

Different trends in Antarctic temperature and atmospheric CO₂ during the last glacial

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Key Points:

- The Antarctic warming rate during Antarctic Isotope Maxima significantly decreased as the climate cooled toward the glacial maximum.
- The Antarctic warming rate during Antarctic Isotope Maxima is insensitive to whether the warming coincides with a Heinrich event.
- In contrast, the millennial-scale CO₂ rise is insensitive to long-term glacial cooling but sensitive to the presence of Heinrich events.

Abstract

Using Antarctic ice-core records, we determine for each Antarctic Isotope Maximum (AIM) of Marine Isotope Stage 3 (MIS-3: ca. 28,000 to 59,000 years before present) the rates and durations of warming and atmospheric CO₂ rise. We find that the AIM warming rates significantly decrease as the climate cools from early to late MIS-3. In contrast, the rate of CO₂ rise during AIMs shows no significant trend across this interval. We further find that the AIM warming rate is not sensitive to Heinrich (H) events, contrasting with CO₂, which rises for significantly longer time (compared to the temperature rise) during AIMs which coincide with H events. These distinct Antarctic temperature and CO₂ responses to varying background climate and H events challenge the view that millennial-scale CO₂ and Antarctic temperature changes are dominated by the same physical processes, suggesting an important contribution of low-to-mid-latitude processes to the CO₂ rises.

Plain Language Summary

Glacial climate is characterized by millennial-scale variations in polar temperature and atmospheric CO₂ concentration. Over the last two decades, a consistent explanation of the temperature changes has emerged, but no such consensus exists with regard to CO₂. However, due to the similarity of their records it is frequently proposed that CO₂ and Antarctic temperature were controlled by the same processes. Here we present a new analysis of millennial-scale Antarctic warming and CO₂ rise based on ice-core data. Our results show that during the latter half of the glacial period, the Antarctica warming rate decreased as the climate cooled, but it was not affected by occasional massive iceberg discharges (known as Heinrich events) that had a dramatic impact on northern hemisphere climate. On the other hand, the rate of CO₂ change was

insensitive to the glacial cooling trend, but the CO₂ rise was sensitive to the occurrence of Heinrich events. This suggests that on top of the processes that control millennial-scale Antarctic temperature variations and also play a role for CO₂ levels, there are other processes (possibly of extra-polar or terrestrial origin) that are important for the CO₂ dynamics.

1 Introduction

Antarctic Isotope Maximums have a systematic relationship with the Dansgaard-Oeschger (DO) oscillations recorded in Greenland ice cores during the last glacial period (EPICA, 2006). The cold (Greenland stadial: GS) phase of the DO oscillation in Greenland coincides with AIM warming in Antarctica and the warm (Greenland interstadial: GI) phase of the DO coincides with AIM cooling (Figure 1). According to the bipolar ocean seesaw hypothesis (Stocker & Johnsen, 2003), this interhemispheric coupling results from changes in northward heat transport by the Atlantic Meridional Overturning Circulation (AMOC): GSs are associated with a weak AMOC, reduced northward heat transport in the Atlantic and warming in the South Atlantic. AIM warming in Antarctica follows, after a centennial-scale lag, as the South Atlantic warm anomaly spreads through the upper to intermediate-depth ocean and is gradually mixed across the Antarctic Circumpolar Current, in turn melting back Southern Ocean sea ice and increasing atmospheric heat transport to Antarctica (see Pedro et al., 2018; for a view that places more emphasis on buoyancy forcing see Thompson et al., 2019).

Previous work, based on the EPICA Dronning Maud Land (EDML) ice core, identified a strong linear relationship ($r^2 = 0.85$) between the duration of GSs and the amplitude of corresponding AIMs during MIS-3 (EPICA, 2006). The linear relationship suggested that the amplitude of AIM warming depends only on the duration of Greenland stadials and that the operation of the bipolar seesaw is not significantly influenced either by the evolution of the

climate state during the glacial, nor by the presence of H events. However, only eleven of the sixteen MIS-3 AIMs were considered in that analysis and more recent studies propose that the bipolar seesaw is sensitive to background climate (Capron et al., 2010; Margari et al., 2010), and the presence of H events (Margari et al., 2010). However, these sensitivities have not previously been quantified or statistically assessed.

The apparent similarity between atmospheric CO₂ concentrations (hereafter CO₂) and Antarctic temperature change in the Antarctic ice-core record (Figure 1) has been hypothesized to result from a common cause, such as changes in wind-driven upwelling (Anderson et al., 2009; Anderson & Carr, 2010; Menviel et al., 2018; Toggweiler et al., 2006), or changes in the formation of Antarctica Bottom Water (AABW, Menviel et al., 2015). In both cases, the release of CO₂ and heat involve coupled Southern Ocean processes: increased ventilation of relatively warm and carbon-rich sub-surface waters in the Southern Ocean (Bauska et al., 2018; Gottschalk et al., 2019; Jaccard et al., 2016; Skinner et al., 2020); and/or increased poleward heat transport accompanying the AABW formation and CO₂ ventilation (Menviel et al., 2015). However, there are two obvious differences between the millennial-scale CO₂ and Antarctic temperature changes: a) CO₂ maximums associated with Heinrich stadials (HS, defined here as GS containing a H event) visibly lag the peak of the corresponding AIM events (Figure 1; Bereiter et al., 2012); and, b) CO₂ rises during non-Heinrich stadials (nHS), if resolved at all, are thought to be of lower amplitude than those during HS (Ahn & Brook, 2014). However, as with the sensitivity of AIM warming to background climate and H events, a robust statistical assessment of these CO₂ features is lacking. Thus, we constrain for each MIS-3 AIM the amplitude and duration of the AIM warming phase and CO₂ rise. We then test the sensitivities of AIM warming and CO₂ rise to the changes in background climate state and the presence of H events.

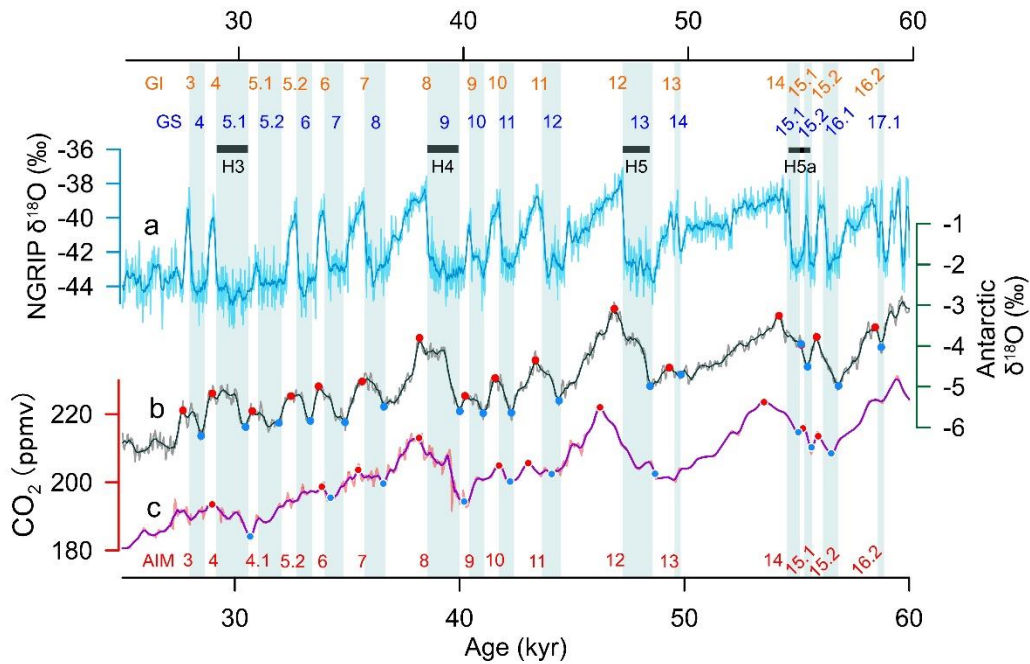


Figure 1. Millennial-scale climate events and atmospheric CO₂ variations for the MIS-3. a.

The oxygen isotope ratio ($\delta^{18}\text{O}$) from the North Greenland Ice Core Project (NGRIP) ice core, as a proxy for Greenland temperature (NGRIP, 2004; Svensson et al., 2008). **b.** The Antarctic five-core averaged $\delta^{18}\text{O}$ record (Buizert et al., 2018). **c.** The composite Antarctic ice core CO₂ record (Bereiter et al., 2012). The red and blue dots in **b** and **c** show the identified AIM and CO₂ maximums and minimums (see text for definitions). **b** is on the WD2014 timescale, all other records were transferred to the WD2014 time scale by stretching the GICC05/AICC2012 age to 1.0063 times (Buizert et al., 2015). The Greenland stadials and interstadials events are numbered following Rasmussen et al. (2014), the timing of Greenland climate transition is from Buizert et al. (2015) and WAIS (2015). Timespan of the H events are marked by the horizontal gray bars. The AIM event number is following EPICA (2006) and extended to AIM 16.2.

2 Methods and results

2.1. AIM temperature rise

We use the recently published Antarctic five-core averaged $\delta^{18}\text{O}$ stack (in WD2014 timescale; Buizert et al., 2014) to determine the individual MIS-3 AIM warming amplitudes. To exclude uncertainties associated with $\delta^{18}\text{O}$ -temperature transformation (Jouzel et al., 2013; Landais et al., 2015), we use per mille (‰) $\delta^{18}\text{O}$ change as our unit of measure for AIM warming amplitude. EPICA (2006) concludes that AIM warming amplitude is only controlled by the duration of the corresponding GS; i.e. that the AIM warming rate is constant. To test this result, we go a step further than EPICA (2006) and define here the AIM amplitude divided by the GS duration as the bipolar seesaw (BPS) warming. The GS durations are obtained from previous work (Buizert et al., 2015).

To obtain a robust estimate for the bipolar seesaw warming rate and its uncertainty, a Monte Carlo (MC) Methods is used: in every iteration, the maximums and minimums of each AIM are determined from a randomly perturbed version of the five-core averaged $\delta^{18}\text{O}$ record which is smoothed with a 200-yr moving average before calculating the BPS warming rate of each AIM. The randomly perturbed $\delta^{18}\text{O}$ record is created by drawing values from the normal distribution of the five-core averaged $\delta^{18}\text{O}$ record with standard deviation set at 0.12‰ (found as the standard deviation of the residual of the smoothed data relative to the unsmoothed one). Considering the 100 to 200 yr lag of the onsets and ends of AIM events relative to Greenland GS and GI transitions (Svensson et al., 2020; WAIS, 2015), we conservatively search for the isotope maximums and minimums in a 300 yr window starting at the time of the corresponding Greenland climate transition (the window is not allowed to cross the adjacent Greenland climate

transition). After 100,000 iterations the median BPS warming rate for each AIM is used as the output (Figure 2a), and the 95% confidence interval (CI) of the MC generated rates of each AIM is used as the estimate of uncertainty range.

In contrast to the results of EPICA (2006), we do not observe a strong linear correlation between GS duration and AIM amplitude (Figure 2a). Instead, our data suggest a lower BPS warming rate during late MIS-3 than during early MIS-3 (Figure 2a, S1). To test if this reduction is systematic, we plot the BPS warming rates against the AIM ages (defined as the onset of the corresponding Greenland stadial). A clear decline in the warming rate across MIS-3 is observed (Figure 2b). The significance of the observed slope is evaluated by comparing it with the slopes generated from randomly permuting the BPS warming rate data points (keeping the AIM ages) and re-calculating the resulting slopes 100,000 times. Only 0.06% of the absolute value of the randomly generated slopes are larger than that obtained from the actual data (Figure 2c), demonstrating the declining trend of BPS warming rate during MIS3 is significant.

Following EPICA (2006), our analysis above assumes that the BPS warming rate is linear. If the AIM warming instead rises asymptotically, as predicted by the “minimum thermodynamic seesaw model” (Margari et al., 2010; Stocker & Johnsen, 2003), then our definition of the BPS warming rate could be biased low for longer GS. To test whether our result holds under the assumption of an asymptotic rather than linear temperature increase during AIMs we repeat our analysis fitting the warming phase of each AIM using the minimum thermodynamic seesaw model (Stocker & Johnsen, 2003; Figure S2): $\Delta T_S = \Delta T_N(e^{-t/\tau} - 1)$, where ΔT_S and ΔT_N are the amplitude of southern warming and northern cooling respectively, t is the time since the start of the stadial, and τ is the equilibration timescale of the seesaw system, estimated as 1120 yr (Margari et al., 2010; Stocker & Johnsen, 2003). We use isotope units

rather than temperature, and introduce the constant $-T_G$ to replace ΔT_N , representing the amplitude of northern temperature (isotope) drop that gives the best fit to the observed AIM isotope growth rate (note the $-T_G$ and ΔT_N have opposite sign). With τ kept constant, the shape of the fitted curve is determined by $-T_G$, with smaller $-T_G$ corresponding to slower rises (Figure S3). Similar to the linear case, Monte Carlo Methods is applied, and $-T_G$ of each AIM is determined as the median of 100,000 iterations (Figure 2d).

Like the linear assessment of BPS warming rate, a slope showing the decline of $-T_G$ toward the glacial maximum is observed (Figure 2d). To test the significance of the slope, we randomly permute the $-T_G$ values and re-calculate the slope 100,000 times and find that the absolute value of the random slopes is larger than the value found from data in only 4.9% of the cases, (Figure 2e), suggesting that $-T_G$ has a significant trend during MIS-3. This could be due to a weakening of efficiency of the bipolar seesaw or a reduction of the driving ΔT_N , but as the latter is not found in the reconstruction of northern temperature (Kindler et al., 2014), we conclude that the bipolar seesaw was indeed weakening significantly during MIS-3 both under the Stocker and Johnsen (2003) minimum thermodynamic seesaw model and that of linear AIM warming rates (EPICA, 2006).

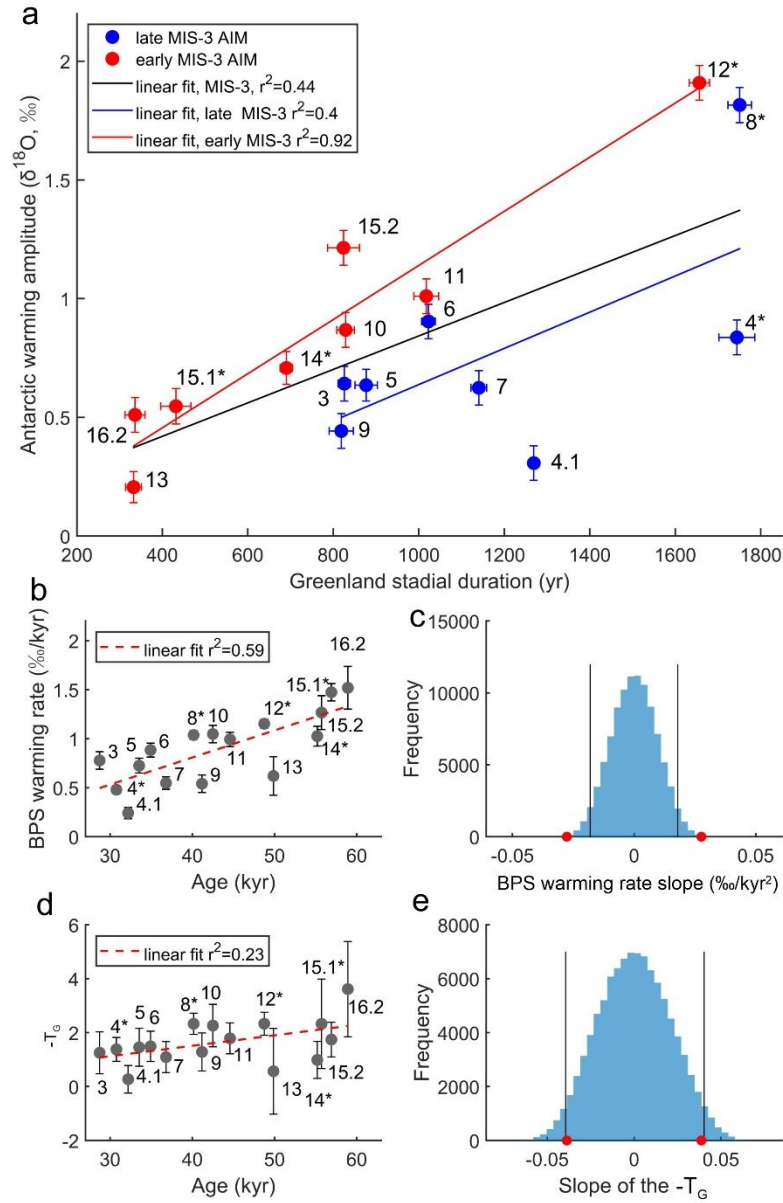


Figure 2. The BPS warming rate results. a. Greenland stadal duration and Antarctic warming amplitude, the AIM nomenclature is following EPICA (2006) and has been extended to AIM 16.2. The error bar shows the uncertainty of GS duration (Buizert et al., 2015; WAIS, 2015) and the 95% CI of the amplitude. **b.** The BPS warming rate plotted against the age of the AIM. The error bar shows the range of 95% CI. **c.** The distribution of the slopes of the randomly permuted BPS warming rates, the red circle marks ± 1 time the slope observed from **b**, the 2.5% and 97.5%

fractile of the randomly generated slope are marked by vertical black lines. **d, e.** The same as **b, c** for $-T_G$ instead of the BPS warming rate.

To test the sensitivity of AIM warming to H events we group the BPS warming rates and $-T_G$ values for each AIM into HS and nHS categories (Figure S1). A Student's t-test shows no significant difference (at 95% significance level) between HS and nHS BPS warming rates or $-T_G$. These results suggest that the processes controlling Antarctic warming are insensitive to the H events.

Our result of a weakening bipolar seesaw during MIS-3 is robust to replacing Greenland stadial durations with the duration of the corresponding Antarctic isotope rise (the time between the identified isotope minimum and maximum, Table S1). Removing the long-term signal represented by 20,000 yr smoothing of the five-core averaged isotope data also does not change the conclusions (Table S1, Figure S4). We also carried out the analysis of the BPS warming rate and $-T_G$ on the individual ice-core records going into the five-core averaged data set: West Antarctic Ice Sheet Divide (WDC; WAIS, 2013, 2015), European Project for Ice Coring in Antarctica (EPICA) in the interior of Dronning Maud Land (EDML; EPICA, 2006), Talos Dome (Landais et al., 2015; Stenni et al., 2011), EPICA Dome C (EDC; EPICA 2004), and Dome Fuji (Fuji; Kawamura et al., 2007; Watanabe et al., 2003). Similar results are obtained (Table S1, Figure S4). Although the EDML record may reflect more low-latitude and atmospheric signals (Landais et al., 2015), our results show that in terms of the overall bipolar seesaw response, it is not systematically different from other Antarctic ice cores. This is consistent with the conclusion that a spatially homogeneous oceanic component of the AIM events related to the

thermodynamic seesaw response is the first and dominant principle component of the AIM variability (Buizert et al., 2018).

2.2. Millennial CO₂ rise:

We determine the amplitude of MIS-3 millennial-scale CO₂ rise using the composite CO₂ record (on the AICC2012 timescale; Bereiter et al., 2015; Figure 1, S5). The MIS-3 section of the composite CO₂ record consists of the Siple Dome (20 to 40 kyr BP; Ahn & Brook, 2014) and Talos Dome (40 to 60 kyr BP; Bereiter et al., 2012) records. To avoid the influence of different gas smoothing between ice cores, the CO₂ data is smoothed using a 300 yr moving average, which is larger than the mean resolution of both CO₂ records (96 and 246 for Siple and Talos Dome respectively). We note that there is a systematic offset between the two CO₂ records (Bereiter et al., 2012) with a junction point (at about 40 kyr BP, Figure S5), this offset does not fall within a period of assessed CO₂ rise, so it does not affect any of our derived CO₂ amplitudes.

Considering the CO₂ peak could lag the abrupt DO warming by several hundred to a thousand years (Bereiter et al., 2012; Figure S5), we here extend the search window for the CO₂ maximums and minimums to 100 yr prior and 1000 yr after the Greenland climate transition. We use the timing of Greenland climate transitions defined in Greenland $\delta^{18}\text{O}$ and Ca⁺ data (Rasmussen et al., 2014).

In calculating the amplitudes and rates of MIS-3 CO₂ rise, we assess only those CO₂ rises with amplitudes significantly larger than the measurement uncertainty (at the 95% level, Figure S6). A total ten CO₂ rise are identified, five for the HS (DO 4, 8, 12, 14, 15.1) and five for the nHS periods (DO, 6, 7, 10, 11, 15.2, Figure 1, 3a). We here focused on the overall CO₂ response

and defined the rate of CO₂ rise (Figure 3b) as CO₂ amplitude divided by the duration of the rise. The CO₂ rate and its uncertainty are estimated by similar MC method of temperature rates.

In contrast to AIM warming rates, we do not find evidence for a significant decline in the rate of CO₂ rise across the MIS-3 AIM events (with two-side significance 82.7%, Figure 3c). Note that this result is not a statistical artifact due to the low number of CO₂ events resolved by the data, see a one-to-one comparison using only events both analysed in the CO₂ and temperature records in Figure S7. However, the bubble enclosure characteristics of different ice core sites (Bereiter et al., 2012) and the depth/age-dependent ice diffusion (Ahn et al., 2008) could unevenly smooth the composite CO₂ records and bias the CO₂ amplitude. To test the influence of these effects, we perform additional tests and compared the CO₂ and temperature trends for the Talos section of composite CO₂ (Bereiter et al., 2012) vs the Talos $\delta^{18}\text{O}$ data (Landais et al., 2015; Stenni et al., 2011; Supporting Information, Figure S8), and for the newly recovered EDC CO₂ record (Nehrbass-Ahles et al., 2020) spanning 330 to 440 kyr BP vs the EDC δD data (EPICA, 2004), (Supporting Information, Figures S9 and S9). In both cases, the CO₂ trend is less significant than the Antarctic temperature trend (Figure S8, S9). The above experiments are also designed to consider the influence of using different timescales and changes in delta-age (age difference between gas and ice at the same depth, see Supporting Information). The results show that our conclusions are not sensitive to these factors (Supporting Information, Figure S8, S9).

Our results are cross-validated by another method that compares the CO₂ and temperature trend: we let the CO₂/temperature rates randomly varying within their uncertainty and repeatedly calculate the slope (without permuting the CO₂/temperature rate data points), after 100,000 iterations it turn out the median temperature rate slopes are larger than the CO₂ rate slopes in all

CO₂–temperature comparison cases (Five-core $\delta^{18}\text{O}$ vs composite CO₂; Talos CO₂ vs the Talos $\delta^{18}\text{O}$; new EDC CO₂ vs the EDC δD , Supporting Information, Figure 3e, 3f, S11). Overall, these results suggest the rate of CO₂ rise is less sensitive to varying background climate across the MIS-3 than the Antarctic temperature trend.

We also do not find a significant difference between the rate of CO₂ rise during HS and nHS (t-test with 95% significance level). However, we do find, consistent with previous work (Bereiter et al., 2012), that CO₂ rise ‘overshoots’ (the excess CO₂ rise time relative to corresponding GS duration) are longer during HS than during nHS (median HS overshoot: 307 yr, nHS: -126 yr, Figure 3d). A t-test confirms that the longer overshoots for HS CO₂ rise relative to nHS are significant (at the 95% level). For comparison, we determine the Antarctic warming overshoot in the same way, and no significant difference between HS and nHS is detected (at 95% significance level, Figure S1), consistent with a stable north to south lag (Svensson et al., 2020; WAIS, 2015).

The CO₂ sensitivity to background climate and H events are similar when detrended CO₂ data (removing the long-term signal represented by 20,000 yr smoothing) are used (Supporting Information, Figure S5, S12). To compare with previous research, we also determined the maximums/minimums from a spline fitted CO₂ data with cut off period of 500 yr (Bereiter et al., 2012), again the results hold (Supporting Information, Figure S5, S12).

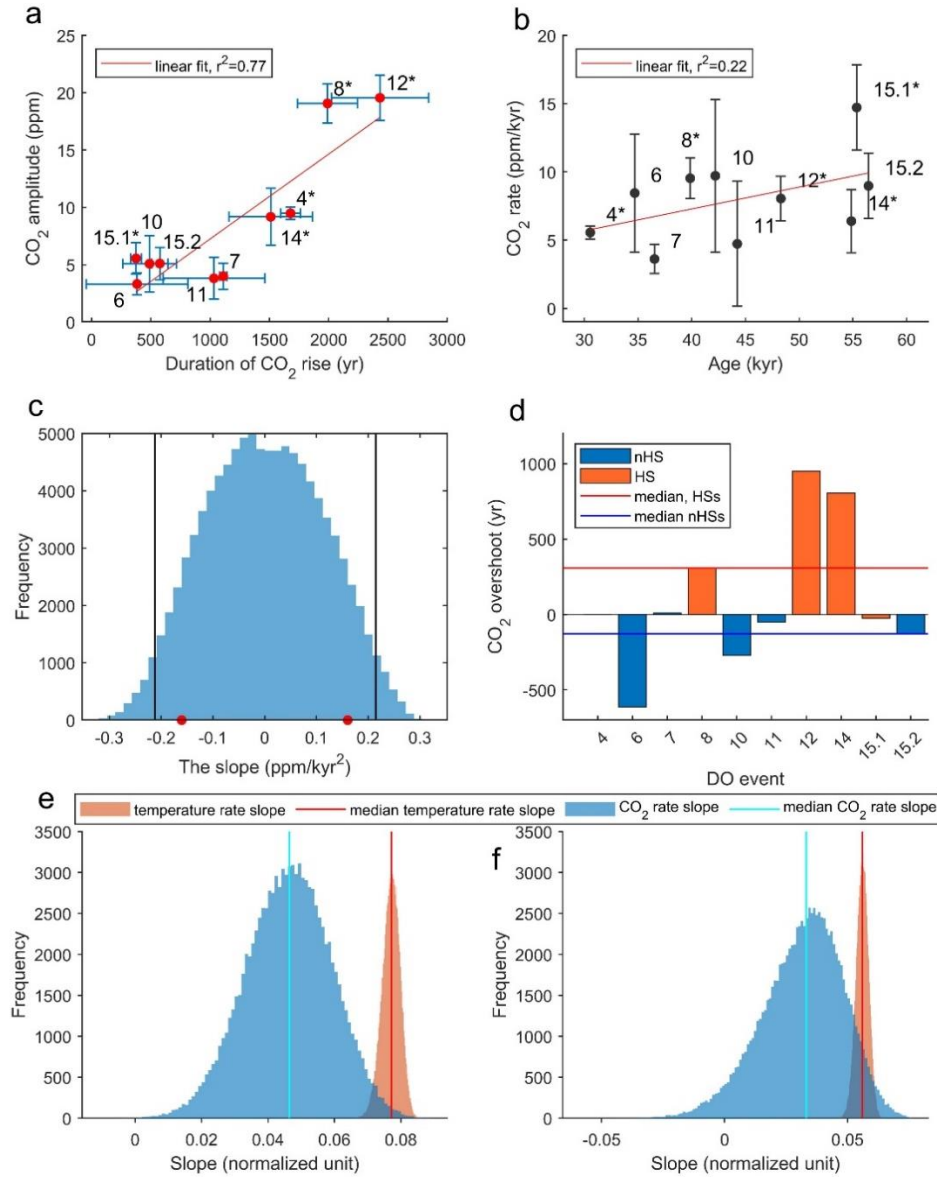


Figure 3. The CO₂ rates and CO₂ overshoot. **a.** The CO₂ amplitude vs the duration of CO₂ rises, the error bar marks the 95% CI of the amplitude and duration. **b.** CO₂ rate vs the age of the CO₂ rise, which is defined by the initiation of the corresponding Greenland stadial. The error bar marks the 95% CI. **c.** The comparison between the ± 1 time observed CO₂ slope in **b** (red dots) with the slopes generated by randomly permuting the CO₂ rate. The 2.5% and 97.5% fractile of the randomly generated slope is marked by vertical black lines. **d.** The CO₂ overshoot. **e.** The

distribution of the randomly generated CO₂/temperature rate slopes of five-core averaged $\delta^{18}\text{O}$ data (Buizert et al., 2018) and the composite CO₂ data (Bereiter et al., 2015). Note the rates are normalized before fitting the slope. **f.** same as **e** but for the new EDC CO₂ (Nehrbass-Ahles et al., 2020) and EDC δD data (EPICA, 2004).

3 Discussion

3.1. Interpreting the temperature sensitivity

Our statistical analysis suggests the AIM temperature response to the DO cycles is gradually weakened throughout the MIS-3. This result provides a firm quantitative basis for previous suggestions of weakened bipolar connections when climate approaches the glacial maximum (Margari et al., 2010; Mcmanus et al., 1999; Wolff et al., 2009).

Recent work indicates that the majority of AIM variability (about 83%) can be explained by a spatially homogeneous oceanic mode (Buizert et al., 2018) that is well captured by the minimum thermodynamic seesaw model (Stocker & Johnsen, 2003). The reduction of the BPS warming rate and $-T_G$ as the climate state cools suggest either a weakening of the mechanisms coupling Antarctic temperature to this oceanic mode and/or a weakening influence of the DO cycles on the Southern Ocean itself. Performing a similar analysis to ours on AIM warming rates observed on surface temperature reconstructions from Southern Ocean marine sediment cores may resolve which is the case.

A recent coupled-model investigation of the bipolar seesaw mechanism (Pedro et al., 2018) suggested that during Greenland stadials, the weakened AMOC drives Antarctic warming

through the following chain of events: Reduced northward advection of heat in the Atlantic Ocean results in heat accumulation in the South Atlantic. This heat then spreads east around the globe along the northern edge of the Antarctica Circumpolar Current (ACC). As a result, the temperature gradient across the ACC increases, driving an increase in the cross-ACC heat flux carried by ocean eddies. Temperature anomalies south of the ACC are amplified by the retreat of sea ice and the resulting ice-albedo feedback. Finally, heat from the Southern Ocean sea-ice zone is transported to Antarctica by atmospheric eddies, i.e., storms.

A cooler background climate state of late MIS-3 (as shown by Antarctic water isotope records (Buizert et al., 2018), or similar trend in ice core noble gas estimates of mean ocean temperature; Bereiter et al., 2018), with thicker Antarctic sea ice, would be expected to reduce the BPS warming rate, because thicker ice would inhibit the ice-albedo feedback (Levermann et al., 2007). Moreover, the concurrent expansion of the sea ice expected in a colder climate would lower the efficiency of atmospheric heat transportation from the warm sea-ice zone to Antarctica, due to the greater distance between Antarctic and the sea-ice-free area. Indeed, marine sediment-core (Collins et al., 2012; Gersonde et al., 2005; Stuut et al., 2004), WDC sea-salt Na record (WAIS, 2015), and climate model (Ferrari et al., 2014), show evidence for Southern Ocean sea-ice expansion at the end of glacial periods.

3.2. Interpreting the CO₂ sensitivity

Our results suggest that in contrast to changes in Antarctic temperature the millennial-scale CO₂ rises show little, if any, sensitivity to varying background climate during MIS-3. Our results also show that this less background state-dependent response of CO₂ is also observed in

earlier glacial periods (Nehrbass-Ahles et al., 2020; Figure S9), which suggests it is a robust feature of millennial climate variability during glacial periods. Recent studies have suggested that increased Southern Ocean deep convection can jointly explain AIM warming and CO₂ trends via the ventilation of heat and CO₂ from ocean depths (Menviel et al., 2018; Skinner et al., 2020). Carbon reservoir age and deep-water temperature reconstructions from the South Atlantic appear to support this ‘Southern Ocean hypothesis’ during HS-4 (Skinner et al., 2020), but do not quantify the scale of its contribution. Furthermore, although strengthened winds are often invoked as the forcing for increased upwelling, Southern Ocean eddies may partially or fully nullify the influence of wind changes on upwelling (Munday et al., 2013). Our results on the different sensitivities of AIM temperature trends and CO₂ to background climate state challenge that a single physical process dominates the millennial signals in both throughout MIS-3. Instead, a dominant influence of Southern Ocean processes on the AIM temperature evolution (Buizert et al., 2018; Pedro et al., 2018) and an important contribution of low- and mid-latitude processes to CO₂, for example by reduction of the biological pump (Nielsen et al., 2019), or CO₂ release from terrestrial sources (Bauska et al., 2016; Marcott et al., 2014; Rhodes et al., 2015), would appear consistent with our results.

Alternately, the Southern Ocean hypothesis could sidestep our observational constraint if the ventilation of ocean heat is reduced relative to the CO₂ as MIS-3 progresses and the climate cools, or if heat ventilated by deep convection has less influence on Antarctic temperature due, for example, to expanded Antarctic sea ice (Collins et al., 2012; Gersonde et al., 2005; Stuut et al., 2004). In either case, our observations provide a new target for model studies seeking to replicate the millennial-scale variability of temperature and CO₂ and their sensitivities to climate state and H events.

The most obvious decoupling of CO₂ and Antarctic temperature occurs during the centennial-scale CO₂ overshoots accompanying H events (Figure 3d). While the CO₂ overshoots are still unexplained, there are a number of candidate mechanisms that would be expected to express a signal in CO₂ and not in Antarctic temperature. Notably, rapid ventilation of accumulated respired carbon from intermediate-depth Atlantic (Chen et al., 2015; Jaccard et al., 2016) and again reduction of the biological pump (Nielsen et al., 2019), or terrestrial CO₂ release (Bauska et al., 2016; Marcott et al., 2014; Rhodes et al., 2015).

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Nehrbass-Ahles, C., Shin, J., Schmitt, J., Bereiter, B., Joos, F., Schilt, A., . . . Grilli, R. (2020). Abrupt CO₂ release to the atmosphere under glacial and early interglacial climate conditions. *Science*, 369(6506), 1000-1005. The EDC δ D data is included in this paper (and its supplementary information files): EPICA. (2004). Eight glacial cycles from an Antarctic ice core. *Nature*, 429(6992), 623-628.

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