

Estimating satellite orbital drag during historical magnetic superstorms

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Key Points:

- Historical magnetic superstorms (minimum Dst ≤ -500 nT) occurred before CHAMP's and GRACE's operation times
- Standard Dst and equivalent Dst data are used by the JB2008 empirical model to estimate extreme orbital drag effects
- Interplay between storm-time duration and minimum Dst and Dst-like values determine the severity of satellite drag effects in low-Earth orbit

Abstract

Understanding extreme space weather events is of paramount importance in efforts to protect technological systems in space and on the ground. Particularly in the thermosphere, the subsequent extreme magnetic storms can pose serious threats to low-Earth orbit (LEO) spacecraft by intensifying errors in orbit predictions. Extreme magnetic storms (minimum Dst ≤ -250 nT) are extremely rare: only 7 events occurred during the era of spacecraft with high-level accelerometers such as CHAMP (CHALLENGE Mini-satellite Payload) and GRACE (Gravity Recovery And Climate experiment), and none with minimum Dst ≤ -500 nT, here termed magnetic superstorms. Therefore, current knowledge of thermospheric mass density response to magnetic superstorms is very limited. Thus, in order to advance this knowledge, 4 historical magnetic superstorms, i.e., events occurring before CHAMPs and GRACEs commission times, are used to empirically estimate density enhancements and subsequent orbital drag. The November 2003 magnetic storm (minimum Dst = -422 nT), the most extreme event observed by both satellites, is used as the benchmark event. Results show that, as expected, orbital degradation is more severe for the most intense storms. Additionally, results clearly point out that the time duration of the storm is strongly associated with storm-time orbital drag effects, being as important as or even more important than storm intensity itself. The most extreme storm-time decays during CHAMP/GRACE-like sample satellite orbits estimated for the March 1989 magnetic superstorm show that long-lasting superstorms can have highly detrimental consequences for the orbital dynamics of satellites in LEO.

Plain Language Summary

We investigate drag effects on satellites orbiting Earth in its upper atmosphere during magnetic storms caused by the impacts of solar superstorms. During magnetic storms, the upper atmosphere is heated and expands upwards, resulting in increased drag forces on satellites flying in those regions. Enhanced drag effects directly impact operations of such spacecraft, for instance, orbital tracking and predictions, maneuvers, and lifetime maintenance. The U.S. Federal Government has recognized space weather phenomena as natural hazards, and the understanding of their consequences, particularly during extreme circumstances, is of paramount importance. The very extreme events, here termed magnetic superstorms, occurred before the space era when no in-situ observations of the atmospheric density are available. Therefore, we use an empirical model to estimate drag from these historical events. Results generally show that the most extreme events drive the most severe effects. Additionally, we show that another storm feature, its time duration, can play a significant role in enhancing drag. Therefore, we argue that space weather forecasters should be aware of events with long duration, particularly the ones caused by sequential impacts of solar disturbances on the Earth's magnetic field, when predicting and forecasting the subsequent drag effects on satellites in the upper atmosphere.

1 Introduction

Magnetic storms are global phenomena that occur due to the interaction of solar perturbations with the Earth's magnetosphere (Gonzalez et al., 1994). The most intense and severe magnetic storms are commonly caused by coronal mass ejections (CMEs) (Gonzalez et al., 1994; Daglis et al., 1999; Balan et al., 2014). CMEs usually have a shock at their leading edge that is promptly followed by a sheath and a magnetic cloud (Gonzalez et al., 1994; Balan et al., 2014; Kilpua et al., 2019). Extreme magnetic storms are caused by the impact of extremely fast CMEs on the Earth's magnetosphere (Tsurutani & Lakhina, 2014), usually associated with highly depressed values of the southward component of the interplanetary magnetic field (Gonzalez et al., 1994; Daglis et al., 1999; Balan et al., 2014; Tsurutani & Lakhina, 2014; Kilpua et al., 2019).

68 Extreme space weather events like severe magnetic storms have been recognized
 69 by the U.S. Federal Government through the National Space Weather Strategy and Ac-
 70 tion Plan (National Science and Technology Council, 2015b, 2015a) as a natural hazard,
 71 and the need to establish benchmarks for extreme space weather events has also been
 72 recognized by the scientific community (e.g., Lanzerotti, 2015; Jonas et al., 2017; Riley
 73 et al., 2018). The intensity of magnetic storms is usually measured by depletions of the
 74 ground horizontal magnetic field component recorded by magnetometers located at mid-
 75 and low-latitudes by means of the disturbance storm time (Dst) index (section 2.1). Ex-
 76 tremely severe events, here termed magnetic superstorms, with minimum Dst ≤ -500
 77 nT, are notably rare (Cliver & Dietrich, 2013; Riley et al., 2018; Vennerstrøm et al., 2016;
 78 Hayakawa, Ebihara, Willis, et al., 2019; Chapman et al., 2020). For instance, the March
 79 1989 event, the only superstorm occurring during the space age (Meng et al., 2019), is
 80 well-known for the occurrence of low-latitude aurorae (Allen et al., 1989; Rich & Denig,
 81 1992; Pulkkinen et al., 2012) and intense geomagnetically induced currents (GICs) which
 82 caused the blackout of the Hydro-Québec system in Canada for several hours, leading
 83 to serious economic losses (Bolduc, 2002; Kappenman, 2006; Pulkkinen et al., 2017). How-
 84 ever, though arguably, the most extreme ground horizontal magnetic field perturbation
 85 (~ -1600 nT) on record was recorded by the Colaba station during the Carrington event
 86 of September 1859 (Tsurutani et al., 2003; Siscoe et al., 2006; Hayakawa, Ebihara, Willis,
 87 et al., 2019). Since that is the only known low-latitude data set available to date, a global
 88 analysis of that storm cannot be performed (Siscoe et al., 2006; Cliver & Dietrich, 2013;
 89 Hayakawa, Ebihara, Willis, et al., 2019; Blake et al., 2019). For this reason, the Carring-
 90 ton event is not addressed in this paper.

91 During active times, large amounts of electromagnetic energy enter the ionosphere-
 92 thermosphere system causing the prompt thermosphere heating and upward extension
 93 due to mechanical collisions between ions and neutrals (e.g., Prölss, 2011; Emmert, 2015).
 94 This energy has access to the thermosphere primarily through high latitudes (Fuller-Rowell
 95 et al., 1994; Liu & Lühr, 2005; Huang et al., 2014; Connor et al., 2016; Lu et al., 2016;
 96 Kalafatoglu Eyiguler et al., 2018), and propagates equatorward due to the occurrence
 97 of gravity waves and wind surges (Fuller-Rowell et al., 1994; Hocke & Schlegel, 1996; Bru-
 98 insma & Forbes, 2007; Sutton et al., 2009). Therefore, the heating and upwelling of the
 99 thermosphere are global phenomena (Richmond & Lu, 2000; Liu et al., 2005; Sutton et
 100 al., 2009). As a result, satellites that happen to fly in those regions experience increased
 101 effects of drag forces leading to their subsequent orbital degradation (Prölss, 2011; Pri-
 102 eto et al., 2014; Zesta & Huang, 2016). The understanding and control of orbital drag
 103 effects during active times can enhance predictability and forecasting of satellite track-
 104 ing, reentry processes, and maintenance of satellite life times (Prölss, 2011; Zesta & Huang,
 105 2016; Berger et al., 2020), particularly during extreme magnetic storms (Oliveira & Zesta,
 106 2019). Most of these studies have used data obtained from state-of-the-art accelerom-
 107 eters onboard two low-Earth orbit (LEO) satellites, namely CHAMP (CHALLENGE Min-
 108 isatellite Payload; Reigber et al., 2002) and GRACE (Gravity Recovery And Climate Ex-
 109 periment; Tapley et al., 2004). These spacecraft were launched after 2001 (section 2.2).

110 The most extreme magnetic storm experienced by CHAMP and GRACE took place
 111 in November 2003 with minimum Dst = -422 nT. Consequently, there are no assessments
 112 of satellite drag in LEO during magnetic superstorms inferred from high-accuracy ac-
 113 celerometer data. The orbital degradations of CHAMP and GRACE associated with the
 114 November 2003 event during stormy times were, respectively, ~ -285 m and ~ -71 m
 115 (Krauss et al., 2015; Oliveira & Zesta, 2019), much more severe than the natural drag
 116 caused by the quiet-time background density estimated by Oliveira and Zesta (2019),
 117 namely -34.45 m and -6.86 m, respectively. Hence, these are the most extreme orbital
 118 decays measured with high-quality accelerometer data. In order to empirically estimate
 119 drag effects during magnetic superstorms, standard Dst data and ground magnetome-
 120 ter data of historical superstorms reconstructed from historical archives are used by a
 121 thermospheric empirical model (section 2.3) for density computations (section 2.4). These

122 events occurred in March 1989 (Allen et al., 1989; Boteler, 2019), with the traditional
 123 Dst index available, September 1909 (Silverman, 1995; Hayakawa, Ebihara, Cliver, et al.,
 124 2019), May 1921 (Silverman & Cliver, 2001; Hapgood, 2019), and October/November
 125 1903 (Lockyer, 1903; Ribeiro et al., 2016), with an alternative version to the Dst index
 126 available. Effects of storm time duration associated with minimum values of Dst and Dst-
 127 like data will be estimated and compared. As a result, this effort will improve our un-
 128 derstanding of severe satellite orbital drag effects in LEO caused by magnetic superstorms.

129 2 Data, model, and a framework for orbital drag estimations

130 2.1 Disturbance storm time indices

131 In this study, magnetic activity is represented by the Dst index provided by the
 132 World Data Center for Geomagnetism, Kyoto et al. (2015). This 1-hr-resolution index
 133 was defined in 1957, the International Geophysical Year (IGY), as described by Sugiura
 134 (1964). Specifically, Dst is computed by averaging latitudinally weighted horizontal mag-
 135 netic field perturbations, with a background removal scheme, recorded by mid- and low-
 136 latitude stations with reasonably even longitudinal separation according to the expres-
 137 sion

$$Dst = \frac{1}{4} \sum_{i=1}^4 \frac{\Delta H_i}{\cos \Lambda_i}, \quad i \text{ in } [\text{HON, SJG, HER, KAK}] \quad (1)$$

138 where ΔH_i is the horizontal magnetic perturbation of the i -th station, and Λ_i is the con-
 139 temporary magnetic latitude of the i -th station. The colored stars in Figure 1 show the
 140 stations, with their corresponding names, abbreviations, and geographic locations, used
 141 to compute standard Dst after the IGY.

142 Additionally, recent efforts have been undertaken to provide alternative (but sim-
 143 ilar) versions to the standard Dst index for historical magnetic superstorms with archival
 144 material. The events took place in October/November 1903 (Hayakawa et al., 2020), Septem-
 145 ber 1909 (Love et al., 2019b), and May 1921 (Love et al., 2019a). This alternative in-
 146 dex, also with resolution of 1 hr, was reconstructed with data obtained from 4 low/mid-
 147 latitude stations, with the best possible longitudinal separation, and is represented here
 148 by Dst^\dagger . The corresponding contemporary magnetic latitudes were computed by the au-
 149 thors. A background removal scheme similar to the one used to calculate Dst is used in
 150 the source papers as well. The stations used to compute Dst^\dagger used in this study are shown
 151 by the colored crosses in Figure 1. Therefore, the Dst^\dagger index is given by

$$Dst^\dagger = \frac{1}{4} \sum_{j=1}^4 \frac{\Delta H_j}{\cos \Lambda_j}, \quad j \text{ in } \begin{cases} [\text{CLA, COI, CUA, ZKW}] & \text{for Oct/Nov 1903} \\ [\text{API, MRI, SFS, VQS}] & \text{for Sep 1909} \\ [\text{API, SFS, VSS, WAT}] & \text{for May 1921} \end{cases} \quad (2)$$

152 The Dst^\dagger data for the magnetic superstorms used here are available as supporting
 153 information provided by the respective references (Hayakawa et al., 2020; Love et al., 2019b,
 154 2019a). Details of individual stations and magnetograms for each corresponding Dst^\dagger
 155 network is provided in the source articles.

156 2.2 Neutral mass density data

157 CHAMP and GRACE neutral mass density (ρ) data obtained from their respec-
 158 tive high-accuracy accelerometers are used in this work. CHAMP was launched in 2001
 159 at the initial altitude 456 km and orbital inclination 87.25° . It covered each 1 hr local
 160 time in 5.5 days with orbital period 90 min. The GRACE-A and -B spacecraft were launched

Figure 1. Geographic locations of the ground magnetometer stations that compose the standard Dst network that has been used by the World Data Center for Geomagnetism, Kyoto et al. (2015) since 1957 (colored stars), and the alternative Dst[†] network used by Hayakawa et al. (2020), Love et al. (2019b), and Love et al. (2019a) for the historic events of October/November 1903, September 1909 and May 1921 (colored crosses), respectively. Magnetic latitudes (solid cyan lines) and magnetic equator (solid orange line) were computed by the Altitude-Adjusted Corrected Geomagnetic Coordinates Model (Shepherd, 2014) for 1957. Note that the SJG and VQS stations are very close to each other.

161 in 2002 at the initial altitude 500 km and orbital inclination 89.5°. The GRACE con-
 162 stellations covered each 1 hr local time in 6.7 days with orbital period 95 min. GRACE-
 163 A flew ~220 km ahead of GRACE-B. As discussed in Oliveira and Zesta (2019), only
 164 GRACE-A data are used, henceforth GRACE data, because GRACE-A data show higher
 165 quality than GRACE-B data. CHAMP re-entered in 2010, while GRACE re-entered in
 166 2018. Uncertainties and calibration techniques of both missions have been discussed by
 167 many papers (e.g., Bruinsma et al., 2004; Doornbos & Klinkrad, 2006; Flury et al., 2008).

168 The density data used in this study are normalized and intercalibrated as described
 169 in Oliveira et al. (2017) and Zesta and Oliveira (2019). Basically, the Jacchia-Bowman
 170 2008 (hereafter JB2008, Bowman et al., 2008, see below) empirical model computes quiet-
 171 time densities (ρ_0) in order to obtain the background state for the quiet thermosphere.
 172 This approach ensures that the ratio and the difference between the storm-time and quiet-
 173 time densities are as close to one ($\rho/\rho_0 \approx 1$) and zero ($\rho - \rho_0 \approx 0$) as possible, respec-
 174 tively. As a result, storm-time density enhancements can be extracted more effectively
 175 (Oliveira et al., 2017; Oliveira & Zesta, 2019; Zesta & Oliveira, 2019).

176 2.3 The Jacchia-Bowman 2008 (JB2008) empirical model

177 The JB2008 empirical model computes thermospheric density from a single param-
 178 eter, the exospheric temperature (see equation 2 in Oliveira & Zesta, 2019). This tem-
 179 perature depends on several satellite parameters such as latitude, local time, and alti-
 180 tude. Additionally, this model uses the solar radio flux at wavelength 10.7 cm, indicated
 181 by the F10.7 index, to account for thermospheric heating due to solar UV radiation (Bowman
 182 et al., 2008). Finally, a term that depends on Dst in the exospheric temperature repre-
 183 sents the magnetic activity contribution. Dst and Dst[†] data of the historical magnetic
 184 superstorms recorded by the stations shown in Figure 1 will be used along with LEO satel-
 185 lite orbital data during the event of November 2003 to estimate drag effects. A descrip-
 186 tion of the JB2008 model along with other popular thermospheric empirical models has
 187 recently been provided by He et al. (2018).

188 2.4 Orbital drag computations

189 Neutral mass densities are derived by high-accuracy accelerometers according to
 190 the drag equation (Prieto et al., 2014):

$$a_d = -\frac{1}{2}\rho C_D \frac{S}{m} V^2 \quad V = |\vec{V}_{s/c} - \vec{V}_{wind}|, \quad (3)$$

191 where a_d is the spacecraft acceleration caused by drag forces; ρ is the local ther-
 192 mospheric neutral mass density; C_D is the drag coefficient; S/m is the area-to-mass ra-
 193 tio; and V is the relative velocity between the spacecraft velocity ($\vec{V}_{s/c}$) and the ambi-

Figure 2. CHAMP (left-hand-side column) and GRACE (right-hand-side column) orbits, in magnetic coordinates, for the northern hemisphere (top row) and southern hemisphere (bottom row). The colorbars represent the corresponding altitudes during the time interval 19-23 November 2003, the benchmark event chosen for this study. The grey arrows in all panels indicate CHAMP's and GRACE's trajectories in both hemispheres.

ent neutral wind velocity (\vec{V}_{wind}). In this equation, all quantities are presumably known, and therefore it is solved for ρ in order to yield density. However, these parameters (particularly C_D) can introduce significant errors in density computations (Moe & Moe, 2005; Prieto et al., 2014; Zesta & Huang, 2016). In this study, drag coefficients computed with error mitigation methods by Sutton (2009) were used.

Chen et al. (2012) provide the following expression for the computation of storm-time orbital decay rate:

$$\frac{da}{dt} = -C_D \frac{S}{m} \sqrt{GM\langle a \rangle} \Delta\rho, \quad (4)$$

with a being the semi-major axis of the satellite orbit (temporal Earth's radius plus satellite altitude, Oliveira & Zesta, 2019), $G = 6.67 \times 10^{-11} \text{ m}^3 \cdot \text{kg}^{-1} \cdot \text{s}^{-2}$ the gravitational constant, $M = 5.972 \times 10^{24} \text{ kg}$ the Earth's mass, and $\Delta\rho$ the difference between the modeled storm-time and quiet-time densities. The daily average of the semi-major axis a is represented by $\langle a \rangle$.

Finally, the storm-time orbital decay is computed by the sum over all da/dt values along the satellite's path for any (t_1, t_2) interval:

$$d(t) = \int_{t_1}^{t_2} a(t) dt \quad (5)$$

3 Results

3.1 The selected magnetic superstorms

The benchmark event for the current study occurred in November 2003. That storm had minimum Dst = -422 nT, the most intense magnetic storm event with both CHAMP and GRACE neutral mass density data available. Ground magnetometer data and neutral mass density data for the GRACE satellite are shown in Figure 1 of Zesta and Oliveira (2019). The solar flux F10.7 index increased from 151 sfu (solar flux units) on 19 November to 175 sfu on 23 November (Liu & Lühr, 2005).

Figure 2 documents the orbits of CHAMP and GRACE in the time interval from 19 to 23 November 2003. The dial plots show orbits as a function of magnetic latitudes (MLATs) and magnetic local times (MLTs). The left column shows altitudes for CHAMP, while the right column shows altitudes for GRACE. The top row indicates data for the northern hemisphere, while the bottom row indicates data for the southern hemisphere. The colorbars indicate altitudes for both satellites in the same periods.

CHAMP is in a near noon-midnight orbit. The orbit altitudes of CHAMP increased at high latitudes and at the poles of both hemispheres and decreased at mid- and low-latitudes. Similar behavior is shown by GRACE whose orbits were confined within the

Figure 3. Ground magnetometer Dst and Dst[†] time series, with resolution of 1 hr, for the storms of (a) October/November 1903 (Dst[†], Hayakawa et al., 2020); (b) May 1921 (Dst[†], Love et al., 2019a); (c) March 1989 (Dst, World Data Center for Geomagnetism, Kyoto et al., 2015); and (d) September 1909 (Dst[†], Love et al., 2019b). The highlighted regions correspond to the time span between storm sudden commencement (SSC, vertical dashed lines) and the beginning of the storm recovery phases (minimum Dst or Dst[†]), or time duration of storm development.

225 mid-noon/dusk and mid-midnight/dawn sectors. Therefore, both spacecraft provide rea-
 226asonable coverage between the day and night sectors. The altitude variations shown in
 227Figure 2 caused by density variations at different MLATs and MLTs are mitigated by
 228the density intercalibration method introduced by Oliveira et al. (2017).

229 CME leading edges are usually associated with the occurrence of positive jumps
 230in the Dst index, while its sudden depression is associated with the arrival of CME mag-
 231netic material or sheaths (e.g., Gonzalez et al., 1994; Kilpua et al., 2019). The first per-
 232turbation, termed storm sudden commencement (SSC), is caused by the shock compres-
 233sion (e.g., Oliveira et al., 2018; Shi et al., 2019), while the second event, termed storm
 234main phase, is associated with strong driving of the magnetosphere via magnetic recon-
 235nection (e.g., Gonzalez et al., 1994; Daglis et al., 1999; Kilpua et al., 2019). Examples
 236of SSCs and storm main phases represented by the Dst and Dst[†] indices during magnetic
 237superstorms caused by fast CMEs are illustrated in Figure 3.

238 Figure 3 shows ground magnetometer time series for the magnetic superstorms of
 239(a) October/November 1903 (Dst[†]); (b) May 1921 (Dst[†]); (c) March 1989 (Dst); and (d)
 240September 1909 (Dst[†]). Data are plotted 12 hr and 72 hr around each respective SSC
 241(dashed vertical black lines). Times are shown as Greenwich Mean Time (GMT) for all
 242events, except as Universal Time (UT) for the 1989 event because UTs were introduced
 243only in 1928 (Hapgood, 2019). Given the similarities of UTs and GMTs, here they will
 244be used interchangeably (Hapgood, 2019). The highlighted areas of each panel correspond
 245to the time interval between SSC and minimum Dst/Dst[†] occurrences, which also marks
 246the beginning of the storm recovery phase. This time interval will henceforth be referred
 247to as the storm development duration time in this paper.

248 Panels (a) and (b) show that the 1903 event is the weakest (minimum Dst[†] = -513
 249nT), whilst the 1921 event is the strongest (minimum Dst[†] = -907 nT) amongst all events.
 250In contrast, the development duration times of both events are almost the same, ~ 14
 251hr and ~ 12 hr, respectively. Storm strengths can be estimated by computing how fast
 252Dst (or Dst[†]) is depressed during storm development. The average slope of Dst/Dst[†] dur-
 253ing the development phase is quantified by the difference of Dst/Dst[†] minimum minus
 254Dst/Dst[†] peak at SSC compression by the development time. This provides a quantifi-
 255able measure of the impactfulness of the storm, meaning that storms with very low am-
 256plitude rates are commonly associated with high geomagnetic activity (e.g., Gonzalez
 257et al., 1994). The estimated amplitude rates are -44.8 nT/hr and -80.0 nT/hr for the
 258October/November 1903 and May 1921 events, respectively. These numbers explain why
 259the effects of the 1921 event, such as equatorial extent of low-latitude aurorae (Chree,
 2601921; Silverman & Cliver, 2001), and GIC impacts on telegraph systems (Kappenman,
 2612006; Hapgood, 2019) were more severe than the effects of the 1903 event, mostly rep-
 262resented by mid-latitude aurorae (Page, 1903; Hayakawa et al., 2020), and local GIC im-
 263pacts on contemporary telegraph systems in the United States and in the Iberian Penin-
 264sula (Ribeiro et al., 2016; Hayakawa et al., 2020).

265 On the other hand, the superstorms of March 1989 and September 1909 (panels
 266c and d) had very similar minimum values for Dst and Dst[†], around -590 nT. However,

Figure 4. JB2008 satellite orbital drag estimating results for the selected events for CHAMP’s orbit (a1-d1) and GRACE’s orbit (a2-d2) during the November 2003 event, but with hypothetical Dst/Dst^\dagger values. Panels a1/b1 and a2/b2: da/dt and d for the events in October/November 1903 (yellow lines) and May 1921 (green lines). Panels c1/d1 and c2/d2 indicate the same, but for the events in March 1989 (red lines) and September 1909 (blue lines). The highlighted areas correspond to the storm development duration, or the time interval between SSC occurrence and the end of the storm main phase (minimum Dst or Dst^\dagger occurrence).

267 the storm development duration of the 1989 event (24 hr) was 3 times longer than that
 268 of the 1909 event (8 hr). Consequently, the development amplitude rates of both super-
 269 storms were -23.8 nT/hr and -75.0 nT/hr, respectively. With respect to the aurorae of
 270 these events, Hayakawa, Ebihara, Cliver, et al. (2019) estimated, based on contempo-
 271 rary observations, that their equatorward extent reached $\sim 32^\circ$ MLAT during the 1909
 272 superstorm, as opposed to 40° MLAT estimated from particle precipitation measurements
 273 by satellites during the 1989 superstorm (Rich & Denig, 1992; Pulkkinen et al., 2012).
 274 Intense GICs occurred during both events, with several reports of geophysical disturbances
 275 on telegraph systems in 1909 (Silverman, 1995; Hayakawa, Ebihara, Cliver, et al., 2019;
 276 Hapgood, 2019; Love et al., 2019b), and on power transmission lines in 1989, particu-
 277 larly the power blackout in Québec, Canada (Allen et al., 1989; Kappenman, 2006; Oliveira
 278 & Ngwira, 2017; Boteler, 2019). During the 1989 event, the only event with satellite-based
 279 data amongst the four superstorms, the number of space objects “lost” in LEO increased
 280 dramatically around periods of maximum intensity due to errors introduced by storm
 281 heating effects into tracking systems (Allen et al., 1989; Joselyn, 1990; Burke, 2018). The
 282 left part of Table 1 summarizes these storm properties.

283 A comprehensive comparison of GIC effects caused by the superstorms on the con-
 284 temporary ground infrastructure, i.e., telegraph systems and power grids, is a difficult
 285 task to be accomplished. However, the comparisons above show that the latitudinal ex-
 286 tent of the auroral oval was more equatorward for the events with lower amplitude rates
 287 (May 1921 and September 1909 events). Next, the effects of these amplitude rates on
 288 storm-time orbital drag will be evaluated and compared for the 4 historical magnetic su-
 289 perstorms studied in this paper.

290 3.2 Storm-time orbital drag effects

291 Figure 4 shows results of storm-time satellite orbital drag effects estimated accord-
 292 ing to the framework presented in section 2.4. The computations are performed for the
 293 orbits of CHAMP and GRACE (Figure 2), with the orbital parameters the satellites had
 294 during the November 2003 storm. The sample CHAMP- and GRACE-like satellites are
 295 flown through an upper atmosphere produced by the JB2008 model for Dst/Dst^\dagger of the
 296 superstorms of Figure 3. All solar indices are kept the same, as those of the benchmark
 297 storm. For the sake of comparisons, results are plotted as a function of arbitrary times
 298 (GMT/UT) 12 hr before and 72 hr after the SSC onset as seen in Figure 3. The dashed
 299 vertical black lines ($t = 0$) indicate the times of SSC occurrence, while the highlighted
 300 areas correspond to the storm development duration as shown in Figure 3 for each cor-
 301 responding storm.

302 The top 4 panels of Figure 4 (a1-d1) show results for CHAMP’s orbit, while the
 303 bottom 4 panels (a2-d2) show results for GRACE’s orbit. Panels a1 and a2 show storm-
 304 time orbital decay rates (equation 4) computed for the October/November 1903 super-
 305 storm (yellow line) and May 1921 superstorm (green line) for CHAMP and GRACE, re-
 306 spectively. Both events had approximately the same development times and very differ-

Table 1. Summary of the properties of the magnetic superstorms and subsequent orbital drag results shown in Figures 3 and 4, respectively.

Storm Month and year	Magnetic superstorm properties				Orbital drag effects		
	SSC	Min	Development	Amplitude	Satellite	Min	Min
	GMT/UT (Day) ^a	Dst/Dst [†] [nT]	duration ^b [hr]	Rate ^c [nT/hr]	Name	da/dt [m/day]	d [m]
Oct/Nov 1903	0100(31)	-513	14	-44.8	CHAMP	-252.67	-102.65
					GRACE	-171.33	-65.39
May 1921	2300(14)	-907	12	-80.0	CHAMP	-432.98	-196.24
					GRACE	-319.43	-142.09
Mar 1989	0200(13)	-589	24	-23.8	CHAMP	-621.29	-388.59
					GRACE	-469.95	-305.58
Sep 1909	1200(25)	-595	8	-75.0	CHAMP	-285.14	-96.61
					GRACE	-191.25	-62.14

^a Greenwich Mean Time or Universal Time and Day of Storm Sudden Commencement (SSC).

^b Time between SSC and minimum Dst/Dst[†] occurrence.

^c $d(\text{Dst}/\text{Dst}^\dagger)/dt$

ent intensities (Table 1). The same is shown in panels c1 (CHAMP) and c2 (GRACE) for the superstorms of March 1989 (red line) and September 1909 (blue line). In this case, the storms had very similar intensities, but different development durations (Table 1). The storm-time orbital degradation (equation 5), is shown for CHAMP (panels b1 and d1) and GRACE (panels b2 and d2). The same colors used to represent da/dt results in panels a1/c1 and a1/c2 above are used to represent d results in panels b1/d1 and b2/d2.

Figure 4a1 shows that da/dt values during October/November 1903 for CHAMP were very close to zero before CME impact. On the other hand, da/dt values preceding the stormy period of May 1921 shows some oscillatory behavior, presumably linked to a similar behavior shown by ground magnetometer data during the same pre-storm period (Love et al., 2019b; Hapgood, 2019). CHAMP da/dt values for the 1921 event decreased faster in comparison to minimum da/dt values for the 1903 event. Similar orbital drag dynamics is observed for GRACE (a2), but the absolute values of the drag response are smaller (Table 1) because GRACE operated at higher altitudes in comparison to CHAMP (Krauss et al., 2018; Oliveira & Zesta, 2019). The da/dt results for CHAMP and GRACE are summarized in Table 1.

For the same pair of storms, the storm-time orbital degradations of CHAMP (panel b1) at the end of 72 hr after CME impact were -102.65 m and -196.24 m for both events, respectively. The same estimated results for GRACE (b2) are -65.39 m (1903) and -142.09 m (1921). Comparatively, the percentual difference between drag effects during both superstorms for CHAMP (91.17%) are higher than the percentual difference of the superstorm intensities (76.80%) most likely because the magnetosphere was hit by another CME on 16 May 1921 (Figure 3; Love et al., 2019a), leading to an additional magnetosphere energization during its recovery, which in turn impacted drag effects. Similarly, the orbital drag relative difference is higher in the case of GRACE (117.30%), when compared with the case of CHAMP. As suggested by Oliveira and Zesta (2019, Figure 10), this is presumably due to the interplay between heating propagation from auroral-to-equatorial latitudes and (possibly) the direct uplift of neutrals at low and equatorial latitudes more evident at altitudes higher than 400 km (Tsurutani et al., 2007).

In summary, the main features that arise from the comparison between these events are: (i) CHAMP and GRACE decayed faster during the most intense event (1921) due to its sharper negative excursion of the Dst[†] index and lower amplitude rate (Figure 3a

Table 2. Comparisons between magnetic superstorm intensity and satellite orbital drag severity for the magnetic superstorms in this study.

Magnetic Superstorm Month/Year	Comparisons between Superstorm intensities and durations	Relative differences of drag effects [%]			
		CHAMP		GRACE	
		da/dt	d	da/dt	d
Oct/Nov 1903	May 1921 is 76.80% stronger	71.36	91.17	86.44	117.30
May 1921	Nearly the same durations				
Sep 1909	March 1989 is 3 times longer	117.30	302.22	145.73	391.76
Mar 1989	Nearly the same intensities				
May 1921	March 1989 is 2 times longer	43.49 ^a	98.02	47.12	115.06
Mar 1989	May 1921 is 53.98% stronger				

^a Percentual differences between more severe (March 1989) with respect to less severe (May 1921) drag effects

339 and b; Table 1); and (ii) the relative differences between d for both events do not closely
 340 follow the relative differences between minimum Dst^\dagger values. This is likely the case be-
 341 cause the magnetosphere was struck by another CME during its recovery, increasing the
 342 magnetospheric activity which in turn affected the subsequent orbital drag effects. Ta-
 343 bles 1 and 2 summarize these results.

344 The comparisons between estimated drag effects for the March 1989 and Septem-
 345 ber 1909 superstorms are remarkably different. These events had very similar strengths
 346 (similar minimum Dst and Dst^\dagger values), but their development times were quite distinct.
 347 Figure 4c1 shows that 1909 CHAMP da/dt values had a very sharp negative excursion
 348 after CME impact, which follows very closely the same feature in the Dst^\dagger index (Fig-
 349 ure 3d). The minimum da/dt value (-285.14 m/day) for the September 1909 superstorm
 350 was reached shortly before minimum Dst^\dagger . On the other hand, the March 1989 drag ef-
 351 fects are quite different, since da/dt decreased more slowly in comparison to the former
 352 case due to the differences in storm development amplitude rates. This is explained by
 353 the fact that the magnetosphere was most likely struck by multiple CMEs while the storm
 354 main phase was developing (Fujii et al., 1992; Lakhina & Tsurutani, 2016; Boteler, 2019).
 355 Similarly to the 1909 case, the minimum da/dt value (-621.29 m/day) occurred shortly
 356 before minimum Dst occurrence. The thermosphere recovery of the 1989 superstorm took
 357 longer than the thermosphere recovery of the 1909 superstorm, most likely because the
 358 magnetosphere was yet hit by more CMEs shortly after the beginning of the magneto-
 359 sphere recovery (Figure 3c). A similar behavior is shown by the GRACE results, panel
 360 c2, but with smaller absolute values due to higher GRACE altitudes. The relative dif-
 361 ferences between da/dt peak values of CHAMP and GRACE for both superstorms are
 362 117.30% and 145.73%, even though both events had approximately the same minimum
 363 Dst and Dst^\dagger values and very different storm development durations and amplitude rates.

364 Now the storm-time orbital degradations in both cases are evaluated. Figure 4d1
 365 shows that CHAMP d decreased faster during the main phase of the 1909 event, reach-
 366 ing values near its minimum value around the beginning of storm recovery. This is a typ-
 367 ical feature of drag effects triggered by a storm caused by an isolated CME (Krauss et
 368 al., 2015, 2018; Oliveira & Zesta, 2019). Conversely, CHAMP's orbital degradation de-
 369 creased more dramatically during the recovery of the 1989 superstorm. These drag ef-
 370 fects correlate well with a very sharp negative excursion presented by the Dst index, which
 371 is also directly related with the occurrence of low-latitude aurorae and very intense GICs
 372 around the world (Allen et al., 1989; Kappenman, 2006; Hayakawa, Ebihara, Cliver, et
 373 al., 2019). This time also coincides with the loss of orbital control of several objects in
 374 LEO as shown by satellite-based data (Allen et al., 1989; Joselyn, 1990; Burke, 2018).
 375 The storm-time orbital decays for the 1909 and 1989 events are -96.61 m and -388.59

376 m for CHAMP and -62.14 m and -305.58 m for GRACE. These orbital degradation lev-
 377 els were never observed by LEO spacecraft with high-precision accelerometers during mag-
 378 netic superstorms, and therefore set a new base for drag effects under such circumstances.
 379 Their relative difference is 302.22% and 391.76%, closely following the proportion of storm
 380 time developments in the case of CHAMP. Taking into consideration that both super-
 381 storms were almost equally intense, these results show that the storm time duration can
 382 play a major role in driving orbital drag effects. Note also that relative differences are
 383 higher in the case of GRACE, most likely explained by the reasons suggested by Oliveira
 384 and Zesta (2019) as mentioned before.

385 Another striking difference concerning minimum Dst and Dst[†] values, storm de-
 386 velopment duration and subsequent amplitude rate impacts arises from the comparison
 387 between the May 1921 and March 1989 superstorms. The 1921 event was more than 50%
 388 stronger than the 1989 event, but active times during the latter lasted twice longer. The
 389 storm-time orbital decay for the March 1989 event was nearly twice more severe than
 390 the May 1921 event in both CHAMP’s and GRACE’s cases (Figure 4 and Tables 1 and
 391 2). These results clearly reveal that a long-lasting magnetic superstorm can drive much
 392 more severe drag effects in comparison to a short-lasting, even stronger, superstorm. Ta-
 393 bles 1 and 2 summarize the main results discussed in sections 3.1 and 3.2.

394 There are no solar wind nor interplanetary magnetic field data available for the mag-
 395 netic superstorms discussed in this paper. Furthermore, it is important to emphasize that
 396 our statements concerning CME impacts are supported by our current knowledge of the
 397 underlying science: intense magnetic storms, particularly extreme events, are usually caused
 398 by CMEs (Gonzalez et al., 1994; Daglis et al., 1999; Balan et al., 2014; Tsurutani & Lakhina,
 399 2014; Lakhina & Tsurutani, 2016; Kilpua et al., 2019).

400 4 Discussion and conclusion

401 Extreme magnetic storms (minimum Dst ≤ -250 nT) are very rare. Only 39 ex-
 402 treme events have taken place since the beginning of the space era (Meng et al., 2019),
 403 while only 7 extreme events were observed by CHAMP and GRACE (Oliveira & Zesta,
 404 2019; Zesta & Oliveira, 2019). Additionally, only one magnetic superstorm (minimum
 405 Dst ≤ -500 nT) occurred since 1957, while none were ever observed by either CHAMP
 406 or GRACE. Therefore, current knowledge of thermospheric mass density response to mag-
 407 netic superstorms and the subsequent storm-time drag effects are very limited. Then,
 408 in order to estimate these effects, 4 historical magnetic superstorms were selected: one
 409 with standard Dst data (March 1989), and 3 with Dst[†] (Dst-like) data occurring on Oc-
 410 tober/November 1903 (Hayakawa et al., 2020), September 1909 (Love et al., 2019b), and
 411 May 1921 (Love et al., 2019a). These Dst and Dst[†] data were used as input data for the
 412 JB2008 thermospheric empirical model for density computations. The extreme magnetic
 413 storm of November 2003 (minimum Dst = -422 nT), the most extreme event during CHAMP’s
 414 and GRACE’s commission times, was used as the benchmark event. The orbital drag
 415 framework provided by Oliveira and Zesta (2019) was used for drag estimations.

416 First, two events with different intensities but with approximately the same storm
 417 development times were compared (October/November 1903 and May 1921). Although
 418 the 1921 superstorm was $\sim 80\%$ stronger than the 1903 superstorm, the drag effects in
 419 the former were up to 120% more severe than the effects in the latter (GRACE’s case).
 420 This is attributed to the likely impact of another CME during the recovery phase of the
 421 1921 superstorm. Second, the other pair of superstorms, with very similar strengths, but
 422 with the September 1909 storm development being 3 times shorter than the March 1989
 423 storm development, were compared. Results show that the relative difference of the storm-
 424 time orbital degradation for the 1989 event was about 400% higher than the 1909 event
 425 (GRACE’s case). This is explained by the likely impacts of several CMEs on the mag-
 426 netosphere during the main and recovery phases of the March 1989 superstorm (Fujii

et al., 1992; Lakhina & Tsurutani, 2016; Boteler, 2019). Therefore, as opposed to latitudinal extent of aurorae, a superstorm with a smaller amplitude rate (absolute value) can cause more detrimental effects on orbital drag in comparison to an even stronger superstorms that develops faster (larger absolute value of amplitude rate). All orbital degradations shown in Table 1 are much more severe than the orbital degradation due to the background density estimated by Oliveira and Zesta (2019) during November 2003 for CHAMP (-34.45 m) and GRACE (-6.86 m); therefore, these results set a new basis for these effects. Despite the fact that these effects can have significant error levels particularly during the storm recovery phases due to the lack of nitric oxide cooling effects in the model (Mlynczak et al., 2003; Bowman et al., 2008; Knipp et al., 2017; Oliveira & Zesta, 2019; Zesta & Oliveira, 2019), these results reveal the comparative roles of time durations and strengths of magnetic superstorms in controlling drag effects.

The results of this work clearly show that multiple CME impacts on the Earth’s magnetosphere (as in the March 1989 superstorm), particularly occurring during active times, can largely enhance satellite orbital drag due to long and sustained storm times. These drag effects can be more severe when compared to drag effects during storms caused by a single CME leading to even more intense storms, but lasting shorter. Therefore, orbital drag forecasters should be aware of potential impacts of several CMEs on the terrestrial magnetosphere during ongoing magnetic storms (e.g., Zhao & Dryer, 2014, and many references therein). In a future work, simulation results of tens of historical severe and extreme magnetic storms, with minimum Dst ≤ -250 nT excluding superstorms (Meng et al., 2019; Oliveira & Zesta, 2019; Zesta & Oliveira, 2019; Chapman et al., 2020), will be statistically studied.

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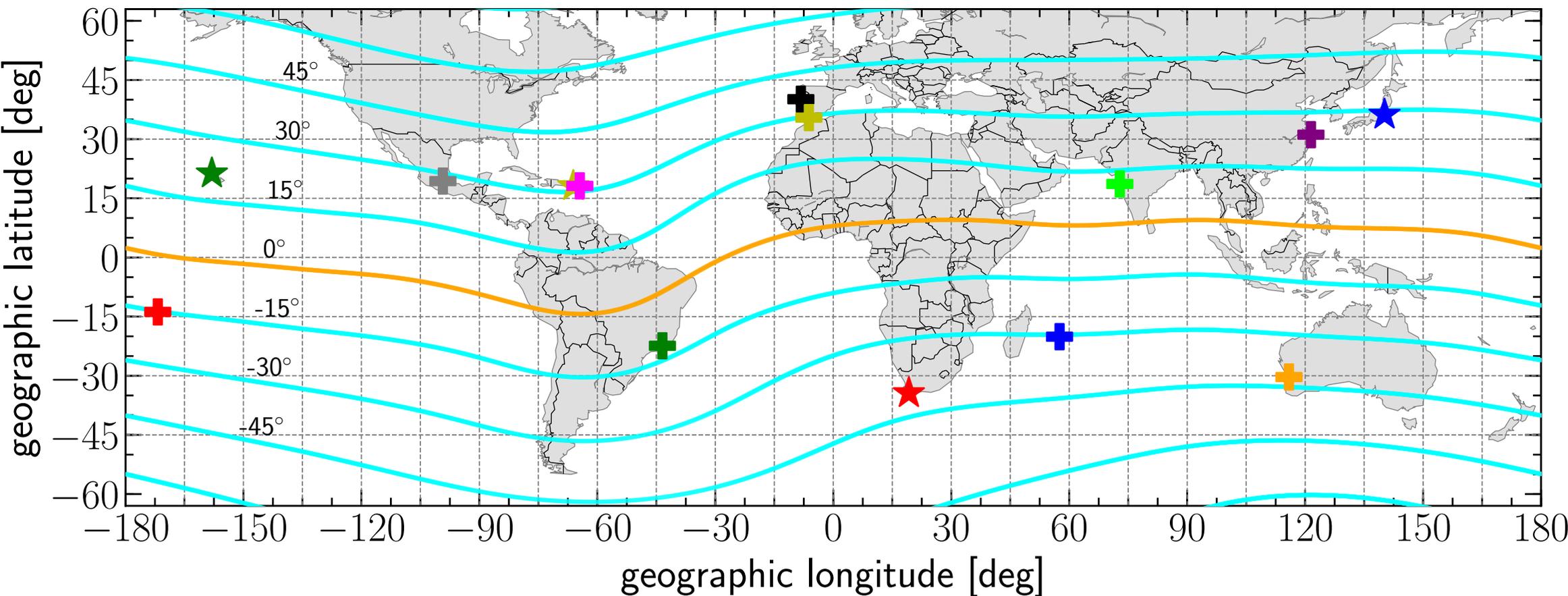
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Figure 1.

Standard (Dst) and alternative (Dst[†]) disturbance storm time stations



Standard Dst stations

- ★ Honolulu [HON], United States
- ★ San Juan [SJG], Puerto Rico
- ★ Hermanus [HER], South Africa
- ★ Kakioka, [KAK], Japan

Magnetic coordinates for 1957

- Magnetic latitudes
- Magnetic equator

Alternative Dst[†] stations

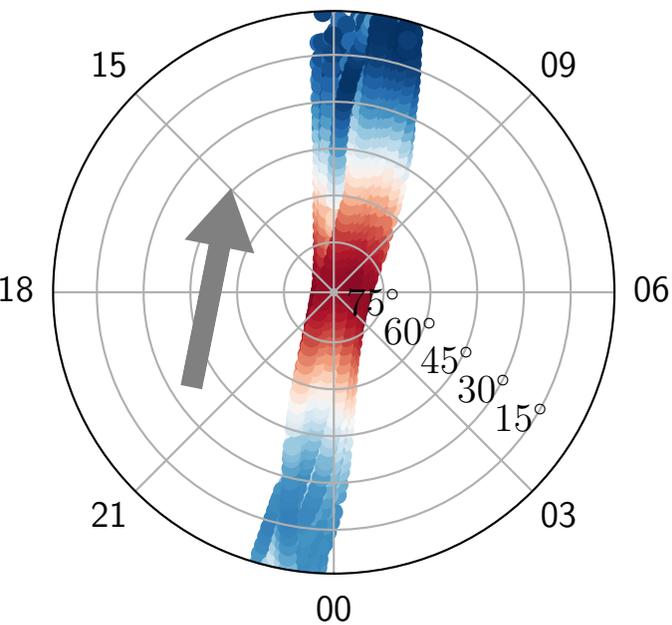
- ✚ Apia [API], Western Samoa (1909, 1921)
- ✚ Coimbra [COI], Portugal (1903)
- ✚ Colaba [CLA], India (1903)
- ✚ Cuajimalpa [CUA], Mexico (1903)
- ✚ Mauritius [MRI], Mauritius (1909)
- ✚ San Fernando [SFS], Spain (1909, 1921)
- ✚ Vassouras [VSS], Brazil (1921)
- ✚ Vieques [VQS], Puerto Rico (1909)
- ✚ Watheroo [WAT], Australia (1921)
- ✚ Zi-Ka-Wei [ZKW], China (1903)

Figure 2.

CHAMP

Northern Hemisphere

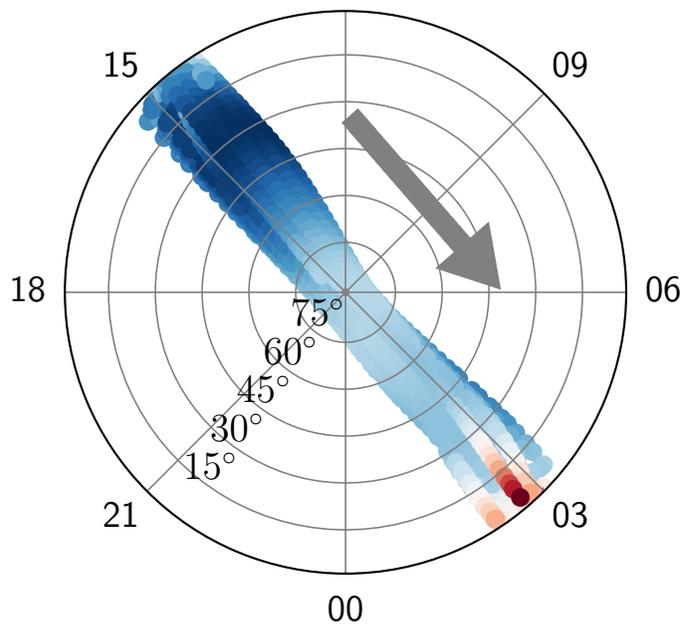
12 MLT



GRACE

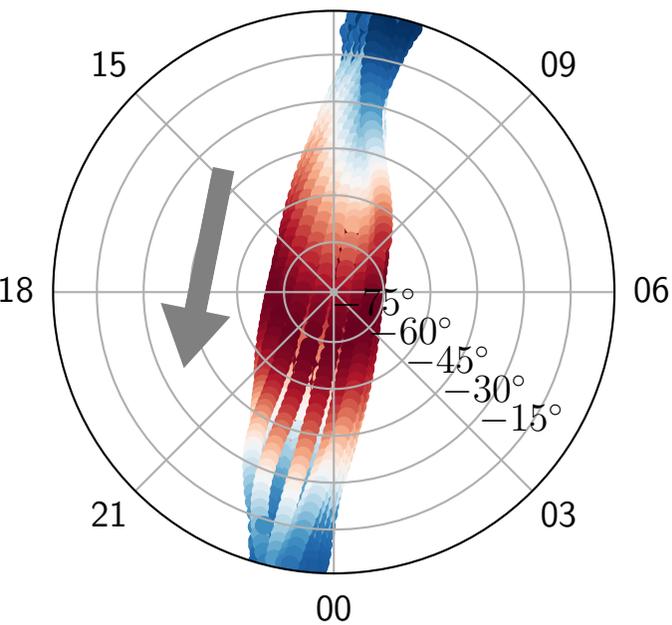
Northern Hemisphere

12 MLT



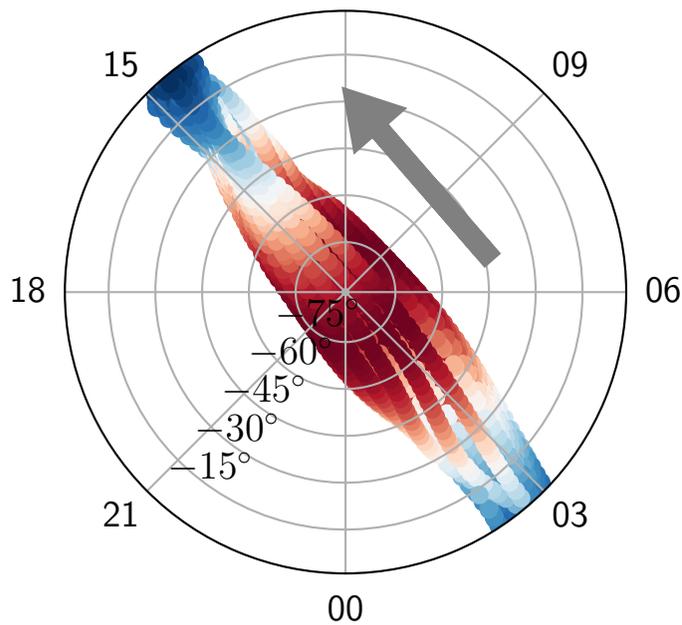
Southern Hemisphere

12 MLT



Southern Hemisphere

12 MLT



390 395 400 405 410 415

CHAMP altitude [km]

480 490 500 510 520

GRACE altitude [km]

Figure 3.

Dst (standard) and Dst[†] (alternative) indices for the selected magnetic superstorms

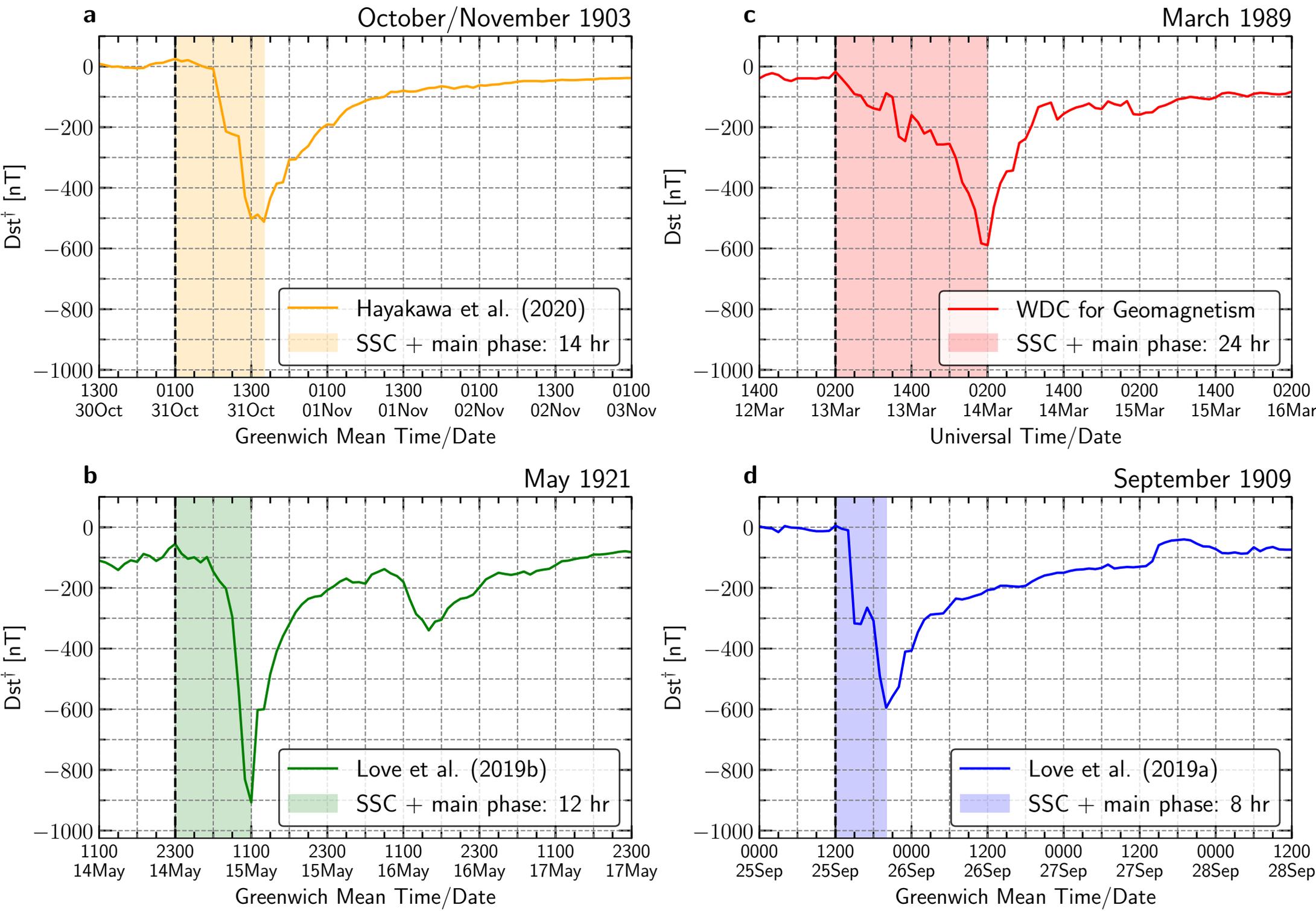
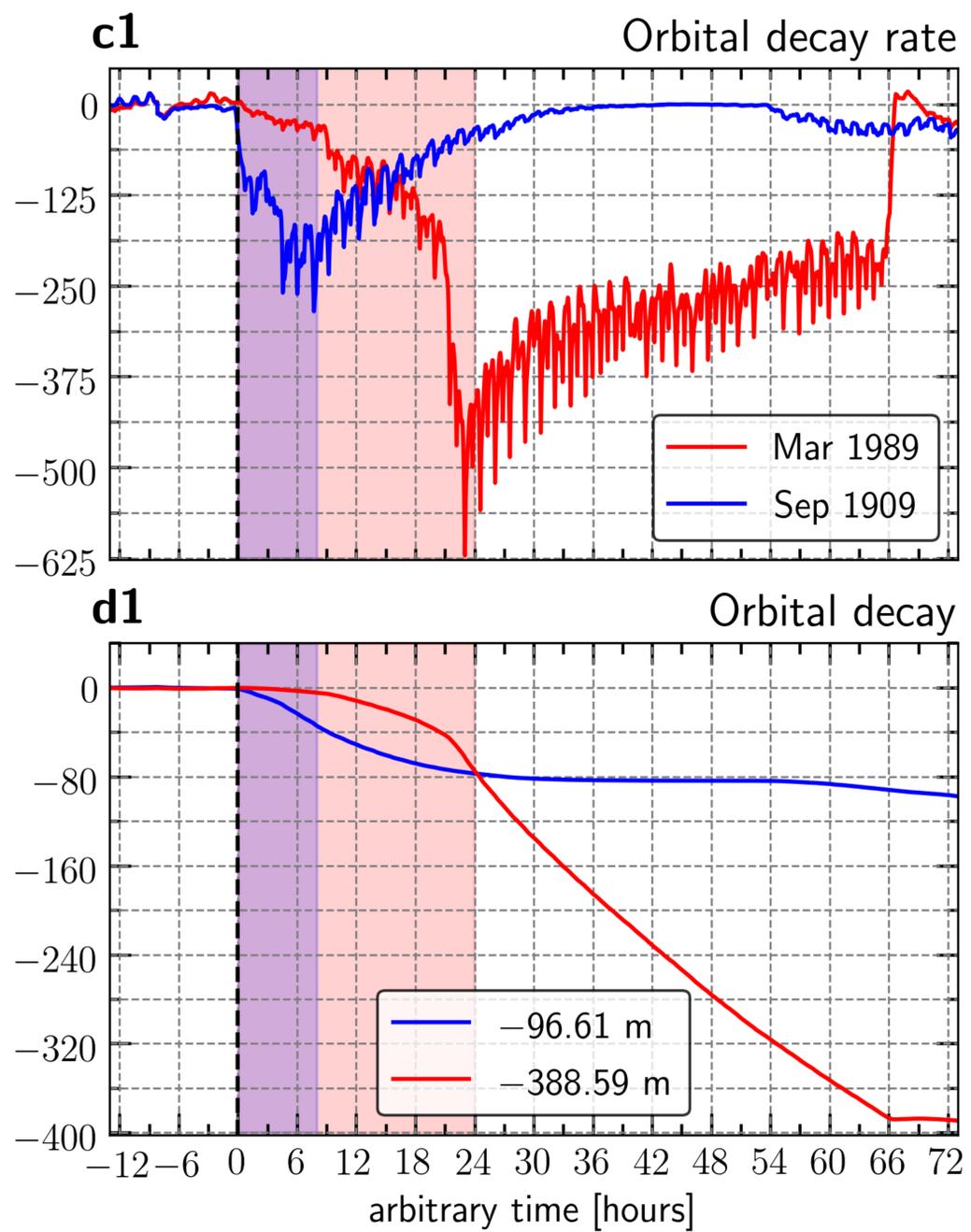
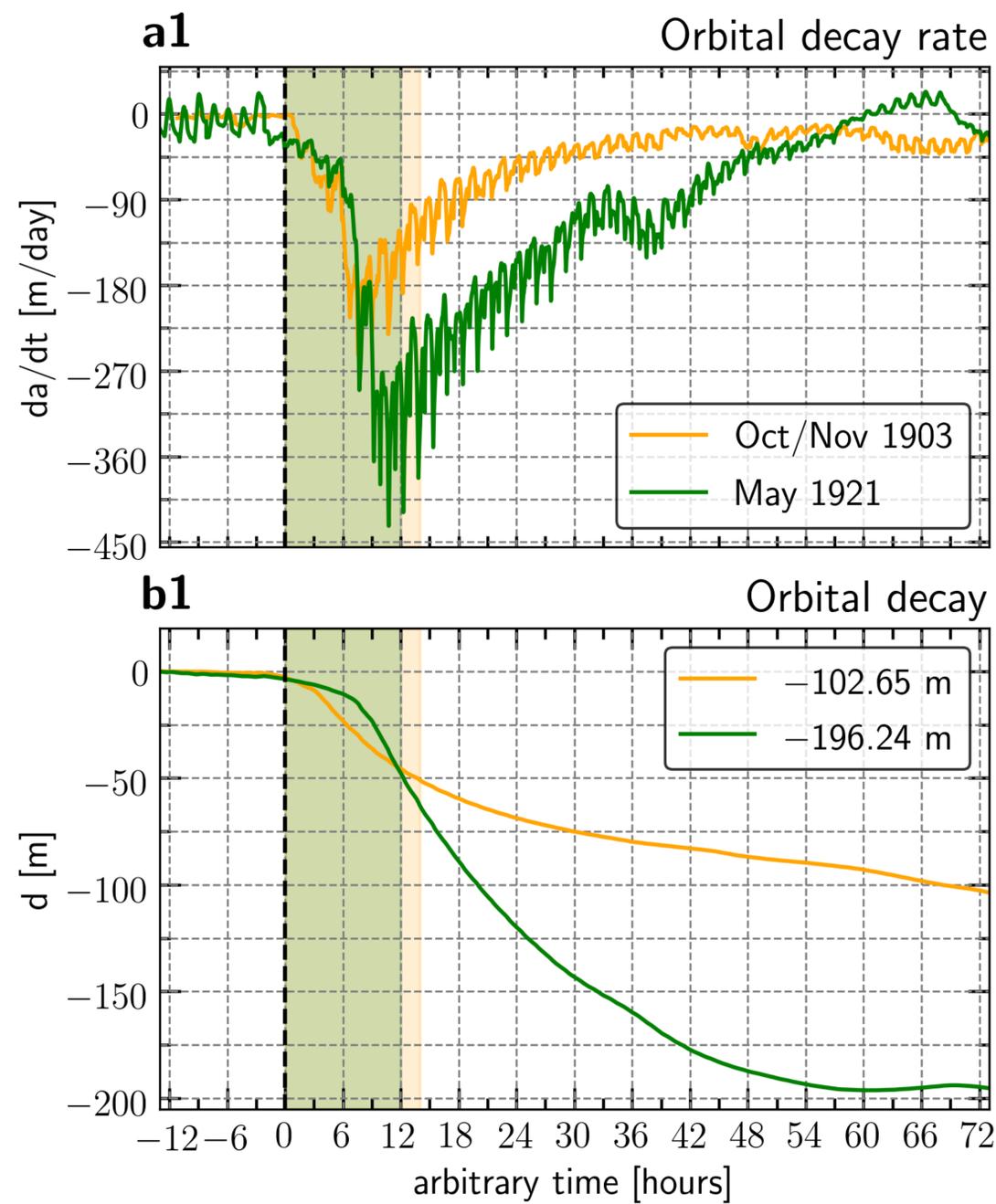


Figure 4.

Storm-time orbital drag effects for CHAMP's orbit



Storm-time orbital drag effects for GRACE's orbit

