

# **Experimental Investigation on the Transport of Sulfide Driven by Melt-rock Reaction in Partially Molten Peridotite**

**Zhenjiang Wang<sup>1,3</sup>, Zhuo-sen Yao<sup>2</sup>, Zhenmin Jin<sup>3</sup>, Yannan Wang<sup>1</sup>**

<sup>1</sup>Key Laboratory of Resource Exploration Research of Hebei Province, and School of Earth Science and Engineering, Hebei University of Engineering, 056038 Handan, China.

<sup>2</sup>State Key Laboratory of Geological Processes and Mineral Resources, and School of Earth Resources, China University of Geosciences, 430074 Wuhan, China.

<sup>3</sup>State Key Laboratory of Geological Processes and Mineral Resources, and School of Earth Sciences, China University of Geosciences, 430074 Wuhan, China.

Corresponding author: Zhuo-sen Yao and Zhenmin Jin ([yaozhuosen@cug.edu.cn](mailto:yaozhuosen@cug.edu.cn);  
[zmjin@cug.edu.cn](mailto:zmjin@cug.edu.cn))

## **Key Points:**

- Systematical variations are produced by the melt-rock reaction in lithologies and compositions of silicate melt and minerals
- Melt-rock reaction is conducive to the upward transport of sulfide droplets in a crystal framework based on experimental observation
- Focused melt flow with high velocity can drive upward extrusion and entrainment of  $\mu\text{m}$ - to mm-scale sulfides through narrow melt tubes

## **Abstract**

Extraction of sulfides from the partially molten mantle is vital to elucidate the cycling of metal and sulfur elements between different geochemical circles but has not been investigated systematically. Using laboratory experiments and theoretical calculations, this study documents systematical variations in lithologies and compositions of silicate minerals and melts, which are approximately consistent with the results of the thermodynamically-constrained model. During a melt-peridotite reaction, the dissolution of olivine and precipitation of new orthopyroxene generate an orthopyroxene-rich layer between the melt source and peridotite. With increasing reaction degree, more melt is infiltrated into and reacts with upper peridotite, which potentially enhances the concomitant upward transport of dense sulfide droplets. Theoretical analyses suggest an energetically focused melt flow with a high velocity ( $\sim 170.9 \mu\text{m/h}$ ) around sulfide droplets through the pore throat. In this energetic melt flow, we, for the first time, observed the mechanical coalescence of sulfide droplets, and the associated drag force was likely driving upward entrainment of fine  $\mu\text{m}$ -scale sulfide. For coarse sulfide droplets whose sizes are larger than the pore throat in the peridotite, their entrainment through narrow constrictions in crystal framework seems to be physically possible only when high-degree melt-peridotite reaction drives high porosity of peridotite and channelized melt flows with extremely high velocity. Hence, the melt-rock reaction could drive and enhance upward entrainment of  $\mu\text{m}$ - to  $\text{mm}$ -scale sulfide in the partially molten mantle, potentially contributing to the fertilization of the subcontinental lithospheric mantle and the endowment of metal-bearing sulfide for the formation of magmatic sulfide deposits.

## **Plain Language Summary**

Sulfides are a pivotal potential reservoir for sulfur and economically important metals. Their transport in the Earth's mantle plays a vital role in understanding many crucial geological and environmental processes, especially the formation of mineral deposits, and the environmental damage and health hazards related to volcanic eruptions. This work proposes a new driving force for the upward transport of dense sulfide drops in the upper mantle that experiences partial melting. The reaction between melt and rock potentially leads to focused melt flow in new-forming channels with three orders of magnitude higher velocity than that of melt flowing among the crystal framework of peridotite. This energetic melt flow drives the upward transport of tiny

μm-scale sulfide droplets in peridotite and may also facilitate the amalgamation of droplets contacting each other. Coarse sulfide droplets could be possibly entrained upward through narrow pore throats, especially when a high-degree melt-peridotite reaction drives fast-flowing melt in the mantle with high porosity.

## 1 Introduction

Sulfide is a ubiquitous phase in the mantle (Alard et al., 2011) and an important repository for sulfur and geochemically and economically important chalcophile metals, which plays a pivotal role in the partitioning behaviors of PGE, Cu, and Ni (Mungall & Brenan, 2014; Patten et al., 2013). Understanding the factors that control the fate of sulfide phases in the partially molten mantle is of fundamental importance in exploring the recycling of sulfur and chalcophile elements among different geochemical reservoirs (Ding & Dasgupta, 2017; Farquhar et al., 2002; Yao et al., 2018; Chen et al., 2022) and identifying the re-fertilization of the depleted lithospheric mantle, which potentially provides the metal endowment for the formation of Cu-rich porphyry and/or Ni-rich magmatic ore system (Holwell et al., 2022; Lee & Tang, 2020; Mungall et al., 2015; Zhao et al., 2022). However, it remains highly contentious for the driving forces of the transport of metal-bearing sulfide liquid, which severely blocks our understanding of the details of fertilization processes that occurred in the sources of sulfide-related magmatic and hydrothermal deposits.

Conventionally, the removal of sulfide from a partially molten peridotite requires progressive dissolution of sulfide into the departing silicate melt (Holzheid & Grove, 2002; Mungall & Brenan, 2014; Yao et al., 2018). Nevertheless, the mechanical entrainment of sulfide could also potentially be an efficient process driving the redistribution and local enrichment of sulfur and chalcophile elements in the mantle (Iacono-Marziano et al., 2022; Heinrich & Connolly, 2022; Z. Wang et al., 2020; Yao & Mungall, 2020). Although a small amount (< percolation thresholds) of sulfide under hydrostatic conditions may not be entrained through the slow porous flow of silicate melt due to high surface tension (Bagdassarov et al., 2009; Holzheid et al., 2000; Yoshino et al., 2003, 2004), it has been argued that the extraction of sulfide driven by stress in a partially molten peridotite could be further facilitated by the directional porous flow of silicate melt, strongly affecting the fractionation and abundances of chalcophile metals in the mantle melting products, *i.e.*, basalts (Ballhaus et al., 2006; Bockrath et al., 2004; Z. Wang et

al., 2020). Except for the deviatoric stress, a preliminary experimental study (Wang & Jin, 2020) proposed that during the melt-rock reaction, entrainment of sulfide liquid may be enhanced by the reaction infiltration instability (RII) in the partially molten peridotite when a reacting melt percolates through a dissolvable, porous, melt-mineral mixture. The mechanical transport and enrichment of sulfide during this melt-rock reaction have been extensively demonstrated in mantle peridotites, which has clarified some geochemical paradoxes of the chalcophile/siderophile elements (Ciazela et al., 2018; Lorand & Luguet, 2016). The RII theory proposed by Chadam et al. (1986) also applies to magmatic systems by introducing compaction of crystal mush and driving a solubility gradient along the flow direction instead of the propagating reaction front (Aharonov, 1995), resulting in positive feedback among the increasing permeability driven by melt-rock reaction and the associated reinforcement of melt flux in reacting regions. This melt-rock reaction is well known as a pervasive process even forming melt-rich channels with an extremely speedy magma flow due to the positive feedback in partially molten regions of the mantle, which has been broadly demonstrated by geological investigations (e.g., Kelemen et al., 1995; Sundberg et al., 2010), high-temperature, high-pressure (HTHP) laboratory works (e.g., Daines & Kohlstedt, 1994; Pec et al., 2015) and theoretical analyses (Aharonov, 1995; Chadam et al., 1986; Spiegelman et al., 2001). During a reaction of melt with a depleted mantle, melt re-fertilization processes potentially occur, thereby impregnating the mantle with sulfides, which have been widely found in mantle xenoliths, orogenic and abyssal peridotites (e.g., Ciazela et al., 2018; Luguet et al., 2003; Niu, 2004; Wang et al., 2009).

Although it has been tentatively demonstrated that the rapid enough ascending magma flow in these melt-rich channels could hold upward entrainment of large olivine phenocrysts or even dense sulfide droplets during a melt-rock reaction (Pec et al., 2017; Wang & Jin, 2020), extraction of sulfide driven by the RII has received little attention, and the precise physical and chemical constraints about these processes appear to be worth studying thoroughly. Here we conducted systematically two-layer reaction experimental studies in which a partially molten peridotite was placed on a sulfide-bearing silicate melt source to explore the entrainment of sulfide liquids driven by the RII and trace the physical and chemical variations of liquid and solid phases during these processes, which sheds lights on the fertilization of lithospheric mantle, as well as the formations of magmatic and/or hydrothermal sulfide deposits.



## 2 Experimental Methods

### 2.1 Starting materials

As a melt source, the starting materials were a mixture of the powdered calc-alkaline tholeiitic basalt (70 wt.%) from the East Pacific Rise (102.7044 °W, 2.64961 °S), olivine crystals (20 wt.%) from fresh spinel lherzolite xenoliths at Damaping (Hannuoba region, North China), and sulfide aggregates (10 wt.%) from Jinchuan Ni-Cu sulfide deposit, NW China. Additionally, 0 or 2 wt.% oceanic sediments (mainly carbonate) were mixed into the melt source. The sulfide aggregate (59 wt.% pyrrhotite, 36 wt.% pentlandite, and 5 wt.% chalcopyrite) adopted here is similar to that in massif peridotites (Lorand et al., 2010), and was grounded in alcohol for about 6h to < 10 μm measured by using the scanning electron microscope (SEM).

To prepare a partially molten rock, the olivine and clinopyroxene crystals from Damaping spinel lherzolite xenoliths were mixed in a 50:50 ratio by weight. Then, ~ 0 or ~ 5 wt.% calc-alkaline tholeiitic basalt was added to change the initial porosity of the crystal framework. All silicate mineral grains were ground to about 10-20 μm grain size in an agate mortar. The partially molten rock was placed on the melt source to explore the possible upward transport of dense sulfides from the melt source to the peridotite. The length ratio of these two parts was slightly larger than ~ 1:1 to avoid chemical equilibration among them over the experimental time scales. Compositions of these starting materials were described by Z. Wang et al. (2020) (Table S3), and Table 1 lists the experimental conditions.

### 2.2 Experimental methods and analytical techniques

The two-layer reaction experiments were conducted via a 150 Ton non-end-loaded type piston-cylinder press at China University of Geosciences. Starting materials mentioned above are loaded into a cylindrical platinum (Pt) capsule (3.75 mm diameter; 5-6 mm high) with a graphite inner sleeve (Figure 1-inset in the lower left). Due to the graphite inner sleeve, the oxygen fugacity of this experimental system was maintained at about CCO-0.8 (the graphite-CO<sub>2</sub> buffer), which corresponds to a log  $fO_2$  < FMQ-2 log unit at ~ 1,360 °C and 1.5 GPa (e.g., Médard et al., 2008). Information on experimental assembly can be found in our previous work (Z. Wang et al. 2020). In order to remove absorbed water, a 12-hour heating (120 °C) in a vacuum oven is conducted for all capsules before being sealed, and we did not observe any

apparent sulfur loss during this drying. The experiment conditions in Table 1 should be held for 12-72 h, and then these materials were quenched to room temperature.

Polished sections were prepared from the recovered experimental specimens cut parallel to the specimen axis. The observation of microstructure in these sections was conducted via a Quanta 450 field-emission scanning electron microscope (FE-SEM) at China University of Geosciences, with a 6.0  $\mu\text{m}$  spot size, 20 kV accelerating voltage, and 12 mm working distance. The backscattered electron (BSE) images were adopted to measure the morphological characteristics of experimental products, such as the dissolution distance, the thickness of reactive boundary layers, and the size of sulfide droplets. Quantitative compositional analyses for olivine and clinopyroxene were obtained via a JEOL JXA-8230 electron probe microanalyzer with four wavelength-dispersive spectrometers (WDS) at China University of Geosciences, with an accelerating voltage of 15 kV, a probe current of 20 nA and a beam diameter of 1  $\mu\text{m}$ . Sulfides were analyzed with a higher accelerating voltage (20 kV), and the compositions of quenched silicate melts were measured with a smaller probe current (10 nA). A defocused beam of 20-30  $\mu\text{m}$  diameter was adopted here for all the standardizations, quenched melts, and sulfides. Dwell times were 10s on element peaks and half that on background locations adjacent to peaks. Raw X-ray intensities were corrected using a ZAF (atomic number, absorption, fluorescence) correction procedure. A series of natural and synthetic SPI standards were utilized and changed based on the analyzing minerals. The following standards were used: sanidine (for K), pyrope garnet (for Al), almandine garnet (for Fe), diopside (for Ca, Si, Mg), jadeite (for Na), rhodonite (for Mn), olivine (for Si, Mg), rutile (for Ti), apatite (for P), chromium oxide (for Cr). The analytical standards for sulfides were pentlandite (for Fe, Ni, S) and native copper (for Cu) (Pan et al., 2023).

### 3 Results

In order to investigate the physical migration of sulfide during reaction infiltration of silicate melt in the partially molten mantle, we have conducted eight reaction experiments at 0.5-1.5 GPa, 800-1,300  $^{\circ}\text{C}$ , and 12-72 h (Table 1). Generally, several distinct lithological regions were developed and, from bottom to top, respectively separated by a diffuse-dominated, mineralogical interface in these two-layer reaction experiments (Figure 1; Figures S2-S8). The

freeware ImageJ (<http://imagej.nih.gov/ij/>) from the National Institute of Health was adopted to analyze all two-dimensional BSE images obtained here.

### 3.1 Phase assemblages and textures

Under the low-temperature, high-pressure conditions (800 °C and 1.5 GPa), both minerals and sulfides were generally polygonal and not molten, and sulfide droplets in the lower part of the sample were uniformly distributed in pores among silicate grains. At 1.5 GPa, when the experimental temperatures exceed ~1,250 °C, the basaltic material was completely molten, and the diameter of sulfide droplets ranges from several  $\mu\text{m}$  to  $> 100 \mu\text{m}$  (Figure 1). Gravitational segregation of sulfide was insignificantly observed in all sections, but obviously, some huge sulfide droplets were attached to the wall of the graphite capsule, especially at high temperature and long experiment durations (Figure 1). These coarse droplets were excluded from the collected data to calculate the average sulfide size. Tiny exsolution lamellae with higher backscattered electron (BSE) intensity were occasionally observed in some quenched large sulfide droplets. Sometimes silicate melts were pooled at the edge of the sample (Figure 1).

In all high-pressure experiments (1.5 GPa), an orthopyroxene (Opx)-rich reaction layer (ORL) was formed at the boundary of the lower melt source and upper partially molten peridotite due to the reaction between them, and chiefly composed of Opx grains, silicate melt, and sulfide droplets (Figure 1). With increasing the annealing time and/or run temperature, the morphology of ORLs became more irregular (Figure 1b; Figures 2a, b, and d), and small olivine inclusions were observed in a few large Opx grains (Figure 2a). The ORL at 1,300°C was even ruptured mainly due to the growth of several Opx grains to enormous size ( $> 100 \mu\text{m}$ ) and possible gravitational settling (Figure 2d). We also observed the bulge of ORL toward the melt source (Figure 1b), which can be attributed to the slight gravitational settling of Opx and volumetric shrinkage of the melt source after the upward reactive percolation of silicate melt. Thus, the original interface between the melt source and peridotite before the experimental reaction cannot be marked by the offset of the graphite inner sleeve. At 1,250 °C, the thickness of ORL increases from  $\sim 118 \pm 11 \mu\text{m}$  ( $1\sigma$ ) at 12 h to  $\sim 416 \pm 119 \mu\text{m}$  at 72 h, which is likely a linear function of the square root of the run duration (slope =  $0.78 \pm 0.05$ ) (Figure 3a-blue line). With increasing the temperature to 1,300 °C, the thickness increases considerably to  $500 \pm 128 \mu\text{m}$  at 12 h (Figure 3a), implying the contribution of a higher temperature to the growth of ORL.

199 Additionally, based on the linear fit, the addition of 5 wt.% basalts into the partially molten  
200 peridotite (run PC548) or 2 wt.% oceanic sedimentary into the melt source (run PC560) seems to  
201 have no significant effect on the thickness of ORL (Figure 3a).

202 On the other hand, peridotites in the upper part of our samples were in disequilibrium  
203 with silicate melt at these run conditions, thereby should be dissolved essentially into the melt.  
204 The degree of dissolution is indirectly quantified via the dissolution distance (Figure 3b), which  
205 is the distance between the current boundary (Figure 1-blue dotted lines) and the original  
206 interface (Figure 1-red dotted lines) that is approximately represented by a sharp decrease in the  
207 size of sulfide droplets in the melt source. All these dissolution distances are measured at least 5  
208 times around the central part of experimental charges. Similarly, dissolution distance linearly  
209 increases with the square root of time (Figure 3b-blue line) at 1,250 °C, and the increase of  
210 temperature to 1,300 °C significantly enlarges the dissolution distance (Figure 3b). The addition  
211 of oceanic sedimentary (PC560) or basalt (PC548) has no substantial effect on the dissolution  
212 distance (Figure 3b).

213 Moreover, based on analyses of SEM images, with increasing run duration or  
214 temperature, more silicate melt penetrated the Opx reaction layer into the partially molten  
215 peridotite ( $\sim 4.65 \pm 0.78$  area% at 12 h;  $\sim 7.13 \pm 0.58$  area% at 72 h) (Figure 4a), resulting in the  
216 formations of melt junctions and channels among silicate minerals, and even some large melt  
217 pools just above the ORL (Figure 2c; Figures S2-S8). At 1.5 GPa, with increasing annealing time  
218 from 12 h to 72 h, the grain size of silicate minerals increases from  $8.1 \pm 3.4$   $\mu\text{m}$  to  $11.0 \pm 4.9$   
219  $\mu\text{m}$  in the partially molten peridotite (Figure 4b; Supporting information; Table S1). Under  
220 conditions of high temperature (1,300 °C) and long-run duration (72 h), an olivine-melt layer,  
221 consisting of sulfide, olivine, and silicate melt, is present above the ORL (Figures 2a and 2d,  
222 Figures S3 and S5), which may be attributed to the dissolution of clinopyroxene and  
223 reprecipitation of olivine during the melt-peridotite reaction. Additionally, crystal faces between  
224 olivine grains were open and full of silicate melt, forming abundant melt channels among the  
225 crystal framework (Figure 2c, Figures S2-S8).

226 In the melt source, an amount of Opx grains were produced by the reaction between  
227 olivine in starting materials and convecting silicate melt (Figure 1), and there is commonly no  
228 significant boundary between these Opx grains and the ORL. The mean area-weighted diameter

of sulfide droplets was  $\sim 4.1 \pm 0.1 \mu\text{m}$  in a low-temperature hot-press experiment (PC537, 800 °C), which potentially denotes the initial size of sulfide droplets in the starting material. At a higher temperature (1,250 °C), sulfide size increases from  $\sim 14.7 \pm 0.2 \mu\text{m}$  to  $\sim 49.4 \pm 1.6 \mu\text{m}$  with increasing the annealing time from 12 h to 72 h (Figure 5a; Table S2), in which the larger standard deviations from the long annealing time (72 h) experiments may be partly attributed to the presence of some enormous ( $> 50 \mu\text{m}$ ) and tiny ( $< 1 \mu\text{m}$ ) sulfide droplets. Under the same annealing time (48 h), there is no conspicuous increase in sulfide size when 5 wt.% basalt was added to the peridotite (PC548), but a slight increase in sulfide size was observed if the experimental temperature increased to 1,300 °C (PC527) (Figure 5a). The addition of carbonate into the melt source (PC560) resulted in a significant increase in sulfide size to  $\sim 54.0 \pm 1.6 \mu\text{m}$  at 24 h (Figure 5a). Additionally, we also observed that the number density of sulfide droplets uniformly decreases with increasing annealing time in the lower melt source (Figure 5a-inset). The area fraction of sulfide in the melt source of the low-temperature experiment (800 °C) was  $\sim 3.26 \pm 0.12 \text{ area}\%$ . At 1,250 °C, the area fraction was essentially constant ( $3.44 \pm 0.14 \text{ area}\%$  at 12h;  $3.50 \pm 0.15 \text{ area}\%$  at 48h) when the annealing time was less than 48 h, but a visible increase ( $5.46 \pm 0.22 \text{ area}\%$ ) can be observed for a longer annealing time, 72 h (Figure 5b). A similar trend was also shown in the upper part of samples (Figures 5c and 5d), which may be due to more silicate melts infiltrating into the upper part of samples with increasing annealing time to 72 h (Figure 4). As previously observed (Yoshino & Watson, 2005), the diffusion of Fe and Ni components from silicate phases into sulfide liquids may be one of the reasons for the increase of sulfide area fraction in the melt source. In contrast, with increasing temperature to 1,300 °C, a lower area fraction ( $2.93 \pm 0.07 \text{ area}\%$ ) of sulfide was observed even in the experiment with a short annealing time (12 h) (Figure 5b), potentially due to partial dissolution of sulfide driven by the higher sulfur content at sulfide saturation (SCSS) at high temperature (*e.g.*, Liu et al., 2007; Mavrogenes & O'Neill, 1999). Similarly, the additions of basalt and carbonate enhance the permeation of more silicate melt into the upper part of samples, thereby increasing the area fraction of sulfide in the melt source (Figure 5b).

Based on the SEM images of products from these high-temperature experiments (1,250 °C), it becomes evident that sulfide droplets in the melt source were entrained into the partially molten peridotite by a porous flow of silicate melt during the melt-peridotite reaction (Figure 2). With permeating more silicate melt upwards, more and larger sulfide droplets can be observed in

the partially molten peridotite, and meanwhile their sizes (area-weighted diameter  $\sim 3.1 \pm 0.1 \mu\text{m}$  at 12 h to  $\sim 11.1 \pm 0.4 \mu\text{m}$  at 72 h) and area fractions ( $\sim 0.15 \pm 0.05 \text{ area}\%$  at 12 h to  $\sim 0.49 \pm 0.03 \text{ area}\%$  at 72 h) increase with the increasing annealing time at 1,250 °C (Figures 5c and 5d; Table S2). Apparently, the addition of carbonate and 5 wt.% basalt into melt source and partially molten peridotite, respectively, drives a higher area fraction and larger size of sulfide droplet in the upper part of samples (Figures 5c and 5d). Hence, these observations propose that the ORL may not efficiently prevents the upward transport of silicate melt and sulfide droplets from the lower melt source to the upper molten peridotite.

Notably, under the condition of lower pressure (0.5 GPa), some vapor bubbles have been found in the partially molten peridotite due to the exsolution of volatile (mainly  $\text{H}_2\text{O}$ ) in the starting basaltic material, and they mostly absorb on sulfide droplets to form compound drops, which has been proposed to potentially enhance the upward transport of sulfide (Figure S7; Mungall et al., 2015; Yao & Mungall, 2020).

### 3.2 Phase compositions

#### 3.2.1 Mineral compositions

Figure 6 and Figure 7 respectively showed the compositional variations of olivine and clinopyroxene as a function of distance away from the final melt-peridotite interface (Table S4). At 1.5 GPa and 1,250-1,300 °C, from the far-field region to the interface, olivine grains became gradually lower in the Mg# (defined as molar  $\text{Mg}/(\text{Mg}+\text{Fe}) \times 100$ ) and concentrations of  $\text{SiO}_2$ , MgO, and NiO, and meanwhile shown the increases in the FeO and MnO contents (Figure 6). Compared with the low-temperature, hot-press experiment (PC537) (Figure 6 purple dotted lines), olivine grains from high-temperature experiments contain lower MgO and NiO, and higher FeO and CaO concentrations. By contrast, olivine compositions in experiments PC548 (5 wt.% basalts) and PC559 (0.5 GPa) were roughly constant across the partially molten peridotite, whereas olivine grains crystallized in the melt source at lower-pressure experiment (0.5 GPa, PC559) had lower Mg#, MgO,  $\text{SiO}_2$  and higher FeO than those in the peridotite region (Figure 6).

In contrast to olivine, more scatter compositions were observed in the clinopyroxene (Cpx) grains from the partially molten peridotite (Figure 7). Only near the melt-rock interface,

the trends of decreasing Mg# and CaO and increasing FeO were present in Cpx grains from all high-pressure experiments, whereas the Cpx compositions in the low-pressure and carbonate-bearing experiment (PC559 and PC560) were essentially constant across the peridotite region (Figure 7). On the other hand, the Opx grains in the reaction layer and melt source were the products of silicate melt-olivine reaction. No distinct compositional difference was observed in these Opx grains (Figure S9; Table S4).

### 3.2.2 Silicate melt and sulfide liquid compositions

The variations of the reacted melt compositions in these experiments mostly depend on the extent of the melt-rock reaction and the formation of a new major phase. In the melt source, with increasing run time and/or temperature, the compositions of silicate melt became higher in CaO and MgO concentrations, and lower in SiO<sub>2</sub>, Al<sub>2</sub>O<sub>3</sub>, and FeO (Figure 8; Table S5), implying the effect of high-degree melt-peridotite reaction on the melt composition. No distinct variation of the melt sulfur contents was observed under the conditions of 1,250 °C and 1.5 GPa. Based on two new SCSS models (Chowdhury and Dasgupta, 2020; Smythe et al., 2017), the calculated ranges of SCSS in silicate melt are almost the same or slightly lower than our results within the error of measurement (Figure 8f-shadow gray region), potentially indicative of sulfur-saturated silicate melt in these experiments. Moreover, the calculated SCSS in the melt source is nearly the same as those in the interface within the error, which is consistent with our measurements (Figure 8f). The decreases in pressure and temperature respectively to 0.5 GPa and 1,200 °C caused the obvious increases in SiO<sub>2</sub> and Al<sub>2</sub>O<sub>3</sub> concentrations and decreases in MgO and S concentrations in silicate melt (Figures 8a, 8b, 8e, and 8f). No significant variations of melt compositions were observed between the melt source and melt-peridotite interface within the error of measurement (Figure 8), suggesting the convective flow is strong enough to drive the chemical equilibrium of silicate melt. The composition of primary silicate melt infiltrating into the peridotite cannot be analyzed due to the small scale. In addition,  $\sim 1.57 \pm 0.15$  wt.% H<sub>2</sub>O was detected in the reacted melts from experiments PC520 and PC528 using the Fourier Transform Infrared Microscopy (Supporting Information, Mercier et al., 2010), which is slightly higher than that of basalt ( $\sim 1.19$  wt.%) added in the starting materials. On the basis of these melt compositions, the viscosity ( $\mu_m$ ) of the starting basalt ( $\sim 1.19$  wt.% H<sub>2</sub>O) before the melt-rock reaction can be calculated as  $\sim 10.8$  Pa·s at 1.5 GPa and 1,250 °C by using the model from Giordano et al. (2008), whereas the  $\mu_m$  of reacting melt after the reaction decreases to  $\sim 1.2$ -2.7



Pa·s at 1.5 GPa and 1,250-1,300 °C assuming the water content of ~ 1.5 wt.% (Figure 8g). When the pressure and temperature respectively decrease to 0.5 GPa and 1,200 °C, the reacting melt has a higher viscosity of ~10.5 Pa·s. Additionally, there is no evident difference in melt viscosity between the melt source and the melt-peridotite interface. However, the addition of carbonate (PC560) potentially leads to a peculiarly lower melt viscosity (~ 0.25 Pa·s) estimated by using a simple model (Di Genova et al., 2014) under the assumption of the 2 wt.% carbonate dissolved in silicate melt, which represents an order-of-magnitude evaluation of the melt viscosity due to the uncertainty of extrapolation (Figure 8g).

Except for several large sulfide droplets containing the Ni-rich quenched phases, the composition of sulfide was homogeneous in the lower melt source of all experiments. In the 12 h experiment, the S content, Ni/S and Fe/S ratios of sulfide liquid were  $\sim 35.16 \pm 0.72$  wt.%,  $\sim 0.39 \pm 0.03$ , and  $\sim 1.43 \pm 0.04$ , respectively, and these values remained constant within the error of measurement when the annealing time increases to 72 h (S  $\sim 35.74 \pm 0.74$  wt.%, Ni/S  $\sim 0.37 \pm 0.09$ , Fe/S  $\sim 1.42 \pm 0.06$ ). However, the Cu/S ratio of sulfide liquid slightly decreased from  $0.017 \pm 0.004$  at the 12 h to  $0.009 \pm 0.002$  at the 72 h at 1,250 °C, while an obvious increase in the Cu/S ratio ( $\sim 0.046 \pm 0.010$ ) was observed at 1,300 °C (Table S4).

## 4 Discussion

### 4.1 Grain-scale processes in melt-peridotite reaction

Firstly, if the upper peridotite reaches equilibrium at 1.5 GPa and 1,250-1,300 °C, the corresponding melt fraction will be less than 0.04-0.08 wt.% based on the thermodynamic model via the pMELTS (Ghiorso et al., 2002), implying that the self-partial melting of peridotite at the experimental conditions is negligible here. The observed variations of melt area fraction, phase proportions, and minerals' sizes and compositions in our experiments can be primarily attributed to the melt-peridotite reaction. On the other hand, the upper peridotites are not in equilibrium with the starting melts at the experimental conditions, and hence olivine and clinopyroxene in the peridotite part should partially or completely dissolve in the reacting and upward percolating melt (Liang, 1999; C. Wang et al., 2020). In multicomponent partially molten systems, the melt-peridotite reaction is suggested to occur via multi-scale mass transfer processes that inevitably



involve both the grain-scale dissolution-precipitation-reprecipitation and diffusion, in conjunction with large-scale advective transport (Cascio et al., 2008; Morgan & Liang, 2005).

In our high-temperature, high-pressure experiments (1.5 GPa and 1,250-1,300 °C) here, due to the disequilibrium between melt and peridotite, the consequent grain-scale processes can mostly occur via the dissolution, precipitation, and reprecipitation of mineral grains, which is evident from the formation of ORL. In the lower melt source, the basaltic melt in starting materials is olivine-undersaturated, and thus the melt-rock reaction between basaltic melt and olivine produces new Opx grains (Figure 1), which is consistent with the following reaction:



where subscripts 0 and 1 designate melt and mineral grains from starting materials and the reaction products, respectively. This reaction should occur via the dissolution of olivine, and precipitation of new orthopyroxene, largely contributing to the formation of ORL in the interface between the peridotite and melt source (Figure 1a). It is well known that dissolution and precipitation occur simultaneously during the development of the ORL under high pressure (> 1 GPa) (Cascio et al., 2008; Morgan & Liang, 2003). With increasing the degree of melt-peridotite reaction (a longer run duration and/or higher temperature), an olivine-melt layer above the ORL is formed by the further dissolution of olivine and few clinopyroxene grains, with concomitant reprecipitation of fresh olivine grains in the melt reaction flow (Figure 2) (Kelemen et al., 1995).

The grain size of silicate minerals plays a vital role in affecting the physical properties of upper partially molten peridotite, such as permeability (Faul, 2001). Here, the growth of silicate mineral sizes in the upper partially molten peridotite with increasing annealing time (Figure 4b) is mostly driven by a reduction of grain boundary energy (Faul & Scott, 2006) and demonstrates that textural coarsening by which small-size crystals are consumed to support the growth of coarser grains (Higgins, 1998, 2011; Yao et al., 2017), may outweigh the decreasing-size effect related to reprecipitation of new, fine-grained crystals. The transference of material from dissolved olivine to other grains occurs via diffusion through grain boundaries, crystals themselves, and interstitial melt (Yao et al., 2017), while the last one is much faster and becomes the main pathway of these grain-scales processes. Finally, large-scale advective transport of

reacting melt is evident from the variations of melt area fraction in the partially molten peridotite and the formations of melt junctions and channels (Figure 1; Figure 4a).

Based on the parabolic law of diffusive dissolution (Liang, 1999; Zhang et al., 1989), the slope (*i.e.*, the diffusive dissolution constant) of the fitted linear regression line for our experiments is used to estimate the growth rate of ORL ( $k_{ORL}$ ), which is about  $\sim 0.78 \pm 0.05 \mu\text{m/s}^{0.5}$  at 1.5 GPa and 1,250 °C. This growth rate is far lower than those of some previous experiments conducted at higher temperatures and/or pressures (Figure 3a) (Morgan & Liang, 2005; C. Wang et al., 2020). The thickness of ORL in experiment PC527 at 1.5 GPa and a higher temperature ( $\sim 1,300$  °C) has a faster-increasing tendency ( $k_{ORL} = 2.37 \pm 0.14 \mu\text{m/s}^{0.5}$ ), consistent with the experiments under 2 GPa and 1,375 °C (C. Wang et al., 2020). This tendency of ORL growth is likely unchanged in the experiments with the carbonate-bearing melt source (PC560) and silicate melt-bearing peridotite (PC548) (Figure 3a), implying an insignificant effect of melt composition and porosity on the growth of ORL. Our results indicate that the increase in temperature will potentially enhance the growth rate of ORL thickness during the melt-rock reaction, which can be mostly attributed to the higher diffusivities of elements in silicate melts at a higher temperature (Mungall, 2002; Zhang et al., 2010).

Similarly, we observed the linear increase of dissolution distance with the square root of time (Figure 3b), suggesting that a substantial amount of peridotite was dissolved in the percolating melt. At 1.5 GPa and 1,250 °C, the dissolution rate of peridotite evaluated by the slope of the linear regression line ( $k_{diss} = 0.58 \pm 0.08 \mu\text{m/s}^{0.5}$ ) is lower than those in previous experiments with higher temperatures and pressures (Figure 3b, C. Wang et al., 2013, 2020), indicating that relatively high temperature also increases the dissolution rate. Moreover, the higher dissolution rate of peridotite at experiment PC527 (1,300 °C) implies that an increase in temperature significantly enhances melt-rock reaction (Mitchell & Grove, 2016). The high content of water ( $\sim 1.57 \pm 0.15$  wt.%) in silicate melt could enhance the formation of several enormous Opx grains in the ORL (Figure 2) (Wang et al., 2016), partly because the addition of water can strongly depress the peridotite solidus. This may be also the reason for the large standard deviations of the dissolution rate in experiments PC527 and PC545 with a high temperature (1,300 °C) and long annealing time (72 h). Overall, the growth rate of the ORL is generally higher than the peridotite dissolution rate, which is potentially attributed to the volume-increasing reaction during the replacement of olivine by orthopyroxene (Milke et al., 2009).

## 4.2 Compositional variations of melt and minerals

During a melt-peridotite reaction, besides the systematic changes in mineralogy and texture of peridotite, we also observed the compositional variations of melt and minerals in the reaction couple (Figures 6-8), which have been widely used to outline the grain-scale processes governing this melt-peridotite reaction (Mallik & Dasgupta, 2012; C. Wang et al., 2020). On the other hand, a thermodynamically-constrained mixing model has been recently used to examine the variations of major element compositions of minerals during the melt-peridotite interaction (Lambart et al., 2012; Pin et al., 2022; Shaw et al., 2018), and this forward model may offer a key to testing and understanding the compositional evolution of our experimental products.

The melt-peridotite reaction here is modeled as a simplified, thermodynamic process in which the peridotite is continually impregnated by a finite amount of basaltic melt from the lower melt source, which is the same as the assumption in previous works (Lambart et al., 2012; Pin et al., 2022; Shaw et al., 2018). After each increment of melt impregnation, the infiltrated melt will eventually be equilibrated with surrounding peridotite, and the thermodynamic properties of the whole system are adopted as a reference for the next increment. In each increment, the proportions and compositions of melt and solid phases after chemical re-equilibrium can be obtained by finding the minimum Gibbs energy of the whole system (Pin et al., 2022; Yao et al., 2018). This process is modeled by constantly adding up to 200 g of the lower basaltic melt by increments of 0.5 g, to 100 g of the upper peridotite, using the pMELTS of alphaMELTS (Ghiorso et al., 2002; Ghiorso & Sack, 1995; Smith & Asimow, 2005) in isenthalpic mode at 1.5 GPa, 1,250-1,300 °C and  $\Delta QFM$ -2.0. Here, the incremental addition of basaltic melt corresponds to the increase of melt: peridotite ratio from 0.005 to 2.

Our simulation shows that the mass fraction of olivine decreases from ~ 49 wt.% to ~ 22 wt.% with the increasing melt: peridotite ratio from 0.005 to 2 (Figure 8h). The silicate melt becomes saturated with orthopyroxene when the melt: peridotite ratio reaches ~ 0.22. After this point, the mass fraction of new orthopyroxene precipitated from silicate melt quickly increases to ~ 32 wt.% at the melt: peridotite ratio of 2 (Figure 8h). The mass fraction of clinopyroxene in peridotite has a quick increase from ~ 51 wt.% to ~ 57 wt.% before the precipitation of orthopyroxene, but then slowly decreases to ~ 46 wt.% at melt: peridotite ratio = 2 (Figure 8h). Hence, the increase in orthopyroxene proportion is mostly due to the dissolution of olivine,

accompanied by a limited contribution from the consumption of less than  $\sim 5$  wt.% clinopyroxene, which approximatively coincides with the hypothetical Reaction 1 mentioned above.

Concurrently, the melt penetrates upwards and further reacts with the peridotite to result in systematic variations of the compositions of olivine and Cpx grains towards the melt-rock interface (Figure 6 and Figure 7). The modeled compositional evolutions of olivine show a decreasing trend of MgO content from  $\sim 49.3$  wt.% to  $\sim 44.1$  wt.%, an increase of FeO content from  $\sim 9.0$  to  $\sim 14.5$  wt.%, and the associated decrease of Mg# from  $\sim 90.7$  to  $\sim 84.4$ , when the melt: peridotite ratio increases from  $\sim 0.005$  to  $\sim 1.5$  (Figure 6h). In the same range of melt: peridotite ratio, the clinopyroxene has an obvious increase in its MgO content (from  $\sim 17.5$  to  $\sim 20.0$  wt.%) and approximately constant  $\text{Al}_2\text{O}_3$  content (Figure 7h). The trends of these compositional variations of minerals are nearly in agreement with those of our experiments, except for the low-pressure experiment (PC559, Figure 6h and Figure 7h). Therefore, the measured compositional variations of olivine and clinopyroxene as a function of distance away from the final melt-peridotite interface may be connected to the melt: peridotite ratio during the reaction. The measured compositions of olivine and Cpx grains just above the final melt-peridotite interface are approximately consistent with those results simulated thermodynamically (Figures 6h and 7h-gray shadow regions) in the melt-peridotite ratio of 1.0 to 1.5, which implies that they have obtained locally chemical re-equilibrium with the reacting melt that is about  $\sim 1.0$ - $1.5$  times the initial mass of peridotite in the same region. As the distance away from the final melt-peridotite interface grows, the growing MgO and Mg# contents of olivine reflect the decrease of melt: peridotite ratio, and consequently imply the gradual weakening of melt-peridotite reaction. Therefore, when the upward-flowing melt encounters the partially molten peridotite that initially has a low permeability, melt convection, backflow, and the associated high-degree melt-peridotite reaction mostly occur at the bottom of peridotite, leading to an increase of porosity of the reaction region and more motivation of upward melt flow. This potentially produces positive feedback between the reaction and melt flow, contributing to upward infiltration of more silicate melt into the partially molten peridotite.

Along with the increasing melt: peridotite ratio from  $\sim 1.0$  to  $\sim 1.5$ , the modeled silicate melt that is reacted with the peridotite shows a decrease of  $\text{SiO}_2$  content from  $\sim 47.1$  wt.% to  $\sim 46.7$  wt.%, which is partly less than the measured ranges in our experiments with long annealing

time (Figure 8a). The  $\text{Al}_2\text{O}_3$  content of silicate melt varies from  $\sim 13.6$  wt.% to  $\sim 13.1$  wt.%, while the MgO content of melt shows a narrow variation range of  $\sim 11.1$  wt.% to  $\sim 11.4$  wt.%, which both roughly match our experiments (Figures 8b and 8e). The possible iron loss to Pt capsule in these experiments due to the incomplete isolation of separate upper and lower graphite capsules with different diameters for silicate melt from Pt capsule (J. Wang et al., 2020) is likely to be the main reason for the lower FeO content (Figure 8c) measured in these experiments with long annealing time ( $> 48$  h). The model suggests that silicate melt contains  $\sim 7.6$ - $7.9$  wt.% CaO at the melt-peridotite of 1.0-1.5, but this range is underestimated in contrast to our measurements (Figure 8d), partly because of the slight differences in the Gibbs free energy among various compositional models of pyroxenes (Yao et al., 2021) may drive large errors in the CaO content of clinopyroxene in the AlphaMELTS. Although there still have some weaknesses in modeling the compositional evolution of melt, the use of AlphaMELTS is encouraging here and helps to examine how mineral modes and chemical compositions evolve as silicate melts are added into the upper peridotite part.

#### 4.3 Coarsening processes of sulfide droplets

In a partially molten system, the nucleation/growth rate and associated size distribution of sulfide droplets are vital constraints on their dynamics (Yoshino & Watson, 2005). Along with the increasing annealing time, the size of sulfide droplets shows a larger increase (from  $\sim 3.2 \pm 0.1$   $\mu\text{m}$  to  $\sim 11.1 \pm 0.4$   $\mu\text{m}$ ) (Figure 5c) compared with that of the surrounding silicate minerals (from  $\sim 8.1 \pm 3.4$   $\mu\text{m}$  to  $\sim 11.0 \pm 4.9$   $\mu\text{m}$ ) (Figure 4b). The growth of silicate minerals has been suggested to be mostly controlled by the melt-enhanced diffusion of material, but sulfur that is the one of major elements in the sulfide has a much lower diffusivity than those of other elements (Freda et al., 2005; Yao & Mungall, 2021). Given that the diffusion of S rather than Fe limits the growth rate of sulfide droplets (Zhang, 2015), slow growth is expected for sulfide if its growth completely occurred via element diffusion. However, the estimated growth rate of sulfide droplets is larger than those of silicate minerals in our measurements, suggesting that there must be other mechanisms, besides the diffusion-driven growth, contributing to the fast growth of sulfide droplets.

On the other hand, the actual 3D size distribution of sulfide droplets in the melt source shows a uniform, nearly log-normal size distribution with a skewed peak to a slightly smaller

size than the average one (Figure 9), which can be attributed to the possible contribution of the Lifshitz-Slyozov-Wagner (LSW) coarsening and the occurrence of normal grain growth (Honour et al., 2019; Yoshino & Watson, 2005). Based on the classical LSW theory (Lifshitz & Slyozov, 1961; Wagner, 1961), the kinetics of crystal growth can be approximately described as  $r \sim t^{1/m}$ , where  $r$  is the grain size, with the exponent  $m = 3$  for a diffusion-controlled process and  $m = 2$  under the control of an interface reaction. At 1,250 °C and 1.5 GPa, growth rates of sulfide droplets are slightly slower than that predicted ( $m = 2$ ) by the interface-reaction-controlled process via a transient regime ripening before static-state ripening (Figure 5a), implying the likely combination of multiple coarsening processes here (Lautze et al., 2011; Sun, 2007). High run temperature (1,300 °C) can contribute to a higher degree of interface-controlled reaction (Yoshino & Watson, 2005), which potentially facilitates the coarsening processes of sulfide in the melt source (Figure 5a). A dramatical decrease in the number density of sulfides (Figure 5a-inset) and the presence of sulfide droplets in contact with each other (Figure 2b and Figure S1), both imply that the contribution of mechanical coalescence of sulfides (agglomeration and successive coalescence) to their coarsening processes may be more important at droplet contacts (Yoshino & Watson, 2005). Moreover, the coalescence of sulfide droplets is enhanced at lower viscous melts (Holzheid, 2010), potentially addressing a significant increase in growth rate (Figure 5a) from the carbonate-bearing experiment (PC560) with a lower melt viscosity (Figure 8g) (Di Genova et al., 2014). Although the previous analog experiment (de Bremond d'Ars et al., 2001) and static high-temperature, high-pressure experiments (Bockrath et al., 2004; Holzheid, 2010) did not observe the aggregation of sulfide droplets, their mechanical coalescence, for the first time, has been demonstrated to be feasible in our experiments with strong melt convection and melt-peridotite reaction. This mechanical coalescence may occur via the impaction of sulfide droplets with each other, and further drainage and rupture of the infinitesimally liquid film separating them in the reacting melt flow, as suggested by theoretical analyses (Yao et al., 2019).

## 5 Implications

### 5.1 Porous melt flow in the partially molten peridotite

These dissolution-reprecipitation processes not only change the composition and proportion of solid phases, but also increase the porosity of peridotite in the upper part of samples, thereby enhancing upward mobilities of silicate melt and sulfide droplets away from the

lower melt source (Figure 4 and Figure 5). Moreover, the porosity of ORL is enough high for the entrainment of sulfide droplets into peridotite (Figure 1 and Figure 2). In addition, with increasing annealing time to 72 h, more basaltic melt infiltrating into the peridotite region from the melt source (Figure 4) could potentially enhance the reaction between basaltic melt and peridotite. This melt-rock reaction may produce a high flow capacity of reacting melt through a dissolvable peridotite (Chadam et al., 1986; Jackson et al., 2018), resulting in the formation of the finger- or tree-like melt-rich channels in which a fast-ascending magma flow can even drive the antigravitational migration of coarse olivine grains (Pec et al., 2017). Hence, the development of positive feedback between the enhanced permeability and the associated increase of melt flux in reacting regions produces a higher flow velocity of melt and the wider melt channels among crystal framework (Aharonov, 1995; Pec et al., 2017), thereby potentially facilitating the upward entrainment of sulfide droplets (Wang & Jin, 2020).

To quantitatively outline the upward physical migration of sulfide droplets along with the reactive melt flows in the dissolvable peridotite, some theoretical parameters of physical properties of the partially molten peridotite should be considered first (*e.g.*, von Bagen & Waff, 1986; McKenzie, 1984). For an ideal porous crystal framework, the permeability ( $k$ ) can be cast by using a power-law relationship between grain size ( $d$ ) and porosity ( $\Phi$ ) of the crystal matrix (von Bagen & Waff, 1986):

$$k = \frac{\Phi^n d^2}{C} \quad (2)$$

where  $C$  and  $n$  are constant parameters depending on the topology of the melt phase and the geometry of an individual melt channel. Based on the previous works (Miller et al., 2014; Z. Wang et al., 2020), some suitable ranges of  $C$  (36-94) and  $n$  (2.4-2.8) are used for our experiments that contain the multiphase-mineral assemblage.

On the basis of the numerical model from McKenzie (1989), the extraction velocity of melt relative to the stationary crystal framework can be estimated by:

$$v = \frac{k \Delta \rho_m g}{\Phi \mu_m}$$

(3)



where  $\Delta\rho_m$  - the density contrast between silicate melt and solid mineral,  $g$  - the gravitational acceleration, and  $\mu_m$  - the melt viscosity, which decreases with increasing temperature and volatile content (mainly CO<sub>2</sub> and H<sub>2</sub>O) (Dingwell, 1996; Di Genova et al., 2014). Based on equation (2) and (3), the porous flow velocity of melt through grain-scale percolation is  $\sim 0.74$   $\mu\text{m/h}$  in the partially molten peridotite of experiment PC545 ( $C \sim 36$ ,  $n \sim 2.6$ , Miller et al., 2014;  $\Phi \sim 7.13$  %;  $\mu_m \sim 2.2$  Pa·s;  $d \sim 11.0$   $\mu\text{m}$ ;  $\Delta\rho_m \sim 600$  kg/m<sup>3</sup>,  $g \sim 9.8$  m/s<sup>2</sup>). Apparently, the transport distance ( $\sim 53.3$   $\mu\text{m}$ ) of this slow-extracting ( $\sim 0.74$   $\mu\text{m/h}$ ) melt within the longest experimental time (72 h) is much shorter than the region of melt-peridotite reaction observed here (Figure 1a), which implies that other mechanisms of melt extraction occur concurrently to drive a higher flow velocity. In addition to the melt junctions among minerals, we also observed thick melt channels (Figure 2) between some opening mineral grain boundaries in the ORL and olivine + melt layer where a new crystal framework may have not been built, indicating the possible existence of melt channelized flow, which has been demonstrated by some experimental studies (e.g., Pec et al., 2017; Zhu et al., 2011) and the field investigations on dunite channels in MORB (e.g., Kelemen et al., 1995). Melt velocity in these channels is suggested to be much higher relative to other regions (Pec et al., 2017; Wang & Jin, 2020).

We then extrapolate our experimental results to the partially molten mantle. Generally, the porosity ( $\Phi$ ) is approximately estimated to be  $\sim 1 - 15$  % in the partially molten peridotite (e.g., Mei et al., 2002; Yoshino et al., 2010; Zhu et al., 2011), and the regions having high-degree melt-rock reaction can obtain a higher porosity (Pec et al., 2017). Hence, the  $\Phi$  in our model of the molten mantle is properly set as  $3 - 20$  %. The observed grain size of upper mantle peridotite is about mm-scale (Ave Lallemant et al., 1980; Karato, 1984), and thus the crystal radius in our model is set as  $\sim 1.5$  mm, which is about two orders of magnitude higher than the grain size ( $\sim 10$   $\mu\text{m}$ ) in the experiments here. Therefore, the permeability of upper mantle peridotite should be much better than the estimated situation of our experiments, and consequently, the extraction velocity of melt relative to stationary crystal framework in the residual mantle increases to  $\sim 4.78 \times 10^{-6} - 5.62 \times 10^{-5}$  m/s ( $\sim 150.81 - 1,771.06$  m/year) with the growth of  $\Phi$  from 3 to 20 %, if other parameters in equation (3) remain constant. On the other hand, the melt transport velocities within the mantle were estimated to have a wide range of 1 to 1,000 m/year by analyses of uranium-series (U-series) disequilibria in lavas and Icelandic deglaciation (Katz et al., 2022; Rees Jones & Rudge, 2020; Rubin et al., 2005). The channelized melt flow potentially



has a higher melt extraction velocity, which can even reach  $\sim 10,000$  m/year (McKenzie, 2000). In contrast, recent laboratory experiments further estimate some conservative velocities of the porous melt flow ( $\sim 2.4 - 29.7$  m/year, Z. Wang et al., 2020;  $\sim 2 - 150$  m/year, Connolly et al., 2009;  $\sim 0.5$  m/year, Zhu et al., 2011). The estimated melt velocity ( $\sim 150.81 - 1,771.06$  m/year) here is higher than the previous experimental results ( $\sim 0.5 - 150$  m/year), but falls into the potential range of melt flowing velocity ( $\sim 1 - 10,000$  m/year) in the natural mantle. The wide range of melt flowing velocity can be attributed to the heterogeneous grain size, various mineral phases, and associated heterogeneous melt flow in natural peridotite relative to these in experimental conditions.

The microstructure of a porous crystal framework is usually simplified as a “throat and chamber” model, where the interstitial chamber pores are connected by many narrow throat pores. The size distribution of these throat pores will have a great effect on the permeability of the crystal framework, which in turn determines the percolation of silicate melt and the migration of sulfide. Previous numerical models had demonstrated that the ratio of mean throat pore radius to mineral radius decreases from 0.24 to 0.09 along with the decreasing porosity of crystal framework from 35.4% to 8.3% (Sweeney & Martin, 2003). Based on this modeling result, we can estimate the variations in the mean throat pore radius of the partially molten peridotite with constant crystal size ( $\sim 1.5$  mm) and variable porosity due to the melt-rock reaction, which can help us to trace the dynamic migration of sulfide in crystal framework.

## 5.2 Upward entrainment of fine sulfide droplets by porous melt flow

The physical mobility of sulfide liquids along with a sulfur-saturated silicate melt within the crystal framework of partially molten peridotite largely depends on the size of sulfide droplets, the permeability of crystal framework, the connectivity of melt network, and the processing degree of melt-rock reaction (Z. Wang et al., 2020; Yao & Mungall, 2020). Chung & Mungall (2009) simplified the downward sinking of sulfide droplets in the crystal framework, and proposed that this process is mostly controlled by the relative sizes of sulfide droplets and surrounding crystals within partially molten peridotite. In this regard, for the upward transport of a sulfide droplet, its terminal settling velocity should be also less than the upward velocity of porous melt flow, although its size has been smaller than the size of the most constricted part of the melt channel or pore throat in the crystal framework. The terminal settling velocity of the

spherical droplet will achieve when its gravity-driven buoyancy force ( $F_B$ ) is balanced by the vertical drag force ( $F_D$ ).

The melt flows passing through sulfide droplets that are nearly stranded in the melt channel or pore throat, as a first approximation, is regarded as the flow of an incompressible fluid around a circular cylindrical particle (*i.e.*, the sulfide droplet) confined by two parallel flat plates (*i.e.*, the melt channel) (Figure 10a) (Lee & Fung, 1969). The drag force ( $F_D$ ) acting on this sulfide droplet can be expressed by a dimensionless coefficient ( $f_D$ ):

$$F_D = f_D \times 4\pi\mu_m UR \quad (4)$$

where  $R$  is half the distance between the two plates, *i.e.*, the radius of the melt channel, and  $U$  is the Stokes flow velocity of melt around sulfide droplet. Based on the modeling results from Lee & Fung (1969), the dimensionless coefficient,  $f_D$ , can be estimated via the ratio between  $R$  and  $r_s$  which is the radius of an initial undeformed sulfide droplet before entering the melt channel (Figure 10b). For simplification, here we assume that the spherical droplet has a radius ( $r_s$ ) equaling to or being slightly higher than  $R$ , and thus the corresponding values of  $f_D$  should exceed  $\sim 4.8$  (Figure 10b) (Lee & Fung, 1969).

On the other hand, dense sulfide droplet always has a trend of settling to the base of the capsule due to gravity, and the corresponding buoyancy force ( $F_B$ ) can be expressed as:

$$F_B = \frac{4}{3}\pi r_s^3 g \Delta\rho_s \quad (5)$$

where  $\Delta\rho_s = 1,700 \text{ kg/m}^3$  is the density difference between sulfide and silicate melt (Kress et al., 2008; Z. Wang et al., 2020). Hence, the velocity of melt flow that just successfully entrains sulfide droplets whose sizes are close to that of melt channel/pore throat can be calculated via:

$$U = \frac{F_D}{4\pi\mu_m R \cdot f_D} = \frac{F_B}{4\pi\mu_m R \cdot f_D} = \frac{r_s^3 g \Delta\rho_s}{3\mu_m R \cdot f_D} \quad (6)$$

Because the measured maximum radius of sulfide droplets in the upper peridotite part is  $\sim 9.5 \text{ }\mu\text{m}$ , the maximum value of  $U$  is estimated as  $\sim 170.9 \text{ }\mu\text{m/h}$  ( $\sim 1.5 \text{ m/year}$ ), when the  $R/r_s$  equals 1. This value can represent the localized velocity of channelized melt flow in the regions

of melt-rock reaction and is nearly three orders of magnitude higher than the velocity of melt porous flow ( $\sim 0.74 \mu\text{m/h}$ ) driven by compaction of crystal framework. As the melt flow velocity is changeable due to the complex morphology of interconnected melt channels and variable fluid dynamic environments in porous peridotite (McKenzie, 1989; Miller et al., 2014; Z. Wang et al., 2020; Zhu et al., 2011), our calculation ( $\sim 1.5 \text{ m/year}$ ) here is used as an order of magnitude estimate for the velocity of melt flow around sulfide droplet through narrow pore throat in these experiments.

When the porous melt flow in crystal mush has a small velocity, it may not have enough kinetic energy to overcome the gravitational settling of large sulfide droplets. Hence, in this case (the region of No. ① in Figure 11), the maximum size of the upward migrating sulfide droplet is mostly controlled by the low flow velocity when the droplet size is smaller than the pore throat at the high porosity of crystal mush. For the porous melt flow with a velocity of  $\sim 5.1 \times 10^{-6} \text{ m/s}$  ( $\sim 161 \text{ m/year}$ ), due to the balance between  $F_B$  and  $F_D$ , the corresponding sulfide droplet radius is about  $98.2 \mu\text{m}$ , which is less than the mean radius of pore throat in crystal mush with increasing  $\Phi$  from 3 to 20% (Figure 11). Here, the maximum radius ( $r_s$ ) of sulfide droplets that can pass through the crystal mush will remain at  $\sim 98.2 \mu\text{m}$  even though the permeability of crystal mush has been highly improved due to the increasing porosity. When the velocity of porous melt flow increases to  $\sim 11.6 \times 10^{-6} \text{ m/s}$  ( $\sim 366 \text{ m/year}$ ), the  $r_s$  is limited by the radius of the narrow pore throat at low porosity (about  $\sim 9\%$ ), but keeps a constant value ( $\sim 148.5 \mu\text{m}$ ) at a higher porosity (Figure 11). A further increase in the porous melt flow velocity (e.g.,  $\sim 2.1 \times 10^{-5} \text{ m/s}$ ,  $\sim 662 \text{ m/year}$ ) will further reduce the porosity range ( $> \sim 15\%$ ) of which the  $r_s$  is just related to the melt flow velocity (Figure 11).

Once the melt velocity increases to  $\sim 3.2 \times 10^{-5} \text{ m/s}$  ( $\sim 1,009 \text{ m/year}$ ), the value of  $r_s$  due to the balance between  $F_B$  and  $F_D$  will increase to  $\sim 246.7 \mu\text{m}$ , which is equal to the radius of pore throat at  $\Phi = 20\%$  (Figure 11). Hence, if the melt velocity is beyond this value, the motivation of porous melt flow is no longer the restriction on the upward transport of fine sulfide droplets whose sizes are less than the radius of pore throat in crystal mush with the porosity range of 3-20%. Of course, this moderate flow velocity cannot provide enough power to overcome the capillary forces preventing droplet deformation as the coarse sulfide droplets attempt to pass into pore throats narrower than themselves. Therefore, in this situation (the No. ② in Figure 11), the sizes of upward transporting sulfide droplets must be strictly no larger than

the pore throat size in the crystal mush. Because the mean size of the pore throat is just related to the grain size and porosity of molten peridotite based on the modeling results from Sweeney & Martin (2003), the value of  $r_s$  should be equal to the estimated pore throat size that increases from 99.7 to 246.7  $\mu\text{m}$  with the increasing porosity from 3 to 20%, which is irrelevant to the melt flow velocity. Thus, the No. ② situation is represented by a black line in Figure 11, although the corresponding flow velocity may vary from  $\sim 3.2 \times 10^{-5}$  m/s ( $\sim 1,009$  m/year) to a much higher value (e.g.,  $\sim 0.08$  m/s, as explained below). Hence, during the melt-peridotite reaction, the accompanying melt flow in the dissolved peridotite is energetic, and its flow velocity seems to be high enough to drive the upward transport of fine  $\mu\text{m}$ -scale sulfide droplets (Figure 10c-right side; Figure 11).

### 5.3 Potentially upward transport of coarse sulfide droplet

When a coarse sulfide that is sufficiently larger than the minimum constriction of the melt channel is gradually squeezed into the small “throat” part along with the ascending reactive melt flow, the sulfide droplet will invade this throat and be changed to a pear shape having a smaller radius of curvature (Figure 10c-left side). Obviously, the upward migration of this coarse sulfide droplet in the partially molten peridotite becomes complicated, and the referred dynamic process simplistically involves a competition between buoyancy, pressure gradient, viscous, capillary, and drag forces (Yao & Mungall, 2020). This squeezing process introduces the additional capillary pressure  $P_c$ , which is a measure of the pressure discontinuity existing at the interface of two immiscible phases (Chung & Mungall, 2009). The additional pressure imposed on this deformed sulfide droplet within the squeezing process equals the difference in capillary pressure between the top and bottom interfaces of the sulfide droplet and can be calculated by:

$$\Delta P_c = 2\gamma_{ms} \times \left( \frac{1}{r_t} - \frac{1}{r_d} \right) \quad (7)$$

where  $\gamma_{ms} = 0.21$  N/m is the melt-sulfide surface tension (Mungall et al., 2015),  $r_t$  is the radius of the upper spheric cap, and  $r_d$  is the radius of the lower part of the pear-shaped droplet (Figure 10c). The droplet cannot be forced through the pore throat unless this excess pressure  $\Delta P_c$  is balanced by an equal or greater pressure ( $\Delta P$ ) exerted by ascending melt flow on the trailing edge of sulfide droplet, which tends to push the droplet to rise through the melt channel.

Because the silicate melt flows at a low Mach number, the simple form of Bernoulli's principle is valid for the quasi-incompressible flows in melt-peridotite reaction, and hence pressure of flowing melt around a sulfide droplet can be highly related to the peripheral flow velocity via the Bernoulli's equation (Bauman & Schwaneberg, 1994):

$$P_1 + \frac{\rho_m U_1^2}{2} + \rho_m g h_1 = P_2 + \frac{\rho_m U_2^2}{2} + \rho_m g h_2$$

(8)

where the subscripts 1 and 2 correspond to the top and bottom outside points of deformed sulfide droplets along the axis of symmetry, respectively;  $P$  is the pressure caused by the flow of silicate melt;  $\rho_m$  is the density of melt;  $h$  is the height of 1 and 2 relative to a reference point beneath sulfide droplet. Hence, the flowing-melt-driven external pressure imposed on the deformed sulfide droplet equals the difference ( $\Delta P$ ) between  $P_1$  and  $P_2$ . In an extreme case, the coarse, deformed sulfide droplet completely blocks the ascending melt flow, and the flow velocity at bottom of droplets ( $U_2$ ) is close enough to zero, leading to a maximum value of external pressure difference ( $\Delta P$ ) that can be calculated as:

$$\Delta P = P_2 - P_1 = \frac{\rho_m U_1^2}{2} + \rho_m g (h_1 - h_2) = \frac{\rho_m U_1^2}{2} + \rho_m g \Delta h$$

(9)

where  $U_1$  is the flow velocity of melt at the top of the sulfide droplet, and  $\Delta h$  is the vertical distance between the top and bottom points of the droplet. Assuming that the volume of a sulfide droplet is unchanged during the squeezing process, the height of this invading sulfide droplet ( $\Delta h$ ) can be described by Chung & Mungall (2009):

$$\Delta h = r_t + r_d + \sqrt{r_d^2 - R^2} - \sqrt{r_t^2 - R^2}$$

(10)

For the deformed sulfide droplet, it bears the upward pressure driven by flow melt (*i.e.*,  $\Delta P$ ), which should balance or even exceed the total of downward capillary pressure ( $\Delta P_c$ ) and gravity-driven pressure ( $P_G$ ) for the droplet to rise through the constriction:

$$\Delta P \geq P_G + \Delta P_c = \rho_s g \Delta h + 2\gamma_{ms} \times \left( \frac{1}{r_t} - \frac{1}{r_d} \right)$$

(11)

where  $\rho_s$  is the density of sulfide liquid. Thus, combining the above equations (7-10), under the critical condition for migration of sulfide droplet into the pore throat to be the case ( $r_t = R$ ) (Chung & Mungall, 2009), the equation 11 substituted with the appropriate values for  $\rho_m = 2,600 \text{ kg/m}^3$  (Robertson et al., 2016) is simplified into:

$$1,300 \times U_1^2 \geq 16,660 \times \left( R + r_d + \sqrt{r_d^2 - R^2} \right) + 0.42 \times \left( \frac{1}{R} - \frac{1}{r_d} \right) \quad (12)$$

And finally, the initial radius ( $r_s$ ) of this sulfide droplet before its invasion into the pore throat can be calculated by Chung & Mungall (2009):

$$r_s = \left( 0.5 \times (r_d^3 + R^3 + (r_d^2 + \frac{1}{2}R^2) \times \sqrt{r_d^2 - R^2}) \right)^{1/3} \quad (13)$$

When the melt flows around a sulfide droplet through a narrow pore throat in the partially molten peridotite, the associated velocity field is suggested to be not homogeneous. In addition, it is well known that exactly deciphering the flow velocity field of silicate melt infiltrating through a multi-mineral phase system plays a critically important role in understanding the geodynamic properties of the partially molten upper mantle, but it is challenging and beyond the scope of this study. The increases in grain size and permeability may also strongly enhance the melt flow velocity in the new channel of the crystal framework. Here, the flow velocity in the channel has been confirmed to be about three orders of magnitude higher than that of porous flow through grain-scale percolation in our experiments, which may be extrapolated to the mm-scale upper mantle peridotite. Therefore, the velocity of melt flow around a sulfide droplet in the channel in upper mantle peridotite may be in the magnitude of 0.001-0.1 m/s, and thus the local velocity ( $U_l$ ) of focused melt flow around the top of a sulfide droplet was analogously assumed as  $\sim 0.3 \text{ m/s}$  and  $\sim 0.08 \text{ m/s}$  here (Figure 11- region ③) for the localized high-degree melt-rock reaction regions of the partially molten mantle. This range of flow velocity is also consistent with the estimates from previous studies discussed in Section 5.1.

Equations 12 and 13 are numerically solved for the above value of  $U_l$  by programming a loop using Maple<sup>®</sup>. Based on the above numerical calculation, when the velocity of focused melt flow around a sulfide droplet is less than  $\sim 0.08 \text{ m/s}$ , the sulfide droplet is difficult to overcome the capillary pressure imposed by melt channel/pore throat with a smaller size, and will be

761 stranded in inter-grain pores. However, these flow velocities ( $\sim 3.2 \times 10^{-5}$  m/s – 0.08 m/s) are  
762 still large enough to drive upward transport of sulfide droplets whose radii are smaller than or  
763 almost the same as the radius ( $R$ ) of pore throat (Figure 11-region ②).

764 For a powerful melt flow with extremely high velocity (e.g.,  $\sim 0.3$  m/s in Figure 11-  
765 region ③), coarse-grained sulfide droplet overcomes the capillary pressure driven by the pore  
766 throat, and its maximum radius for upward migration increases from  $\sim 106.8$  to  $\sim 276.8$   $\mu\text{m}$  with  
767 increasing the porosity from 3 % to 20 % (Figure 11-region ③), which is larger than the  
768 corresponding radius of the pore throat (from  $\sim 99.7$  to  $246.7$   $\mu\text{m}$ ). In this condition (region ③ in  
769 Figure 11), the coarse sulfide droplets can pass through the pore throat narrower than themselves.  
770 Although the estimates of the maximum sulfide size that are capable of upward migrating  
771 through the porous molten peridotite are divided into three conditions along with the increasing  
772 melt flow velocity (the regions ①-③ in Figure 11), all of them are mostly dominated by the melt  
773 flow velocity, porosity, and grain size of peridotite, which would be strongly enhanced by the  
774 high-degree melt-rock reaction.

775 On the other hand, based on these previous studies on the abyssal and orogenic  
776 peridotites, and peridotite xenoliths (e.g., Lorand & Luguet, 2016; Lorand et al., 2010; Luguet et  
777 al., 2003), highly variable grain size ( $< 20$   $\mu\text{m}$  to  $500$   $\mu\text{m}$ ) of sulfides is widely observed in the  
778 upper mantle, while most sulfides are likely smaller than about  $100$   $\mu\text{m}$  in diameter. Hence, it is  
779 plausible that almost all sulfide droplets are stranded among the crystal framework of the upper  
780 peridotite mantle with low porosity, while the high-degree partial melting and/or melt-peridotite  
781 reaction will lead to high porosity, large flow velocity, and even the new-forming melt channel,  
782 which highly proposes that efficient entrainment of most sulfide droplets into ascending magma  
783 flow among porous peridotite is physically possible, especially among channelized extraction of  
784 silicate melt in melt-peridotite reaction. Conceivably, mechanical entrainment of sulfide droplets  
785 during the melt-rock reaction would be more efficient than the transport of sulfide liquid by  
786 dissolving in departing silicate melt, especially given the slow S diffusion in the basaltic melt  
787 (Freda et al., 2005). The sub-continental lithospheric mantle that undergoes large-scale partial  
788 melting and melt-peridotite reaction under tectonically active craton margins may be favorable  
789 conjunction of these factors and potentially conducive to the entrainment of large sulfide droplets,  
790 thereby contributing to the fertilization of sub-continental lithospheric mantle and the primary  
791 enrichment of metal-bearing sulfides for the formation of magmatic sulfide deposits (e.g., Griffin



et al., 2013). This may be one of the reasons to account for the issue of why the spatial and temporal distribution of magmatic Ni-Cu-(PGE) sulfide deposits are genetically related to these regions (Maier & Groves, 2011).

## 6 Conclusions

Reaction-infiltration of silicate melt and concomitant transport of sulfide droplets in the partially molten peridotite are examined experimentally and thermodynamical-quantitatively in this study, which provides important insights into the fertilization of sub-continental lithospheric mantle and the efficient recycling of sulfur and metal elements from the partially molten mantle. The reaction between peridotite and basalt leads to the preferential dissolution of olivine and precipitation of orthopyroxene at high temperature (1,250-1,300 °C) and pressure (1.5 GPa), forming an orthopyroxene-rich reaction layer (ORL) with a high growth rate ( $\sim 0.78 \pm 0.05 \mu\text{m/s}^{0.5}$  at 1,250 °C) in the melt-rock interface. With increasing the degree of melt-peridotite reaction, more silicate melt infiltrates through the ORL into the upper partially molten peridotite, forming an olivine-melt layer above the ORL, and coarse sulfide droplets could also be entrained along with the infiltration of silicate melt. Systematic variations observed in the compositions of minerals and reactive melt are approximately consistent with the thermodynamically-constrained mixing model. Meanwhile, the sizes of sulfide droplets have fast growth in the melt source, which is partly attributed to the multiple coarsening processes (including the mechanical coalescence) of sulfide droplets in the dynamic reactive melt flow.

Suspension and entrainment of sulfide droplets in the partially molten peridotite indicate a fast-flowing velocity of reactive melt ( $\sim 170.9 \mu\text{m/h}$ ) in our experiments, which demonstrates the occurrence of focused melt flows with high velocity in the reactive-forming melt channels, besides the melt porous flow driven by the density difference between silicate melt and minerals forming the crystal framework. It is convincible that the melt flow velocity within the melt-rock reaction is potentially high enough to drive upward transport of fine sulfide droplets with smaller diameters than that of the pore throat in the partially molten peridotite. In this condition, the maximum size of sulfide for upward entrainment is limited by the radius of the pore throat in the crystal framework of peridotite and increases with the growth of porosity during partial melting and melt-peridotite reaction. Only in the channelized melt flow with extremely high velocity, coarse-grained sulfide droplets could upward migrate through the pore throats narrower than



themselves. Hence, the high-degree melt-rock reaction regions in the upper mantle, such as the sub-continental lithospheric mantle of craton margins, are likely to have wide pore throats, high porosity, and even channelized melt flow with high velocity, which can potentially drive upward entrainment and/or extrusion of sulfide droplets in the partially molten mantle, and thereby fertilize the lithospheric mantle and lead to the endowment of Cu- and Ni-bearing sulfide for the formation of associated deposits.

## Acknowledgments

We are grateful to two anonymous reviewers for their useful suggestions and comments. We thank Prof. Zhang Jinsen and Dr. Wenlong Liu for their help with EPMA and EBSD analysis. Especially, I appreciate my wife Shasha Guo for the support and help in my life, so that I could do my research without any worries. This research was supported by the National Natural Science Foundation of China (No. 42102057, 42272084, and 41902088), the Natural Science Foundation of Hebei Province (No. D2021402019), the fund from SinoProbe Laboratory (No. SinoProbe Lab 202219), and the Education Department Foundation of Hebei Province of China (BJ2020023).

## Data Availability Statement

Microprobe images used for the composition measurements and all data are shown in figures and the supplemental materials are freely available online (<https://doi.org/10.6084/m9.figshare.22352632>).

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**Figure 1.** Back-scattered electron images (BSE) of sections from experiments PC520 (a) (annealing time ~ 12 h) and PC545 (b) (annealing time ~ 72 h) under the conditions of 1.5 GPa and 1,250 °C. Dashed red and blue lines respectively denote the original and current interfaces between the melt source and peridotite before and after the melt-rock reaction. The dashed yellow line marks the top of the Opx-rich reaction layer. The position where the melt fraction at the top of the melt source decreases below 10 – 15 vol% is set as the possible lower boundary of ORL for the measurement (the dashed white line), which is approximately consistent with the original interface between the melt source and peridotite at short annealing time (< 48 h). The width of these quenched cracks needs to be subtracted. The specimen setup of these experiments

is shown as the inset in the left lower part. Mineral abbreviations: Ol-olivine; Opx-orthopyroxene; SM-silicate melt; SL-sulfide liquid.

**Figure 2.** Microstructures of the interface of melt-rock reaction in experiments PC545 (72 h, 1,250 °C) (a), PC528 (48 h, 1,250 °C) (b, c), and PC527 (12 h, 1,300 °C) (d). Detailed microstructure of the distribution of sulfide droplets and silicate melt in the partially molten peridotite is displayed in (c). Symbols are the same as in Figure 1. Mineral abbreviations: Ol-olivine; Opx-orthopyroxene; SM-silicate melt; SL-sulfide liquid. Note that some huge Opx grains are present in the Opx-rich reaction layer of experiment PC527 with high reaction temperature, potentially leading to the destruction of the layer.

**Figure 3.** Plots of thicknesses of Opx-rich reaction layers (ORL) (a) and dissolution distance (b) as a function of the square root of run time. The lines show the best fit for these data.  $k$  is the diffusive dissolution or growth rate which is the slope of the linear regression line. The grey symbols and lines denoted as 3 and 4 are from Morgan & Liang (2005), and those denoted as 1, 2, 5, and 6 are from C. Wang et al. (2020). The blue symbols and lines denoted as 7 are from this study. Note that the dissolution distance and the thickness of ORL may be slightly underestimated and overestimated, respectively, due to the possible gravitational settling of minerals in our experiments.

**Figure 4.** Plots of melt fraction (a) and grain size of silicate minerals (b) as a function of the square root of run time in the partially molten peridotite from our experiments.

**Figure 5.** Plots of grain size ( $\mu\text{m}$ ) of sulfide droplets weighted by the area (a and c) and the area fraction of sulfide droplets (b and d) as a function of the square root of run time in the melt source (lower part) and partially molten peridotite (upper part). Inset in (a) shows the relationship between the number density of sulfide droplets in the melt source and the square root of run time. All symbols representing samples are the same as those in (c).

**Figure 6.** Plots (a-g) of measured oxide abundance (in wt.%) and Mg# in olivine as a function of distance (in  $\mu\text{m}$ ) away from the melt-rock interface. The compositional variations of MgO, Mg#, and FeO are also simulated thermodynamically as a function of the melt-peridotite ratio in (h). These measured compositions of olivine approaching to the interface are roughly consistent with

those results simulated thermodynamically (the gray shadow region in h), when the melt-rock ratio ranges from 1 to 1.5 (Table S6). At low pressure (PC559, 0.5 GPa), the compositions of olivine recrystallized in the melt source during the melt-rock reaction are present as PC559-Lower (hollow black square). The compositions of olivine grains within the low-temperature hot-press experiment (PC537) were almost constant across the partially molten peridotite region and thus can be used as the reference line (purple dotted lines) to show the composition variations of olivine in those high-temperature reaction experiments. The gray shadow regions in (a)-(g) denote the region that has a vigorous melt-peridotite reaction, whose width is not in accordance with the actual distance.

**Figure 7.** Plots (a-g) of oxide abundance (in wt.%) and Mg# in clinopyroxene as a function of distance (in  $\mu\text{m}$ ) away from the melt-rock interface. The compositional variations of MgO,  $\text{Al}_2\text{O}_3$ , and CaO are also simulated thermodynamically as a function of the melt-peridotite ratio in (h). The legends are the same as in Figure 6.

**Figure 8.** Plots of oxide abundance (in wt.%) (a-e), sulfur content (f), and viscosity ( $\mu_m$ ) (g) of silicate melt versus the run time in melt-rock interface and lower melt source of these reaction experiments. The ranges of composition variations of  $\text{SiO}_2$ ,  $\text{Al}_2\text{O}_3$ , FeO, CaO, and MgO simulated thermodynamically are denoted as shadow gray regions in (a)-(e) when the melt-peridotite ratio varies from 1 to 1.5. The shadow gray region in (f) is the calculated SCSS based on the models from Chowdhury & Dasgupta (2020) and Smythe et al. (2017), and the SCSS value at  $t = 0$  is calculated by using the compositions of starting materials under conditions of 1,250 °C and 1.5 GPa. The composition of sulfides in starting materials (Table S3) was used to estimate the SCSS in experiments PC559 and PC560. Produced phase proportions of different silicate minerals during the melt-rock reaction are simulated thermodynamically as a function of melt-peridotite ratio in (h). A lower melt viscosity ( $\sim 0.25 \text{ Pa}\cdot\text{s}$ ) of the carbonate-bearing experiment (PC560) is roughly estimated by using a model from Di Genova et al. (2014) considering the effect of  $\text{CO}_3^{2-}$  ( $\sim 2 \text{ wt.}\%$ ) and denoted as a hollow inverted triangle in (g). Cpx-clinopyroxene, Opx-orthopyroxene.

**Figure 9.** Normalized 3D size distributions of sulfide droplets in melt source. The size of sulfide droplets ( $r$ ) and frequency ( $f$ ) have been normalized to average size ( $r^*$ ) and maximum frequency ( $f_{max}$ ), respectively.

**Figure 10.** Schematic map illustrating the fluid flow around a circular cylindrical particle confined by two parallel flat plates (a), the dimensionless coefficients  $f_D$  is shown as a function of the  $R/r_s$  ratio (b) (modified from Lee & Fung, 1969), and the extrusion and entrainment of large (left side) and small (right side) sulfide droplets driven by focused melt flow (c). In (a), the two parallel flat plates and a circular cylindrical particle are considered as the walls of the pore throat/melt channel and sulfide droplet, respectively.

**Figure 11.** Plots of the radius ( $r_s$ ) of initial undeformed sulfide droplets that could extrude the pore throat as a function of the porosity ( $\Phi$ ) of the partially molten mantle. The relationship between  $r_s$  and  $\Phi$  is divided into three parts according to the relative size of sulfide droplets' radius ( $r_s$ ) and the pore throat's minimum constriction ( $R$ ). The description of these three regions can be found in the main text.

**Table 1.** Experimental conditions and the observed lithologies.

Figure 1.



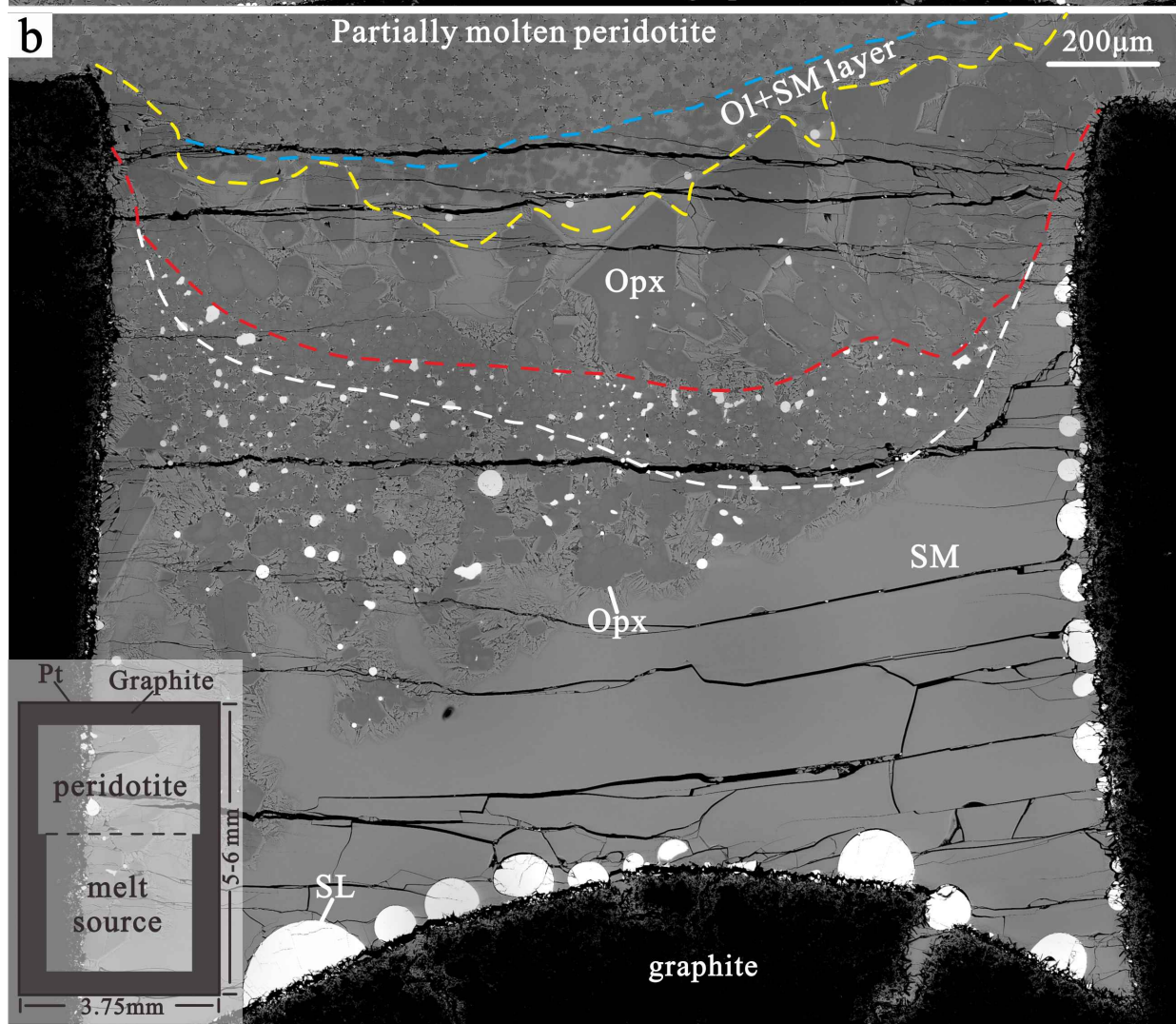
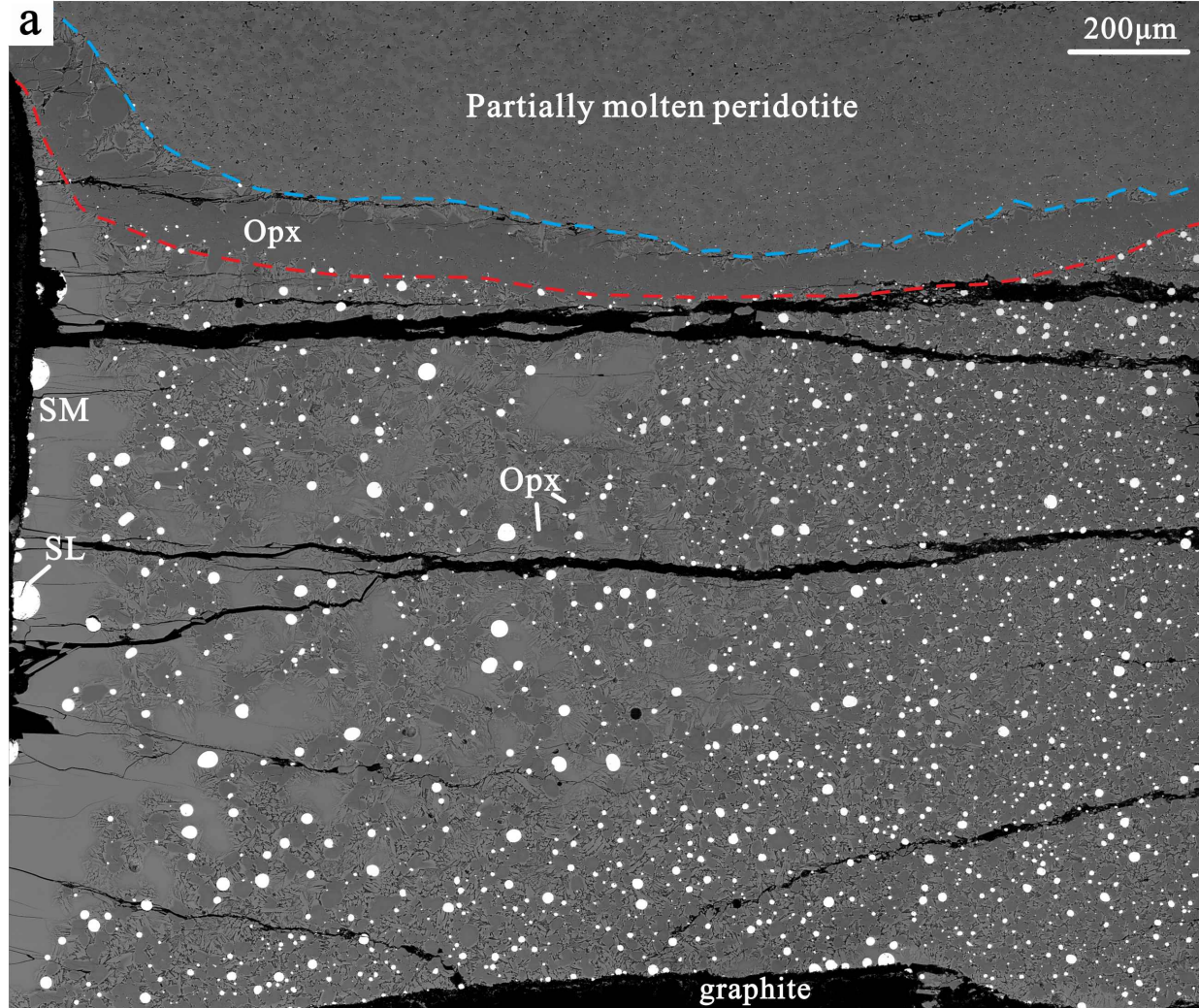


Figure 2.



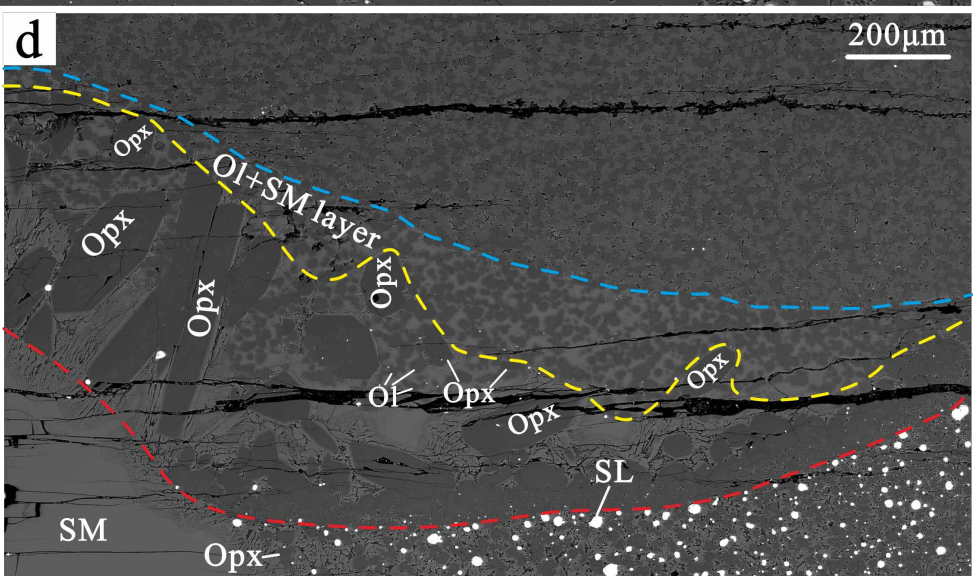
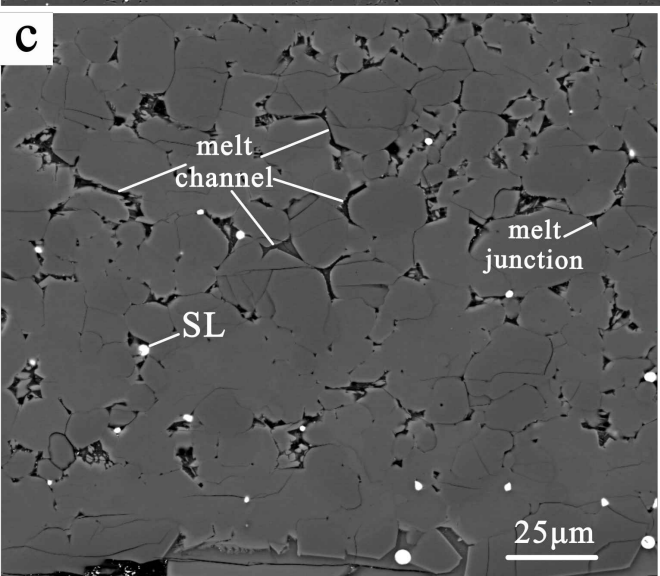
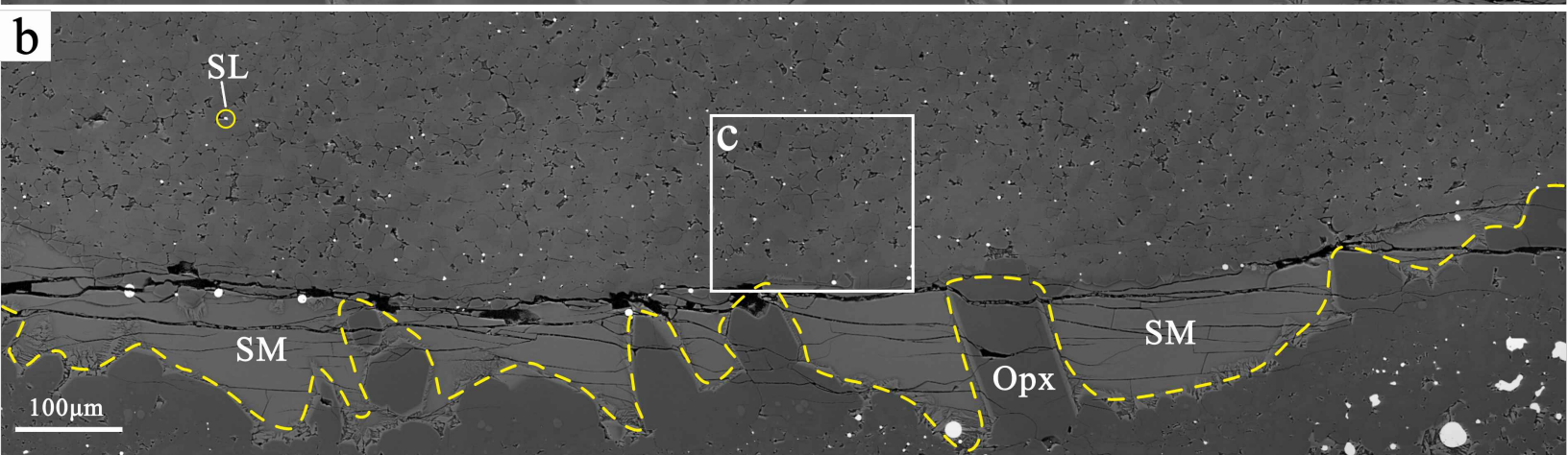
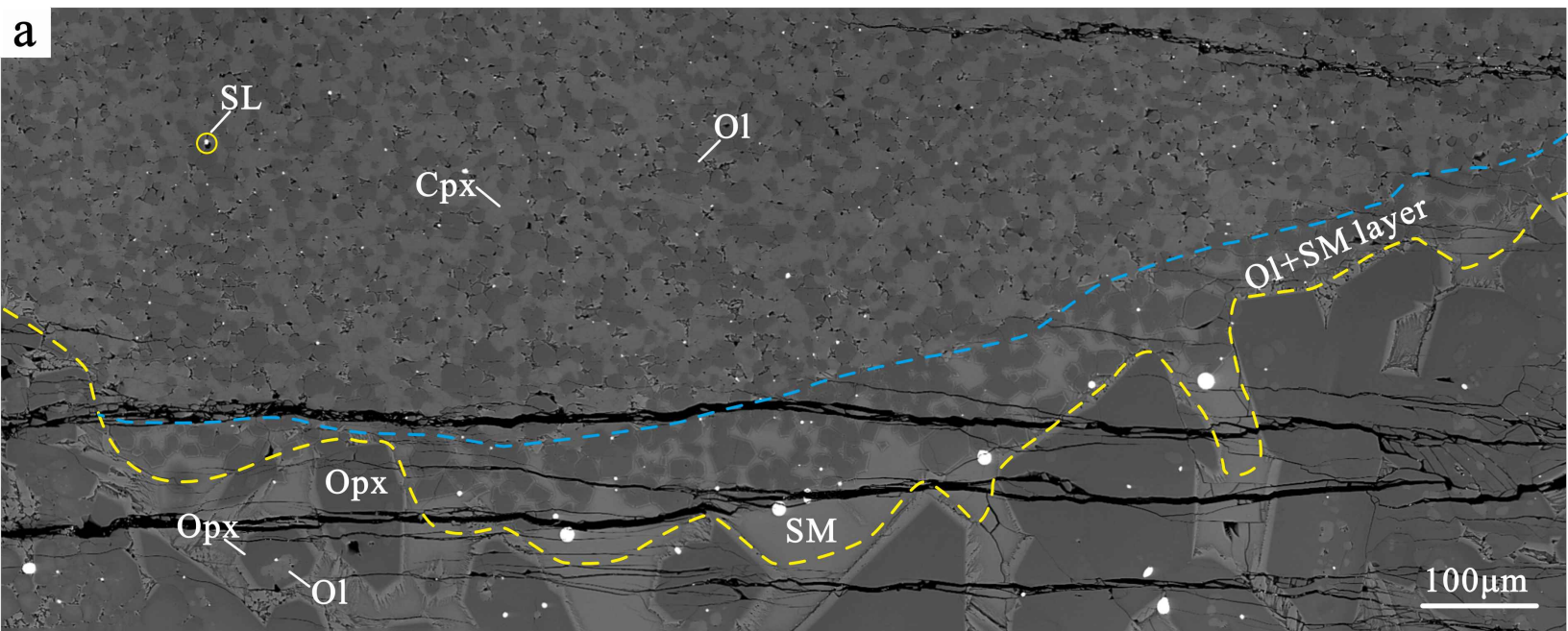


Figure 3.

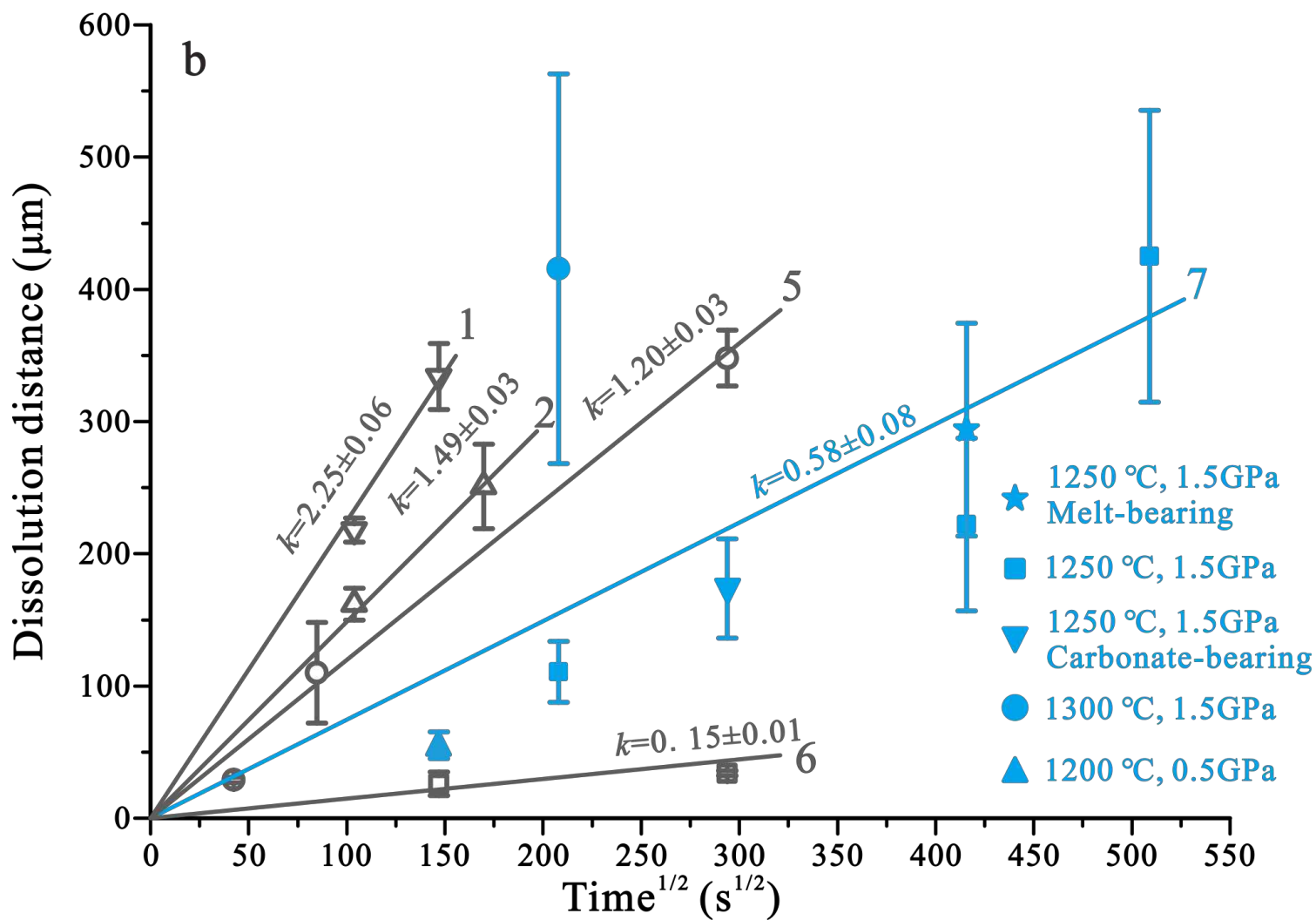
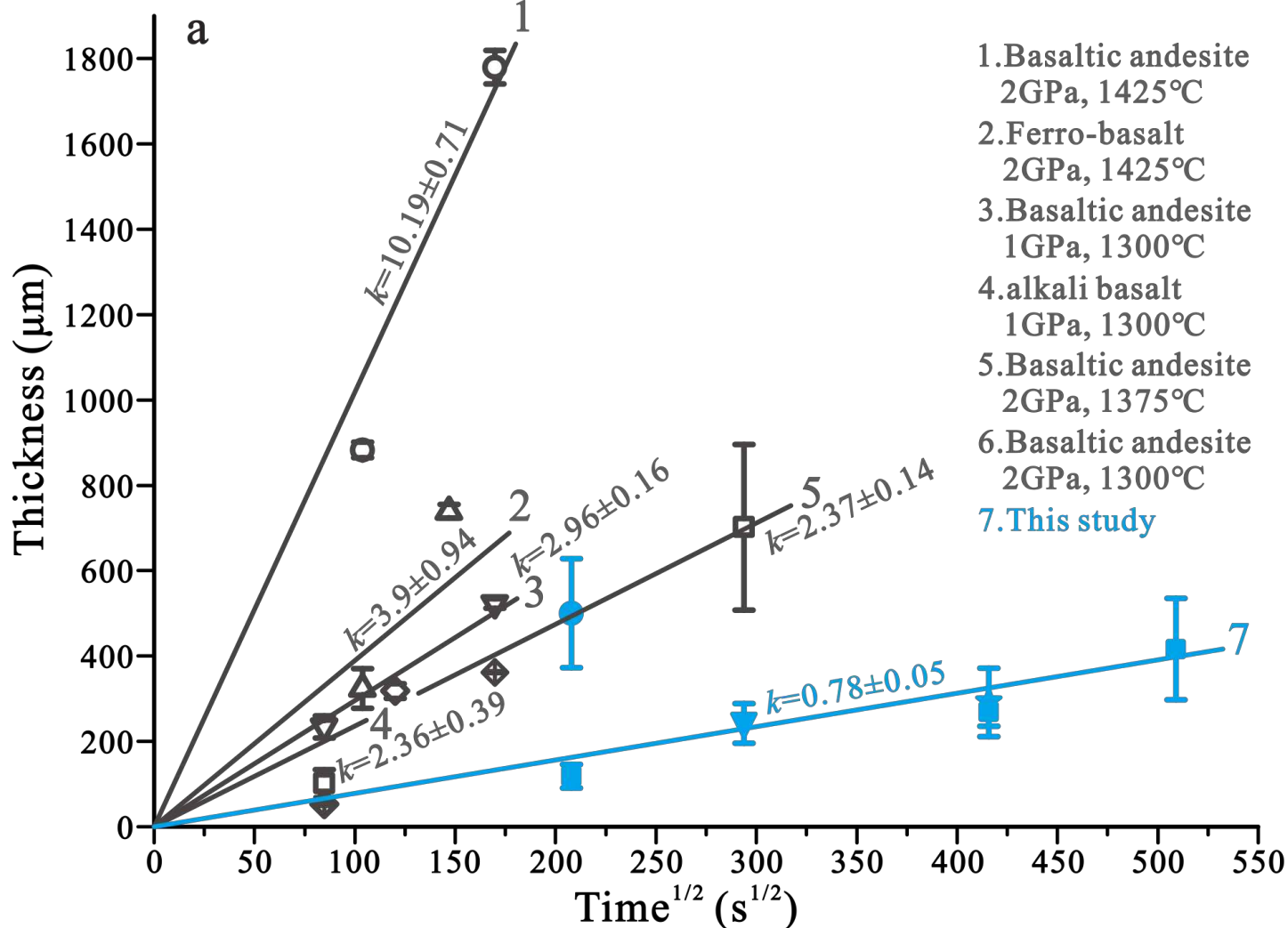


Figure 4.

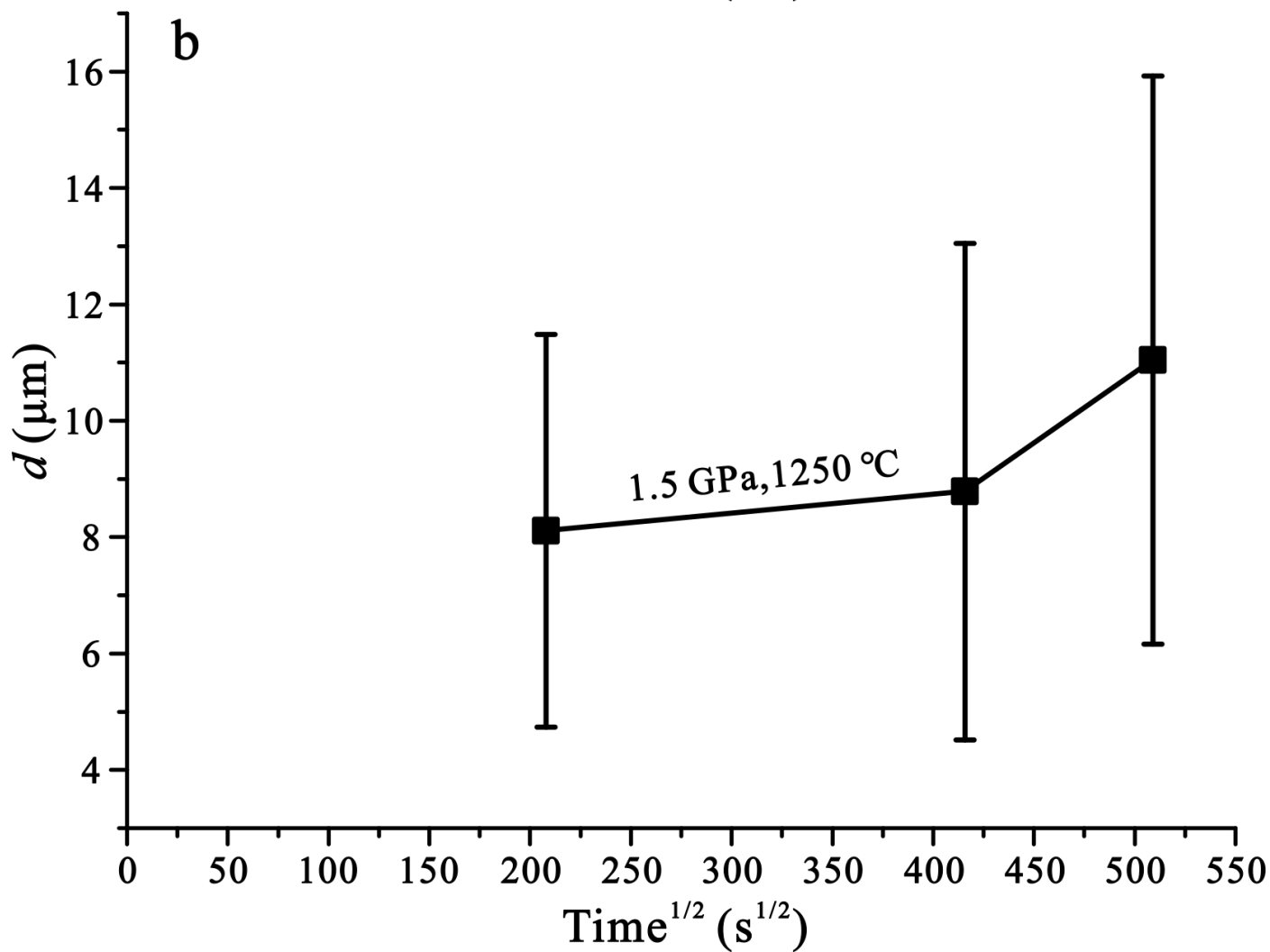
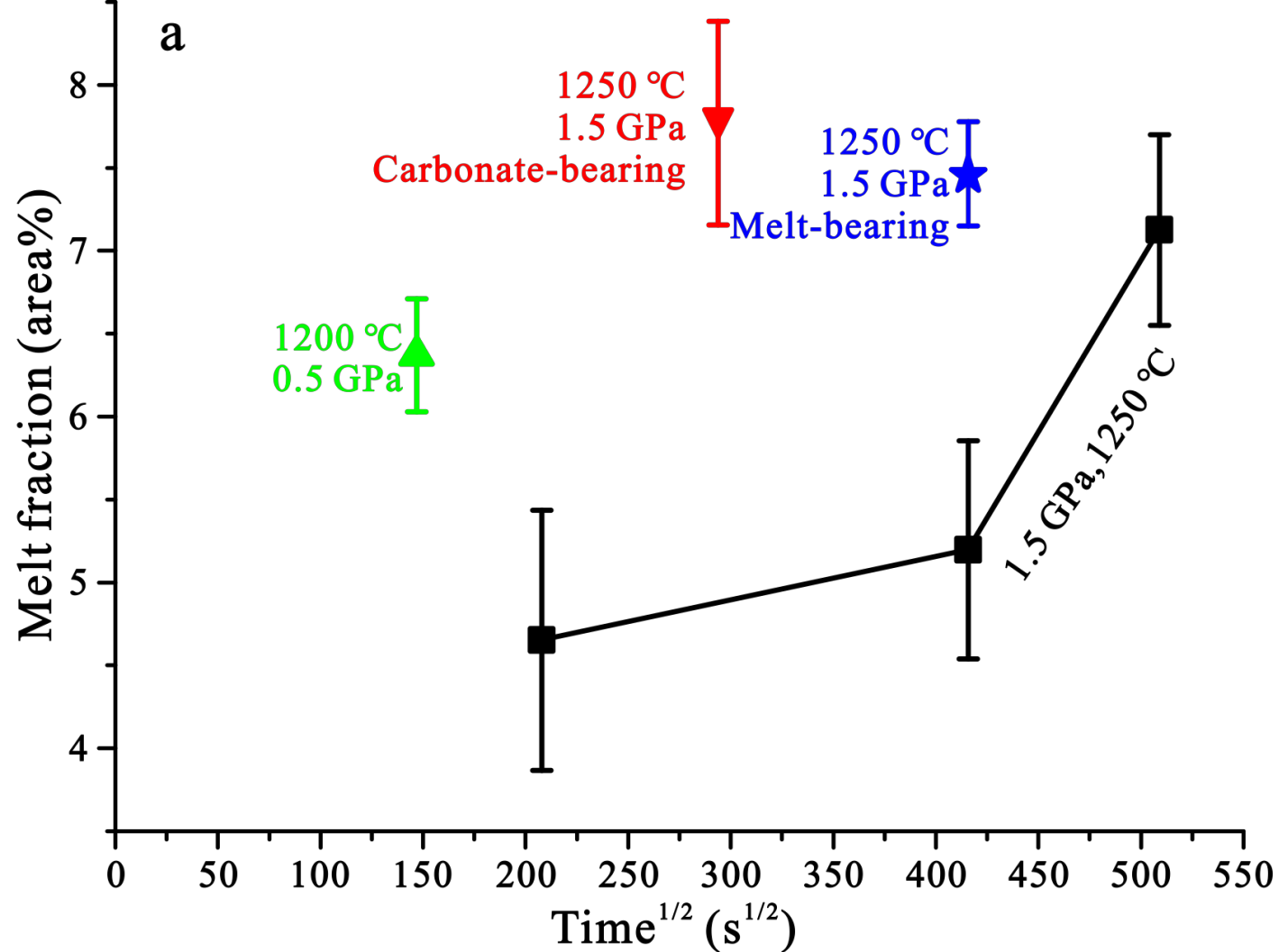


Figure 5.



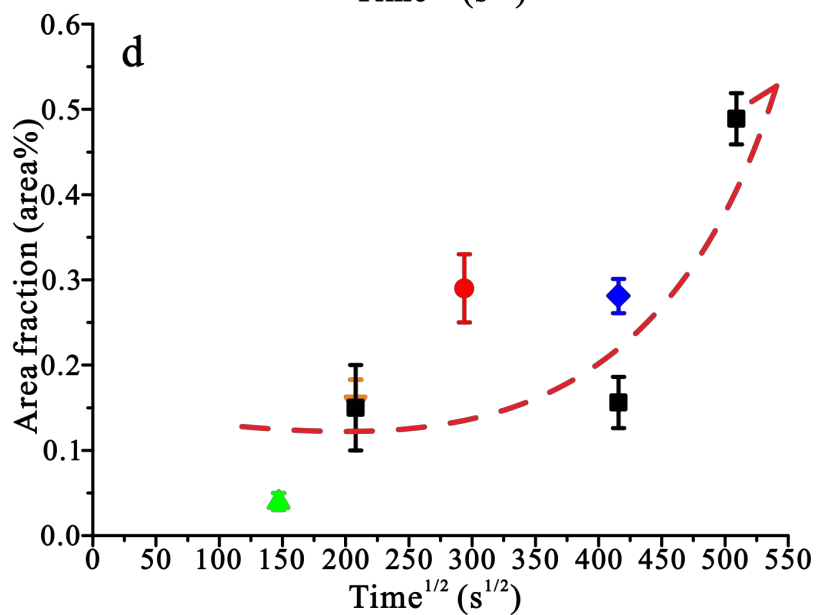
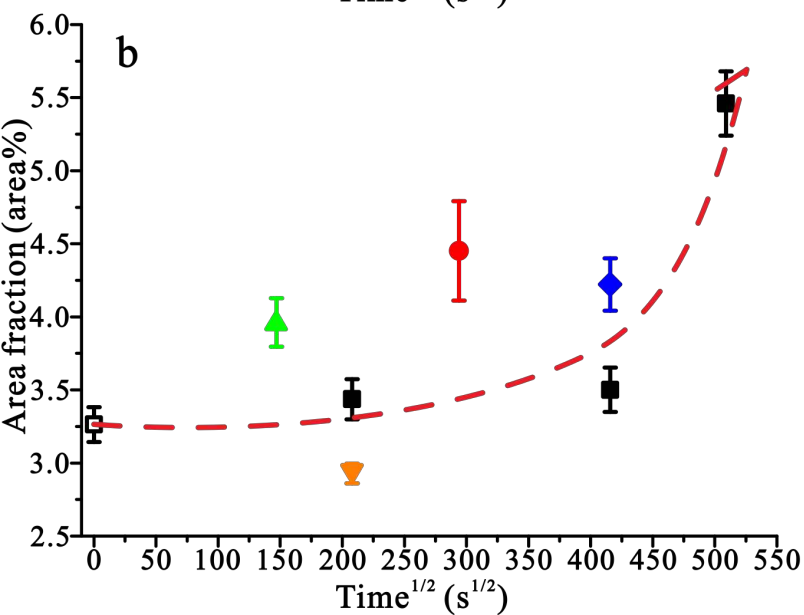
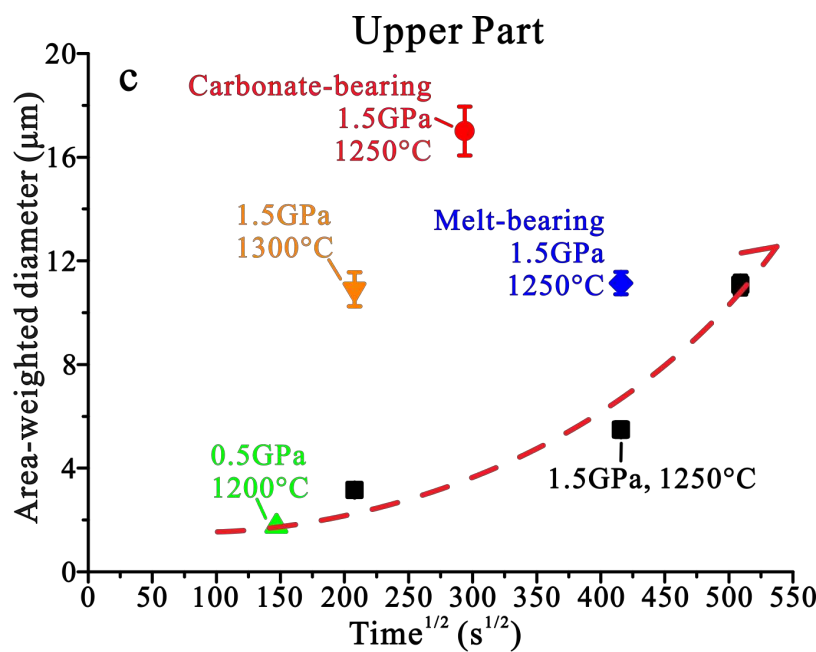
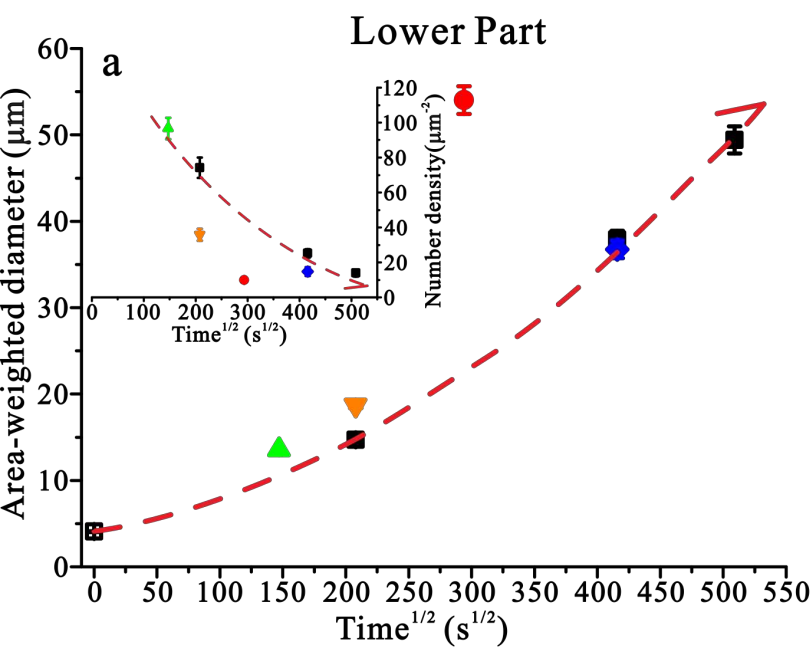


Figure 6.

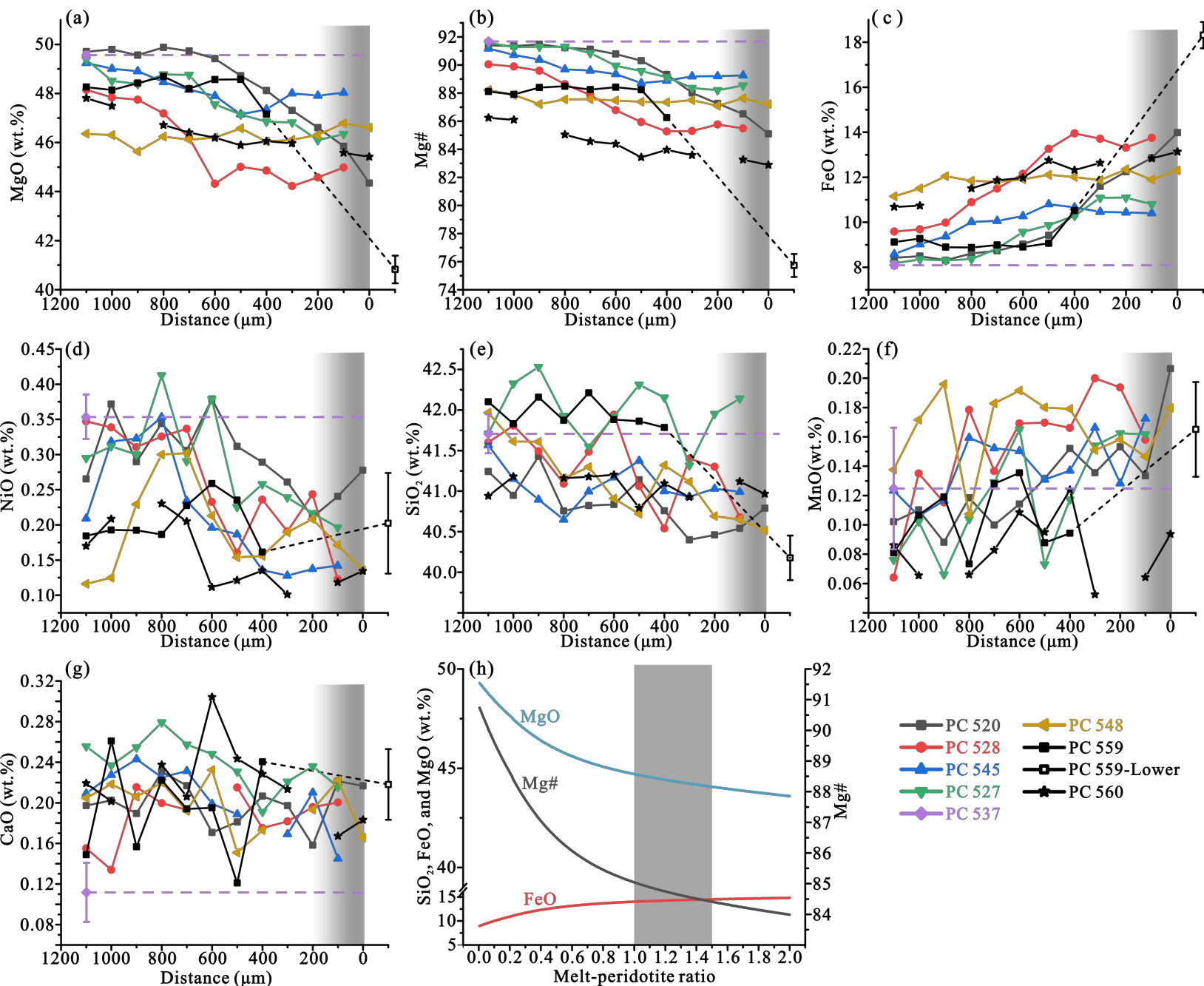


Figure 7.

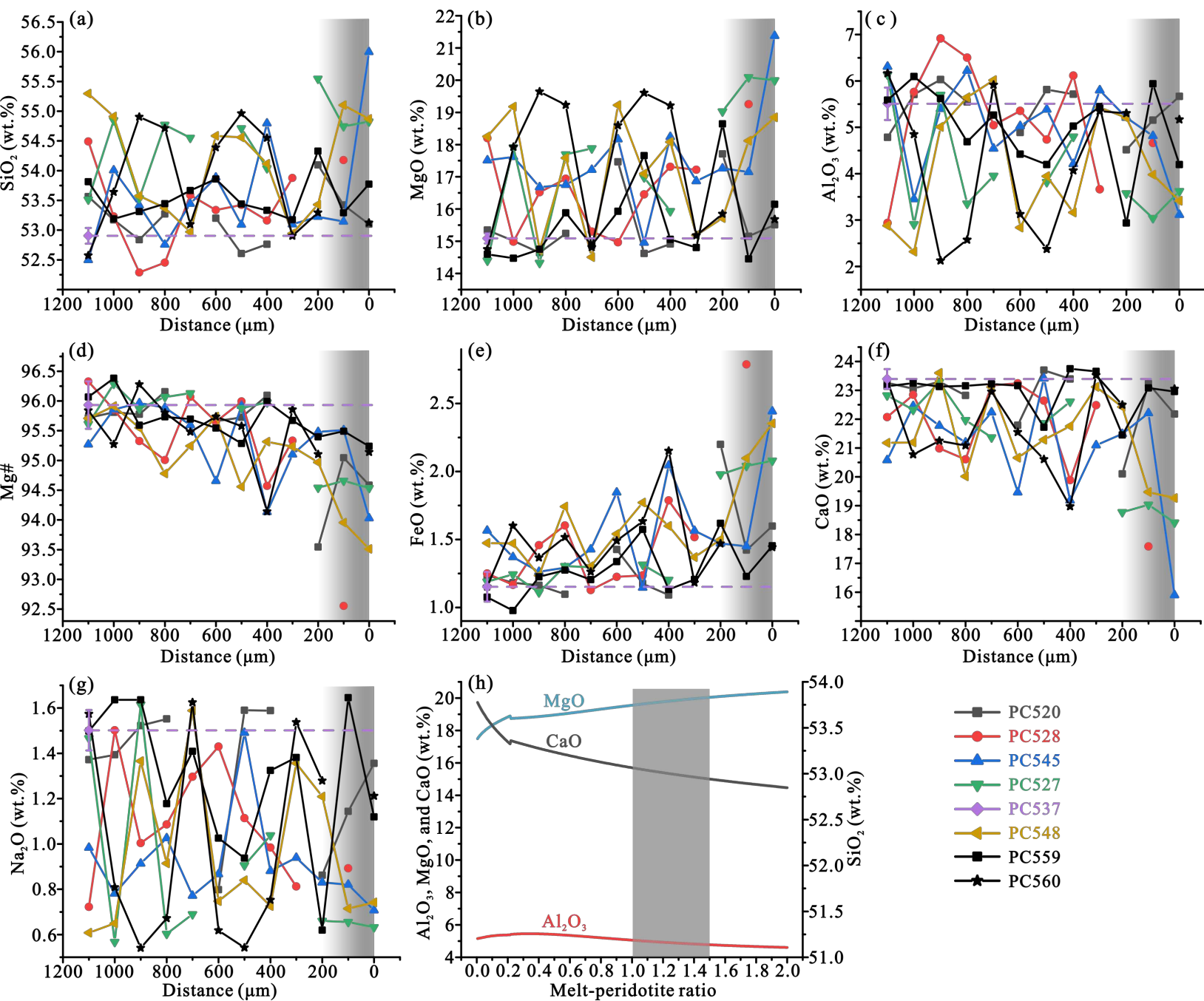


Figure 8.

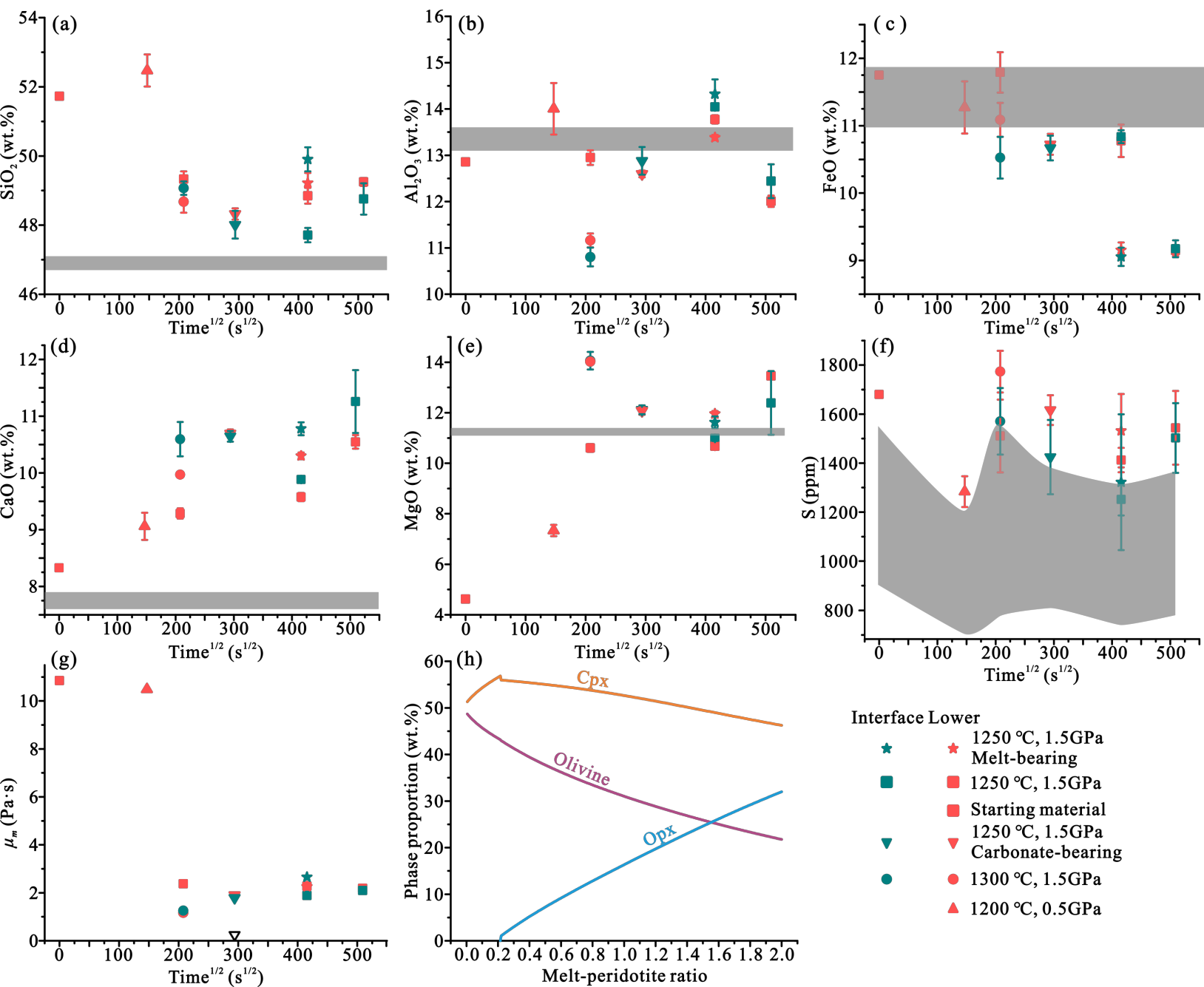


Figure 9.



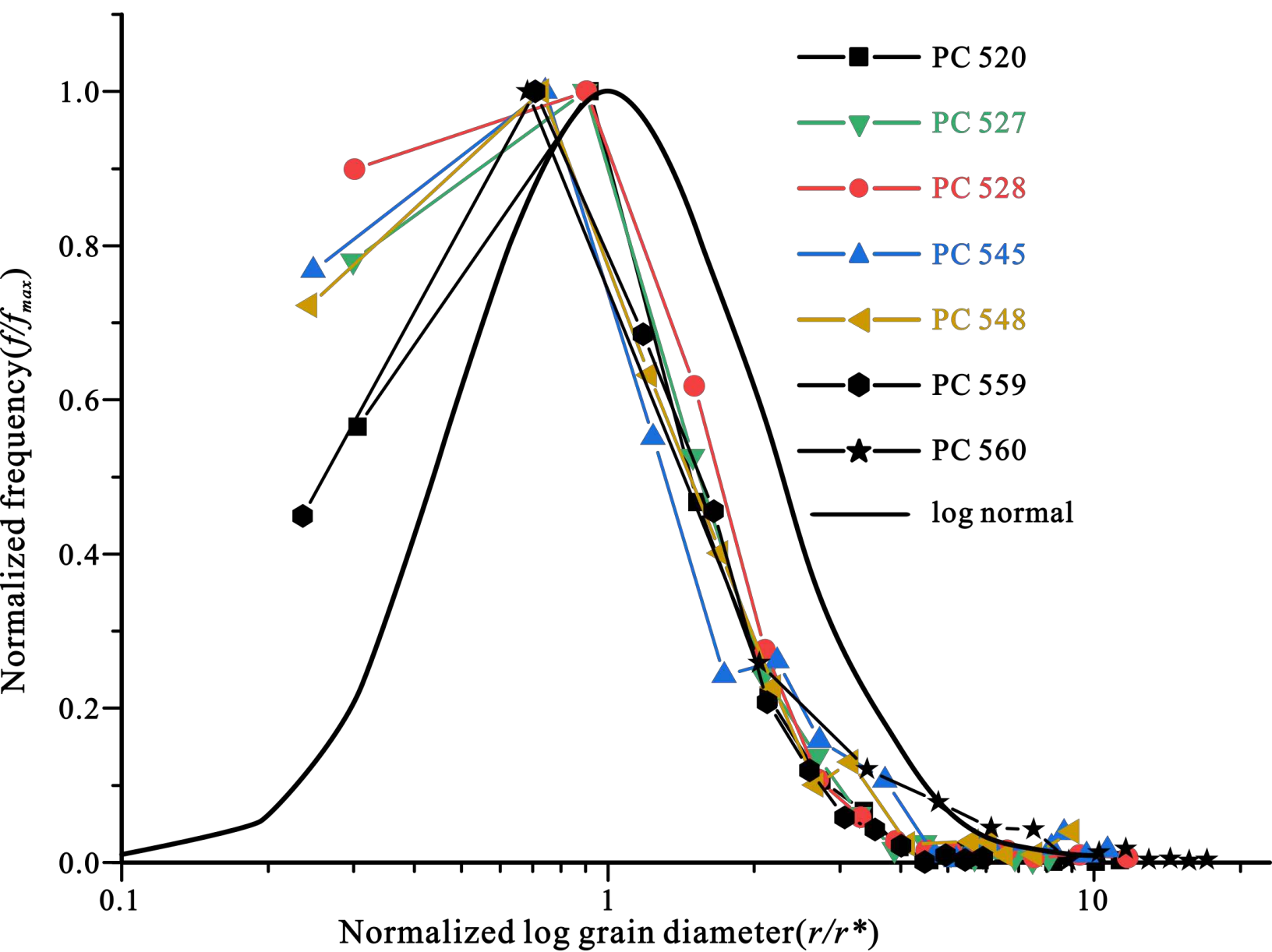


Figure 10.

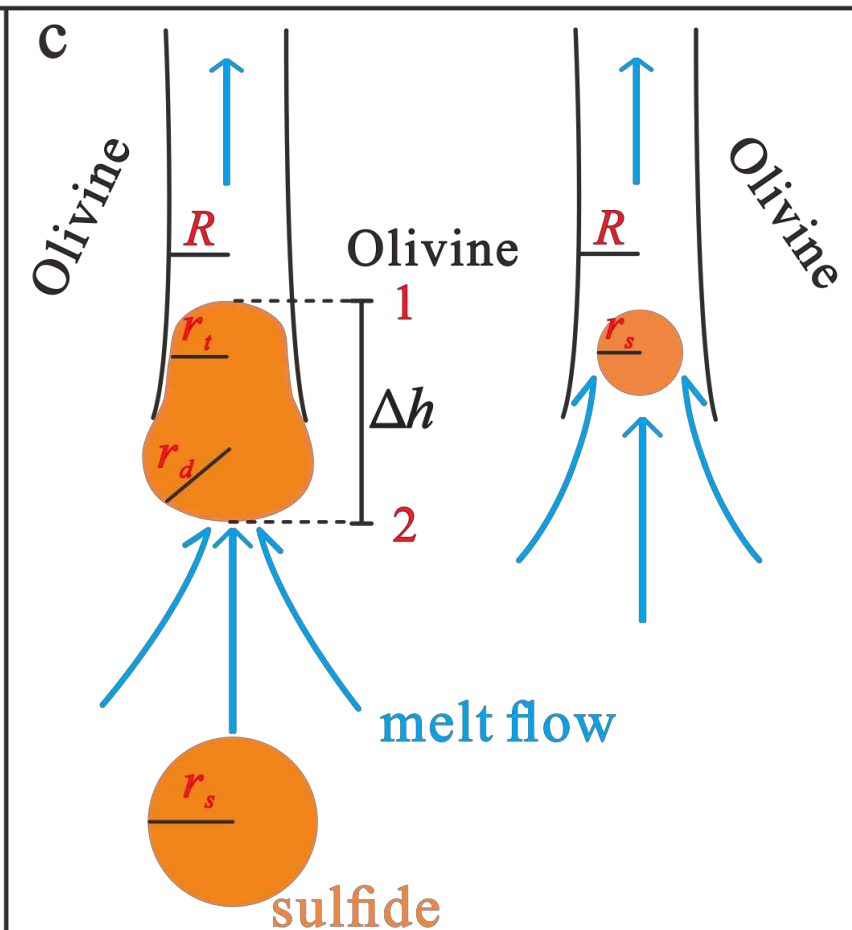
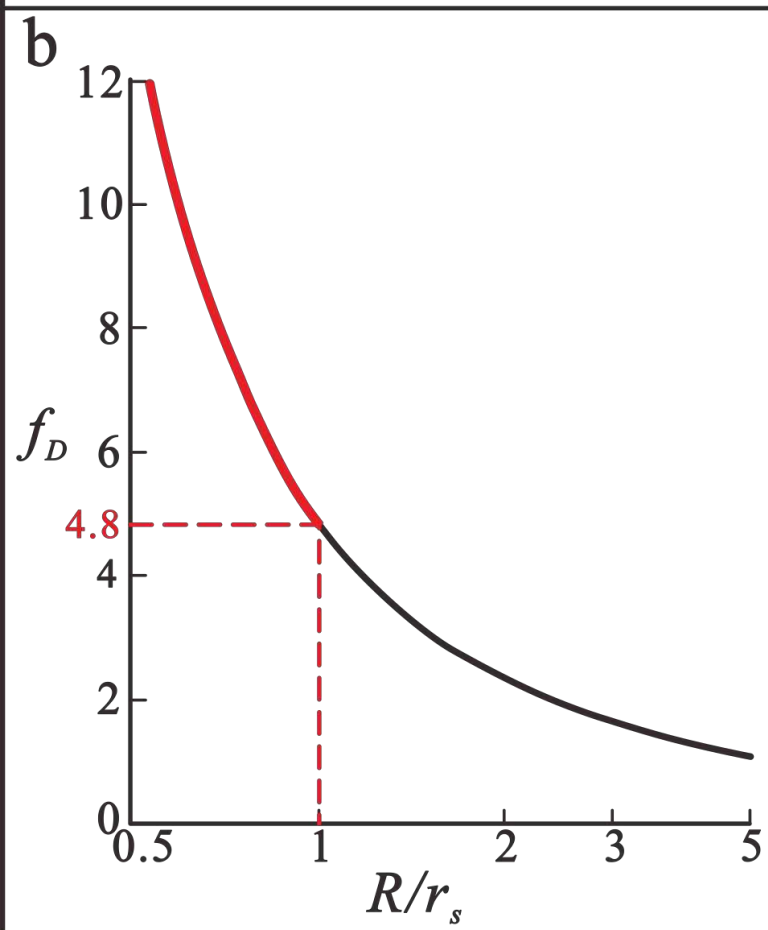
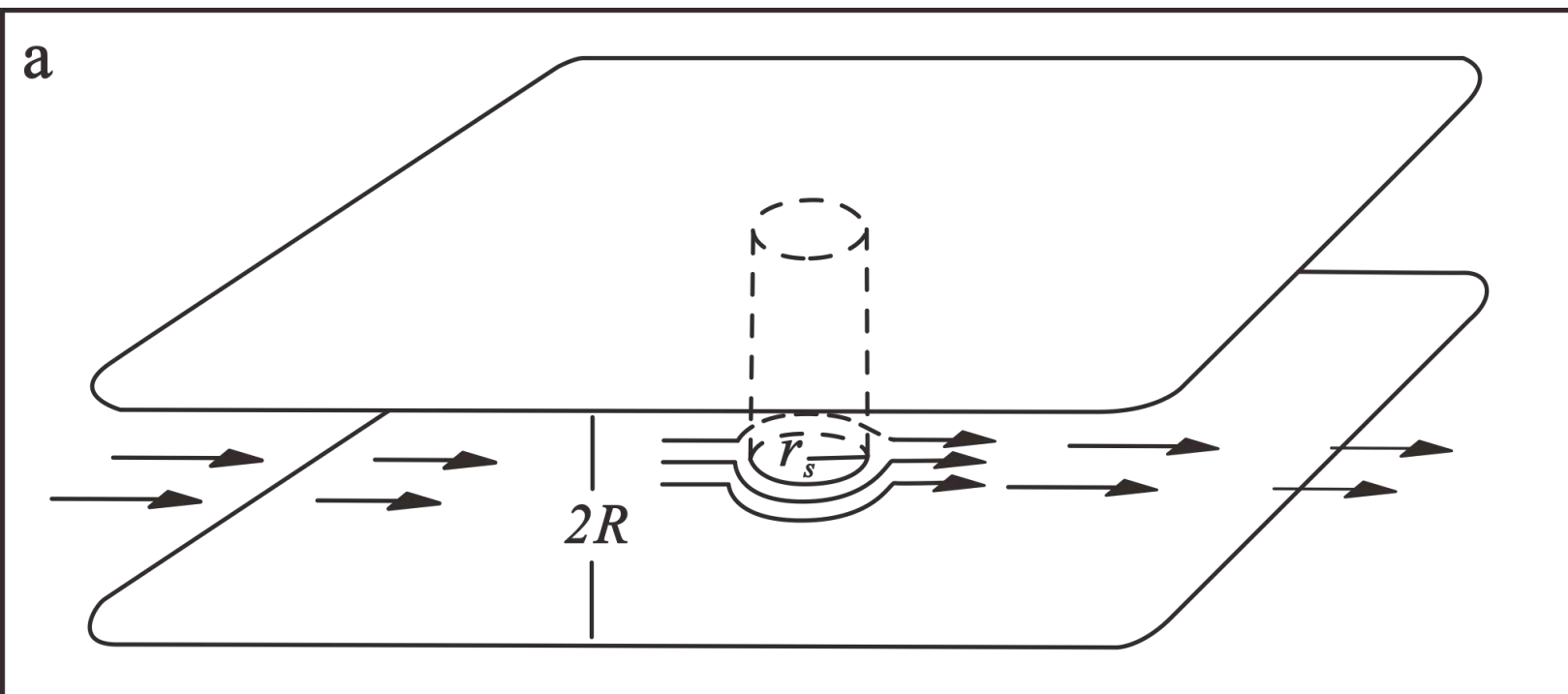


Figure 11.

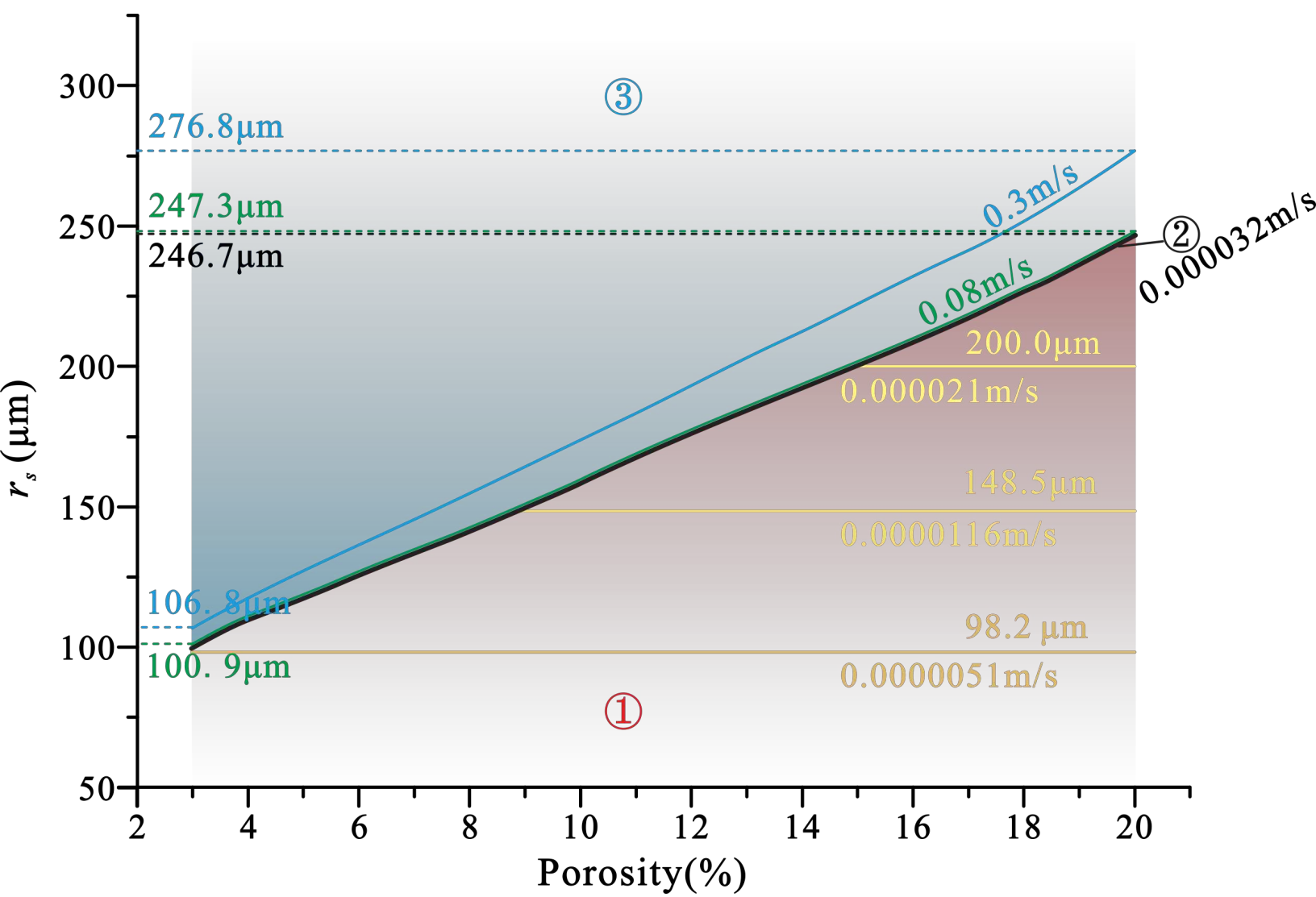


Table 1. Experimental conditions and the observed lithologies

NO.	Starting material (upper peridotite/lower melt source)	Experimental conditions (P/GPa, T/°C, t/h)	Capsule	Lithologies (upper peridotite/reaction interface/melt source)
PC520		1.5, 1250, 12		Ol+Cpx+SM+SL/ ORL/ Opx+SM+SL
PC527		1.5, 1300, 12		Ol+Cpx+SM+SL/ Ol-SM layer/ ORL/ Opx+SM+SL
PC528	Ol+Cpx/ Ol+basalt+sulfide	1.5, 1250, 48		Ol+Cpx+SM+SL/ ORL/ Opx+SM+SL
PC537		1.5, 800, 12		Ol+Cpx/ Ol+basalt+sulfide
PC545		1.5, 1250, 72	Pt-graphite	Ol+Cpx+SM+SL/ Ol-SM layer/ ORL/ Opx+SM+SL
PC548	Ol+Cpx+basalt (5wt.%)/ Ol+basalt+sulfide	1.5, 1250, 48		Ol+Cpx+SM+SL/ ORL/ Opx+SM+SL
PC559		0.5, 1200, 6		Ol+Cpx+SM+SL+V/ Ol+SM+SL
PC560	Ol+Cpx+basalt (5wt.%)/ Ol+basalt+carbonate (2wt.%)+sulfide	1.5, 1250, 24		Ol+Cpx+SM+SL/ Ol-SM layer/ ORL/ Opx+SM+SL

Note: Ol-olivine; Cpx-clinopyroxene; SM-silicate melt; SL-sulfide liquid; V-vapor bubble; ORL-orthopyroxene-rich reaction layer.