

1 **Fusion of MISR Stereo Cloud Heights and Terra-MODIS Thermal Infrared Radiances to**
2 **Estimate Multi-layered Cloud Properties**

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8
9 **Key Points:**

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11 1. Accurate, high-precision MISR low cloud heights are employed in a physics-based
12 correction to MODIS CO₂-slicing in multi-layered scenes.
13 2. Cloud-top pressure bias drops from 65 hPa to 5 hPa, resulting in a quartering of cloud-
14 height and emissivity bias for cirrus over low cloud.
15 3. Up to 88% of cloud-top pressure retrieval errors are bound by theoretical estimates,
16 resulting in near-closure of CO₂-slicing error budget.
17

18 Abstract

19 Our longest, stable record of cloud-top pressure (CTP) and cloud-top height (CTH) are derived
20 from the Moderate Resolution Imaging Spectroradiometer (MODIS) and Multi-Angle Imaging
21 Spectroradiometer (MISR) on Terra. Because of single cloud-layer assumptions in their standard
22 algorithms, they provide only single CTP/CTH retrievals in multi-layered situations. In the
23 predominant multi-layered regime of thin cirrus over low clouds, MODIS significantly
24 overestimates cirrus CTP and emissivity, while MISR accurately retrieves low-cloud CTH.
25 Utilizing these complementary capabilities, we develop a retrieval algorithm for accurately
26 determining both-layer CTP and cirrus emissivity for such 2-layered clouds, by applying the MISR
27 low-cloud CTH as a boundary condition to a modified MODIS CO₂-slicing retrieval.

28 We evaluate our 2-layered retrievals against collocated Cloud-Aerosol Transport System (CATS)
29 lidar observations. Relative to CATS, the mean bias of the upper cloud CTP and emissivity are
30 reduced by ~90% and ~75% respectively in the new technique, compared to standard MODIS
31 products. We develop an error model for the 2-layered retrieval accounting for systematic and
32 random errors. We find up to 88% of all residuals lie within modeled 95% confidence intervals,
33 indicating a near-closure of error budget. This reduction in error leads to a reduction in modeled
34 atmospheric longwave radiative flux biases ranging between 5-40 Wm⁻², depending on the position
35 and optical properties of the layers. Given this large radiative impact, we recommend that the
36 pixel-level 2-layered MODIS+MISR fusion algorithm be applied over the entire MISR swath for
37 the 22-year Terra record, leading to a first-of-its-kind 2-layered cloud climatology from Terra's
38 morning orbit.

39 Plain Language Abstract

40 Our longest climate-quality record of global cloud-top heights (CTH) comes from the Moderate
41 Resolution Imaging Spectroradiometer (MODIS) and Multi-Angle Imaging Spectroradiometer
42 (MISR) on the Terra satellite. These sensors assume a single cloud-layer in retrieving CTH, even
43 though ~30% of global cloud cover is multi-layered. Multi-layered clouds predominantly consist
44 of thin ice clouds over low clouds. Under such conditions, MISR accurately retrieves low-cloud
45 CTH, while MODIS systematically underestimates upper-cloud-layer CTH. Here, we have
46 developed a 2-layered MODIS+MISR fusion CTH retrieval by using MISR's accurate low-cloud
47 CTH as an input to a modified MODIS algorithm. This algorithm combines the complementary
48 capabilities of MISR and MODIS in distinguishing higher and lower clouds and estimates both-
49 layer cloud heights and high-cloud emissivity.

50 Through comparisons against coincident Cloud-Aerosol Transport System (CATS) lidar
51 observations, we find that the new algorithm improves the accuracies in retrieved CTH and cloud
52 emissivities by ~75% over standard MODIS products. We further demonstrate significant
53 improvements in estimates of simulated atmospheric longwave radiation from our implementation.
54 Owing to its large radiative impact, we suggest that the pixel-level fusion algorithm be applied to
55 all 22 years of Terra record to facilitate public dissemination of the first 2-layered cloud record
56 from its morning orbit.

57 1. Introduction

58 The vertical and horizontal distribution of clouds induces gradients in 3D radiative and latent
59 heating rates (McFarlane et al., 2008; Cesana et al., 2019; Athreyas et al., 2020), affecting
60 atmospheric circulation and precipitation patterns (Li et al., 2015; Voigt et al., 2021). As such,
61 clouds play an important role in the Earth's climate – yet, even after decades of research, they
62 remain the key source of uncertainty in predicting future climate change under any given climate
63 change scenario (Boucher et al., 2013). The cloud component of the uncertainty in climate model
64 predictions arises, in part, from approximate sub-grid parametrization of cloud processes in those
65 models (McFarlane, 2011). The sub-grid scale parameterizations are applied to microphysical
66 (hydrometeor size and content) and macrophysical cloud properties (amount-by-altitude and cloud
67 overlap), which together govern the radiative and hydrological properties of clouds. Accurate
68 satellite records of these micro- and macro-physical properties, and their diurnal to long-term
69 variability, are essential to provide empirical constraints on sub-grid cloud parameterizations and
70 climate predictions (e.g., Zhou et al., 2013; Terai et al., 2016; Mace & Berry, 2017).

71 Our longest record of cloud properties that are stable over multiple decades (features of a desirable
72 climate record) and from a single satellite platform comes from NASA's flagship Earth Observing
73 System (EOS) mission, Terra. It maintained a stable equator-crossing time (ECT; 10:30 am \pm 15
74 minutes) for >20 years (2000-2022), with remarkable radiometric stability in its instruments. This
75 long-term stability in Terra's ECT makes it a unique climate record, since diurnal variability has
76 not been aliased into the patterns of long-term variability.

77 Two of the instruments on Terra – the Multiangle Imaging Spectroradiometer (MISR) and the
78 Moderate Resolution Imaging Spectroradiometer (MODIS) – employ independent cloud-top
79 height (CTH) retrieval algorithms. MISR retrieves CTHs through visible-channel stereoscopy
80 (Moroney et al., 2002; Muller et al., 2002; Mueller et al. 2013), whereas MODIS employs infrared
81 (IR) techniques, namely the CO₂-slicing and 11 μ m brightness temperature techniques (Menzel et
82 al., 2008; Baum et al., 2012). Both MODIS and MISR CTH retrieval algorithms assume a single
83 cloud layer in the scene. This assumption is often not met in nature as multi-layered clouds occur
84 frequently, with CALIPSO/CloudSat showing that >30% of all clouds occur under various degrees
85 of overlap (Sassen et al., 2008; Joiner et al., 2010; Yuan & Oreopoulos, 2013; Li et al., 2015;
86 Oreopoulos et al., 2017; Hong & Di Girolamo, 2020). By far the most dominant multi-layered
87 cloud regime is a 2-layered system with thin cirrus overlying water clouds, followed by thin cirrus
88 overlying mixed-phase clouds (Wang & Dessler, 2006; Oreopoulos et al., 2017; Hong and Di
89 Girolamo 2020). Numerous validation studies against ground and space-based active sensors have
90 shown that the presence of optically thin cirrus overlying low clouds leads to the most significant
91 disagreements in retrieved CTH between MISR and MODIS (Naud et al., 2007; Marchand et al.,
92 2010; Mitra et al., 2021), suggestive of the presence of independent information of the upper and
93 lower cloud layers in the two datasets.

94 The path to improving the Terra record relies on exploiting the distinctiveness of the MODIS and
95 MISR CTH techniques to estimate the properties of multi-layered clouds more accurately, as
96 previously suggested (Naud et al., 2007; Mitra et al., 2021). CTH errors in multi-layered cloud
97 regimes have been comprehensively studied for the Terra MODIS and MISR records by Mitra et
98 al. (2021) using collocated Cloud-Aerosol Transport System (CATS) lidar observations (McGill
99 et al., 2015; Yorks et al., 2016) that operated aboard the International Space Station (ISS) from
100 2015-2017. Comparison of MODIS Collection 6.1 CTH with CATS showed that the CTHs of thin
101 cirrus in these multi-layered regimes were underestimated by more than 1 km on average. 42% of
102 the MODIS CTH retrievals occurred below the cloud base detected by the lidar in these conditions.

103 Such biases are common in thermal CTH retrievals and are due to the radiative influence of the
 104 lower cloud layer reaching the sensor through the optically thin cirrus at infrared wavelengths. On
 105 the other hand, the stereoscopic technique employed by MISR tended to retrieve the height of the
 106 lower layer when cirrus visible optical depths were less than ~ 0.4 , and with a high degree of
 107 precision and accuracy (-280 ± 300 m). However, MISR failed to detect the higher layer in favor of
 108 the lower layer $>80\%$ of the time in these multi-layered conditions. This is due to the greater
 109 contribution of the optically thicker, more textured low clouds to the overall image texture that is
 110 used in stereoscopic retrieval. The distinct error characteristics of MISR and MODIS CTH
 111 retrievals indicate that there is information about multi-layering of clouds that can be extracted
 112 through fusion of the two retrieval methodologies. Here, we present a retrieval algorithm that
 113 makes use of the strengths of MISR's sensitivity to low clouds and MODIS CO₂-slicing
 114 technique's sensitivity to high clouds to retrieve the coincident heights of up to two cloud layers,
 115 which also improves the CO₂-slicing technique's estimate of the cirrus emissivity. We carry out a
 116 detailed error budget analysis and validate the retrievals using CATS.

117 The remainder of the paper is organized as follows. Section 2 describes the theoretical
 118 underpinnings of the CO₂-slicing algorithm for retrieving CTH and emissivity of thin ice clouds,
 119 and how it has been updated here to account for the presence of an optically thick low cloud
 120 measured by MISR. Section 3 describes the datasets used and the method of implementation of a
 121 variant of the MODIS single-layered CO₂-slicing, along with the implementation of our 2-layer
 122 CO₂-slicing technique. Section 4 documents the validation of the 2-layer CO₂-slicing against
 123 coincident CATS lidar observations, along with an error budget analysis for the same. Since cloud
 124 radiative effect depends strongly on cloud overlap (e.g., Li et al., 2011; L'Ecuyer et al., 2019, Kang
 125 et al. 2020), Section 5 demonstrates significant improvements in modeled cloud radiative effects
 126 when using inputs from the 2-layer algorithm compared to the 1-layer algorithm. Concluding
 127 remarks follow in Section 6.

128 **2. Theoretical Foundation**

129 CO₂-slicing (Smith & Platt, 1978; Wielicki & Coakley, 1981), as used in MODIS (Menzel et al.,
 130 2008), makes use of the difference of clear- and cloudy-sky radiances from closely separated
 131 channels in the 13-15 μm CO₂ absorption band, where the emissivity for ice clouds (such as cirrus)
 132 remain invariant across wavelengths within the band. Clear-sky radiance are estimated through
 133 infrared radiative transfer to account for the radiance reaching MODIS that originates from below
 134 thin ice clouds. The spectral clear-sky IR radiance, I_{cs} (neglecting scattering) at wavelength λ ,
 135 reaching a satellite sensor viewing at nadir over a black surface (for simplicity here) is given by:

$$136 \quad I_{cs}(\lambda) = \mathcal{T}(\lambda, P_s)B(\lambda, T(P_s)) - \int_0^{P_s} B(\lambda, P) \frac{d\mathcal{T}(\lambda, T(P))}{dP} dP \quad \dots (1)$$

137 where, P_s denotes the surface pressure, $B(\lambda, T)$ denotes the Planck radiance at temperature T and
 138 wavelength λ , with temperature defined as a function of pressure, P . $\mathcal{T}(\lambda, P)$ denotes the
 139 atmospheric transmittance between P and the satellite. For a completely opaque cloud covering
 140 the instantaneous field of view (IFOV) of the sensor, the effective emissivity, which is the product
 141 of cloud fraction (A_c) within the IFOV and the cloud layer emissivity (ϵ_c), is unity. In this case,
 142 provided the opaque cloud is geometrically infinitesimally thin, the nadir radiance observed by the
 143 satellite, I_c , is devoid of all emissions from below the cloud-top pressure (P_c), and is given by:

$$144 \quad I_c(\lambda, P_c) = \mathcal{T}(\lambda, P_c) B(\lambda, T(P_c)) - \int_0^{P_c} B(\lambda, P) \frac{dT(\lambda, T(P))}{dP} dP \quad \dots (2)$$

145 In reality, cirrus are often transmissive ($\epsilon_c A_c < 1$). Then, the observed nadir top-of-atmosphere
146 (TOA) radiance is:

$$147 \quad I(\lambda) = I_{cs}(\lambda) + \epsilon_c(\lambda) A_c [I_c(\lambda, P_c) - I_{cs}(\lambda)] \quad \dots (3)$$

148 where, A_c is the cloud fraction, and $\epsilon_c A_c$ is often interchangeably referred to as the effective cloud
149 amount or effective emissivity. As effective emissivity for ice clouds is nearly equal for any two
150 wavelengths (say λ_1 and λ_2) in the 15 μ m CO₂-absorption band, we set them equal to each other,
151 which, from Eq. 3, leads to

$$152 \quad \frac{I(\lambda_1) - I_{cs}(\lambda_1)}{I(\lambda_2) - I_{cs}(\lambda_2)} = \frac{I_c(\lambda_1, P_c) - I_{cs}(\lambda_1)}{I_c(\lambda_2, P_c) - I_{cs}(\lambda_2)} \quad \dots (4)$$

153 Cloudy-sky radiances are calculated for a number of discrete P_c values, and the value of P_c for
154 which the right-hand side (RHS) and the left-hand side (LHS) have the least absolute difference is
155 taken as the retrieved P_c . Using this value of P_c , we can solve for the cloud effective emissivity
156 from Eq. 3, for either band, by:

$$157 \quad \epsilon_c(\lambda) A_c = \frac{I(\lambda) - I_{cs}(\lambda)}{I_c(\lambda, P_c) - I_{cs}(\lambda)} \quad \dots (5)$$

158 For a 2-layer cloud system, with lower altitude cloud at P_l of effective amount $\epsilon_l(\lambda) A_l$ and an
159 upper altitude cloud at P_u of effective amount $\epsilon_u(\lambda) A_u$, Eq. 3 misrepresents the observed TOA IR
160 radiation at the satellite sensor as it does not consider the emission from the lower cloud layer
161 when the upper-layer is thin (i.e., $\epsilon_u(\lambda) A_u < 1$). In reality, for such a 2-layered system, the
162 background emission (equivalent to the clear-cloudy sky radiance difference in a single-layered
163 case) comes not only from the surface but also from the lower-layer, and hence, $I_{cs}(\lambda)$ in Eq. 3 is
164 modified to be $I'_{cs}(\lambda) = \epsilon_l(\lambda) A_l I_c(P_l) + (1 - \epsilon_l(\lambda) A_l) I_{cs}(\lambda)$, and the TOA IR radiance is:

$$165 \quad I'(\lambda) = I_{cs}(\lambda) + \epsilon_l(\lambda) A_l [1 - \epsilon_u(\lambda) A_u] \int_{P_l}^{P_s} B(\lambda, P) \frac{dT(\lambda, T(P))}{dP} dP \\ 166 \quad + \epsilon_u(\lambda) A_u [I_c(\lambda, P_u) - I_{cs}(\lambda)] \quad \dots (6)$$

167 Since $I'(\lambda)$ is usually less than $I(\lambda)$, the cloudy-clear radiance differences on the LHS of Eq. 4 are
168 typically reduced when a second layer is present. Hence, simply using the single-layer strategy of
169 Eq. 4 results in a CTP solution that is numerically greater than the true P_u . Comparing Eq. 3 and
170 Eq. 6, we note that the second term of Eq. 6 must be accounted for in the CO₂-slicing of 2-layered
171 clouds, and hence, Equations 4 and 5 must be updated accordingly. Since the number of unknown
172 variables in Eq. 6 would make solving the equation intractable, we make the simplifying
173 assumption that the lower cloud is black [i.e., $\epsilon_l(\lambda) A_l = 1$], and define the following term:

$$174 \quad \Delta I(\lambda) = \int_{P_l}^{P_s} B(\lambda, P) \frac{dT(\lambda, T(P))}{dP} dP \quad \dots (7)$$

175 As in a 1-layered CO₂-slicing, we assume $\epsilon_u^1 A_u^1 = \epsilon_u^1 A_u^1$ (but now strictly for the upper cloud
176 marked by 'u'). With all these modifications, Eq. 4 for multi-layered cases is recast as:

$$177 \quad \frac{I(\lambda_1) - I_{cs}(\lambda_1) - \Delta I(\lambda_1)}{I(\lambda_2) - I_{cs}(\lambda_2) - \Delta I(\lambda_2)} = \frac{I(\lambda_1, P_u) - I_{cs}(\lambda_1) - \Delta I(\lambda_1)}{I(\lambda_2, P_u) - I_{cs}(\lambda_2) - \Delta I(\lambda_2)} \quad \dots (8)$$

178 Similarly, Eq. 5 is adjusted to account for $\Delta I(\lambda)$, and is recast from Eq. 6, as:

$$179 \quad \epsilon_u(\lambda)A_u = \frac{I(\lambda) - I_{cs}(\lambda) - \Delta I(\lambda)}{I_c(\lambda, P_u) - I_{cs}(\lambda) - \Delta I(\lambda)} \quad \dots (9)$$

180 **3. Methodology**

181
182 Section 3.1 briefly describes the datasets used in this study to both implement and validate our
183 CO₂-slicing algorithm. Section 3.2 describes the method of implementation of this algorithm.

184 **3.1. Data**

186 The operational MODIS Cloud Top Property algorithm [detailed in the MODIS Algorithm
187 Theoretical Basis Document or ATBD (Menzel et al. 2015)], which produces the 1 km-resolution
188 Collection 6.1 MOD06 product, uses gridded model output from the National Center of
189 Environmental Prediction Global Data Assimilation System (GDAS) (Derber et al., 1991) for
190 temperature and moisture fields and Reynolds Sea Surface Temperatures (Reynolds et al., 2007)
191 to set up the forward model atmosphere. In our implementation, we have instead used gridded
192 ERA5 Reanalysis products (Hersbach et al., 2020) at 0.25°-resolution, at 4 times a day (i.e., 0, 6,
193 12 and 18 UTC), to do the same. ERA5 is chosen over other reanalyses because it has been
194 demonstrated to compare better against observations than older reanalyses (Tetzner et al., 2019;
195 Tegtmeier et al., 2020), as well as to use its publicly available modeling error estimates for error
196 budget analysis (see Section 4.2). ERA5 temperatures, specific humidity, and geopotential heights
197 from all 37 ERA5 pressure levels are linearly interpolated as a function of the logarithm of pressure
198 to arrive at the atmospheric state for the 101 pressure levels employed by the MOD06 algorithm.
199 Surface pressures, temperatures (2m temperature over land and sea-surface temperature over
200 oceans) and 2m dewpoint temperatures (to calculate surface moisture) are also used from ERA5
201 reanalysis, 4 times daily, to define surface temperature and near-surface humidity.

202 Well-mixed and trace gases (except ozone) are taken from standard atmospheric profiles
203 (Northern/Southern Midlatitude Summer/Winter, Tropical) (Anderson et al., 1986); as are
204 temperatures, specific humidity, and geopotential heights in the uppermost reaches of the
205 atmosphere (i.e., pressures < 1 hPa; ERA5 reanalyses are not available at these altitudes). Between
206 April-September, we assume a Northern Midlatitude Summer; while, between October-March, we
207 assume a Northern Midlatitude Winter. The opposite is true for the Southern Hemisphere. The
208 tropical profile remains invariant for all times of the year and is applied between 30°N-30°S,
209 whereas the midlatitude profiles are chosen for latitudes poleward of ±30°. From Collection 6
210 MOD06, ozone profiles are taken from gridded GDAS output; however, for simplicity, we
211 obtained ozone profiles similar to legacy MOD06 products – climatological ozone mixing-ratio
212 profiles were estimated by linear interpolation in latitude and month among model atmospheres
213 (Tropical, Midlatitude Summer/Winter). Surface emissivity is taken from the same global surface
214 emissivity database used in MOD06 (Seemann et al., 2008).

215 The observed infrared radiances used in Equations 4/5 and 8/9 are taken from the Collection 6.1
 216 MODIS Level 2 geocalibrated radiance product (MOD021KM). Terra MODIS uses Bands 33, 35
 217 and 36 (13.3, 13.9 and 14.2 μm , respectively) for CO₂-slicing CTP estimation [Band 34 (13.6 μm),
 218 also a CO₂ absorption channel, is unused due to high noise]. Hence, the band-pairs 36/35 and 35/33
 219 are used for estimating CTP (Equations 4 and 8). Band 31 (11.2 μm) radiances are used to calculate
 220 effective cloud amounts (Equations 5 and 9).

221 The low-cloud pressure, P_l , is taken from MISR Level 2 CTH (in pressure coordinates). We use
 222 the 1.1 km-resolution MISR “wind-corrected” cloud height, from the TC_CLOUD Version
 223 F01_0001 product. The low cloud CTH is transformed to pressure coordinates through a linear
 224 interpolation between multi-level ERA5 geopotential height and the logarithm of pressure. MISR
 225 CTH is reported on the 1984 World Geodetic System (WGS84) ellipsoid, and hence, 0.25°-
 226 resolution nearest neighbor geoid heights were added to MISR CTH to obtain low cloud heights
 227 above mean sea level, before calculating CTP from it.

228 We validate our CO₂-slicing technique by comparing against coincident observations from the
 229 CATS lidar. Thus, our validation is restricted to latitudes traversed by the ISS orbit ($\pm 52^\circ$ in either
 230 hemisphere). The CATS data is taken from the CATS Version 2.01 Level 2 Product, that reports
 231 lidar observations such as 1064 nm cloud-masked lidar backscatter at an along-track resolution of
 232 5 km and a vertical resolution of 60 m. We use the same dataset of CATS CTH, layer depth and
 233 layer-integrated backscatter used in Mitra et al. (2021) for this study. As in Mitra et al. (2021),
 234 CATS, MISR and MODIS samples were selected only if they are collocated (< 1 km) and
 235 coincident (< 5 minutes), for robust statistical analysis. Note that the filtering of multi-layered
 236 scenes in our study must be based solely on MISR and MODIS retrievals. Based on the discussion
 237 in Section 2, our algorithm is best suited for scenes with a thin ice-phase cloud overlying a
 238 vertically well-separated low cloud layer (further discussed in Section 3.2). To ensure application
 239 only on ice-phase clouds, we apply our algorithm only on scenes where the MOD06 product had
 240 used CO₂-slicing for cloud-top detection (since CO₂-slicing is only applied on ice-phase clouds).
 241 To ensure that our algorithm is applied on scenes with well-separated cloud layers, we choose only
 242 those scenes where MODIS-MISR CTH difference > 1 km [suggestive of well-separated cloud
 243 layers, based on Mitra et., al (2021)]. Upon imposing these conditions, it is found that all scenes
 244 in the remaining dataset are indeed multi-layered according to CATS. 95% are likely 2-layered
 245 (for 92% of such cases, the CATS signal completely attenuates in the second layer). The remaining
 246 5% pixels show attenuation in a third cloud layer. The final dataset constitutes 2790 pixels from
 247 501 independent scenes (i.e., unique MISR and MODIS granules and CATS orbits), hence ~ 6
 248 samples per scene (Figure S1). Out of these, 305 ($\sim 11\%$) pixels are no-retrievals. This is largely
 249 due to the presence of radiance artifacts, such as striping within the MODIS data. In the current
 250 study, such bad pixels are discarded from the analyses, but can be dealt with in future
 251 implementations by established procedures of MODIS radiance de-striping (Weinreb et al., 1989;
 252 Bouali & Ladjal, 2011).

253 IR emissivity (ϵ_{IR}) of a cloud layer is related to visible optical depth (τ_{VIS}) over the layer, as

$$254 \quad \tau_{VIS} = -\zeta \ln(1 - \epsilon_{IR}) \quad \dots (10)$$

255 where, $-\ln(1 - \epsilon_{IR})$ equals the thermal IR optical depth (τ_{IR}). The constant ζ is taken to be 2.13
 256 for ice clouds (Minnis et al., 1993; Rossow & Schiffer 1999). Estimates of visible optical depth

257 (τ_{VIS}) of the topmost cloud layer from CATS comes from a linear regression between layer-
 258 averaged integrated backscatter and layer-integrated optical depth for high clouds (CTH > 7 km)
 259 [detailed in (Mitra et al., 2021)]. These estimates of high cloud τ_{VIS} are converted to infrared
 260 effective emissivity ($A_c \epsilon_c$, assuming $A_c = 1$) using Eq. 10 for validation. MODIS 1 km-resolution
 261 CTP, CTH, effective emissivity ($A_c \epsilon_c$) and visible optical depth (τ_{VIS}) from the MOD06 product
 262 are also used in comparison to CATS and our 2-layered solution.

263 **3.2. Implementation of the CO₂-slicing Algorithm**

264 For our implementation of the CO₂-slicing algorithm, we have modified the original MOD06
 265 Fortran Cloud-Top Property code obtained from the MODIS Adaptive Processing System (see
 266 Section 7) and wrapped it in Python. Salient features of the operational code and the modifications
 267 for our implementation are hereby discussed.

268 The MOD06 algorithm simulates clear- and cloudy-sky radiances using Equations 1 and 2, on 101
 269 vertical pressure levels between 0.05 to 1100 hPa, taking gaseous absorption, surface emissivity
 270 and satellite zenith angle into account. These radiances are calculated for the channels centered on
 271 11.2, 13.3, 13.6, 13.9 and 14.2 μm , using a transmittance model named Pressure layer Fast
 272 Algorithm for Atmospheric Transmissions (PFAAST) (Hannon et al., 1996), and further corrected
 273 for increased path-length along off-nadir viewing zenith angles. The usage of these modeled
 274 radiances along with the observed radiances from MODIS, in Eq. 4, requires that the cloud
 275 emissivity for pairs in the CO₂-slicing spectral bands be nearly equal, which is more satisfied by
 276 ice clouds than water or mixed-phase (Zhang & Menzel, 2002). Hence in generating the standard
 277 MOD06 product, the MODIS cloud phase detection algorithm is run ahead of the cloud-top
 278 algorithm. The CO₂-slicing technique is applied only on such scenes with ice phase detection (11.2
 279 μm brightness temperature technique is applied elsewhere).

280 In our implementation we use the same PFAAST model and we account for cloud phase by
 281 selectively working only on those pixels where the Collection 6.1 MODIS CO₂-slicing had been
 282 previously used. Global comparison of Aqua-MODIS cloud phase with CLOUDSAT-CALIPSO
 283 data had shown that the MODIS cloud phase algorithm mischaracterizes multi-layered clouds with
 284 an upper ice layer as liquid or mixed in <1% of all cases (Marchant et al., 2016). This ensures
 285 confidence that pixels flagged as confidently ice by the Terra MODIS cloud phase algorithm is
 286 nearly always ice topped and hence, suitable for the implementation of our algorithm.

287 **3.2.1. Implementation of a Single-layered CO₂-slicing and its Bias**

288 To obtain solutions for CTP and emissivity, Eq. 4 is solved iteratively between the surface and the
 289 tropopause, to obtain the value of P_c that best reduces the difference between LHS and RHS of
 290 Eq.4. The tropopause is chosen as the upper limit of CTP solution, because the temperature profile
 291 is nearly flat across the tropopause, leading to non-unique solutions. The tropopause is taken to be
 292 the level of the highest altitude inflection point in the reanalysis temperature profile for pressures
 293 > 100 hPa. If many points satisfy such a condition, the lowest altitude point is chosen to be the
 294 tropopause. The solution of P_c from Eq. 4 is then used in Eq. 5 using 11.2 μm radiances to estimate
 295 effective cloud amounts ($A_c \epsilon_c$).

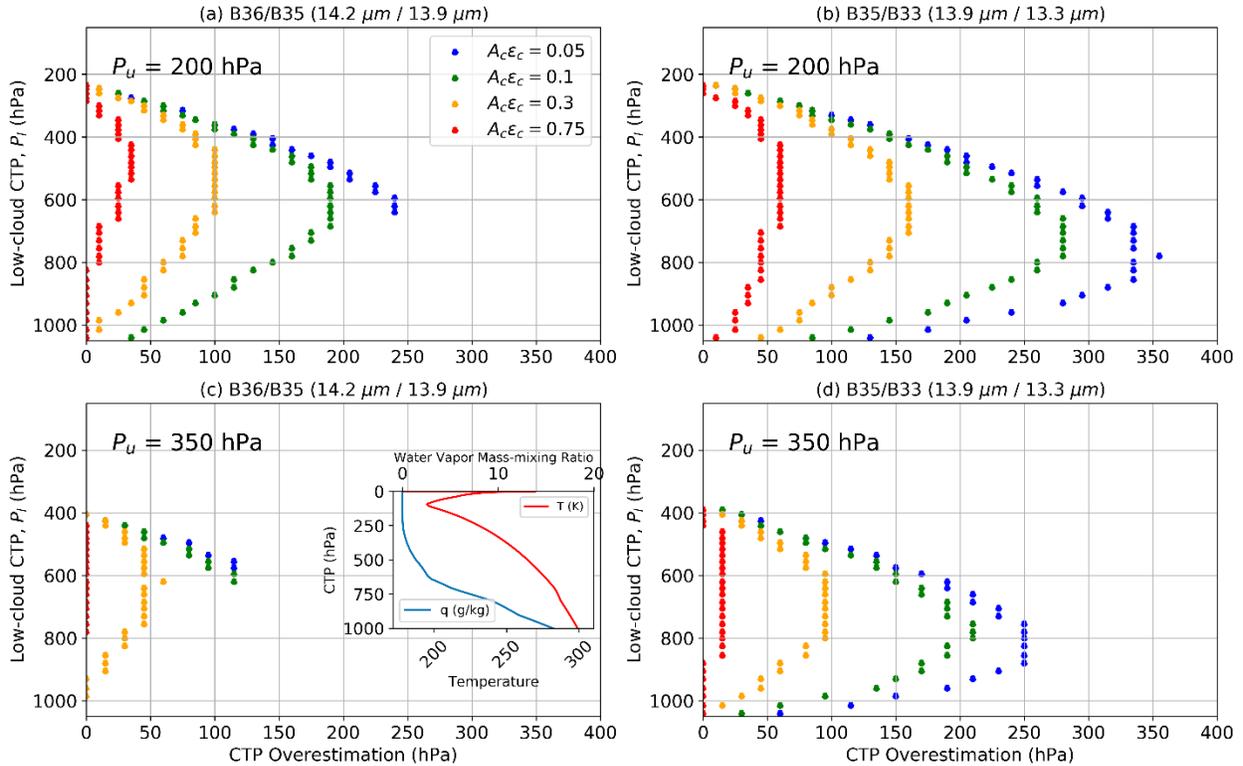
296 The standard MOD06 algorithm calculates all possible CTP solutions, before only reporting a
 297 “best” solution through a “top-down” method that checks for the possibility of a higher wavelength
 298 solution before a lower wavelength or brightness temperature solution (i.e., 36/35 solution over

299 35/33 solution, over an IR BT solution) (Menzel et al., 2008). For a solution to be viable, the clear-
300 cloudy radiance difference must exceed noise levels for each particular channel in that spectral
301 band pair (designated to be 1.25, 1.0, 1.0 and 0.75 W m⁻² sr⁻¹ for Bands 36-33, respectively), and
302 the solution from that channel must lie within a specific portion of the troposphere where the
303 atmosphere is emissive for that spectral channel (i.e., for 36/35 pair, CTP solutions must be < 450
304 hPa; for the 35/33 pair, CTP solutions must be < 650 hPa) (Baum et al., 2012).

305 To verify the implementation of our algorithm, we compared our 1-layer CTP solutions against
306 MOD06 CTP for 500 CATS single-layer high cloud (CTH > 7 km) pixels from 42 independent
307 scenes in January-February 2016. We find a mean (\pm standard deviation) difference in CTP
308 between our implementation and MOD06 to be -5 ± 30 hPa. For these scenes, the mean CTP bias
309 (relative to CATS) for MOD06 is 20 ± 30 hPa, whereas it is 15 ± 35 hPa for our implementation.
310 This provides confidence in our implementation, while also underscoring the fact that moving from
311 GDAS to ERA5 reanalysis had only a minor impact on the single-layer CO₂-slicing retrieval.

312 To estimate the systematic errors accrued from cloud overlap in CO₂-sliced CTP, we conduct an
313 experiment where we apply the 1-layered CO₂-slicing on 2-layered cloud systems. For these
314 experiments, we employ the forward model described in Section 3.2 to calculate synthetic
315 radiances for the 2-layered system, except we include a lower, black cloud layer as in Eq. 6. We
316 then use Equations 4 and 5 to retrieve the CTP under the assumption of a single layer and examine
317 the resulting errors. This experiment is idealized in that it neglects any errors in the forward model.
318 We perform retrievals on the synthetic two-layered systems for a climatological tropical
319 atmosphere for different values of P_u and P_l . We calculate the overestimations of CTP above P_u for
320 four effective cloud amounts between 0.05-0.75 and for each of the spectral band pairs that are
321 used by Terra MODIS, with results shown in Fig. 1. Here we see that the highest overestimation
322 of high-cloud CTP (i.e., an underestimation of high-cloud CTH) occurs in the 35/33 band pair for
323 a combination of very thin high cirrus over a low cloud (provided the low cloud is sufficiently
324 decoupled from the surface). It is unsurprising that the 35/33 band pair is more susceptible to the
325 presence of low clouds, because there is a large reduction in the amount of near-surface radiation
326 that reaches the satellite sensor in going from 13.3 to 14.2 μm due to increasing absorption by
327 CO₂. For the same high-low cloud combination and same spectral band pair, it is also unsurprising
328 that the thinnest of clouds ($\epsilon_c A_c = 0.05$) has the highest errors in CTP determination. As the lower
329 cloud approaches either the high cloud or the surface, the 2-layered system essentially becomes
330 indistinguishable from a single-layered high cloud; hence, in both these extreme conditions, the
331 bias is reduced. These results are similar to the estimates of CTP bias arising from the application
332 of a 1-layered CO₂-slicing for 2-layered cloud systems by the HIRS/2 sounder (Figures 3, 5 and 6
333 in Baum & Wielicki, 1994) and MODIS (Figure 10a of Menzel et al, 2015).

334 Based on these findings, our bias-correction approach (Equations 8 and 9) for two-layered cloud
335 systems will have the largest correction for well-separated cloud layers, particularly when the
336 lower cloud-top is both sufficiently colder than the surface and warmer than the upper-layer cloud.



337
 338 *Figure 1. Bias in CTP from MODIS CO₂-slicing (under single-layer assumption) for Bands 36/35 (left panels) and*
 339 *35/33 (right panels) for a high cloud at pressure = 200 hPa (upper panels) and 350 hPa (bottom panels), given a*
 340 *standard tropical atmosphere profile of water vapor (g/kg) and temperature (K; inset in c). Climatological profiles of*
 341 *ozone and trace gases are used. The lower cloud is assumed opaque, and the surface (1014 hPa) is a dark ocean. For*
 342 *each high-low combination, the experiment is repeated for cloud emissivities of 0.05 (blue), 0.1 (green), 0.3 (orange)*
 343 *and 0.75 (red).*

344 3.2.2. Implementation of the 2-layered CO₂-slicing

345 The modification to the CO₂-slicing solution for a 2-layered system involves replacing Equations
 346 4 and 5 with Equations 8 and 9 in the CO₂-slicing workflow, which, in turn, requires the
 347 computation of the term ΔI , given by Eq. 7. This step requires the value of MISR CTP (Section
 348 3.1). The closest of the 101 MODIS levels to MISR CTP is taken as P_l in Eq. 7. Solutions for P_u
 349 from band pairs 36/35 and 35/33 are recorded. A best solution is also chosen using the “top-down”
 350 method. If no legitimate solution is found (Section 3.2.1), it is a no-retrieval.

351 All 2485 valid CTP retrievals are converted to CTHs, using ERA5 geopotential heights. All such
 352 retrievals are also used to estimate effective cloud amounts (using Eq. 9). MOD06 effective cloud
 353 amounts are also used for comparison. Following Eq. 10, effective cloud amounts are converted
 354 to visible optical depths (τ_{VIS}), assuming $A_c = 1$. Note, the estimates for $A_c \epsilon_c$ and τ_{VIS} are
 355 estimates of the high cloud optical properties retrieved after the radiative contribution of the lower
 356 cloud has been removed. In contrast, the corresponding MOD06 retrievals are effective estimates
 357 of those quantities retrieved using the combined radiation from both upper and lower cloud layers.

358 This aforementioned modification to the CO₂-slicing is rooted in physical theory and makes use
 359 of Terra’s unique design for fusion between instruments, which allows us to improve the MODIS
 360 upper-layer CTP/CTH and emissivity, provided the layer is optically thin for MISR to retrieve

361 CTH of the lower cloud [this is also the regime where MODIS CO₂-slicing CTH errors are
362 maximum (Mitra et al., 2021)]. To distinguish the new high cloud properties from the operational
363 MODIS data variables, we shall refer to the new estimates of cirrus CTP/CTH, $A_c\epsilon_c$ and τ_{VIS} as
364 the *MISR-MODIS Fusion Product for Cloud-Top Height* (MM_CTH).

365 **4. Validation**

366
367 In Section 4.1, MM_CTH and MOD06 estimates of high cloud macrophysics and optical
368 properties will be validated against CATS estimates of those quantities. Section 4.2 provides a
369 detailed error budget analysis of our 2-layered CO₂-slicing CTH retrieval with the goal of closing
370 the total error budget had through a comparison with CATS CTH.

371 **4.1. Comparison with the CATS lidar**

372
373 To validate our new algorithm, we compare the results of high cloud CTP/CTH, high cloud
374 effective emissivity ($A_c\epsilon_c$) and visible optical depths (τ_{VIS}) from MM_CTH against concurrent
375 MOD06 and CATS observations. We divide the validation of MM_CTH along two lines –
376 validation of high cloud macrophysics (CTP, CTH) and high cloud optical properties ($A_c\epsilon_c$, τ_{vis}).

377 **4.1.1. Validation of High-Cloud Macrophysical Properties**

378 As in Mitra et al. (2021), we take CATS CTH/CTP to be an unbiased truth in our analysis. CATS
379 CTH is converted to CATS CTP, using ERA5 geopotential and standard geoid heights, in the same
380 manner as MISR CTH to CTP conversion. Figure 2 shows the distribution of CTP/CTH differences
381 between CO₂-slicing techniques (MOD06 and MM_CTH) and the lidar on the left panels, and the
382 distributions of high cloud CTP/CTH from the 3 techniques (MOD06, MM_CTH and CATS) on
383 the right panels. The mean bias (\pm standard deviation) in retrieved CTP and CTH improves from
384 65 ± 85 hPa and -1.6 ± 2.3 km, respectively, for MOD06 to 5 ± 80 hPa and -0.4 ± 2.4 km for MM_CTH.
385 This represents a $\sim 90\%$ reduction in CTP bias and a $\sim 75\%$ reduction in CTH bias.

386 The reduction in the CTP/CTH bias for high-cloud retrievals results in improved high cloud
387 macrophysical distributions (right panels of Figure 2), with the MM_CTH distributions of
388 CTP/CTH closely mirroring those from CATS. Mitra et al. (2021) showed that for 42% of all
389 scenes with a thin cirrus overlying a low cloud, MODIS CTH lies below the vertical extent of the
390 cirrus (i.e., lower than CATS cloud-layer base). The application of the 2-layered MM_CTH
391 reduces the instances of such below-cloud-base height retrievals to 12%.

392 For the distributions of MM_CTH minus CATS CTP and CTH (Figures 2a and 2c), we note the
 393 existence of a significant number of scenes (~4% of each distribution) where MM_CTH appears
 394 to overestimate the value of CATS CTH by > 4 km (i.e., underestimate CTP > 100 hPa). Previous
 395 studies (Rajapakshe et al., 2017; Mitra et al., 2021) had identified these as scenes where the 1 km-
 396 resolution infrared sensor detects physically tenuous (e.g., broken cirrus) clouds, but the lidar's 5
 397 km-resolution algorithm picks the height of a lower, possibly horizontally continuous, cloud field.
 398 Here, we show that this assertion is indeed true by calculating the mean MISR-CATS CTH for
 399 scenes with MM_CTH – CATS CTH difference > 4 km, and finding a mean difference of -0.5 ± 0.5
 400 km. This is close to MISR's CTH accuracy for low clouds (Mitra et al., 2021). Thus, in these
 401 scenes, MODIS retrieved cirrus heights and CATS retrieved low cloud heights. Such an effect is
 402 noticeably smaller in the corresponding MOD06 distributions because MOD06 estimates of CTH
 403 (CTP) are lower (higher), and hence, closer to the CATS low-cloud retrievals.

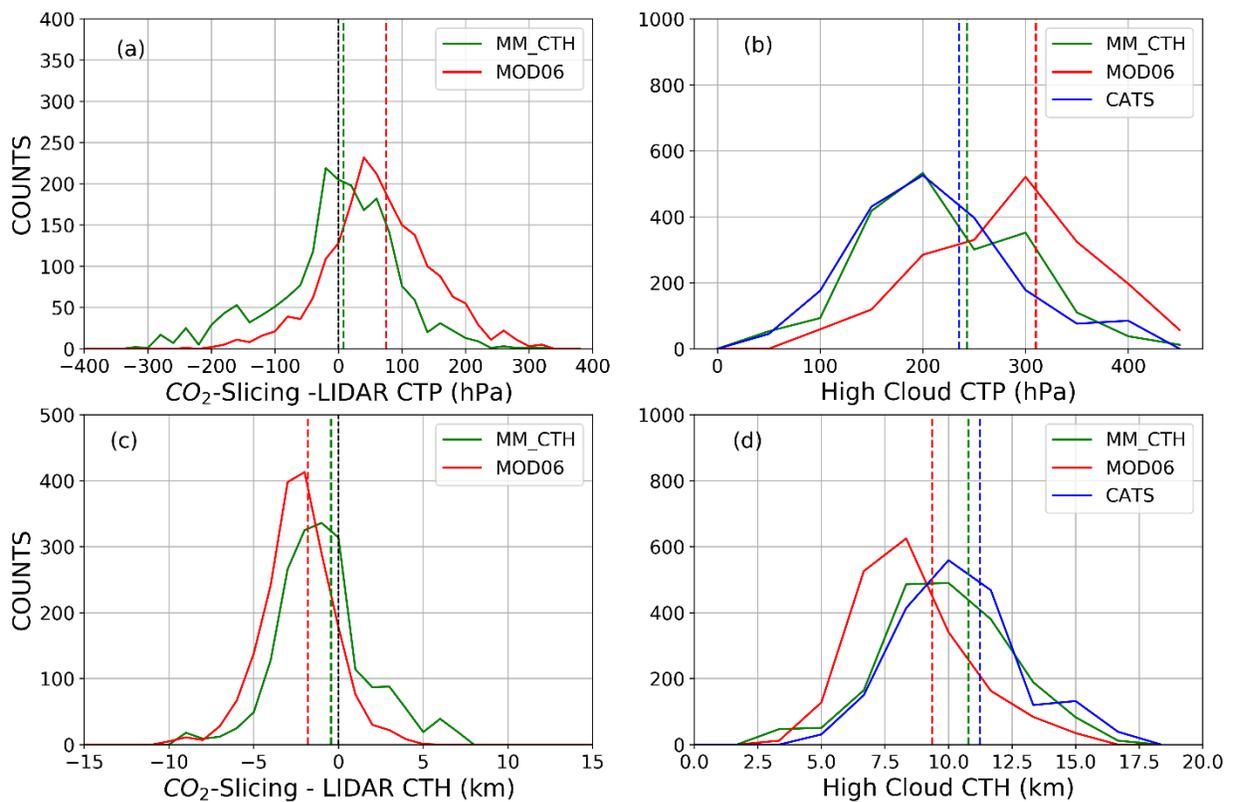


Figure 2. Distribution of errors (left) in CTP (top panels; hPa) and CTH (bottom panels; km) from MOD06 (red) and MM_CTH (green) and the distribution of high cloud macrophysics (right panels) for multi-layered scenes from MOD06 (red), MM_CTH (green) and CATS (blue). The vertical dashed lines in each color represents the mean value of the quantities whose distributions are in that same color.

404 4.1.2. Validation of High-Cloud Optical Properties

405 Unlike cloud-top properties (CTP and CTH), we do not have an unbiased estimate for cloud
 406 effective amount ($A_c \epsilon_c$). As a result, we have converted CATS τ_{VIS} to CATS $A_c \epsilon_c$ by inverting
 407 Eq. 10 (taking $\zeta = 2.13$ and assuming $A_c = 1$). Even though this is not an unbiased estimate of
 408 true $A_c \epsilon_c$, one can reasonably expect the CATS $A_c \epsilon_c$ to be a closer estimate of true cirrus

409 emissivity as compared to MOD06 $A_c\epsilon_c$, because the MOD06 $A_c\epsilon_c$ estimate is impacted by the
 410 lower cloud layer. As shown in Figure 3, we have compared MM_CTH estimates of $A_c\epsilon_c$ and τ_{VIS}
 411 against CATS and MOD06 estimates of those quantities. The improvements in cloud
 412 macrophysical retrievals shown in Section 4.1.1 have propagated to improvements in retrievals of
 413 high cloud optical properties. From Fig. 3, we notice a $\sim 75\%$ increase in accuracy in both $A_c\epsilon_c$
 414 and τ_{VIS} for MM_CTH over MOD06 (assuming that CATS emissivity and τ_{VIS} are unbiased).
 415 These improvements lead to MM_CTH distributions of high cloud emissivity and optical depths
 416 that are comparable to the corresponding distributions from the CATS lidar.

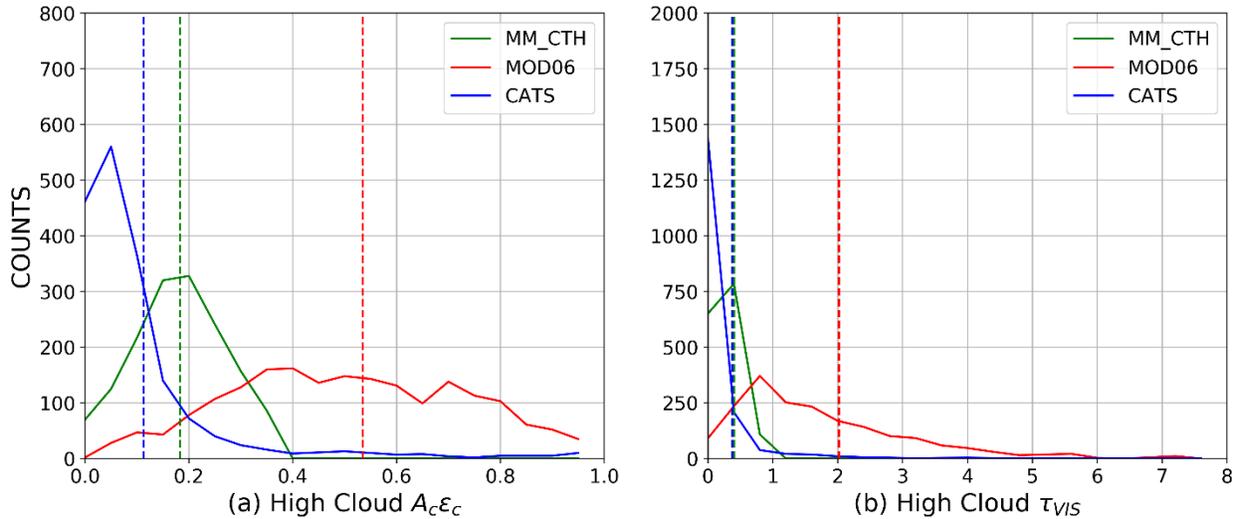


Figure 3. Distribution of effective emissivity (left) and visible optical depth (right) from MOD06 (red), MM_CT (green) and CATS (blue) for high clouds in multi-layered scenes. The dashed lines in each color represents the mean value of the quantities whose distributions are in that same color. On the right plot, the mean values of τ_{vis} from MM_CTH and CATS are visibly indistinguishable.

417 The MOD06 estimates of $A_c\epsilon_c$ and τ_{VIS} are both overestimations of true high-cloud optical
 418 properties because their individual retrieval methods do not remove the radiative contribution of
 419 the lower cloud. As a result, both are effective retrievals over all cloud layers. Improved estimates
 420 of upper-cloud optical properties (especially τ_{VIS}) are crucial in the accurate representation and
 421 tuning of cloud radiative effects in models (which we demonstrate in Section 5). In Fig 4b, the
 422 MOD06 estimates of τ_{VIS} are from the standard MODIS bispectral optical depth retrievals
 423 (Platnick et al., 2017) which use visible channel radiances and separate pre-computed look-up
 424 tables for ice and water clouds. As such, since we are working on scenes where MODIS cloud
 425 phase detected ice, the ice look-up tables had been used to retrieve τ_{VIS} for a 2-layered multi-phase
 426 system. Here, we have improved the retrieval by improving our estimates of only the cirrus τ_{VIS} .
 427 However, with the improvements achieved by MM_CTH in defining the CTP of the ice and water
 428 cloud layers, future work can design 2-layered ice + water/mixed phase cloud look-up tables to
 429 simultaneously retrieve the visible optical depths of both cloud layers present within the scene.

430 In the previous sections, we have presented the validation of MM_CTH retrievals against CATS
 431 lidar. A detailed discussion of the CTP error budget follows in Section 4.2.

432 4.2. The 2-layered CO_2 -slicing Error Budget Analysis

433 In this section, we shall investigate the effect of various sources of systematic and random errors
 434 on MM_CTH CTP with the goal of comparing the total computed error against those shown in
 435 Section 4.1 (note that we do not repeat this exercise for effective emissivity as we do not have a
 436 truth dataset for that quantity). We consider the following sources of errors:

- 437 i. the uncertainty in MISR low-cloud stereo heights,
- 438 ii. the covariance of modelling errors in ERA5 Reanalysis temperature and specific humidity,
- 439 iii. the inherent noise in detected radiances from the MODIS spectral bands,
- 440 iv. the effect of geometric depth of cirrus clouds,
- 441 v. the uncertainty in the geo-collocation of CATS, MISR and MODIS pixels,
- 442 vi. the uncertainty incurred from the application of spatial interpolation to obtain atmospheric
 443 parameters at the 101 MOD06 vertical pressure levels,
- 444 vii. the breakdown of the assumption that the low clouds are perfectly black, and
- 445 viii. the effect of uncertainty in surface emissivity.

446 Empirical error estimates are known (as explained below) for the first six items on the list above.
 447 However, we lack a ‘truth’ dataset for low cloud opacity and surface emissivity. Hence, error
 448 sources vii and viii will be dealt with in a different manner to the others.

449 We run radiative transfer simulations over a range of 2-layered cloud combinations and use the
 450 simulated radiances in MM_CTH retrievals to estimate CTP errors. We compile these errors, E
 451 (bias and standard deviation), in the prescribed functional form: $E =$
 452 $E(P_{high}, depth, \tau_{VIS}, P_{low}, \lambda_{pair}, climate\ zone)$. Here, P_{high} , $depth$, τ_{VIS} are CTP, geometric
 453 depth, and visible optical thickness of the high clouds in the simulations. P_{low} is the CTP of the
 454 low black cloud. λ_{pair} refers to the MODIS band-pair being employed (i.e., either 35/33 or 36/35),
 455 and $climate\ zone$ denotes the 5 climate zones introduced in Section 3.1. For each climate zone
 456 and λ_{pair} , we run the MM_CTH algorithm for every combination of the following:

- 457 (a) 10 values of P_{high} (50 hPa intervals between 150 and 550 hPa)
- 458 (b) 6 values of P_{low} (50 hPa intervals between 700 and 1000 hPa)
- 459 (c) 5 values of geometric depth (25 hPa intervals between 25 and 150 hPa)
- 460 (d) 8 values of τ_{VIS} (0.25 intervals between 0.25 and 2.5)

461 This leads to 2400 cases for each band-pair and climate zone (hence 24000 in total). We choose
 462 the ranges for high cloud properties and P_{low} from the distributions of high cloud properties and
 463 low cloud heights (in units of pressure) that we observed in the CATS and MISR data used in this
 464 study. Out of the variables on which the error function E depends, we expect there to be significant
 465 random variability in estimates of P_{low} , ERA5 reanalysis profiles and MODIS infrared radiances
 466 (error sources i, ii, and iii). To model this expected variability, we perturb these 3 quantities to
 467 derive 200 different realizations of each of the aforementioned 24000 cases. We do this to
 468 propagate the uncertainties in these quantities to uncertainties in simulated radiances and thereby,
 469 to uncertainties in retrieved CTP. This procedure is detailed below.

- 470 a) **Low-cloud CTP:** Mitra et al. (2021) showed that MISR low cloud CTH error is -230 ± 300
 471 m. This error is propagated to CTP error using the formula $\sigma_P = \left| \frac{P}{H} \right| \sigma_z$, where, σ_P is the
 472 pressure uncertainty at a pressure level P corresponding to a height uncertainty of σ_z , for a
 473 pressure profile that varies with height according to the formula $P(z) = P_0 e^{-z/H}$. Here, P_0

474 is the pressure at surface ($z = 0$), z is the altitude of the pressure level and H is the scale
 475 height of the atmosphere, given by the altitude where $P = P_0/e$. For every low-level cloud,
 476 we bias our estimate of low-cloud CTP by taking the pressure equivalent of MISR CTH +
 477 230 m (using the form for $P(z)$, given above) and then perturb the resulting P_{low} by drawing
 478 200 different samples drawn from a normal distribution given by $\mu = P_{low}$, $\sigma = \sigma_p$.

479 b) **ERA5 Reanalysis Error:** To estimate the error-covariances of the ERA5 temperature and
 480 moisture profiles, we used the results of all model ensemble (Hersbach et al., 2020) that
 481 are publicly available along with ERA5 reanalysis (given by the ensemble mean). These
 482 ensemble members provide flow-dependent uncertainties based on propagation of
 483 assimilated measurement uncertainties as well as perturbations to physical tendencies. We
 484 took data from all grid cells over the globe over a day from each month of 2016 and
 485 calculated flow-dependent perturbations by subtracting each ensemble member from the
 486 ensemble mean. We then grouped the perturbations by latitude and season in the 5 pre-
 487 defined climate regimes (Section 3.1). Here, we estimated the error-correlations between
 488 all pressure levels of the profiles of temperature and moisture reanalysis, neglecting error-
 489 correlations between adjacent columns. Horizontal error correlations are neglected, as they
 490 are only relevant for the aggregation of pixel retrievals, not for individual pixel-level
 491 uncertainties. Upon comparing against estimates of ERA5 uncertainty from field studies
 492 (Graham et al., 2019), we found that the ERA5 ensemble variance is similar to observed
 493 uncertainty for specific humidity profiles. However, the ensemble uncertainty
 494 underestimates observed uncertainty of Graham et al. (2019) by a factor ranging between
 495 4-6, depending on pressure level. To correct this discrepancy, temperature profile
 496 perturbations from the ERA5 ensemble data are inflated by a constant value of 5, for all
 497 pressure levels. For each climate regime, we then propagated the resulting errors to errors
 498 in CTP through Monte Carlo sampling. Specifically, we drew 200 perturbed profiles of
 499 temperature and specific humidity assuming multivariate Gaussian distributions. The mean
 500 value of these distributions are given by their climatological profiles and their covariance
 501 matrix is set as described above.

502 c) **Instrument Noise:** We introduced further perturbations to the calculated TOA radiances,
 503 by drawing 200 random samples from a normal distribution with $\mu = 0$, $\sigma = 1 \text{ W m}^{-2}$. Here,
 504 we have set σ as the mean noise level for the Terra MODIS CO₂-slicing channels (as noted
 505 in Section 3.2.1, the noise levels in Bands 33-36 varies between 0.75-1.25 W m^{-2}).

506 To model the error from finite cloud geometric depth (error source iv), we modify the gas-only
 507 model (Section 3.2) for clear-sky radiative transfer to include cloud. We prescribe a cloud optical
 508 depth, cloud-top and bottom pressure (based on our choices of P_{high} , P_{low} , $depth$, τ_{high} listed (a)
 509 to (d)). We assume that cloud extinction is homogeneously distributed in pressure over the cloud
 510 depth. We verify our implementation using the analytic solution for an isothermal, non-scattering
 511 atmosphere. We use this model to simulate radiances in the CO₂-slicing bands for geometrically
 512 thick, non-black clouds and estimate the CTP retrieval errors stemming from the infinitesimally
 513 thin high cloud assumption of the CO₂-slicing technique. Gas optics uncertainties are numerically
 514 insignificant ($\ll 1\%$ of instrument noise) (Hannon et al., 1996) and are hence, ignored.

515 With the major sources of systematic and random errors accounted for, we run the MM_CTH
 516 algorithm for all 200 perturbed instances of each of the 24000 combinations of

517 $(P_{high}, P_{low}, depth, \tau_{high}, \lambda_{pair}, climate\ zone)$. We note the bias and standard deviation in CTP
 518 for each of those instances to construct the error function, E for comparison against observed error.

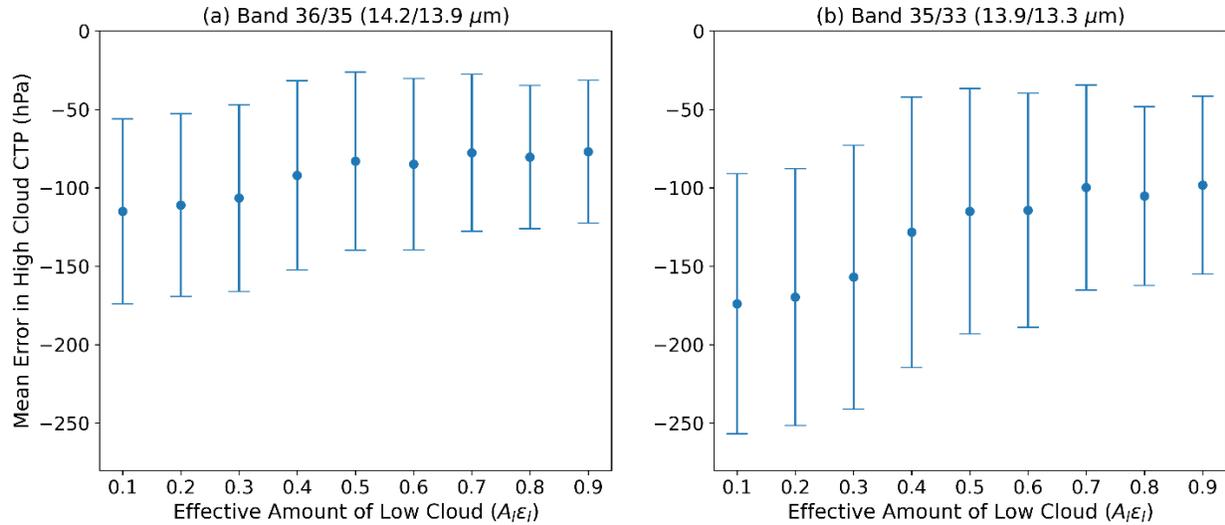
519 To account for further sources of random error (error sources v and vi), we estimated the
 520 uncertainty in CTP introduced by the process of geo-collocation of MODIS and CATS pixels.
 521 Mitra et al. (2021) showed a maximum uncertainty of 900 m in CTH due to the geo-collocation of
 522 MODIS and CATS pixels for CATS retrievals above an altitude of 5 km. Using the equation to
 523 propagate height errors to pressure errors given earlier, we estimate this collocation uncertainty
 524 (given by σ_{coll}) for all pixels. The errors in interpolating our CTP solutions to the discrete grid
 525 employed by the MODIS algorithm also result in an additional source of random error. This error,
 526 which we denote by σ_{grid} is numerically equal to half the grid-spacing between the nearest two
 527 levels of a CTP solution. As in Mitra et al. (2021), the random error in CATS CTH (converted to
 528 a CTP error given by σ_{CATS}) is equal to that associated with an equal probability of successful or
 529 failed retrieval over a 60 m CATS range gate, i.e., a random error of 30 m. Since, these sources of
 530 error are mutually independent, we estimate total random uncertainty (in a pixel-level retrieval) as

531
$$\sigma = \sqrt{\sigma_{modelling}^2 + \sigma_{coll}^2 + \sigma_{grid}^2 + \sigma_{CATS}^2}$$
 where, $\sigma_{modelling}$ is the error incurred from the
 532 various uncertainties in the radiative transfer simulations (sources i to iv), that are accounted by
 533 the standard deviation estimates from the error matrices, E .

534 To ascertain the fraction of pixels that are bound by our calculated total error estimates in E , we
 535 investigated the distribution of bias-corrected errors, normalized by σ , i.e., $\frac{CTP_{MM} - bias - CTP_{CATS}}{\sigma}$,
 536 where CTP_{MM} is the estimated value of CTP from the MM_CTH method, CTP_{CATS} is the observed
 537 (also, the assumed “true”) CTP from CATS, whereas, $bias$ is the closest estimate of theoretical
 538 systematic error for a particular pixel from the error matrices, $E(P_{high}, P_{low}, depth, \tau_{high}, \lambda_{pair})$.
 539 We find 78% of all pixels to be within the bounds of 95% confidence interval (i.e., [-1.96, 1.96]
 540 in units of σ). The remaining 17% (i.e., 95% minus 78%) of errors remain outside the purview of
 541 what can be constrained against empirically observed variables. We suspect that low cloud non-
 542 opacity and uncertainty in surface emissivity are the reasons behind these outliers.

543 We argue that surface emissivity is a less significant source of uncertainty than low clouds because
 544 in most multi-layered cases, the surface remains partly to nearly obscured by an opaque low cloud
 545 and >70% of all retrievals in our dataset are done by the 36/35 band pair (which is nearly
 546 insensitive to surface emissions; Menzel et al., 2015). Moreover, the effect of surface emissivity
 547 only becomes relevant in the very cases where the black low cloud assumption breaks down – e.g.,
 548 for broken low clouds. Hence, we do not investigate surface emissivity separately. To investigate
 549 the effects of low-cloud properties, we first note that the non-opacity of low clouds (i.e., $A_l \epsilon_l \neq 1$)
 550 may arise due to the presence of sub-pixel low clouds (e.g., small trade wind cumuli) or due to the
 551 presence of optically thin low clouds with $\epsilon_c < 1$. To quantify the errors in such scenarios, we
 552 relaxed the condition of a low, black cloud by assuming low cloud effective amounts of 0.1
 553 iterations between 0.1-0.9 for each of the 24000 test cases listed above. Effective IR emissivity of
 554 the low cloud is then converted to cloud optical depth (using Eq. 10 with $\zeta = 2.56$ for liquid water
 555 (Minnis et al. (1993)), and the transmission profile is adjusted accordingly. Surface emissivity is
 556 taken to be 1. In spite of the non-black low cloud, we still solve for the high cloud CTP assuming
 557 $A_c \epsilon_c = 1$. The mean and standard deviation of the resulting errors over all possible cases, for each
 558 value of low-cloud effective amount and MODIS CO₂-slicing band pair, are computed and shown

559 in Figure 4. For the Band 36/35 pair, unsurprisingly (since this pair is less sensitive to surface
 560 emission), low-cloud semitransparency leads to lower and nearly constant error, irrespective of the
 561 low cloud amount (especially, for $A_c \epsilon_c > 0.4$). However, the standard deviations of error for the
 562 Band 35/33 pair drops significantly as low cloud amount increases.



563

Figure 4. Distribution of errors in CTP (in hPa) incurred from the breakdown of the assumption of a black low cloud, from MODIS Band Pair 36/35 (left) and 35/33 (right) for different values of thermal IR effective emissivity ($A_c \epsilon_c$) of the low cloud.

564 Taking the effect of non-opaque low cloud into account, we redefine the bias-corrected errors to
 565 mean $\frac{CTP_{MM} - bias - bias_{low} - CTP_{CATS}}{\sigma}$, where $bias_{low}$ is defined as the mean bias for both band-pairs
 566 in Fig 4, weighted by their relative frequency of usage in our dataset. We calculate distributions of
 567 bias-corrected error (in units of σ) for all values of $A_l \epsilon_l$ and study the percentage of errors which
 568 lie within 95% CI in each case. Taking low clouds into account results in $> 80\%$ of all pixels lying
 569 within 95% CI for all values of $A_l \epsilon_l$. We find that the maximum agreement between theoretical
 570 and observed errors is achieved for $A_l \epsilon_l = 0.3$, resulting in 88% of all bias-corrected errors within
 571 the 95% CI. Here, we note that the expected dominant effect of low-cloud heterogeneity is likely
 572 from sub-pixel clouds. Assuming $\epsilon_l = 1$, this would suggest that the average value of low-cloud
 573 fraction in our dataset is $A_l = 0.3$. For 1-km resolution MODIS pixels, a low-cloud fraction of 0.3
 574 equals an average area-equivalent diameter for low clouds in our dataset of 620 m. This seems
 575 reasonable as our dataset has samples from both trade cumulus regions with typical cloud
 576 diameters of ~ 450 m (Zhao and Di Girolamo, 2007) and from regions with more stratiform clouds
 577 (that would typically cover the entire 1 km MODIS pixel). Thus for $A_l \epsilon_l = 0.3$, only 7% of all pixels
 578 are not constrained by our theoretical estimates (denoted by 95% CI), we can say that a near-
 579 closure of the MM_CTH CO₂-slicing error budget has been achieved. The sources of error that
 580 could potentially explain these outliers are the incomplete modeling of low-cloud uncertainties,
 581 uncertainties in surface emissivity, inaccuracies in MODIS cloud phase detection and the
 582 assumption in CO₂-slicing of equal ice-cloud effective emissivities in closely spaced IR channels.

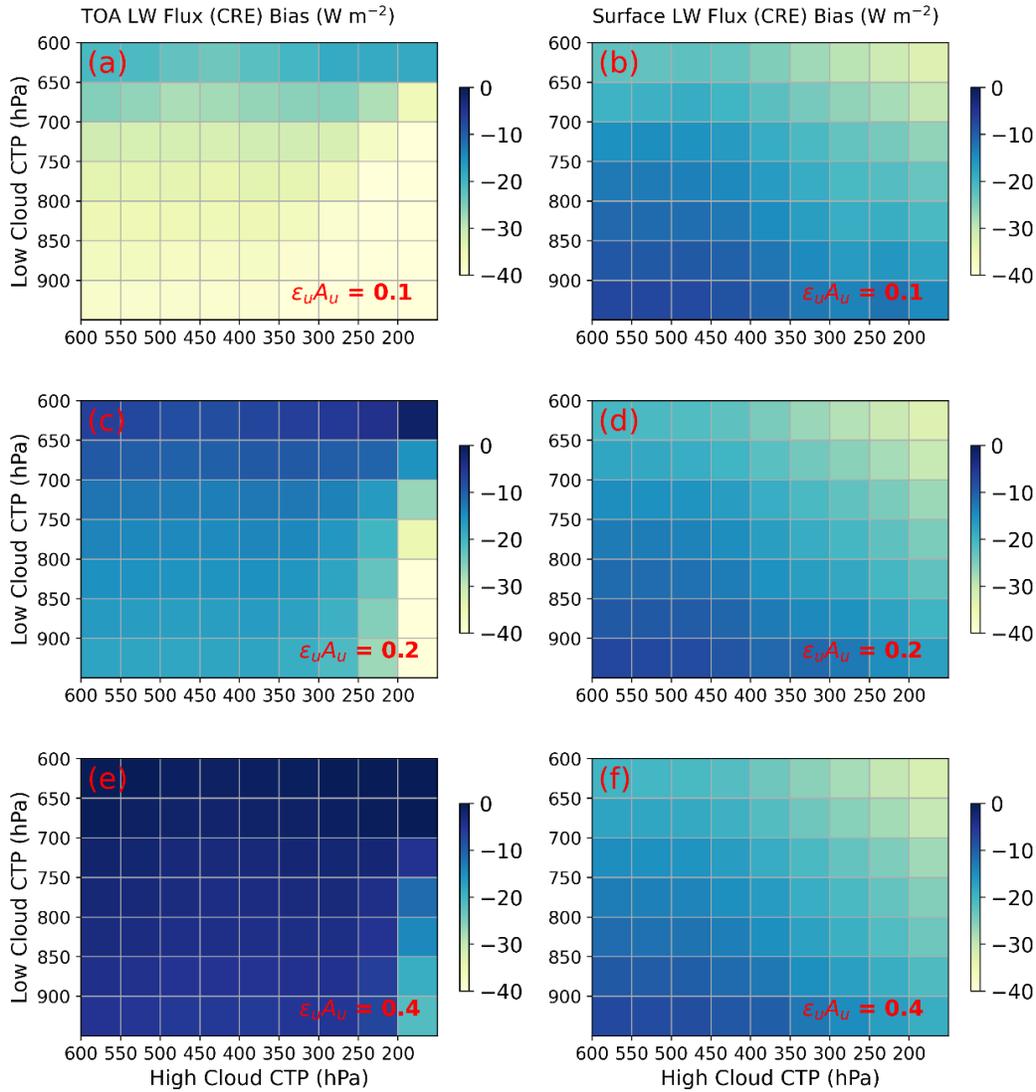
583 **5. Impact on Cloud Radiative Effect**

584 As noted in Section 1, the vertical distribution of cloud properties controls the vertical variation of
585 cloud radiative effect (CRE), defined here as the difference in upwelling cloudy and clear sky
586 radiative fluxes at the top of the atmosphere (and similarly for downwelling radiative fluxes at the
587 surface). When high and low clouds coexist in multi-layered situations, the CRE will depend on
588 the optical properties of the two layers and their geometric locations within the atmosphere; the
589 later controlling their temperature and the extent of absorbing gases above, below and between the
590 cloud layers. Hence, the longwave (LW) or shortwave (SW) CRE due to a 2-layered cloud system
591 cannot simply be expected to equal the corresponding CRE due to the ‘effective’ single-layered
592 ice cloud with CTP and effective emissivity from a 1-layered CO₂-slicing solution. Thus, the
593 accurate representation of the macrophysical and optical properties of both cloud layers in a scene
594 is likely needed for accurate estimation of CRE in radiative transfer simulations. As a result, the
595 accuracies of the MM_CTH method in determining macrophysical and optical cloud properties in
596 2-layered systems (Section 4.1) are expected to improve our estimates of modeled CRE for 2-
597 layered systems. Here, we demonstrate this improvement due to the implementation of MM_CTH.
598 We do this simply by estimating the impact of the 1-layer CO₂-slicing CTP and effective emissivity
599 biases on simulated TOA upwelling and surface downwelling LW CRE. The impact of single-
600 layer retrievals are ostensibly significant for shortwave (SW) CRE as well. Here, we do not study
601 the SW CRE bias as that would be strongly dependent on multiple factors beyond layer-averaged
602 properties (e.g., ice/water single-scattering properties and sun-satellite geometry), which would be
603 beyond a concise explanation for the simple demonstration we are aiming for.

604 To estimate the LW impact of 1-layered CO₂-slicing retrievals applied to a 2-layered cloud system,
605 we run radiative transfer simulations for different combinations of high and low cloud CTP and
606 high cloud effective emissivity. In each of these cases, we calculate TOA upwelling and surface
607 downwelling LW CRE for both the ‘True’ 2-layered cloud configuration (that we prescribe) and
608 the ‘Effective’ single-layered ice cloud parameters (from a 1-layered CO₂-slicing retrieval). We
609 define the CRE bias resulting from the application of a 1-layered CO₂-slicing as the difference
610 between the ‘True’ CRE and ‘Effective’ CRE, defined as follows. ‘True’ LW CRE is defined as
611 the difference between cloudy and clear-sky LW atmospheric radiative fluxes calculated using our
612 pre-defined parameters for higher ice and lower water cloud properties. However, after the
613 application of a 1-layered CO₂-slicing retrieval, we retrieve a single ice cloud layer at a lower
614 altitude than the true altitude of the upper layer (Sections 3.2.1 and 4.1), along with its effective
615 emissivity that is larger than its true emissivity (Section 4.1). We then use this retrieved 1-layer
616 CTP and effective emissivity to calculate the LW CRE and compare it to the true LW CRE to
617 assess the LW CRE bias. Further details of the radiative transfer simulations are in Text S2 of
618 Supporting Information, which are for thin cirrus overlying a lower liquid water cloud that is
619 opaque in the infrared. Figure 5 shows the variation of the surface and TOA LW CRE bias with
620 true high and low cloud CTP and high cloud effective emissivity.

621 The left panels of Fig. 5 shows that the TOA LW CRE bias is sensitive to both the cloud
622 macrophysics and high-cloud emissivity, which the true LW CRE is also sensitive too. The
623 absolute value of the bias decreases with increasing effective emissivity of the upper cloud. As
624 shown in Sections 3.2.1 and 4.1, applying a 1-layered CO₂-slicing retrieval on a 2-layered system
625 results in overestimations in CTP and $A_c \epsilon_c$ for the upper-cloud layer. Since the retrieved cloud is
626 lower in altitude, hence warmer, and more opaque in the infrared, the resultant top-of-atmosphere

627 LW CRE bias is negative, as shown in Figure 5, with the largest absolute bias ($\sim 40 \text{ W m}^{-2}$)
 628 occurring for thin clouds near the tropopause overlying low altitude clouds.



629

630 *Figure 5. Variation of top-of-atmosphere (left panels) and surface (right panels) LW Flux (CRE) bias (W m^{-2}) with*
 631 *variations in high and low CTP, due to a single-layered CO_2 -slicing retrieval on a 2-layer scene. The atmosphere and*
 632 *surface properties are set up similar to Figure 1. CRE bias is defined as true minus modeled LW CRE. High Cloud*
 633 *Effective Emissivity is taken to be 0.1 (top panels), 0.2 (middle panels) and 0.4 (bottom panels).*

634 The right panels of Fig. 5 show the surface LW CRE bias is strongly sensitive to cloud
 635 macrophysics but less sensitive to high-cloud emissivity. Unlike the TOA, the true LW CRE at the
 636 surface is dependent only on the height of the low cloud because its LW emissivity is one in our
 637 simulations. Thus, in order to achieve the little differences that we see between Fig. 5(b), (d) and
 638 (f), the surface LW CRE calculated using the 1-layer CO_2 -slicing solution must also be somewhat
 639 insensitive to the effective emissivity of the upper cloud. This occurs because the 1-layer CO_2 -
 640 slicing solution produces a larger CTP bias, hence warmer cloud, for clouds with smaller
 641 emissivity compared to clouds with larger emissivity. Thus, changes in high cloud effective
 642 emissivity leads to competing changes in the resultant 1-layered retrieval (cloud temperature

643 versus emissivity), thus impacting surface LW CRE bias only weakly. Thus, it is the heights of the
 644 two cloud layers that have the largest effect on the LW CRE bias, with absolute values of the bias
 645 being largest ($\sim 30 \text{ W m}^{-2}$) for high tropospheric clouds overlying mid-level clouds.

646 Based on these findings, application of the MM_CTH algorithm is expected to provide
 647 improvements in modeled LW radiative fluxes that are of a similar order of magnitude ($\sim 10 \text{ W m}^{-2}$)
 648 to the CRE biases calculated here. These improvements to modeled radiative fluxes will be
 649 helpful when estimating, for example, the surface and atmospheric radiation budgets based on
 650 retrieved cloud properties [e.g., Kato et al. (2018)]. They may also provide a set of cloud properties
 651 that have variability that is more consistent with the variability in Earth's radiation budget, thereby
 652 providing improved benchmarks for the evaluation of climate models.

653 **6. Conclusions**

654 Thin cirrus cloud overlying low clouds constitute $>80\%$ of multi-layered clouds globally (multi-
 655 layered clouds themselves constitute $\sim 30\%$ of all cloud cover) (Wang & Dessler, 2006; Oreopoulos
 656 et al., 2017; Hong and Di Girolamo 2020). For 2-layered scenes, MODIS underestimates top-layer
 657 CTH by $>1 \text{ km}$ as the CO_2 -slicing technique converges at a higher CTP solution, when an optically
 658 thin cirrus is present. As a result, MODIS produces more midlevel CTH than MISR and MISR-
 659 MODIS CTH differences have generally low absolute values (Naud et al., 2007; Mitra et al., 2021).
 660 However, MISR often retrieves the lower cloud height in a majority ($>80\%$) of such 2-layered
 661 cases, provided the top-layer optical depth $< \sim 0.4$ (Mitra et al., 2021). In this study, we have
 662 developed an algorithm to retrieve accurate high-cloud properties for 2-layered cloud systems,
 663 named the *MISR-MODIS Fusion Product for Cloud-Top Height (MM_CTH)*. MM_CTH used a
 664 modified version of the standard MODIS CO_2 -slicing algorithm (of the Collection 6.1 MOD06
 665 product), using accurate MISR low-cloud CTH retrievals as an input to account for the presence
 666 of the lower cloud in multi-layer scenes. Using collocated ISS-CATS as a reference, we validate
 667 the MM_CTH retrievals to find a $\sim 90\%$ reduction in cirrus CTP bias over MOD06. This
 668 improvement to CTP accuracy propagates to $\sim 75\%$ improvements in accuracy for cirrus CTH and
 669 effective emissivity over the standard MOD06 products. The MM_CTH algorithm also allows us
 670 to retrieve lidar-like distributions of high cloud macrophysics (Figure 2b and 2d) and optical
 671 properties (Figure 3) in 2-layer cloud systems from passive sensors. Table 1 summarizes the results
 672 of the validation (Section 4.1) of CO_2 -slicing CTP, CTH and thermal IR $A_c \epsilon_c$ (against CATS), and
 673 the distributions of CATS, MOD06 and MM_CTH CTP, CTH and $A_c \epsilon_c$.

Data Source	Mean Errors (with respect to CATS)			Net Distribution for High Clouds		
	CTP (hPa)	CTH (km)	$A_c \epsilon_c$	CTP (hPa)	CTH (km)	$A_c \epsilon_c$
MOD06	65 \pm 85	-1.6 \pm 2.3	0.4 \pm 0.3	300 \pm 85	9.7 \pm 2.3	0.5 \pm 0.3
MM_CTH	5 \pm 80	-0.4 \pm 2.4	0.1 \pm 0.2	235 \pm 70	11.2 \pm 2.0	0.2 \pm 0.2
CATS	N/A	N/A	N/A	225 \pm 80	11.7 \pm 2.5	0.1 \pm 0.2

Table 1. Summary of mean errors in CO_2 -slicing CTP, CTH and effective emissivity for MOD06 and MM_CTH with respect to CATS and the mean value of the retrieved distributions of CTP, CTH and effective emissivity from MOD06, MM_CTH and CATS.

674 We also performed a detailed error budget analysis using CATS high cloud retrievals as reference.
675 CATS high cloud retrievals, ERA5 modeling error estimates, and estimates of MISR CTH and
676 MISR, MODIS, CATS geo-collocation errors from Mitra et al., (2021) are used to model the
677 systematic and random sources of CTP error, which are then compared against empirical estimates
678 of errors (from comparison with CATS). 78% of all observed errors were found to be within
679 theoretical limits (i.e., 95% CI), when non-opacity of low-cloud properties (stemming primarily
680 from sub-pixel clouds) are neglected. However, when the sub-pixel nature of low-cloud is
681 accounted for, up to 88% of observed MM_CTH error estimates fall within the limits of 95% CI
682 – thus providing a near-closure of the MM_CTH error budget. The lack of a truth dataset for low-
683 cloud cloud fraction and emissivity, uncertainties in prescribed surface emissivity, inaccuracies in
684 MODIS cloud phase detection and the assumption in CO₂-slicing technique that ice-cloud effective
685 emissivities in closely spaced IR channels are equal could potentially lead to the existence of the
686 7% outlier pixels. Since the benefit of including an estimate of sub-pixel (i.e., within a 1-km
687 MODIS pixel) low-altitude cloud fraction is significant, it is recommended that MISR's
688 Stereoscopic Derived Cloud Mask (SDCM; Mueller et al. (2013)) be reported at the native
689 resolution of MISR, i.e. 275 m, rather than its current resolution of 1.1 km.

690 We demonstrated that the improvement in high cloud properties from the MM_CTH algorithm
691 may be highly relevant in studies involving Earth's radiation budget. In 2-layered cloud systems,
692 our results show improved estimates of modeled atmospheric fluxes (demonstrated for TOA and
693 surface LW CRE in Figure 5) by ~5 to 40 W m⁻², depending on the 2-layered properties, when
694 using MM_CTH retrievals rather than the standard single-layer CO₂-slicing retrievals. Thus, our
695 algorithm could provide a climatology of CTH and high-cloud optical properties that is more
696 consistent with the fluctuations in the Earth's radiation budget than corresponding estimates from
697 standard MOD06 retrievals for multi-layered scenes.

698 Although this current study is concerned with introducing the pixel-level MM_CTH algorithm and
699 its validation and error budget analysis, we would like to stress its future importance to broader
700 climate science, especially in leveraging the 22-year-long stable Terra record to study long-term
701 climate-scale cloud responses, especially for high cloud populations. Of the many cloud responses
702 to anthropogenic forcing predicted by models, the highest confidence is associated with rising
703 CTHs (Boucher et al. 2013). Rising CTH is predicted to be the first signal of forced change that
704 will emerge above natural variability (Chepfer et al., 2014; Winker et al., 2017). For example,
705 simulations of a uniform 21st century 4K warming had predicted the increase in high cloud amounts
706 by ~5-15%, along with ~25 m/year increase in mean tropical high CTH (Chepfer et al., 2014). In
707 fact, there have been non-significant detection of the expected rising patterns in global high cloud
708 amounts from passive sensors (Norris et al., 2016; Aeronson et al., 2022). For confident detection
709 of such trends, however, we need stable multi-decadal observations (subject to robust uncertainty
710 analysis) of cloud vertical distribution, globally (Shea et al., 2017). While active sensors capable
711 of vertically resolving cloud layers like lidars might seem ideal, the emergence of such trends from
712 lidars are thwarted by their short lifetimes and lack of swath coverage. Hence, multidecadal passive
713 sensor records from stable-orbit satellites like Terra are still the best suited for such a task.

714 However, as demonstrated in Section 1, both stereoscopic and multi-spectral retrievals of cloud
715 macrophysics suffer from issues of sensitivity to different cloud types and accuracy. MISR stereo
716 misses a majority of cirrus in 2-layered cases. On the other hand, unless the cirrus is very thin (OD
717 $\ll 1$), MODIS IR channels detect cirrus emission above the channels' noise levels, but it is the

718 restrictive choice of a 1-layer solution (in the MODIS forward model) that leads to the
 719 misrepresentation of cirrus properties, including its retrieved emissivity. Left unchecked, it would
 720 be difficult to impossible to decouple long-term changes in high cloud heights and emissivity from
 721 true changes in low cloud heights and amount using MODIS data alone. Similarly, it would be
 722 difficult to impossible to decouple long-term changes in low cloud heights and amounts from true
 723 changes in high cloud amount and optical depths from MISR data alone. MM_CTH is a means to
 724 tackle these problems as it can provide lidar-like distributions of high cloud properties over a
 725 passive sensor swath (the MISR swath) over the 22-year stable-orbit satellite record of Terra.

726 Due to its unmatched stability and longevity, the Terra record will remain a unique climate record
 727 of global cloud macro-physical and optical properties between 2000-2022. We are therefore left
 728 with the goal to ensure that the Terra record produces cloud products with well-characterized
 729 uncertainties for future studies on the Earth's climate. Towards this goal, we strongly recommend
 730 that the pixel-level MM_CTH algorithm introduced here be scaled to a fully operational product
 731 over the entire Terra record for public dissemination.

732

733 **7. Acknowledgements, Software and Data Sources**

734

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 738 Software modified to create the MM_CTH software was downloaded from the NASA Goddard
 739 Space Flight Center MODIS Adaptive Processing System (MODAPS) website
 740 (<https://modaps.modaps.eosdis.nasa.gov/software/MODIS/AM1M/PGE06/Collection61/>). The
 741 MISR data was downloaded from NASA Langley Research Center Atmospheric Sciences Data
 742 Center (<https://opendap.larc.nasa.gov/opendap/MISR/MIL2TCSP.001/>). The MODIS data were
 743 obtained through the Level 1 and Atmosphere Archive and Distribution System of NASA Goddard
 744 Space Flight Center (<https://ladsweb.modaps.eosdis.nasa.gov/archive/allData/61/>). The CATS
 745 data was downloaded from the NASA Langley Research Center's ASDC DAAC
 746 (<https://opendap.larc.nasa.gov/opendap/CATS/>). We are thankful to the NASA MODIS, MISR
 747 and CATS teams for supplying the documentation and tools, including the MISR toolkit
 748 (<https://nasa.github.io/MISR-Toolkit/html/index.html>). All ERA5 Reanalyses are downloaded
 749 through the European Center for Medium-Range Weather Forecast (ECMWF) Climate Data Store
 750 (CDS) website (<https://cds.climate.copernicus.eu/cdsapp#!/home>). The geoid data used in this
 751 study was downloaded from the National Geospatial-Intelligence Agency (NGA) WGS84 website
 752 (<https://earth-info.nga.mil/index.php?dir=wgs84&action=wgs84>). Data were stored and
 753 computations were conducted on the computing infrastructure managed by the University of
 754 Illinois at Urbana-Champaign's School of Earth, Society, and Environment (SESE).

755

756 **8. References**

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