

1           **Limitations in one-dimensional (an)elastic Earth**  
2           **models for explaining GPS-observed M<sub>2</sub> Ocean Tide**  
3           **Loading displacements in New Zealand**

4           **B. Matviichuk<sup>1</sup>, M. A. King<sup>1</sup>, C. S. Watson<sup>1</sup> and M. S. Bos<sup>2\*</sup>**

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7           **Key Points:**

- 8           • M<sub>2</sub> ocean tide loading displacements in New Zealand are inferred from GPS ob-  
9           servations  
10          • Estimates for the North Island are not reproduced by models combining ocean tide  
11          loading and a 1D (an)elastic Earth structure.  
12          • Spatially-coherent residual displacements of  $\sim 0.4$  mm (2%) are likely due to lat-  
13          eral Earth structure associated with Pacific Plate subduction

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**Abstract**

[GPS observations of ocean tide loading displacements can help infer the regional anelastic properties of the asthenosphere. We estimate  $M_2$  ocean tide loading displacements at 170 GPS sites in New Zealand and compare these to modeled values using a range of numerical tide and radially symmetric (1D) elastic and anelastic Earth models. Regardless of the model combination we are unable to reduce the strong spatial coherence of the  $M_2$  residuals across the North Island where they reach 0.4 mm (2%). The best fit in the North Island is obtained when combining the FES2014b tide model with spatially-variable ocean density and water compressibility, and the STW105 Earth model. The residuals exhibit a change of  $\sim 0.3$  mm in magnitude between the Taupo Volcanic Zone and the east coast ( $\sim 100$  km), suggesting that this region's laterally-varying, shallow rheological structure may need to be considered to explain these observations.]

**Plain Language Summary**

[The solid Earth changes shape due to the changing weight of the ocean as the ocean tides rise and fall. Measuring this change and comparing it to predictions can yield insights into the interior properties of the Earth, tens to hundreds of kilometers below the surface. We used GPS to measure the changing shape of New Zealand and compared it to predictions based on a range of Earth and tide models. The difference between the observed and modeled displacements revealed a complicated pattern over New Zealand, especially in the North Island and specifically near the Taupo Volcanic Zone. Due to the high accuracy of our GPS analysis and the complex geological structure of the region, we observed the limitations of modeling the loading of the solid Earth by a simple model that varies only with depth.]

**1 Introduction**

The asthenosphere, the weak viscoelastic substrate beneath the lithosphere, is fundamental to the concept of plate tectonics and the earthquake cycle (Hu et al., 2016). The rheological properties of the asthenosphere are, however, not well understood (Karato, 2012). The importance of the asthenosphere is amplified at active convergent boundaries of tectonic plates, specifically subduction systems that initiate forces principal in driving plate tectonics and mantle convection (Stern, 2004). New Zealand is split by the transform Alpine Fault and is locked between two subduction systems: Hikurangi in the north and Puysegur in the south (Lamarche & Lebrun, 2000). These lithospheric discontinuities should produce the largest perturbations observable on the earth tide and perhaps the ocean load tide (Zürn et al., 1976).

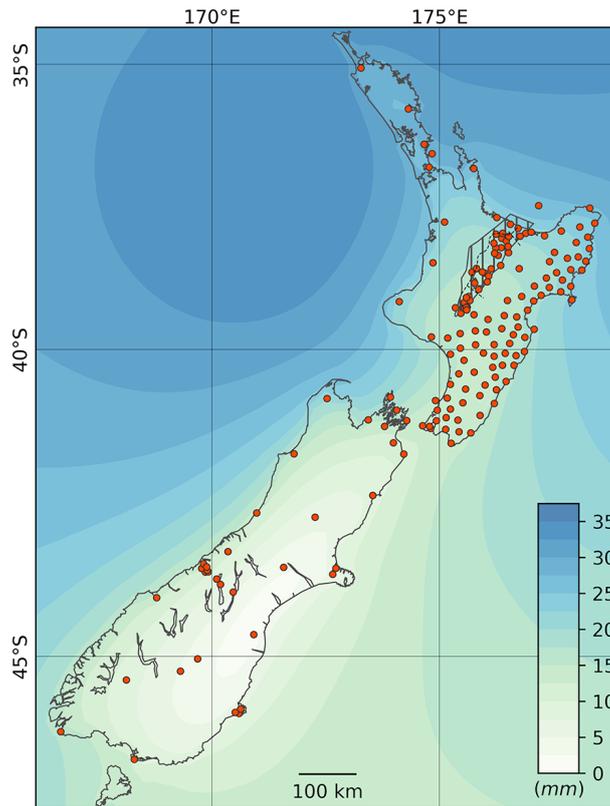
Analysis of Ocean Tide Loading (OTL) displacements, a phenomenon created by the solid Earth's response to tidal-water mass redistribution, can be measured by GPS and such measurements have been used to validate different ocean tide models and elastic Earth models at tidal periods (e.g., Martens et al., 2016; Yuan et al., 2013; Yuan & Chao, 2012). More recently these measurements have been used to constrain the asthenosphere's anelasticity at the period of the major  $M_2$  tidal constituent (period of 12.42 h) by showing improved agreement with deformation modeled using anelastic Earth models. To date, studies of asthenosphere anelasticity have focused on continental settings such as western Europe, western USA, South America, the eastern China Sea region and Alaska (Bos et al., 2015; Ito & Simons, 2011; Martens & Simons, 2020; Wang et al., 2020).

In this paper, we examine the tidal deformation of New Zealand, at the dominant  $M_2$  tidal period, using an array of continuous GPS stations. We combine recent ocean tide models with a range of purely elastic and anelastic Earth models and compare modeled deformation with GPS observed estimates to further understand the anelastic properties of the asthenosphere beneath New Zealand.

## 2 Methods

### 2.1 GPS Data and Analysis

We analyzed all available continuously operating GNSS stations in New Zealand over the period 2013 001 to 2020 153, chosen to maximize the common period of GPS observations. Over this seven-year period, data are available from 170 stations, with all but two (CHTI and RAUL) on mainland New Zealand (see Table S1 for a full list of sites). These stations were designed for nationwide coverage with station spacing in the range 80–100 km to monitor and control the national datum and for geophysical studies (Gale et al., 2015). As shown in Fig. 1, the network provides approximately uniform (but sparse) coverage in the South Island with a substantially higher spatial density of coverage across much of the North Island.



**Figure 1.** Map of New Zealand showing modeled  $M_2$  Up OTL displacement computed with TPXO7.2 ocean tide model and PREM Green’s function. GNSS sites are represented by orange circles. The hatched area in the North Island represents the approximate region of the Taupo Volcanic Zone.

These data were analyzed using GipsyX v1.3 software (Bertiger et al., 2020) using a quite standard kinematic Precise Point Positioning (PPP) approach (Zumberge et al., 1997). The dataset processing was facilitated by a custom wrapper (Matviichuk, 2020). Our approach was described in full by Matviichuk et al. (2020) with the main difference being that we used only the GPS data. Data from other GNSS (e.g. GLONASS) were not logged at all sites over this period hence was excluded from this analysis. We used NASA JPL’s orbit and clock products from their third internal reprocessing campaign (repro 3.0, released March 2018). Ambiguities were fixed to integers where possible (Bertiger et al., 2010). Earth body tides were modeled within GipsyX according to IERS 2010 Con-

83 ventions (Petit & Luzum, 2010). A priori OTL values were removed based on the FES2004  
 84 ocean tide model (Lyard et al., 2006) and Gutenberg-Bullen purely-elastic Earth model  
 85 (Farrell, 1972) in a centre-of-mass of the whole Earth system frame (<http://holt.oso.chalmers.se/loading>)  
 86 – we later restored the OTL component at the coordinate time series level for further  
 87 study; this remove-restore approach is done to reduce the magnitude of companion tides  
 88 and follows approaches adopted previously (e.g., Abbaszadeh et al., 2020; Matviichuk  
 89 et al., 2020; Penna et al., 2015).

90 The GipsyX coordinate and zenith-wet-delay process noise values were chosen based  
 91 on the tests of Penna et al. (2015), Wang et al. (2020) and Matviichuk et al. (2020), us-  
 92 ing values of 3.2 mm/sqrt(s) and 0.1 mm/sqrt(s), respectively. Our parameterization pro-  
 93 duces coordinate estimates every 300 s from which we remove large outliers identified  
 94 with clock bias estimates larger than  $3 \times 10^3$  meters and residuals to a linear trend larger  
 95 than  $\pm 3\sigma$  of each global cartesian coordinate component. These timeseries were converted  
 96 to local topocentric east, north and up components which were then further analyzed.

## 97 2.2 OTL Models

98 We focus here on the difference between the GPS-derived OTL displacements and  
 99 those modeled based on ocean tide models and elastic and anelastic Earth models. For  
 100 the tides we mainly consider three relatively recent global ocean tide models: GOT4.10c  
 101 (Ray, 2013), TPXO9.v1 (Egbert & Erofeeva, 2002) and FES2014b (Carrere et al., 2016),  
 102 although we also explore FES2012 (Carrere et al., 2013) and FES2004 (Lyard et al., 2006).  
 103 We also consider one regional New Zealand tide model (Walters et al., 2001), EEZ, which  
 104 we combine with FES2014b outside the model’s domain for loading computations.

105 The  $M_2$  tide reaches over 1 m near the coast of New Zealand, due to the shallow  
 106 bathymetry, and decreases to 10-20 cm in the open ocean (Stammer et al., 2014). The  
 107 pattern of  $M_2$  between the two islands of New Zealand acts like an amphidromic point  
 108 although the amplitudes are not zero. As a result, the tides to the east and west of New  
 109 Zealand are out of phase and partly cancel out each other’s contribution to the total OTL  
 110 value.

111 All modeled OTL values were computed using the CARGA software (Bos & Baker,  
 112 2005). The coastline was taken from the GMT database (Wessel & Smith, 1996) and has  
 113 a resolution of around 150 m. In most studies a constant sea water density is assumed,  
 114 for example 1030 kg/m<sup>3</sup>. Ray (2013) advocated to take the spatial variation of the den-  
 115 sity into account, and even the fact that water is slightly compressible, which means that  
 116 the mean density of a water column should increase due to the extra density at the bot-  
 117 tom of the column. For the ocean around New Zealand the effect on the resultant de-  
 118 formation is around 1-3%. Assuming a mean 2% effect and a mean OTL amplitude of  
 119 20 mm, this corresponds to a potential error of 0.4 mm which is too large to be ignored.  
 120 We have implemented the equations of Ray (2013) and obtained mean density values from  
 121 the World Ocean Atlas 2013 - WOA13 (Zweng et al., 2013) based on a  $0.25 \times 0.25^\circ$  grid.

122 Three Green’s functions were assessed with this set of ocean tide models: PREM  
 123 (Dziewonski & Anderson, 1981), STW105 (Kustowski et al., 2008) and S362ANI (Kustowski  
 124 et al., 2008). PREM and STW105 provide radial (1D) profiles for the density, and seis-  
 125 mic velocities  $V_p$  and  $V_s$ . These profiles were used to compute load Love numbers which  
 126 were converted into Green’s functions (Bos & Scherneck, 2013). The method is based  
 127 on Alterman et al. (1959) but uses the more recent Chebyshev collocation method to solve  
 128 the differential equations (Guo et al., 2001). These profiles are based on seismic data and  
 129 are only valid at a period of 1 s. To convert them to the period of harmonic  $M_2$ , a con-  
 130 stant absorption band ( $Q=\text{constant}$ , see Table S5) is assumed between these two peri-  
 131 ods (Bos et al., 2015). S362ANI is based on STW105 but has a shear velocity that varies  
 132 horizontally, not just by depth. We averaged the values in a rectangular region between  
 133  $48^\circ\text{S}$  and  $33^\circ\text{S}$  and  $165^\circ\text{E}$  and  $180^\circ\text{E}$  to yield a model representative of the average val-  
 134 ues over the study region. Once converted into a radially symmetric model, the Green’s  
 135 function for S362ANI was computed in similar manner as PREM and STW105.

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### 2.3 OTL Analysis

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Amplitudes and phases of tidal constituents were estimated from the timeseries using the Eterna software v.3.30 (Wenzel, 1996) for 17 tidal constituents, with local phases converted to Greenwich phases with lags positive to compare with the models of OTL displacement. Our focus is solely on the largest loading constituent in New Zealand,  $M_2$ , the major semi-diurnal lunar constituent. To decrease the computation time and measurement noise, the timeseries were first downsampled to 30-min through window averaging.

We follow the naming conventions of Yuan and Chao (2012) with observed and modeled OTL referred to as  $Z_{obs}$  and  $Z_{OTL}$  respectively with  $Z_{res}$  being their vector difference. We refer to the magnitude of the vector difference as  $\|Z_{res}\|$ .

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## 3 Results

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### 3.1 Preliminary analysis of the ocean tide models

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To better understand the variation between recent global ocean tide models at the  $M_2$  period we compute the mean ocean tide model (Fig. S1a) and inter-model standard deviation (Fig. S1b) using the three more recent ocean tide models: FES2014b, GOT4.10c, and TPXO9.v1. We found  $M_2$  SD values of 0.18 cm and 2.68 cm for the deep ocean (>1000 m depth) and the shallow sea ( $\leq$  1000 m depth) respectively. These values are smaller by 40% and 20% than globally derived values reported by Stammer et al. (2014) for  $M_2$ . The largest SD values of up to 0.6 m are located in the Hauraki Gulf in the western North Island, which indicates the region where the largest ocean tides errors are expected.

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We compared FES2004 and EEZ models with a mean ocean tide model, specified above (Fig. S2) by computing a vector difference grids. FES2004 has higher discrepancies (up to 1 m) in the semi-closed water bodies and shallow bights (e.g. South and North Taranaki bights, Fig. S2a), while the EEZ regional tide model shows an approximately constant vector difference in the shallow sea waters ( $\leq$  1000 m) of around 0.1 m (Fig. S2b).

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### 3.2 Comparison of GPS and PREM-based Models

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The GPS-estimated  $M_2$  OTL displacements (with the a priori model restored) are shown in Fig. S3 and listed in Table S2 for each of the up, north and east components. These show a spatially coherent signal across New Zealand with the amplitude ranging from 2 (WAIM) to 32 (KTIA) mm. Using these observations and the modeled  $Z_{OTL}$  based on FES2014b and PREM we computed  $Z_{res}$  as shown in Fig. 2a.  $M_2$  up residuals in the North Island are significant and demonstrate a spatially coherent amplitude of  $\sim$ 1 mm and phase residual of  $\sim$ -10° while residuals in the South Island are small but harder to interpret due to the lower station density and the low OTL amplitude (Fig. 1). This is consistent across different global ocean tide models as indicated by the  $\|Z_{res}\|$  values summarized in the boxplots (Fig. 2c, S4-5).  $\|Z_{res}\|$  variation over the range of tide models with PREM has median value of around 0.7 mm for any of the global tide models while the median for EEZ is  $\sim$ 2 mm. This bias within the EEZ model results in a spatially coherent signal evident from the phasor maps (Fig. S6.2, up component), especially in the North Island.

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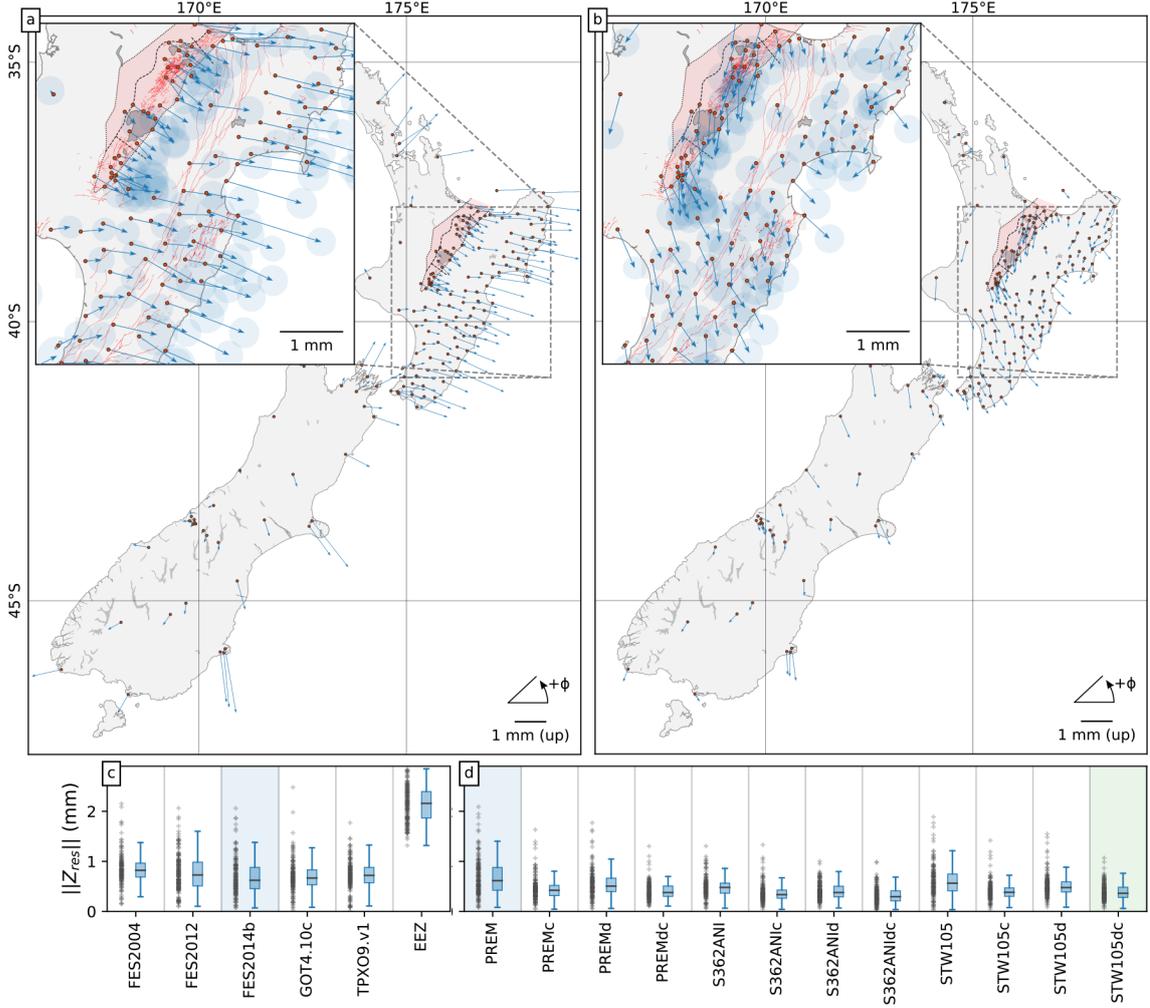
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While all the recent global ocean tide models perform similarly in the horizontal components, FES2014b demonstrates the largest reduction of  $\|Z_{res}\|$  over the set of Green's functions in the up component (Fig. S5). Thus, we continue with FES2014b (Fig. 2c) as a baseline ocean tide model for the subsequent tests.

We considered the impact on the total OTL of specific water bodies by dividing the global oceans into nine separate water areas surrounding New Zealand (Fig. S1). To illustrate the influence of different regions, we selected three sites that experience high, moderate and low  $M_2$  OTL: KTIA, RGMT and MQZG, respectively (Fig. 3). The set of ocean tide models considered consists of the three recent global atlases (FES2014b,



**Figure 2.**  $M_2$  OTL residuals relative to FES2014b.PREM (a), FES2014.STW105dc (b) and residual magnitude ( $\|Z_{res}\|$ ) boxplots for different model setups (c, d). The horizontal line on each box is the median value, the box represents the inter-quartile range (IQR) and the whiskers show an additional  $1.5 \times \text{IQR}$ . Blue and green shading highlights boxplots of (a) and (b) maps, respectively. The Earth model suffixes ‘d’ and ‘c’ in panel (d) refer to the additional treatment of dissipation and compressibility, respectively.

186 TPX09.v1 and GOT4.10c), FES2012 and EEZ. The latter produces  $\sim 2$  mm residual  
 187 amplitude (purple symbols in Fig. 3) and is thus excluded from further analysis. The  
 188 other models show closer agreement but in general outside the observation 95% confi-  
 189 dence interval when using PREM (Fig. 3, bottom panels). However, we note the simi-  
 190 lar magnitude of the variance in  $\|Z_{res}\|$  for all models including EEZ (when the bias is  
 191 ignored) in the up component and complete absence of a  $\|Z_{res}\|$  bias in the horizontal  
 192 components (Fig. S5).

193 Residuals using the purely elastic STW105 show a similar level of variance and median  
 194 as PREM (Fig. 2d) while S362ANI has 50% reduced variance and slightly reduced  
 195 median (0.48 mm compared with 0.61 mm for PREM). However, neither model produces  
 196 consistent agreement within the GPS uncertainty as shown, for example with the three  
 197 sites presented in Fig. 3. We next explore the sensitivity of the modeled OTL to anelas-

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tic dissipation (denoted suffix "d"), and spatially-varying ocean density and compressibility ("c").

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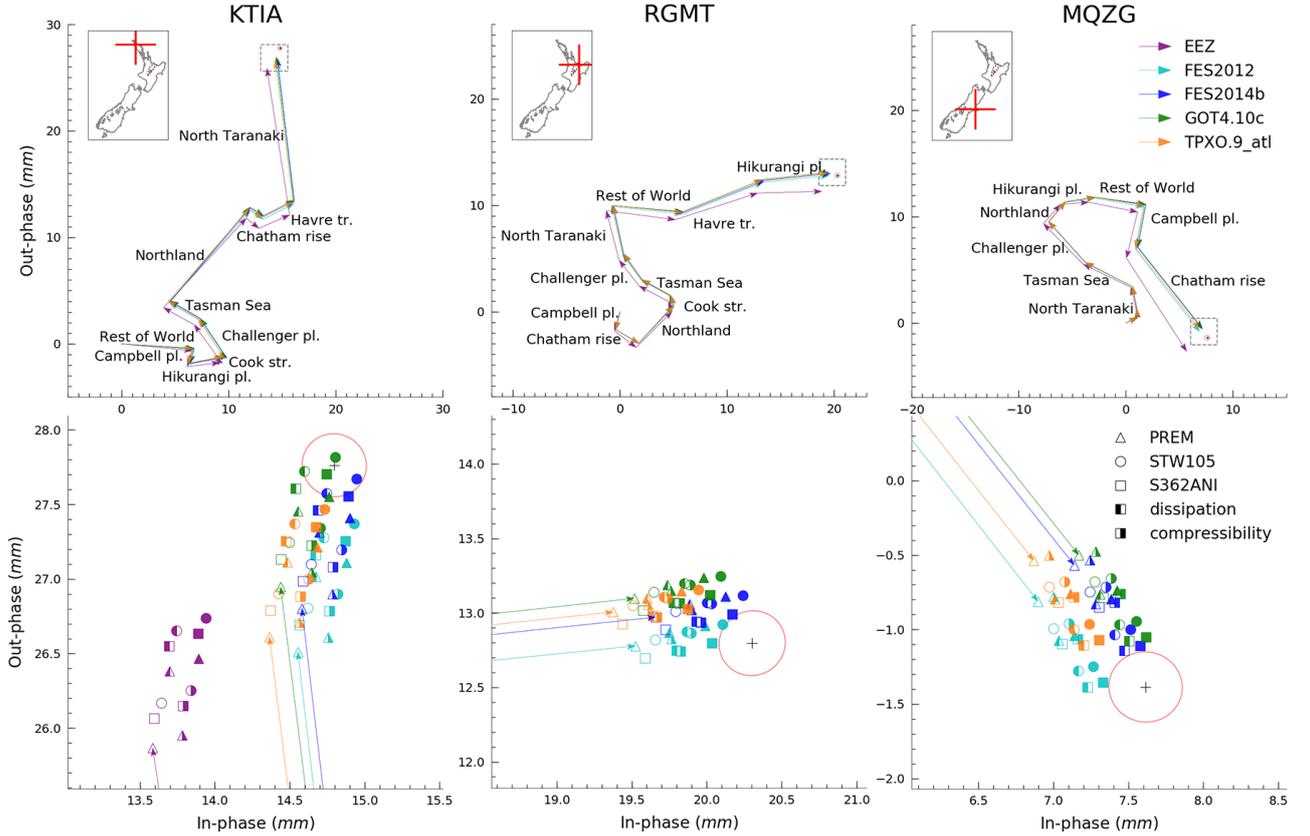
### 3.3 Effect of Considering Anelasticity (Dissipation)

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Bos et al. (2015) demonstrated that accounting for some of the effects of  $M_2$  mantle anelasticity by modifying the Green's functions to include dissipation, decreased OTL residuals in western Europe by up to 0.2 mm. Matviichuk et al. (2020) confirmed these results for the same region but using a different time frame, while similar results have been found by Wang et al. (2020) for south-east Asia.

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For New Zealand, we find a reduction of  $\|Z_{res}\|$  variance and median for all Earth models when dissipation is included (Fig. 2d). The effect is illustrated in Fig. 3 where the models including dissipation (squares with left side only filled) are shown to be closer to the GPS estimates. These do, however, remain outside the GPS 95% confidence interval. At the same time as this improvement, we noticed the introduction of up to 0.2 mm  $\|Z_{res}\|$  bias into the north component with dissipation enabled, independent of the Green's function used; the east component also shows this effect but only with S362ANI (Fig. S4). Enabling sea water compressibility correction partially suppresses the bias. We discuss this further below.



**Figure 3.** Phasor plots of the OTL contributions from different oceanic regions (see Fig. S1) for  $M_2$  Up displacements computed with various Green's functions and ocean tide models. The right bottom panels show the detail for the vector tip as highlighted in the respective left top panels. GPS observations are shown with a black "+" and 95% confidence interval as a red circle. OTL produced by the outside area is titled as "rest of the world".

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### 3.4 Assessment of Water Density and Compressibility Correction

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Enabling the seawater compressibility correction decreases the median  $\|Z_{res}\|$  by a further  $\sim 0.2$  mm in the up component, as shown in Fig. 2d and by example in Fig. 3 (fully filled symbols). In some cases, the application of both dissipation and compressibility eliminates the residual in the up component, although as we discuss in the next section, large, regionally coherent residuals persist. Horizontal components show an increase in variance (Fig. S4) with only compressibility. The dissipation-introduced  $\|Z_{res}\|$  bias in the north component can be partially (S362ANIdc) or completely (PREMdc, STW105dc) removed by additionally applying the compressibility correction (Fig. S4-5, FES2014b). The east component shows a marginal (less than 0.1 mm) increase in both  $\|Z_{res}\|$  median and variance over the solutions with just dissipation included for PREM and STW105, while S362ANI shows further dissipation-introduced increase in  $\|Z_{res}\|$  bias by another 0.1 mm (Fig. S4).

Following Martens and Simons (2020), we constructed Empirical Cumulative Distribution Function (ECDF) plots (Fig. S7.1). The ECDF analysis shows the expected behavior of the corrections in the up component: each correction increases the slope of the ECDF indicating successive improvement with each correction. This is not the case for the horizontal components where both corrections introduce biases using S362ANI, which otherwise demonstrates performance comparable to other models without the corrections. The optimum correction of PREM and STW105 in the north component very much relies on the selection of ocean tide model. The dissipation-introduced bias is suppressed by the compressibility correction in the case of FES2014b and GOT4.10c, which suggests the best performance with both dissipation and compressibility corrections enabled. In the case of TPX09.v1, the bias is too large for compressibility to overcome, effectively repeating the trend as observed for S362ANI.

Removing the respective mean  $Z_{res}$  values from each set of residuals (Fig. S7.2) aligns the ECDFs over all components, fully removing the differences in the horizontal components with exception for S362ANI-based values in the north component. Removing mean  $Z_{res}$  also absorbs any long-wavelength errors in the solid Earth body tide.

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## 4 Discussion

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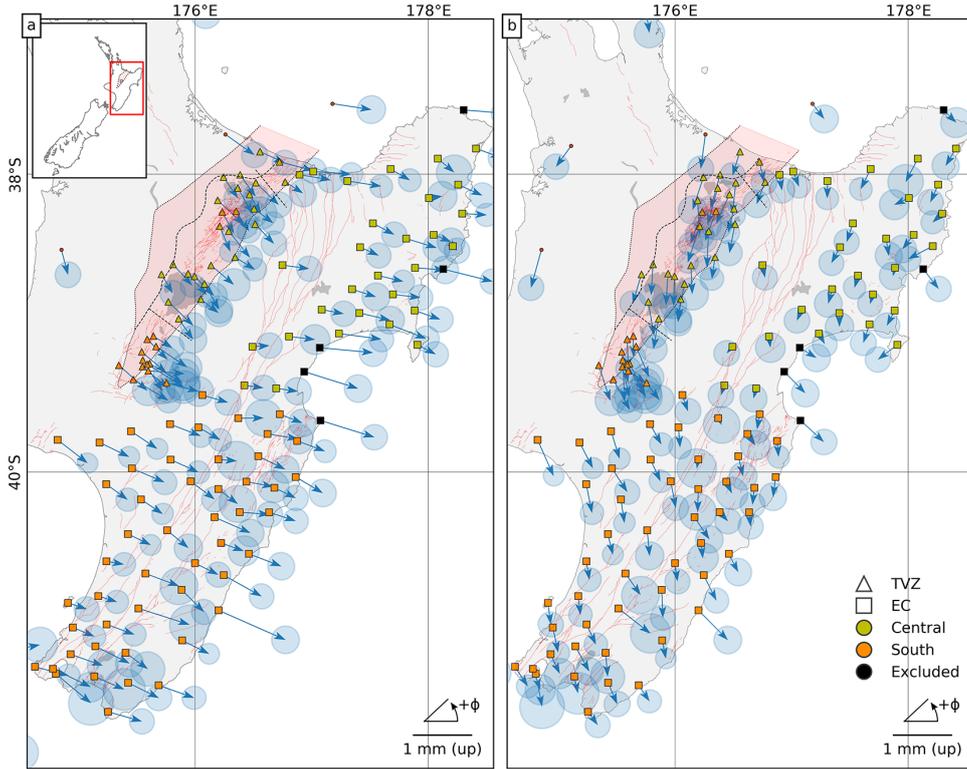
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Following these tests, the optimal agreement between the observed and modeled OTL in all three components occurs when using FES2014b and STW105dc. The spatial distribution of  $Z_{res}$  shows  $\sim 0.5$  mm residual spatially coherent field over the Taupo Volcanic Zone (TVZ)  $Z_{res}$  in the North Island, as shown in Fig. 4. The dense coverage of stations in these regions reveals a distinct change of  $Z_{res}$  between sites in the East Coast (EC) and TVZ that experience the same  $M_2$  OTL (Fig. 1).

To aid discussion, we consider four different regions (blocks) within this region as illustrated by the symbols in Fig. 4: TVZc, TVZs, ECc, ECs, with “c” and “s” subscripts identifying central and southern subareas, respectively. Residual OTL in each block was averaged to provide  $Z_{res}$  summary metrics (per component) relevant to each region (Table S6). Note that several sites along the EC were removed (e.g. Hawke Bay) as they experience a localized signal (Fig. 4, black symbols) which is independent of the ocean tide model or Green’s function used. The sites in both TVZ regions show residual amplitudes of  $\sim 0.5$  mm with phase changing sharply from  $-102^\circ$  to  $-70^\circ$  between TVZc and TVZs. The relative phase change between TVZ and EC within the same central or south area (TVZc/ECc and TVZs/ECs) is found to be approximately constant ( $\sim 35^\circ$ ) while revealing 0.25 mm and 0.15 mm larger amplitudes for TVZc and TVZs, respectively.

The sharp change in residual phase between TVZc and TVZs, and the strong spatial variation in residual amplitude between respective EC and TVZ sub-regions over length-scales of the order of  $\sim 100$  km suggests that the variations are due to localized effects. We discount errors in ocean tides given our previous tests and the spatial distribution



**Figure 4.** GPS-derived  $M_2$  OTL residuals for a section of the North Island relative to FES2014b ocean tide model combined with dissipation corrected STW105d (a) and STW105dc (b). "d" and "c" suffixes stand for dissipation and compressibility corrections. Sites are categorized into Taupo Volcanic Zone (TVZ) and East Coast (EC) regions (symbol shape) with subdivision of each into central and south along the TVZ central/south boundary (symbol color).

266 of the residuals. Also, biases in the adopted deep Earth rheological structure (Lau et al.,  
 267 2017) would be effectively constant over this area.

268 Instead, our assumption is that the residuals result from mismodeled shallow-Earth  
 269 rheological structure. To explore this further, we iterate through a range of alternative  
 270 Earth models with different rheological structure in the upper tens of kilometers based  
 271 on seismic tomography inversions (Eberhart-Phillips & Bannister, 2015; Eberhart-Phillips  
 272 & Fry, 2018). No single one-dimensional (radially-varying) Earth model could explain  
 273 the regional pattern of residuals, with changes generally producing changes that were  
 274 spatially uniform across the region of Fig. 4.

275 Deviations in the shallow rheological structure from that used to compute the Earth  
 276 body tides could produce localized residuals. (Zürn et al., 1976) developed a 2D finite-  
 277 element model of a subduction zone in Alaska, and showed that the subduction zone struc-  
 278 ture can produce an effect up to 0.8% on the solid Earth body tide in the radial direc-  
 279 tion directly above the asthenospheric slab. For  $M_2$  at the latitude of the North Island,  
 280 this equates to 0.7 mm. However, their modeling also showed that the maximum gradi-  
 281 ent over the distance from East coast to TVZ (up to 150 km) should not exceed 0.25%.  
 282 We note that the effect on phase is not described in their work. However, if we consider  
 283 the relative location of the TVZ over the subduction slab (observed by the  $V_p$  anomaly  
 284 at 100-130 km depth (Eberhart-Phillips & Fry, 2018)), the maximum expected change  
 285 becomes close to 0.15%, or 0.13 mm for  $M_2$  at these latitudes. As such, this is well be-  
 286 low the magnitude of the variations seen in Fig. 4.

287 The effect of lateral rheological structure on modeled OTL is unclear. However, mod-  
 288 eling of elastic deformation due to longer-period surface mass displacement indicates that  
 289 consideration of global, localized Earth structure produced differences of the order 10%  
 290 in the vertical and 20% in the horizontal over distances of 10-50 km (Dill et al., 2015).  
 291 The average  $M_2$  OTL in the region of the TVZ shown in Fig. 4 is  $\sim 19$  mm and so even  
 292 a 2% effect due to lateral variation may be relevant to explaining the observed residu-  
 293 als. Given the minor, but non-negligible effect of lateral variation on Earth body tides,  
 294 and likely effects on OTL displacements, our analysis suggests that one-dimensional mod-  
 295 els of this region are unlikely to fully explain GPS observations of OTL at  $M_2$ . To check  
 296 for potential long-wavelength errors that could introduce the observed dissipation-introduced  
 297 biases in the horizontal components, we repeated our analyses for a set of 15 stations in  
 298 the inland Australia (see Table S3 for site list and Table S4 for derived observations) where  
 299 the geological setting is simpler and where a 1D model should produce accurate results.  
 300 For this dataset we needed to adopt a different time period (2015-2018 inclusive) due to  
 301 data availability but checking a subset of sites in New Zealand found time-period was  
 302 inconsequential. Figures S9 and S10 demonstrate that, although the magnitude of the  
 303 OTL is still several mm, for these stations the residuals (observed minus predicted OTL)  
 304 are indeed small and within the uncertainty of the observations. This validates the cor-  
 305 rectness of our analyses and suggests that tidal centre-of-mass errors in this region are  
 306 small.

307 Figures S7.1 and S7.2 show that the OTL residuals for the horizontal components  
 308 suffer from some common mode problem but that modification of the Green's function  
 309 cannot reduce the misfit. For the up component, the influence of the dissipation effect  
 310 within asthenosphere that requires us to modify the elastic properties of the Earth model  
 311 from the reference period of 1s to tidal periods is noticeable. Furthermore, including spa-  
 312 tially varying seawater density and compressibility results in an additional reduction of  
 313 the misfit. The two figures also demonstrate that the difference between the ocean tide  
 314 models used in the loading computations is small. Therefore, the most likely candidate  
 315 to reduce the misfit further reduction is using an advanced (3D) (an)elastic model of the  
 316 region.

317 Similar problems using a 1D Earth modeling OTL loading in Alaska were recently  
 318 described by Martens and Simons (2020). We are unaware of three-dimensional mod-  
 319 els being in use for the computation of OTL, however Latychev et al. (2009) have com-  
 320 puted Earth Body tides with a three-dimensional model. One practical consequence of  
 321 this is that mismodeled tidal deformations in this region will propagate into conventional  
 322 24 hr coordinate solutions (Penna et al., 2007). Such propagation will introduce long-  
 323 period noise in GPS coordinate time series in New Zealand and impact subsequent geo-  
 324 physical interpretation.

## 325 5 Conclusions

326 We estimate  $M_2$  ocean tide loading (OTL) displacements at 170 GPS sites in New  
 327 Zealand between 2013 001 and 2020 153. Comparison with modeled OTL displacements  
 328 using a range of global tide models and elastic PREM shows sub-mm agreement, with  
 329 much larger disagreements using a local New Zealand tide model.

330 However, we find that the no single one-dimensional elastic Earth model, when com-  
 331 bined with modern global tide models, can consistently explain the GPS-derived OTL  
 332 within uncertainties. Of the tested ocean tide models, FES2014b produced the best re-  
 333 sults. However, application of an anelastic dissipation correction, and varying water den-  
 334 sity and compressibility substantially improves the agreement between the various mod-  
 335 els and observed OTL. Despite this, some regional spatially-coherent unmodeled resid-  
 336 ual signals remain in the North Island with magnitudes of up to 0.3 mm. These show  
 337 substantial variation in phase over  $\sim 100$  km in the region between the Taupo Volcanic  
 338 Zone and the East coast. We attempted to reproduce the observed signal using a range  
 339 of 1D Earth models with varying shallow Earth structures, including the effects of anelas-

340 ticity, however no single model could explain the residuals. We anticipate that these resid-  
 341 uals are a result of unmodeled lateral variations in Earth rheological structure forced largely  
 342 by ocean tide loading but with a smaller component likely from mismodeled Earth body  
 343 tides.

344 This analysis of residual OTL demonstrates the deficiencies of the 1D Earth mod-  
 345 eling approach that is currently standard practice. This is particularly relevant to GPS  
 346 analysis using 24 hr coordinate solutions, given mismodeled tidal displacements prop-  
 347 agate into long-period signal. Utilizing 3D Earth modeling to compute tidal phenom-  
 348 ena is likely required to explain the observations in regions with major discontinuities  
 349 in Earth’s lateral structure (e.g. subduction margins). Such models, combined with these  
 350 observations, could provide new insights into the shallow rheological structure of these  
 351 regions.

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