

368 and a fire-generated circulation is evident in the lee of the plume with surface inflow, and
 369 outflow at about 3-4 km MSL. The Fuelx8 simulation has the deepest, strongest, and most
 370 upright plume of any scenario with vertical velocities exceeding 40 m s^{-1} and penetrating to 8-10
 371 km MSL (Fig. 10c). The wide (~5km) updraft base is inducing strong inflow at the surface and
 372 strong outflow at 4-6 km MSL in the lee of the plume. Notably, this scenario produced multiple
 373 instances of pyroCb, with a high-based pyroCb occurring at 0200 UTC between 6 and 10 km
 374 MSL (see blue cloud water contour in Fig. 10c), consistent with NEXRAD and photographic
 375 observations during this period (Fig. 10d, see also Fig. 10 in Lareau et al., 2022a). The strong

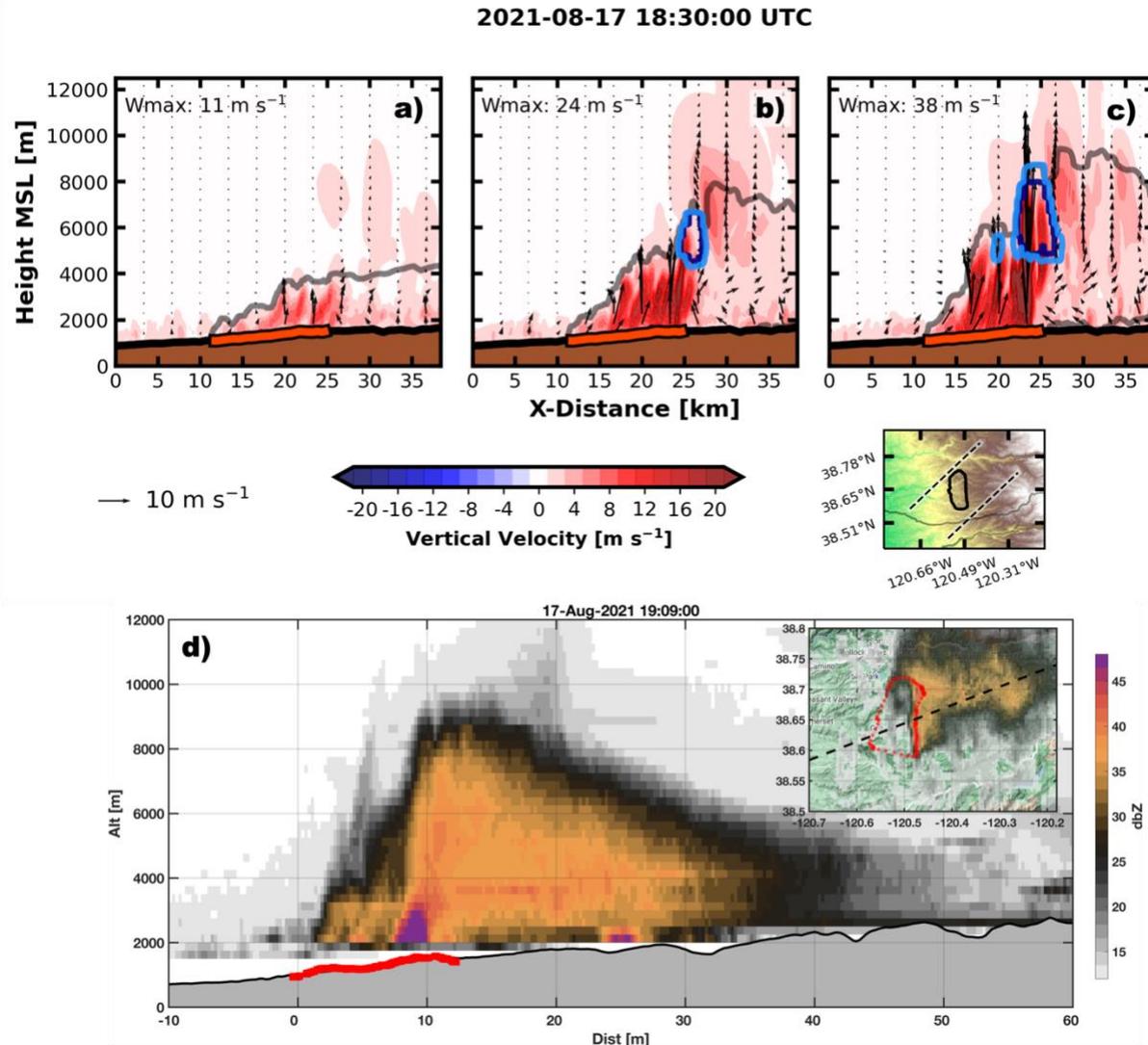


Figure 11. Bulk cross section normal to the Caldor Fire head of fire-generated maximum vertical velocity (shaded), in-plane average fire-generated wind vectors, smoke plume extent (gray contour), cloud water (navy) and ice (light blue) contours in the (a) Fuelx1, (b) Fuelx4, and (c) Fuelx8 simulations, and (d) observed NEXRAD reflectivity (shaded) cross-section.

376 simulated updrafts linked to pyroCb are consistent with observations of other extreme wildfires
 377 (Rodriguez et al. 2020).

378 Whereas the Bear Fire updrafts must compete with very strong ambient winds, the Caldor
 379 Fire's updrafts experience much weaker background flow yet show similar sensitivity to fuel
 380 load. For example, at 1830 UTC, the Caldor Fire Fuelx1 simulation (Fig. 11a) produces
 381 maximum updrafts of $\sim 10 \text{ m s}^{-1}$ reaching $\sim 4 \text{ km}$ MSL with ill-defined updraft cores. The Fuelx4
 382 simulation (Fig. 11b) produces updraft velocities greater than 20 m s^{-1} , penetrating up to $\sim 5 \text{ km}$
 383 MSL and producing shallow pyroCu between 5 and 7 km MSL. This scenario contains a
 384 comparatively wide ($\sim 10 \text{ km}$) updraft region containing several narrow updraft cores from the
 385 surface up to 5 km MSL. There is also a weak fire induced circulation in the lee of the plume,

386 with weak surface inflow vectors and slightly stronger outflow at about 3-4 km MSL. The
387 Fuelx8 scenario (Fig. 11c) contains the most coherent updrafts with a ~5 km wide and ~8 km
388 MSL deep region of vertical velocities just under 40 m s^{-1} . The resulting plume depth in this
389 scenario neared 10 km, with a 4-5 km deep pyroCu/Cb, well-developed surface inflow region,
390 and ~4km MSL outflow in the lee of the plume. This simulated plume and pyroCu/Cb structure
391 compares well with the KDAX NEXRAD data, where plume tops were around 10 km MSL (Fig.
392 11d) and high reflectivity cores suggest upright and vigorous updrafts.

393 The results of both the Bear Fire and Caldor Fire simulations indicate that not only are
394 fuel characteristics important in generating realistic plumes and fire-generated flows, but they
395 also have clear implications for simulating deep pyroCb, which can generate additional
396 feedbacks on the fire environment (e.g., downdrafts, lightning, FGTVs) and may result in
397 stratospheric smoke injection.

398 **5 Summary and Discussion**

399 Our sensitivity analyses indicate that WRF-Fire run with Scott and Burgen 40 fuel
400 categories under-represents fuel quantity and its consumption, and thus underrepresents fire-
401 generated heat fluxes, resulting in deficient simulation of the atmospheric response to landscape-
402 scale wildfire processes. Among these deficiencies are shallow plumes with weak updrafts and
403 little-to-no fire-induced flow perturbations. These deficiencies are also driven by insufficiently
404 wide areas of combustion behind the fire-front (e.g., “deep flaming” in the model), which is
405 linked to both the deficient fuel load and the fire’s residence time. For example, neither the Bear
406 nor Caldor Fire baseline (Fuelx1) simulations produced a broad combustion region with
407 sufficiently large sensible heat fluxes to produce deep updrafts initiating pyroCu/Cb. This stands
408 in stark contrast to radar observations of both fires, which reveal deep, upright convective cores
409 linked to pyroCu/Cb. In contrast, the Fuelx4 and Fuelx8 scenarios generated wider combustion
410 zones and greater total heat fluxes resulting in deep (e.g., 10 km MSL) upright plumes with
411 vigorous updrafts initiating pyroCu/Cb. Since strong inflows, updrafts, and pyroCu/Cb initiation
412 are all vital mechanisms for the development of extreme fire behavior (e.g., FGTVs, long-range
413 spotting) this sensitivity to fuel load underscores current shortcomings in the fuel inputs driving
414 WRF-Fire. Such shortcomings likely apply to other coupled fire-atmosphere models using the
415 Rothermel spread model combined with LANDFIRE-informed fuel data sets (such as Anderson
416 13 or SB40), ultimately limiting their capacity to accurately simulate landscape-scale fires.

417 While some efforts have been made to improve this representation by adjusting fuel
418 categories via machine learning (DeCastro et al., 2022) and accounting for canopy fuels through
419 addition of crown fire heat and improved heat release schemes (Shamsaei et al., 2023b), the
420 foundation of both the Anderson 13 and SB40 fuel data is surface fuels in LANDFIRE which
421 appears to severely underrepresent real-world fuel loads available for consumption in large fires.
422 Such fuel availability is directly linked to wildfire energy release (Goodwin et al., 2021). Thus
423 for accurate, operational simulations of landscape scale fire spread, a methodology that
424 incorporates both surface and canopy fuel loading (e.g., dead and down debris, standing dead,
425 etc.) and landscape-scale fire processes (e.g., spotting, mass-fire, post-frontal combustion) must
426 be incorporated into coupled fire-atmosphere models.

427 In identifying these shortcomings, a potential path forward involves improved
428 representation of fuel inputs (e.g., inclusion of canopy and down woody fuel loading in
429 LANDFIRE) for use in WRF-Fire and other coupled fire-atmosphere models. However, in our

430 simulations we bypassed the large uncertainties in fire spread due to fuel loading by forcing the
431 fire perimeter with observations. This enabled us to change the fuel loads without changing the
432 rate of spread. In freely evolving simulations this is not possible, and simply increasing the fuel
433 load will yield, by formulation, slower rate of spread from the Rothermel model. This issue is
434 compounded in that our forced perimeters include the result of near- and long-range spotting
435 whereas the Rothermel model does not represent long-range spotting. Thus, to achieve high
436 fidelity and freely evolving simulations critical for operational forecasting, the community will
437 need to improve the underlying physical representation of fire spread processes, not just the fuel
438 and its consumption. In the meantime, a combination of assimilating fire perimeter observations
439 (e.g., Farguell et al., 2021 and the approach used herein) and adjusting fuel loads based on
440 machine learning is one approach to bypass uncertainties in the model physics and realize
441 potentially useful simulations not just of the fire spread but also the attendant atmospheric
442 circulations.

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452 **Data Availability Statement**

453 WRF-Fire output was analyzed using Python 3.8. Model output, processing codes, and fire
454 perimeter files (<https://doi.org/10.7910/DVN/FEHPIH>; Roberts and Lareau, 2023) are available
455 on Harvard Dataverse. Ancillary data used in these analyses are free and publicly available
456 through AWS. NEXRAD and GOES-17 data are available at [https://registry.opendata.aws/noaa-](https://registry.opendata.aws/noaa-nexrad/)
457 [nexrad/](https://registry.opendata.aws/noaa-nexrad/) and <https://registry.opendata.aws/noaa-goes/>.

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