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A synthetic spring-neap tidal cycle for long-term morphodynamic models

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Key Points:

- A new approach to devise periodic tidal boundary conditions for long-term morphodynamic simulations is developed
- The new method better represents tidal water level dynamics, bed shear stress, and residual sand transport
- The pros and cons of both the new and existing approaches are evaluated

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Abstract

Existing tidal input reduction approaches applied in accelerated morphodynamic simulations aim to capture the dominant tidal forces in a single or double representative tidal cycle, often referred to as a “morphological tide”. These heavily simplified tidal signals fail to represent the tidal extremes, and hence poorly allow to represent hydrodynamics above the intertidal areas. Here, a generic method is developed to construct a synthetic spring-neap tidal cycle that (1) represents the original signal; (2) is exactly periodic; and (3) is constructed directly from full-complexity boundary information. The starting point is a fortnightly modulation of the semi-diurnal tide to represent spring-neap variation, while conserving periodicity. Diurnal tides and higher harmonics of the semi-diurnal tide are included to represent the asymmetry of the tide. The amplitudes and phases are then adjusted to give a best fit to histograms of water levels and water level gradients. A depth-averaged model of the Ems estuary (The Netherlands) demonstrates the effects of alternative tidal input reduction techniques. Adopting the new approach, the shape of the tidal wave is well-represented over the entire length of the estuary, leading to an improved representation of extreme tidal conditions. In particular, representing intertidal dynamics benefits from the new approach, which is reflected by hydrodynamics and residual sand transport patterns that approach non-schematized tidal dynamics. Future morphodynamic simulations forced with the synthetic signal are expected to show a more realistic exchange of sediment between the channels and tidal flats, likely improving their overall predictive capacity.

Plain Language Summary

The time-scales of erosion and deposition processes in estuaries and tidal basins are several orders of magnitude larger than the time scales of the changing flows (years versus hours, respectively). To efficiently simulate years of erosion and deposition, an acceleration factor is applied to estuarine and coastal models that simulate the long-term bed level developments. Tidal information used to force these accelerated models at the seaward boundary requires an exactly repetitive signal to avoid inconsistencies in the up-scaling approach. A tidal input reduction technique is required to cope with the fact that successive spring-neap cycles are never identical. In this paper, a tidal input reduction method is developed that yields a synthetic, periodic tidal signal representing the variation of amplitudes and asymmetries present in a multiyear tidal signal. These variations are not captured well in existing, more limited, approaches for tidal input reduction. The results from a numerical model forced with the synthetic tidal signal shows that intertidal dynamics and residual sand transports are simulated more realistically, compared to existing approaches. The new tidal input reduction method should improve the exchange between the channels and intertidal areas in long-term estuarine and coastal models, presumably allowing for a more realistic assessment of erosion and deposition in these areas.

1 Introduction

The long-term or multi-decadal evolution of estuaries and tidal basins is largely controlled by the interaction between hydrodynamic forcing and the sediment bed (De Swart & Zimmerman, 2009). This morphodynamic dependence on hydrodynamic controls allows for a quantitative investigation on the evolution of tidal basins using process-based numerical models. Such morphodynamic tools in turn allow to simulate the evolution of deltaic environments, which are increasingly influenced by anthropogenic impacts (Syvitski et al., 2009) jeopardizing ecosystem services and potentially leading to morphological instability (Hoitink et al., 2020). Although numerical bed evolution models are often developed to predict the direct morphological response to engineering measures (De Vriend et al., 1993), they appear to be more realistic when the time scales related to the inves-

64 tigated changes (T_c) and the time-scale at which the model attains dynamic equilibrium
65 (T_e) are longer (Dam et al., 2016). Often, the developments from the initial conditions
66 towards the model’s dynamic equilibrium obscure the morphodynamic impact of the in-
67 terventions where the model was originally designed for. As a consequence, process-based
68 models are increasingly used to investigate not only decadal but also centennial and even
69 millennial morphological evolution of estuarine and tidal environments (e.g.; Dastgheib
70 et al., 2008; Van der Wegen & Roelvink, 2012; Nnafie et al., 2018; Braat et al., 2017).

71 Long-term morphodynamic modelling requires appropriate up-scaling of the effects
72 of hydrodynamic processes that typically fluctuate within hours or days to the time pe-
73 riods relevant for morphological changes. Various techniques exist to reduce the com-
74 putational costs for the slow bed level evolution, while accounting for the shorter hydro-
75 dynamic variability. These techniques range from postponed morphological updating,
76 based on gradients in the tide-averaged residual transport, to constructing simplified sed-
77 iment balances that express bottom change in terms of sediment transport gradients de-
78 pending only on the local water depth (Latteux, 1995; De Vriend et al., 1993; Roelvink,
79 2006; Roelvink & Reniers, 2011). The most commonly used morphological updating tech-
80 nique is the fully coupled approach (Roelvink, 2006), where the bed level is updated ev-
81 ery hydrodynamic time step. Such continuous updating includes short-term interactions
82 between flow, sediment transport, and morphology, resulting in a stable bed evolution,
83 also in intertidal areas which are inundated during high water conditions only. From a
84 physical point of view, the hydrodynamics should be resolved as detailed as possible. For
85 reasons of computational efficiency, long-term morphological evolution is often modelled
86 with the additional use of a so-called morphological timescale factor (or MorFac, MF),
87 essentially a multiplication factor for the depth change (Roelvink & Reniers, 2011). At
88 each hydrodynamic time step, the calculated bed level change is multiplied with this fac-
89 tor, reducing the required simulation time with a factor MF. This accelerated approach
90 resolves morphodynamic processes operating at intratidal timescales, while maintaining
91 the speed, stability, and accuracy of tidally averaged updating approaches (Van der We-
92 gen et al., 2008). In idealized geometric configurations, the MF approach can produce
93 stable bed evolution patterns for values up to $O(1000)$ that do not deviate significantly
94 to the patterns simulated with smaller values of MF. A pre-requisite for stability is that
95 the bed level changes are small compared to water depth, so that no irreversible changes
96 develop within a phase of the tidal cycle (Van der Wegen & Roelvink, 2008).

97 Accelerated long-term simulations require schematized boundary conditions with
98 limited extremes, because the sediment transport fields for bed level adaptation are ex-
99 trapolated with the MF approach (and may not exceed critical values within a compu-
100 tational timestep). The time-series of boundary conditions need to be represented by a
101 reduced number of conditions consisting of a repetitive pattern that includes the dom-
102 inant forcing conditions, but excludes intermittent events (e.g. storms) that may exag-
103 gerate the bed evolution. The goal of input reduction is therefore to derive a limited rep-
104 resentative subset of forcing conditions that approach the residual transport and asso-
105 ciated morphological change patterns compared to a simulation forced with the full time-
106 series over the period of interest (i.e. a ‘brute-force’ simulation).

107 Existing methods for tidal input reduction aim at capturing the dominant tidal dy-
108 namics in a single tide (e.g.: Dastgheib et al., 2008; Van Maanen et al., 2013) or with
109 two representative tidal cycles (e.g.: Latteux, 1995; Lesser, 2009). Such simplified tidal
110 signals have been shown to reasonably reproduce morphological changes of tidal chan-
111 nels (Van der Wegen & Roelvink, 2012; Van Der Wegen et al., 2011; Dissanayake et al.,
112 2009; Dastgheib et al., 2008). However, heavily simplified tidal signals fail to represent
113 the tidal extremes (the tidal elevation above Mean High Water and below Mean Low Wa-
114 ter), because they neglect these variations. They poorly represent intertidal areas, which
115 exert a major impact on the development of tidal asymmetry (Friedrichs & Aubrey, 1988).
116 Although the tide-averaged transport of non-cohesive sediments in the main estuarine

117 channels is captured well with solely a semi-diurnal tide and relevant overtides (Van de
118 Kreeke & Robaczewska, 1993), the long-term morphological development of tidal basins
119 is driven by tidal asymmetries resulting from the combination of multiple tidal constituents
120 (Guo et al., 2016). Preserving asymmetries present in the original tidal signal, as well
121 as providing the hydrodynamic conditions necessary for the development of intertidal
122 areas, is therefore a key requirement for the tidal input reduction approach. Despite its
123 importance for long-term morphological modelling, the impact of tidal input reduction
124 methods has rarely been systematically investigated.

125 A systematic investigation of tidal input reduction techniques preferably correlates
126 such techniques to morphological output. However, morphological models are sensitive
127 to parameterizations (e.g. the sediment transport formula) and settings (grid size, bed
128 slope effect) used in the morphodynamic model (Van Maanen et al., 2011; Baar et al.,
129 2019). Although the morphological output is steered by the simulated hydrodynamics,
130 it is also strongly influenced by morphodynamic calibration parameters, diffusing the ef-
131 fect of the boundary schematization. In this paper we therefore refrain from morpho-
132 dynamic simulations and focus on the effect of the tidal input reduction approach on hy-
133 drodynamic model parameters considered relevant for morphodynamics.

134 The aim of this paper is twofold. First, a tidal input reduction technique is intro-
135 duced that yields a synthetic, periodic spring-neap tidal signal representing the tidal ex-
136 tremes as well as tidal asymmetry. Second, the effects of both existing and the new tidal
137 input reduction approaches are systematically investigated, since such an evaluation is
138 missing in the literature. Simulations forced with the original tidal signal (as a reference)
139 and simulations forced with schematized tides are evaluated in terms of tidal asymme-
140 try, bed shear stress, inundation of intertidal flats, and residual sand transports. For this
141 latter purpose, we develop and apply a morphostatic (i.e. no bed level updating) model
142 of the Ems estuary (The Netherlands).

143 The structure of the remainder of this paper is as follows. We first review exist-
144 ing tidal input reduction techniques and explain the new methodology (Section 2). We
145 then develop a numerical model of a real-world estuary (The Ems estuary, Section 3) and
146 apply this to examine the effect of various types of tidal input reduction techniques on
147 simulated hydrodynamics and sand transport (Section 4). The implications of simpli-
148 fying tidal signals are discussed in Section 5, and conclusions are drawn in Section 6.

149 **2 Tidal input reduction**

150 **2.1 The morphological tide**

151 The goal of tidal input reduction is to create simplified representative tidal bound-
152 ary conditions for up-scaling bed level changes in process-based morphological models.
153 The aim is to represent the original tidal series in a simplified signal in a sense that it
154 produces the same residual transport or initial morphological change patterns for a de-
155 fined period and region of interest. The simplified tide is constructed as a periodic sig-
156 nal, so that a sequence of the same synthetic tidal signals is continuous. Such a simpli-
157 fied tide is often referred to as the “morphological tide” (Latteux, 1995).

158 The most common method to derive a morphological tide can be summarised as
159 follows (Roelvink & Reniers, 2011). The morphological development over a sufficiently
160 long time period (e.g.: several spring-neap cycles) is executed with both full hydrody-
161 namic forcing and with several accelerated simulations, each forced with a single tidal
162 cycle, selected from the time-series. The simulated patterns of residual transport or bed
163 level adaptations resulting from reduced input simulations and from the full forcing sim-
164 ulations are subsequently compared based on a correlation coefficient, and the slope of
165 the regression. The tidal cycle that produces simulated results that best resemble the
166 results from a full forcing simulation is then considered to be most representative.

167 Lesser (2009) demonstrated that such a simplified tide fails to correctly represent
 168 residual transport in some cases, because it neglects the asymmetry resulting from in-
 169 teraction between the main semi-diurnal constituent (M_2) and the main diurnal constituents
 170 (O_1 and K_1). Hoitink et al. (2003) demonstrated that in diurnal, or mixed mainly di-
 171 urnal regimes a residual transport can develop resulting from the tidal asymmetry that
 172 arises from these primary constituents because they have angular frequencies that con-
 173 sist of sums and differences of two of the basic astronomical frequencies (see Pugh, 1987),
 174 leading to substantial residual transport and morphological changes (Van Maren et al.,
 175 2004; Van Maren & Gerritsen, 2012). In these regimes, the residual transport that arises
 176 from the triad interaction of K_1 , O_1 and M_2 can be more important than the residual
 177 transport caused by the non-linear interaction of the main semi-diurnal component (M_2)
 178 with its first overtide (M_4) (Song et al., 2011), often considered to be the dominant mech-
 179 anism for shallow water tides (e.g., Friedrichs & Aubrey, 1988; Van de Kreeke & Robaczewska,
 180 1993). Lesser (2009) therefore included this triad interaction by defining an artificial con-
 181 stituent C_1 with half the frequency of the M_2 tidal constituent. The resulting *double tide*
 182 consists of C_1 , M_2 and its overtones, and may include an additional scaling factor for the
 183 amplitude of M_2 and/or C_1 , to account for the presence of a residual flow.

184 A literature review on publications that apply online-updated accelerated process-
 185 based morphodynamic models in tide-dominated settings was performed to provide an
 186 overview of current tidal forcing approaches (Table 1). The 40 publications reviewed re-
 187 veal that tidal forcing is often reduced to the M_2 tidal constituent (17 publications). All
 188 of these studies comprise idealised model configurations. In modelling studies that give
 189 a more realistic representation of the estuarine environment, the tide is usually repre-
 190 sented by M_2 and its overtones (4 publications), the empirically derived morphological
 191 tide (2 publications), or the morphological *double tide* (4 publications). These studies
 192 aim at capturing the dominant tidal forces in a single or double representative tidal cy-
 193 cle. However, the (1D) simulated long-term morphodynamic development of estuarine
 194 environments is governed by the combined effects of asymmetries resulting from the in-
 195 teraction of multiple tidal constituents and river-tide interaction (Guo et al., 2016). Par-
 196 ticularly, the omission of the S_2 constituent reduces the effects of river-tide interaction
 197 and tidal asymmetry, leading to an underestimation of tide-induced residual transport.
 198 Yet, the effects of ignoring significant constituents in simplified tides are not well stud-
 199 ied for 2D morphodynamics. Presumably because of the unknown effects of oversimpli-
 200 fying tides in 2D morphodynamic simulations, the authors of 13 publications chose to
 201 overcome the considerations for tidal input reduction by forcing the full tide (Table 1).
 202 However, using this approach in accelerated simulations the morphodynamic time can-
 203 not accurately be interpreted, because the sum of the tidal periods imposed lacks peri-
 204 odicity. The interval for integration of the residual transport is inconsistent and there-
 205 fore the transport is not accurately averaged over the tidal periods in the signal. These
 206 studies that did not apply tidal input reduction focused on decadal time-scales and such
 207 an imperfect sediment balance may be acceptable with small acceleration factors. For
 208 long-term (i.e. longer than decadal) simulations using larger acceleration factors, a sim-
 209 plified cyclic tide representing all significant tidal constituents (and therefore their in-
 210 teractions) would be an important advance over earlier simplified tides because (the in-
 211 teraction between) each significant tidal constituent plays a role in driving tidal resid-
 212 ual transport, and therefore in morphodynamic development (Guo et al., 2016).

213 2.2 A synthetic representative signal

214 We aim to develop a generic method to construct a representative tidal signal that
 215 incorporates tidal extremes in a synthetic spring-neap cycle, while remaining periodic.
 216 The target synthetic spring-neap cycle: (1) sufficiently represents the original signal to
 217 preserve asymmetries; (2) is periodic, to ensure consistency in the start and end of con-
 218 secutive cycles and to control the relative phasing with other types of forcings (e.g.: wind,
 219 waves, river discharge, ecology); and (3) is derived directly from the boundary informa-

Table 1. Tidal forcing approaches used in online-updated accelerated morphodynamic simulations.

Tidal forcing	Literature
M ₂	Bolla Pittaluga et al. (2015), Braat et al. (2017), Elmilady et al. (2022, 2020), Geleynse et al. (2011), Guo et al. (2015), Hibma et al. (2003), Leonardi et al. (2013), Marciano et al. (2005), Nahon et al. (2012), Van der Wegen et al. (2010), Van der Wegen and Roelvink (2008), Van der Wegen et al. (2008), Van Maanen et al. (2013), Xie et al. (2017), Yu et al. (2014), Zhou et al. (2014)
M ₂ + M ₂ overtones	Dissanayake et al. (2009), Dastgheib et al. (2008), Nnafie et al. (2018), Nnafie et al. (2019)
Morphological tide	Chen et al. (2022), He et al. (2022)
Morph. double tide	Elmilady et al. (2019), Van Der Wegen et al. (2011), Van der Wegen and Roelvink (2012), Van der Wegen and Jaffe (2014)
Full tidal forcing	Dam et al. (2008), Dam et al. (2016), Ganju and Schoellhamer (2010), Ganju et al. (2011), Ganju et al. (2009), George et al. (2012), Luan et al. (2017), Styles et al. (2016), Van der Wegen and Jaffe (2013), Van der Wegen et al. (2017), Weisscher et al. (2022), Zhang and Mao (2015), Zheng et al. (2021)

220 tion, to avoid the empirical procedure required for the *morphological tide*, which intro-
 221 duces a dependency on the parameters and the locations chosen for the analysis. The
 222 aim for the procedure is to provide a synthetic signal that resembles the original tidal
 223 signal, excluding variations resulting from non-tidal processes.

224 The construction of the synthetic signal starts with a fortnightly modulation of the
 225 amplitude of the semi-diurnal tide to represent spring-neap variations. A synthetic sig-
 226 nal with the duration of a fortnight resembles more accurately the real-world amplitude
 227 and phase variation than a single or double tide. Higher harmonics of the semi-diurnal
 228 tide are included to represent the asymmetry of the tide. Diurnal tides are included, fol-
 229 lowing the method of Lesser (2009) to account for the O₁-K₁-M₂ interaction while main-
 230 taining periodicity of the signal. The synthetic signal is given by:

$$\begin{aligned}
 \zeta(t) = & \\
 & (\overline{A_{D_2}} + A_{D_{sn}} \cos(\omega_{sn}t)) \cos(\omega_{D_2}t - \phi_{D_2}) \\
 & + \overline{A_{D_4}} \cos(\omega_{D_4}t - \phi_{D_4}) \\
 & + \overline{A_{D_6}} \cos(\omega_{D_6}t - \phi_{D_6}) \\
 & + \overline{A_{D_8}} \cos(\omega_{D_8}t - \phi_{D_8}) \\
 & + \overline{A_{C_1}} \cos(\omega_{C_1}t - \phi_{C_1})
 \end{aligned} \tag{1}$$

231 where $A_{D,n}$ is the amplitude, $\omega_{D,n}$ the angular frequency, and $\phi_{D,n}$ the phase of the n^{th}
 232 tidal constituent. The angular frequency ω_{D_2} is taken equal to ω_{M_2} , and all other an-
 233 gular frequencies are an integer product or one over an integer product of this primary
 234 forcing frequency. The diurnal C₁ constituent has an amplitude of $\sqrt{2A_{O_1}A_{K_1}}$ and the
 235 phase average of ϕ_{O_1} and ϕ_{K_1} . The overbar denotes time-averaging and t is time. The
 236 amplitude of D_{sn} modulates $\overline{A_{D_2}}$ and is equal to the amplitude of the second largest peak
 237 in the semi-diurnal frequency band, which corresponds to S₂ or N₂. The length of the

238 “morphological spring-neap cycle” we introduce is given by the closest even number (de-
 239 noted by i) of D_2 cycles that fit into the length of the spring-neap period induced by M₂-
 240 S₂ interaction; exactly 28 semi-diurnal cycles. The angular frequency of the fortnightly
 241 modulation is then given by

$$\omega_{sn} = \frac{2\pi}{28T_{D_2}} \quad (2)$$

242 where T_{D_2} is the period of the D_2 constituent.

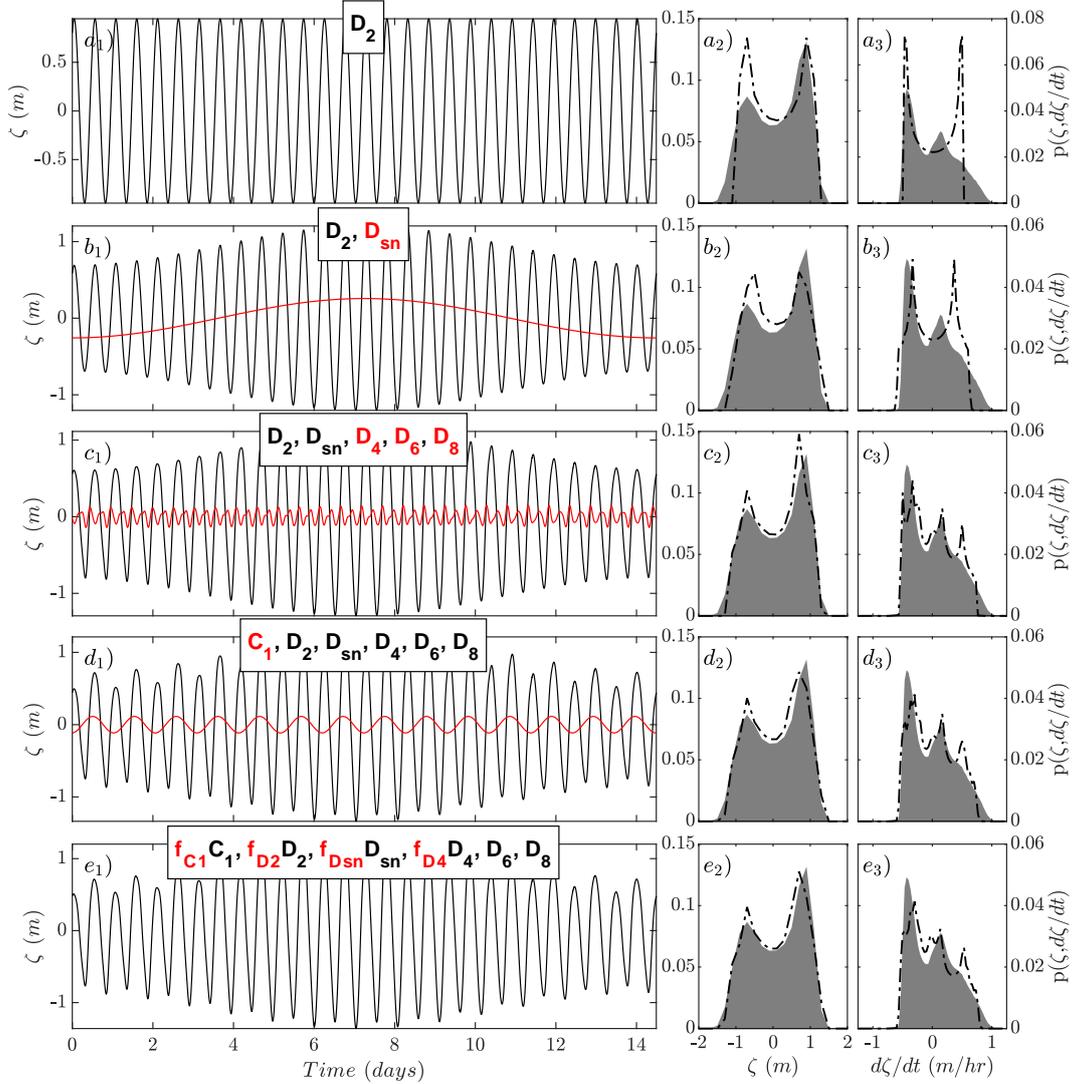


Figure 1. Step-wise construction of the synthetic spring-neap cycle, adding constituents in panel a-d, and scaling in panel e. For each step the resulting time-series (subscripted by 1) are shown in black and the added tidal constituent in red. The panels subscripted by 2 and 3 show the histograms of the synthetic signal (dashed line) and the full tidal signal (gray patch) for ζ and $d\zeta/dt$, respectively.

243 The step-wise construction of the morphological spring-neap cycle is illustrated in
 244 Figure 1, using a 19-year record of water level observations collected in the Dutch North
 245 Sea (monitoring station Wierumergronden). The synthetic signal is compared with the

246 full tidal signal using histograms of the free surface elevation (ζ) and the surface level
 247 gradient ($d\zeta/dt$). The histogram of ζ indicates asymmetry in tidal peaks, i.e. tidal peak
 248 asymmetry, and the histogram of $d\zeta/dt$ indicates asymmetry in the duration of the rising
 249 and falling limbs of the surface elevation time-series. The latter is also referred to
 250 as tidal duration asymmetry and is highly relevant for the direction and magnitude of
 251 residual bed-load transport of non-cohesive sediment (Van de Kreeke & Robaczewska,
 252 1993). This approach based on histograms concisely characterises tidal asymmetry result-
 253 ing from the interaction of all constituents, in contrast to the harmonic method that
 254 characterises the asymmetry resulting from two or more interacting constituents. The
 255 histograms in Figure 1 illustrate how the addition of the individual terms of Equation
 256 1 provide a signal that progressively better resembles the nearly complete tidal signal
 257 (reconstructed with 68 significant constituents resolved through harmonic analysis, see
 258 Pawlowicz et al. (2002)).

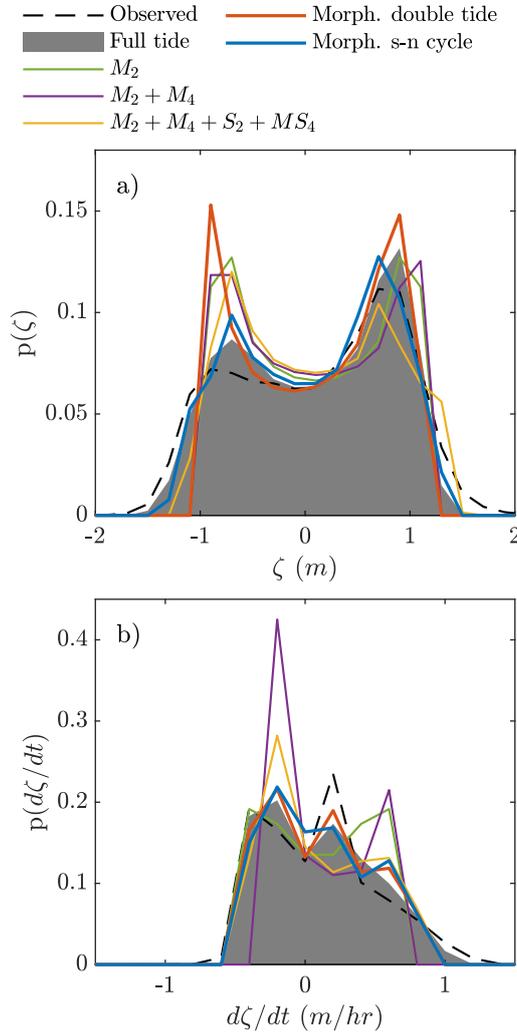


Figure 2. Histograms of ζ (a) and $d\zeta/dt$ (b) for the observed signal (dashed line), a tidal prediction including 68 resolvable tidal constituents (gray patch), and the simplified tidal signals (coloured) previously used for long-term morphological modelling. Histograms are constructed using a bin width of 0.2 m and $\frac{1}{6}$ m/hr for the the histogram of ζ and $d\zeta/dt$, respectively.

259 Applying basic trigonometry, the synthetic signal is rewritten as a linear combi-
 260 nation of sines and cosines with zero phases, which facilitates the optimisation. This equa-
 261 tion is fitted to the full astronomical tidal signal using scale factors to the amplitudes
 262 of the sines and cosines of D_2 , D_{sn} , C_1 , and D_4 (higher harmonics of D_2 are not scaled
 263 because of time efficiency in the algorithm). A combined Root-Mean-Squared-Error (RMSE)
 264 for the histogram of ζ and $d\zeta/dt$ is computed for each individual scaling factor. The er-
 265 ror values are stored in a matrix to optimise the combination of scaling factors for the
 266 amplitudes of each tidal constituent.

267 Histograms of (1) the observed water levels at monitoring station Wierumergronden,
 268 (2) water levels from a full tidal reconstruction, and (3) water levels from the syn-
 269 thetic spring-neap cycle and other simplified tidal signals are shown in Figure 2. Rep-
 270 resenting the full tide with a single M_2 constituent clearly oversimplifies the signal as this
 271 M_2 tide is completely symmetric. Although this is slightly improved by adding an M_4
 272 constituent, tidal extremes are not yet captured. These extremes are better represented
 273 when spring-neap variations ($M_2+M_4+S_2+MS_4$) are included, but the asymmetry of ζ
 274 is reversed. The morphological *double tide* represents the asymmetry of $d\zeta/dt$ well, but
 275 does not capture the extremes and asymmetry of ζ . The synthetic spring-neap cycle bet-
 276 ter approximates the extremes and asymmetries in the full tidal signal than the other
 277 simplified tides do. The synthetic signal does include, however, a third peak in the his-
 278 togram of $d\zeta/dt$, which is not present in the full tide. Apparently, this peak is suppressed
 279 by tidal constituents other than included in the simplified tide.

280 3 Numerical model

281 3.1 Model set-up

282 A numerical model is developed to quantify how various tidal reduction techniques
 283 influence the spatial variation of hydrodynamics and sediment transport. The model is
 284 set up to represent a real-world estuary rather than an idealized case, because the com-
 285 plex topography of a realistic environment introduces tidal asymmetries to be represented
 286 appropriately. For this purpose we have selected the Ems estuary, a meso-tidal system
 287 on the Dutch-German border that is part the Wadden Sea. The tidal prism is predom-
 288 inantly accommodated by a single channel that aligns with the incoming tidal wave prop-
 289 agation direction, as the tidal wave travels from west to east along the North Sea coast.
 290 The discharge of the main river draining into the estuary (the Ems river) varies between
 291 30 - 300 m^3/s , and is small compared to the flood tidal prism ($10^9 m^3$) (De Jonge et al.,
 292 2014). Other rivers discharging in the Ems estuary have a mean annual discharge that
 293 is smaller than 10 m^3/s .

294 The model is developed in the Delft3D Flexible Mesh model suite (Kernkamp et
 295 al., 2011). The numerical domain covers the offshore coastal part in the Wadden Sea,
 296 the estuary, and the river up to an up-estuary weir, with a grid cell size ranging from
 297 1 km (offshore) to 30 m (Figure 3). The model is set-up in 2D depth-averaged (2Dh) mode,
 298 with corrections for spiral motion (secondary flow) applied to the depth-averaged mo-
 299 mentum equations. Water level boundary conditions are derived from a validated hy-
 300 drodynamic model that covers the Northwest European Shelf (Zijl & Groenenboom, 2019)
 301 for the years 2018-2019. Tidal constituents at the boundaries are adjusted according a
 302 comparison between modelled and observed amplitudes and phases, derived through har-
 303 monic analysis (Pawlowicz et al., 2002) at station Wierumergronden (close to the west-
 304 ern boundary of the model - see Figure 3). A time-varying observed river discharge is
 305 prescribed at the upstream end of the Ems river for model calibration and validation,
 306 whereas a constant value (80 m^3/s for the Ems river and less than 10 m^3/s) for the smaller
 307 rivers) is prescribed for various scenario simulations. The bathymetry of the model is
 308 based on echosounding observations collected in 2014, which are made freely available
 309 by the Dutch Directorate-General for Public Works and Water Management.

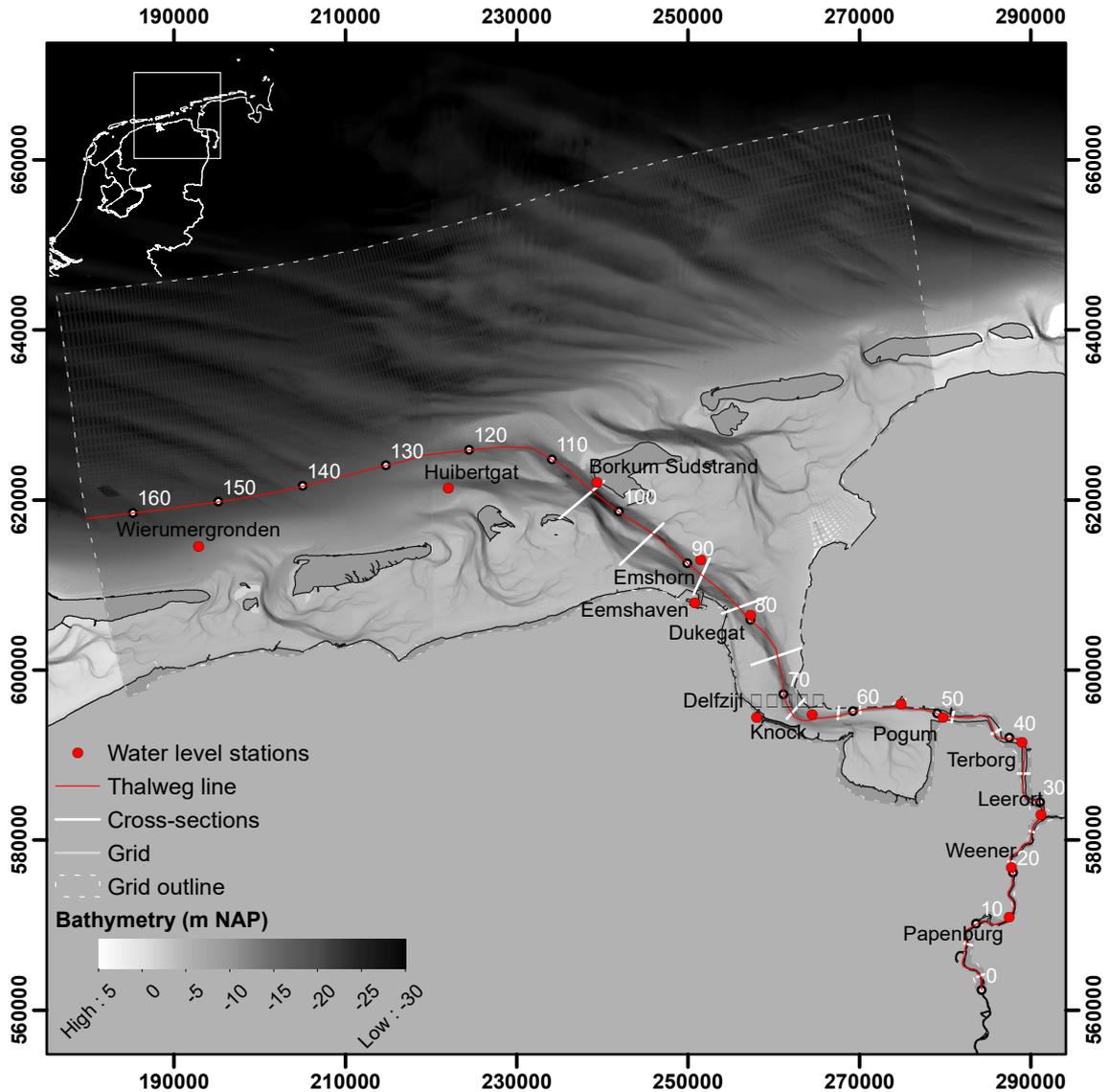


Figure 3. The Ems estuary and numerical model domain (gray lines), with the locations of water level observations (red dots) and a line that follows the main route of tidal propagation (red line) from the western boundary of the model through the thalweg of the estuary and river, with estuary kilometres defined with respect to the point of maximal tidal intrusion at the weir.

310 Sediment transport is computed with the Van Rijn (1993) formula for medium fine
 311 sand ($180 \mu\text{m}$). The model is executed in morphostatic mode (i.e. no bed update) be-
 312 cause the feedback loops initiated by bed level adaptation complicates the analysis on
 313 the direct effects of the boundary schematization on hydrodynamics and residual trans-
 314 port. An equilibrium sand concentration is prescribed at the marine model boundaries,
 315 but no sand enters the model domain through the fluvial boundaries. There is interac-
 316 tion with the bed, which has an unlimited sand supply potential. The simulated sedi-
 317 ment transports in the model can deviate significantly from the natural conditions. How-
 318 ever, the settings for the sand transport model have a limited effect on the results be-
 319 cause the simulations are used for a relative comparison between simulations with var-
 320 ious boundary conditions.

3.2 Hydrodynamic calibration and validation

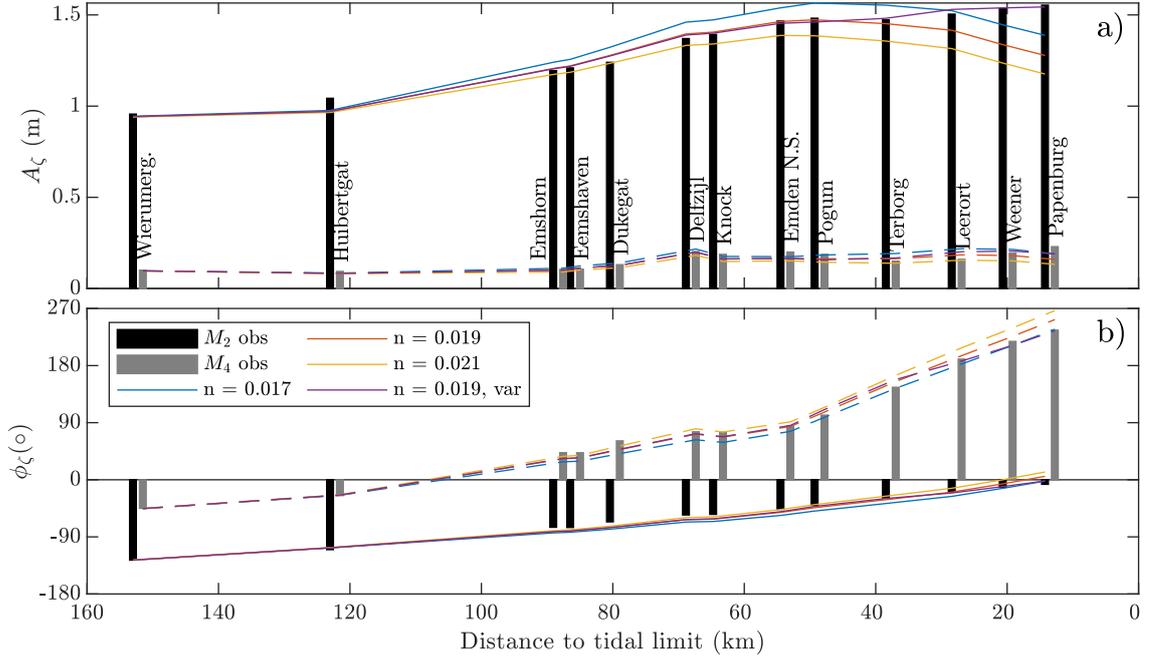


Figure 4. Observed and modelled amplitudes (a) and phases (b) of the M_2 and M_4 tidal constituents, based on the 2018 simulation. Model results (coloured lines) show the effect of different values for a spatially uniform Mannings’ n ($\text{m}^{1/3} \text{s}^{-1}$) and the best calibrated model with a spatially varying roughness in the Ems river.

Water level observations for the years 2018 - 2019 collected throughout the estuary are used to calibrate and validate the model (see Figure 3). The time-series are decomposed into tidal constituent amplitudes and phases using harmonic analysis (Pawlowicz et al., 2002). In the calibration phase, the model simulates the year 2018, using a spatially uniform roughness coefficient, Mannings’ n , amounting to 0.017, 0.019, and 0.021 $\text{m}^{1/3} \text{s}^{-1}$ (Figure 4). Tidal propagation is best represented by a Manning’s n value of 0.019 $\text{m}^{1/3} \text{s}^{-1}$. Such a bed roughness, however, overestimates dampening of the tide in the Ems river. In reality, the tides amplify as a result of extensive fluid mud deposits in the Ems River, resulting in an apparent bed roughness around 0.10 $\text{m}^{1/3} \text{s}^{-1}$ (Van Maren et al., 2015). A linear decrease in bed roughness (from 0.019 $\text{m}^{1/3} \text{s}^{-1}$ at the entrance of the river towards 0.011 $\text{m}^{1/3} \text{s}^{-1}$ at the upstream end at the weir) is therefore employed, which better represents the tidal dynamics.

The model was validated against water level observations over the first five months of 2019. The modelled amplitudes of the four primary tidal constituents (M_2 , S_2 , O_1 , K_1) and M_4 are typically within 15% of the observed amplitudes (Figure 5a). The errors are larger (up to 28%) for the S_2 and M_4 tidal constituents in the landward part of the Ems river (Figure 5b). Modelled phases are typically within 10° of observations (Figure 5c), but the modelled phases of O_1 and especially K_1 differ more than 20° in the tidal river part (Figure 5d).

The calibrated model introduced herein serves to evaluate alternative tidal input reduction approaches for morphodynamic modelling. The non-schematized tidal boundary conditions (full tidal, providing a reference condition) and alternative simplified tidal

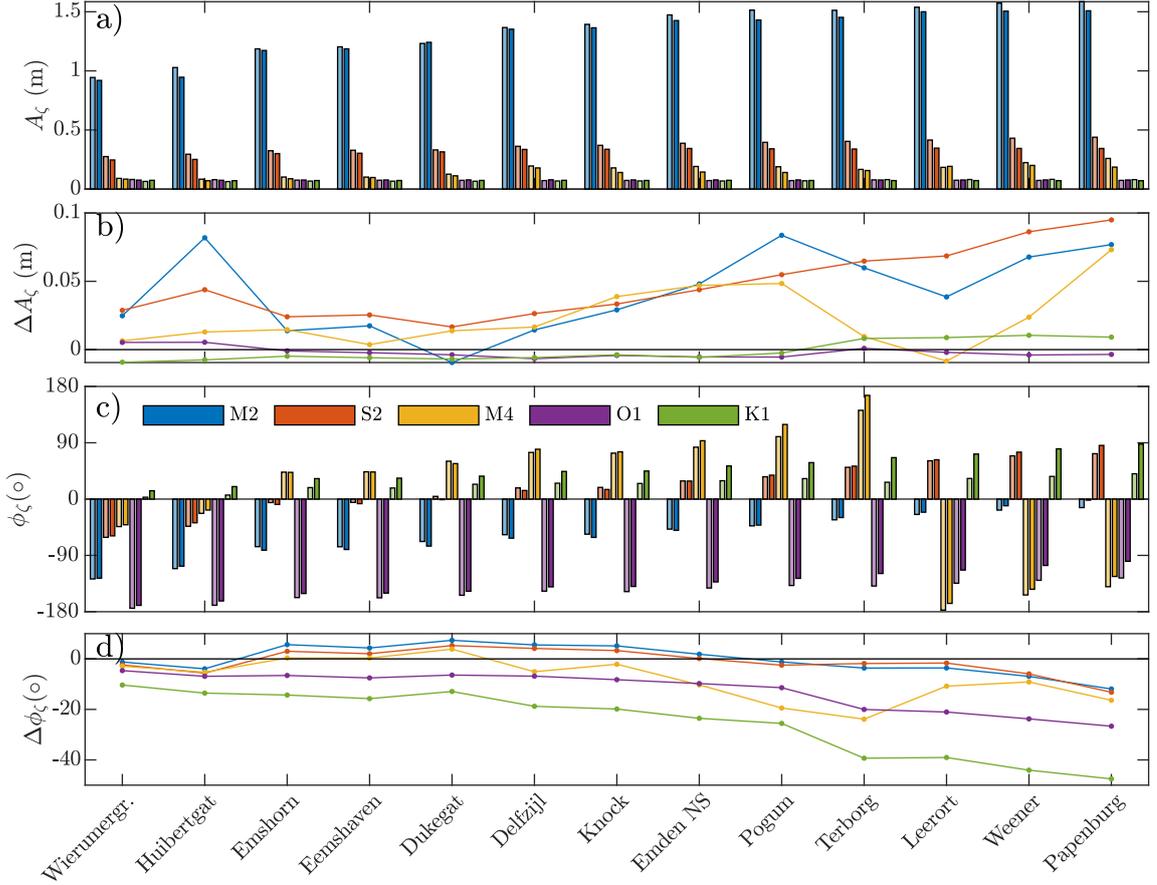


Figure 5. Observed (light coloured bars) and modelled (dark coloured bars) tidal constituent amplitudes (a) and phases (c), based on the 2019 simulation. The difference (observed - modelled) of the amplitudes and phases are shown in panels b and d, respectively.

344 representations (as in Figure 2) are detailed in Table 2. The boundary forcing with the
 345 morphological *double tide* includes an analytically derived scaling factor for M_2 (see Lesser
 346 (2009) for the derivation), to incorporate the total energy of the full tide (the sum of squares
 347 of the amplitudes of all tidal constituents) in the semi-diurnal frequency band. Applying
 348 the scaling factor, residual transports resulting from a mean (residual) flow is conserved
 349 in the simplified tide. The various tidal input reduction scenarios are compared
 350 to the reference in terms of tidal wave shape, bed shear stress, inundation, and sand trans-
 351 port in the following sections. All simulations (Table 2) are preceded by a two-week pe-
 352 riod that is excluded from the analysis to arrive at equilibrium conditions for the hydro-
 353 dynamics and suspended sediment concentrations at the start of the analysis.

354 4 Results

355 4.1 Tidal wave shape

356 The representation of tidal wave shape is a primary indicator for the error made
 357 in the simulations forced with simplified tidal conditions. Figure 6 quantifies the ade-
 358 quacy of the tidal wave shape representation based on the RMSE between the tidal re-
 359 duction scenario and the full tidal signal, for histograms of both ζ and $d\zeta/dt$. The fig-

Table 2. Duration of the simulations forced with simplified tidal signals and the full tidal simulation that serves as the reference. Simulation names are used in the legends of the figures in the results.

Simulation name	Duration
Full tidal	1 year
M_2	24 hr, 50 min
M_2M_4	24 hr, 50 min
$M_2M_4S_2MS_4$	14.77 days
Morph. double tide	24 hr, 50 min
Morph. spring-neap	14.48 days

360 ure clearly shows that only using an M_2 boundary forcing leads to the largest error. In-
 361 cluding more tidal constituents in the boundary information decreases the error and in-
 362 troducing spring-neap variations ($M_2M_4S_2MS_4$) leads to a markedly better representa-
 363 tion of tidal wave shape. The *morphological spring-neap tide* shows the smallest error,
 364 both for ζ and for $d\zeta/dt$. The improvement established by introducing spring-neap vari-
 365 ations is largest in the coastal and central parts of the estuary (km 70 - 160), because
 366 error estimates for all tidal reduction techniques converge to the same value in the up-
 367 per reaches of the estuary.

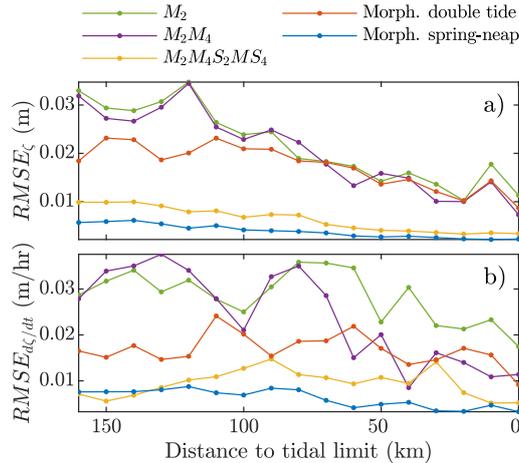


Figure 6. RMSE for the histogram of ζ (a) and $d\zeta/dt$ (b) between the simulations forced with simplified tides and the full tidal simulation, calculated at points in the thalweg along the estuary kilometres defined in Figure 3.

368 4.2 Bed shear stress

369 Maximum bed shear stress magnitudes along the estuary thalweg (Figure 7a) are
 370 most accurately represented when accounting for spring-neap variations, although there
 371 still is an underprediction of 30-40%. Including spring-neap variations gives a better rep-
 372 resentation of tidal wave shape, therefore, asymmetries are better preserved leading to
 373 higher maximum tidal velocities. The mean shear stresses in the thalweg (Figure 7b),
 374 on the other hand, are represented well by all simplified tides (although they are slightly

375 overpredicted using the morphological *double tide*). Maximum shear stresses are largest
 376 in the main tidal channels (Figure 8a) and, consequently, absolute improvements are largest
 377 in the tidal channels when accounting for spring-neap variations (compare in Figure 8
 378 panel d and f to panel b, c, and e). However, maximum shear stresses on the intertidal
 379 areas are also underpredicted, in all simulated scenarios. An analysis on the error made
 380 in representing bed shear stress magnitudes over the complete model domain (presented
 381 in Figure 8) indicates that both the maximum (Figure 9a) and the mean (Figure 9b) shear
 382 stress magnitudes improve by incorporating tidal extremes. A reduction in RMSE is found
 383 in the subtidal (channels) and intertidal parts of the model domain. The consistent over-
 384 prediction of mean bed shear stress magnitudes with the morphological *double tide* in
 385 the thalweg (Figure 7b) is reflected by larger RMSE values in the subtidal domain (Fig-
 386 ure 9b). Possibly, the overprediction is due to the implementation of a scaling factor for
 387 the M_2 tidal amplitude, to account for non-tidal energy in the spectral tidal frequency
 388 band.

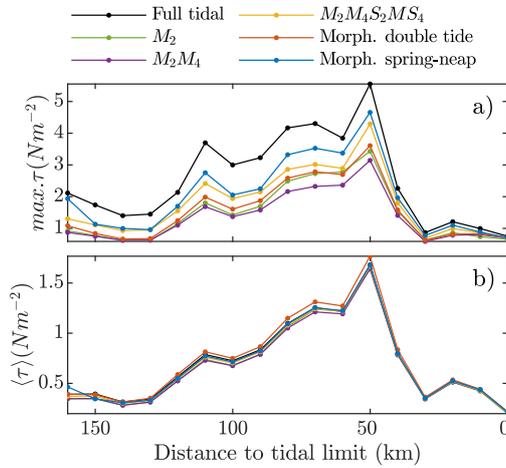


Figure 7. Maximum (a) and mean (b) bed shear stress magnitudes simulated with the full tidal forcing and simplified tides, calculated at points in the thalweg along the estuary kilometres defined in Figure 3.

389 4.3 Inundation

390 The intertidal areas, represented by computational cells that experience regular
 391 flooding and drying, make up $\approx 20\%$ of the model domain. In those areas, the duration
 392 of inundation strongly controls sediment dynamics and therefore, the residence time of
 393 water over the tidal flats (Figure 10) is an important property to capture in morpholog-
 394 ical simulations of tidal environments. Particularly the high littoral zone (Figure 10a,
 395 b) is not captured by the simulations that exclude spring-neap variations, evidenced by
 396 too many computational cells that are permanently dry. Sediment cannot settle or erode
 397 in the higher intertidal parts when those areas never inundate. The bed level height of
 398 tidal flats will not be able to adjust to a height that resembles reality. Similarly, in the
 399 low littoral zone (Figure 10e, f), the simplified signals without spring-neap variations re-
 400 sult in too many computational cells that are permanently inundated such that the lower
 401 intertidal zone becomes a subtidal area. Average conditions in the mid-littoral zone are
 402 well-represented by all simplified tides.

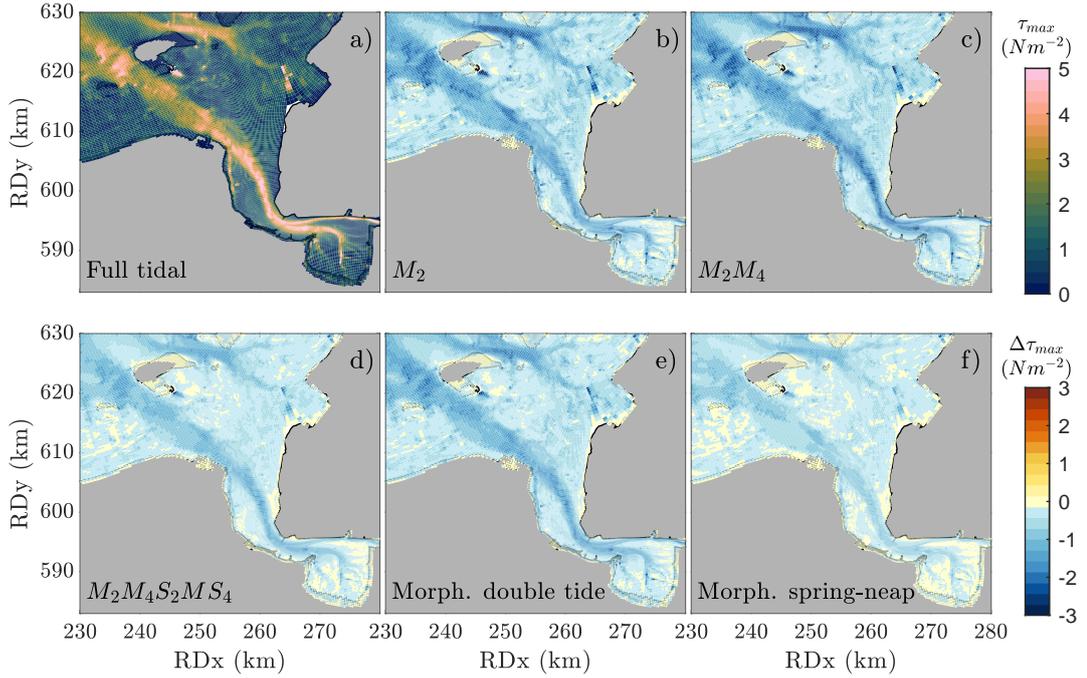


Figure 8. Maximum bed shear stress during the reference simulation (a) and difference in maximum bed shear stress between the scenarios and the reference (b-f)

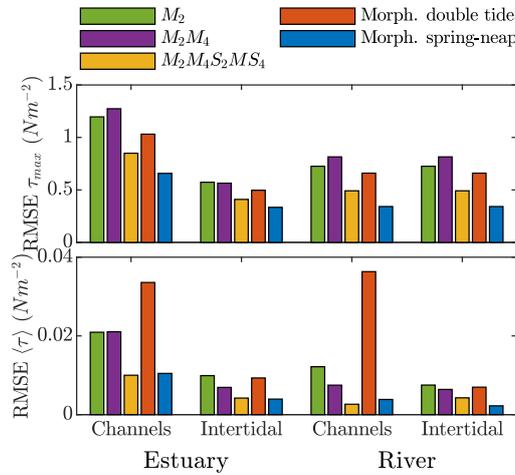


Figure 9. RMSE for the maximum (a) and mean (b) bed shear stress magnitudes between the simulations forced with simplified tides and the full tidal simulation. RMSE values are calculated as mean values for all the computational cells within the specified subregions estuary, river, subtidal channels and intertidal areas.

403

4.4 Sediment transport

404

The gross, cross-section integrated sand transport fluxes vary with each tidal cycle in the *full tidal* simulation. The mean of the range in gross transport flood fluxes (Fig-

405

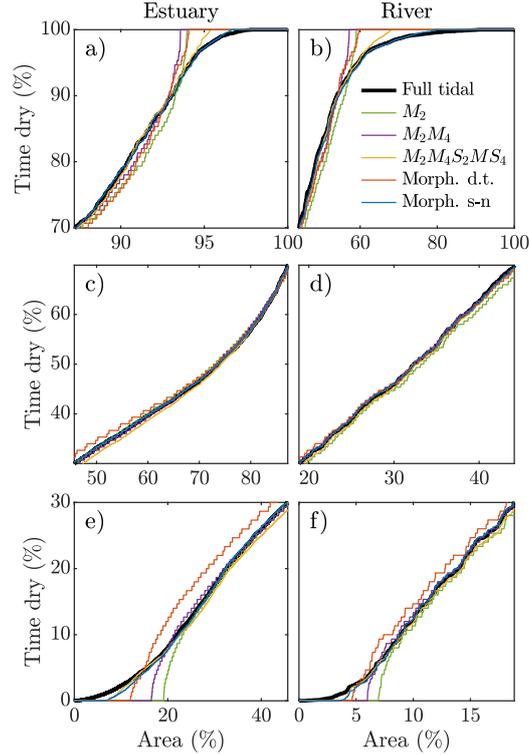


Figure 10. Cumulative distributions of the fraction of time of the total simulation length (in %) that a computational cell is dry (emerged), as a function of the fraction of the total intertidal area in the modelling domain. The distributions are shown for defined subregions; the estuary (a, c, e) and the river (b, d, f), and subdivided in the high (a, b), mid- (c, d), and lower (e, f) littoral zone.

406 ure 11a) is well-captured by the $M_2M_4S_2MS_4$ tide, the morphological *double tide*, and
 407 and the morphological *spring-neap* simulations. The M_2 tide, the *morphological double tide*,
 408 and the *morphological spring-neap* simulations all reproduce the mean gross ebb trans-
 409 ports reasonably well. For the full tidal simulation, the residual transport (Figure 11b)
 410 is flood-dominant at the mouth (km 85 - 108), ebb-dominant in the central part (km 45
 411 - 85) of the estuary; and neither flood nor ebb dominant in the tidal river (km 0 - 45).
 412 This large-scale behaviour is captured well by each of the alternative simplified tides, ex-
 413 cept for the M_2 simulation, which prescribes a perfectly symmetric tide at the sea bound-
 414 aries and therefore leads to an underestimation of the flood directed residual transport
 415 (Figure 11a). The morphological *spring-neap* tidal boundary conditions lead to resid-
 416 ual transport best representing *full tidal* residual transport (Figure 11b). The M_2M_4 and
 417 $M_2M_4S_2MS_4$ tidal boundary conditions lead to an underestimation of the magnitude of
 418 the residual transport fluxes, and the *morphological double tide* generates slightly more
 419 ebb-dominant transport in the entire estuary.

420 The morphological evolution is not only driven by the magnitude of gradients in
 421 the residual transport flux, but also by the directions. An analysis of the error (RMSE)
 422 made in the direction and magnitude of residual transports averaged over all computa-
 423 tional cells (Figure 12) reveals that particularly the error in direction is smaller for the
 424 simulations that include spring-neap variations. The RMSE for the magnitude of the resid-
 425 ual transport shows less scatter, except for the M_2 simulations, which clearly deviates

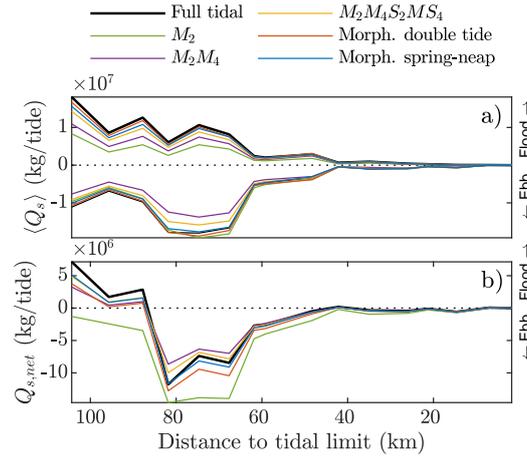


Figure 11. Mean of the total (bed + suspended) load gross transport fluxes (a) and residual transport per tidal cycle (b) in the thalweg (see the cross-sections in Figure 3 for locations).

426 in the channels. In general, including spring-neap variations reduces the error in mag-
 427 nitude and direction of residual transports in the channels and over the intertidal areas.

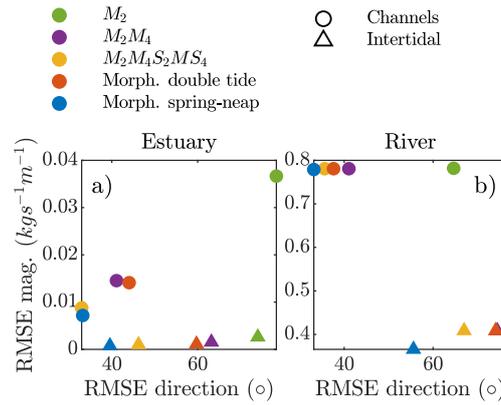


Figure 12. Error (RMSE) in the direction (horizontal axis) and magnitude (vertical axis) of the residual total (bed + suspended) load sand transport in the channels (circles) and on the intertidal areas (triangles).

5 Discussion

A new tidal input reduction method was developed which includes periodic spring-neap variation in a simplified tide. Prescribing this new method as boundary conditions in an estuarine setting improves the representation of tidal wave shape, maximum and mean bed shear stress magnitudes, inundation times, and residual sand transport patterns, compared to existing tidal input reduction methods to represent the non-schematized tidal dynamics. The strong and weak points of the new methodology and existing tidal input reduction techniques are summarised in Figure 13 using normalised scores, with 0 indicating a poor representation of the full tidal signal, and 1 indicating a full representation of the full signal. The scores are calculated as

$$score = \frac{z}{\max(z)}, z = 1 - \frac{x}{\max(x)} \quad (3)$$

The x -values for the parameters *Periodic* and *Deterministic* are binary (0 or 1), and *Cycle length* proceeds directly from Table 2. The other values for x are computed from the model scenario metrics presented in Chapter 4. An averaged RMSE between the output computed with a simplified and a full tide serves as x -value. The new method scores maximal on 10 out of 12 metrics, with lower scores only for the duration of the cycle and the duration of inundation (second-best score). Especially the scores computed from the various model scenarios strongly influence the morphodynamic evolution of a model (bed shear stress parameters, sand transport, inundation, and tidal asymmetry). When converting the model into morphodynamic mode, we therefore expect the new input reduction technique to provide physically more meaningful bed level predictions. However, as elaborated earlier, we do not explore the resulting morphodynamic impacts, which may be very case-specific and therefore cannot easily be generalized. The higher scores in Table 2 therefore motivate to replace traditional approaches for tidal input reduction with the new method.

The main drawback of the synthetic spring-neap cycle, following directly from Figure 13, is the simulation duration. The 28 M_2 cycles (≈ 14.48 days) required in the computations is 14 times longer than the time required to simulate a cycle of the morphological *double tide* (Lesser, 2009). In practice this drawback is minor, because a shorter representative tidal period (e.g. the M_2 period) is usually frequently repeated. Simulating many tidal cycles is preferred because bed elevation changes over a single tidal cycle are small compared to inaccuracy, which are then linearly amplified by a comparatively large morphological upscale factor (MF). For this reason, a single morphological tidal cycle is repeated even more often than 28 times, up to multiple hydrodynamic years (e.g. Dastgheib et al., 2008). The longest acceptable hydrodynamic simulation time is then usually combined with the smallest possible MF because (too) large values for the MF can produce unrealistic bed development (Ranasinghe et al., 2011).

Numerical morphological models may also be forced with non-tidal processes, such as a seasonally varying river discharge (e.g. Van Der Wegen et al., 2011; He et al., 2022) or wave- and wind-driven re-suspension (e.g. Van der Wegen et al., 2017). Such non-tidal conditions are typically accelerated by a factor MF as well (i.e., an annual river flood recurs MF times per year). In these cases, the relative phasing of the various forcing factors with the tide need to be explicitly accounted for as well. Otherwise, for instance, persistently combining seasonal river floods or storm events with spring tide or flood conditions leads to biased bed development. In tidal series the M_2 - S_2 phase differences throughout a spring-neap cycle differ for successive spring-neap cycles, which also holds for the phase differences between semi-diurnal and diurnal tides. In the synthetic spring-neap cycle, the relative phasing is identical for successive spring-neap cycles, which allows to optimize the relative phasing with non-tidal processes.

476 A tide-averaged transport for coarse sediment is generated by a representative tide
 477 consisting of a tide-induced Eulerian mean current (M_0), M_2 and any of its even over-
 478 tides (Van de Kreeke & Robaczewska, 1993). When diurnal components are important,
 479 a similar net residual transport arises from the triad interaction of M_2 - K_1 - O_1 (Hoitink
 480 et al., 2003), which can be captured in a periodic double tide through an artificial di-
 481 urnal component with half the frequency of M_2 (Lesser, 2009). Spring-neap variations
 482 are so far mainly ignored in representative tides (Dastgheib et al., 2008; Roelvink & Re-
 483 niers, 2011). This paper demonstrates that simplified tides consisting of a single or a dou-
 484 ble tide (which are most frequently used for long-term morphological modelling) do per-
 485 form well in representing mean bed shear stress and residual sand transports inside the
 486 estuarine channels. However, they fail to reproduce the full range of asymmetry in the
 487 tide leading to an underestimation of maximum bed shear stresses (controlling the timescales
 488 of bed level adaptation) and to represent the upper and lower intertidal inundation that
 489 steers the development of intertidal flats (Friedrichs, 2011). Representing the variation
 490 in tidal asymmetries is shown to be important to capture the residual sand transports
 491 on the intertidal flats as well. This is because the velocity skew (flood versus ebb dom-
 492 inance) over tidal flats is modulated during the spring-neap cycle (Nidziedo & Ralston,
 493 2012). Therefore, if the intertidal areas are of insignificant importance in the environ-
 494 ment studied and the focus of the study is on the tidal channels, ignoring spring-neap
 495 variations in a representative tide is presumably allowed. However, if the intertidal parts
 496 of the modelling domain are an integral part of the phenomena studied, morphodynamic
 497 models cannot suffice with a representative tide consisting of a *single* or *double tide*.

498 Applying the synthetic spring-neap cycle in a fully coupled morphodynamic model
 499 (including bed level adaptations) leads to much more realistic tidal dynamics. The com-
 500 puted residual sand transport will improve by including the tidal extremes and asym-
 501 metries resulting from the spring-neap modulations, promoting a more realistic chan-
 502 nel transport and channel-shoal exchange. The inclusion of tidal extremes may also have
 503 negative effects, however. The resulting higher maximum bed shear stresses possibly lim-
 504 its the morphological acceleration factor, which can otherwise lead to unrealistic bed level
 505 developments. Such a potential shortcoming depends on various model settings (Reyns
 506 et al., 2014), and requires a case-specific analysis. Furthermore, the gross and net sand
 507 transport presented in this paper was based on simulations with a single fraction sed-
 508 iment bed existing of non-cohesive sediments, calculated with the Van Rijn (1993) for-
 509 mula. Multiple fraction sediment beds, including cohesive sediments, may develop un-
 510 expected interactions in conjunction with the synthetic spring-neap cycle which needs
 511 to be explored in a practical case. Planned long-term morphodynamic modelling will re-
 512 veal the advantages and the challenges of the more realistic representation of tidal dy-
 513 namics advocated in this paper.

514 6 Conclusions

515 Spring-neap variations can be included in simplified tidal signals that are applica-
 516 ble as boundary conditions in long-term morphological models. Compared to a single
 517 or a double tide, often used in morphodynamic simulations, tidal variation in a synthetic
 518 spring-neap cycle is better represented through a fortnightly modulation on the ampli-
 519 tude of the semi-diurnal tide. The tidal input reduction method developed in this pa-
 520 per yields a signal that: (1) resembles the amplitude variation of the full tidal signal and
 521 sufficiently preserves asymmetries to approach non-schematized tidal dynamics and resid-
 522 ual sand transports; (2) is strictly periodic; and (3) can readily be derived from the full
 523 boundary conditions. It does not require a fitting procedure based on modelling results.

524 Process-based numerical models of tidal environments that include the tidal extremes
 525 induced by spring-neap variations represent the shape of the tidal wave through the tidal
 526 basin more realistically. Simulations with simplified tidal signals that neglect the tidal
 527 extremes underestimate maximum bed shear stresses in the channels and simulate a too

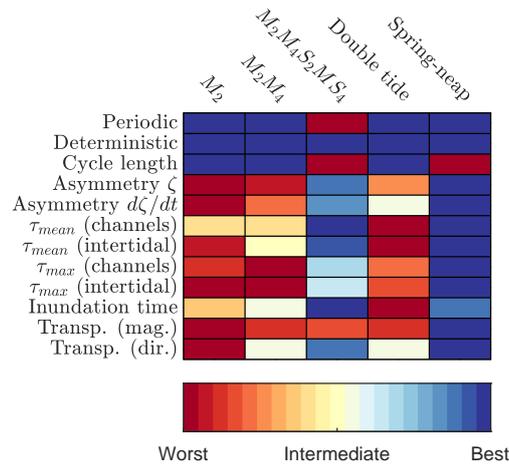


Figure 13. Normalised scores (0-1) for simplified tides to represent (simulated) non-schematized tidal conditions. The calculation of the score values is explained in the main text.

528 limited extent of the tidal flats. Although simulations forced with these signals approx-
 529 imate the tidally averaged residual sand transport patterns in the channels quite rea-
 530 sonably, an appropriate representation of the extremes is required to reproduce the pat-
 531 terns both in the channels and on the intertidal areas. The newly developed tidal input
 532 reduction method provides a signal that resolves non-cohesive sediment transport within
 533 the estuary more accurately, and may improve the simulated exchange of sediment be-
 534 tween the channels and tidal flats.

535 7 Open Research

536 A toolbox is developed that allows to construct a synthetic spring-neap tidal cy-
 537 cle from a time-series of tidal elevations. The toolbox is developed in MATLAB code,
 538 and available for download at [https://github.com/Rschrijvershof/morphoSpringNeap](https://github.com/Rschrijvershof/morphoSpringNeap.git)
 539 [.git](https://github.com/Rschrijvershof/morphoSpringNeap.git).

540 There are no restrictions on the data used in this study. The bathymetry data used
 541 for model set-up was requested through the servicedesk data of Rijkswaterstaat ([https://](https://www.rijkswaterstaat.nl/formulieren/contactformulier-servicedesk-data)
 542 www.rijkswaterstaat.nl/formulieren/contactformulier-servicedesk-data). Ob-
 543 served water level data from Dutch monitoring stations are available at [https://waterinfo](https://waterinfo.rws.nl)
 544 [.rws.nl](https://waterinfo.rws.nl) and the data from the German monitoring stations was requested at WSA Ems-
 545 Norsee (<https://www.wsa-ems-nordsee.wsv.de/>). The configurations of the numer-
 546 ical model simulations used in this article are stored at 4TU.ResearchData ([https://](https://doi.org/10.4121/19845262.v1)
 547 doi.org/10.4121/19845262.v1).

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 554 Wasserwirtschaft, Küsten- und Naturschutz (NLWKN). The modelling software Delft3D
 555 Flexible Mesh was made available by Deltares via the Delft3D Educational Service Pack-
 556 age.

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