

Observations of Electron Vorticity and Phase Space Holes in the Magnetopause Reconnection Separatrix

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Key Points:

- We observe enhanced electron vorticity at an earthward separatrix of magnetopause reconnection.
- The electron vorticity causes perturbations in the magnetic field.
- There is correlation between large amplitude phase space holes and enhanced vorticity in the dayside reconnection separatrix.

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15 **Abstract**

16 We report Magnetospheric MultiScale (MMS) observations of two electron vortices with
 17 a roughly 10 km or 2-3 electron skin depths cross-section across the magnetospheric sep-
 18 aratrix of a dayside magnetopause reconnection. The enhanced electron vorticity, which
 19 is caused by a significant velocity shear in the electron flows streaming in opposite di-
 20 rections along northward-directed magnetic field lines, leads to perturbations perpen-
 21 dicular to the reconnecting magnetic field on the order of 1-2 nT. The high electron vor-
 22 ticity periods are also associated with electron phase space holes having the largest bipo-
 23 lar parallel electric field amplitudes ~ 25 mV/m observed across this exhaust. Less in-
 24 tense ~ 10 mV/m electron phase space holes propagate along a northward magnetic field
 25 away from the X-line within a northward reconnection exhaust. The electron phase space
 26 holes with the largest bipolar parallel electric fields are observed to propagate toward
 27 the X-line within discrete channels of southward and super-Alfvénic 1500 km/s electron
 28 flows in the magnetospheric separatrix layer.

29 **1 Introduction**

30 Magnetic reconnection is a physical process that converts magnetic energy into ki-
 31 netic energy. This process happens throughout the universe where opposing magnetic
 32 field lines exist including at the Earth’s magnetopause. The plasmas on the two sides
 33 of the Earth’s dayside magnetopause current sheet are often characterized by a strong
 34 asymmetry of the magnetic field, plasma number density and temperature. The gradi-
 35 ents of magnetic field and density are known to result in an Earthward displacement of
 36 the magnetosheath flow stagnation point relative to the magnetic reconnection X-line,
 37 where the in-plane reconnecting magnetic field changes direction. (Cassak & Shay, 2007,
 38 2008) This is commonly referred to as asymmetric magnetic reconnection.(Pritchett &
 39 Mozer, 2009)

40 Previous numerical studies have predicted a rotation phenomenon for electrons in
 41 the exhaust region close to the separatrix of asymmetric magnetopause reconnection. Pritchett
 42 and Mozer (2009) used two-dimensional particle-in-cell (PIC) simulations of asymmet-
 43 ric reconnection in the presence of a moderate guide field ($\sim 50\%$) and showed a series
 44 of electron vortices form sunward of the magnetospheric separatrix region of a southward
 45 reconnection jet. Vorticity is the curl of the velocity vector and is a measure of local ro-
 46 tation of the fluid motion. In the simulation, these kinetic-scale electron vortices are pro-

47 posed to be a result of the electron Kelvin-Helmholtz instability with a $\sim 0.3 d_i$ width
 48 ($4 d_e$ with the ion to electron mass ratio of 200) and $\sim 1 d_i$ spacing between separate elec-
 49 tron vorticity structures. Here, $d_i = c/\omega_{pi}$ is the ion inertial length and $\omega_{pi} = \sqrt{n_i e^2 / m_i \epsilon_0}$
 50 is the ion plasma frequency. The simulated electron vortices propagated at a speed of
 51 about $0.3 v_A$ with $v_A = B / \sqrt{\mu_0 n_i m_i}$ being a local ion Alfvén speed and they were as-
 52 sociated with fluctuating dissipation regions and changes in the guide magnetic field mag-
 53 nitude that may impact the reconnection topology and energy dissipation. Fermo et al.
 54 (2012) used PIC codes to suggest that electron Kelvin-Helmholtz vorticity formation near
 55 the separatrix can fold the magnetic field lines and may seed secondary islands.

56 Observations of electron scale vortices was not feasible before the Magnetospheric
 57 MultiScale (MMS) mission due to limited time resolution plasma measurements. (Burch
 58 et al., 2016) Phan et al. (2016) showed enhanced electron vorticity observed within the
 59 reconnection exhaust region and the exhaust boundaries downstream of the X-line us-
 60 ing MMS data. Observed reductions in electron vorticity magnitude toward the inflow
 61 region indicated that the enhancements in vorticity resulted from the reconnection pro-
 62 cesses. In a recent study, Hwang et al. (2019) showed that the electron vorticity can cause
 63 magnetic field fluctuations along the reconnecting magnetic field component near the elec-
 64 tron diffusion region (EDR) in the magnetotail and they proposed that it could be a sig-
 65 nature of spatial proximity to an EDR. They also found that the magnitude of the elec-
 66 tron vorticity increases toward the edge of the electron jet due to a flow shear of the out-
 67 of-plane electron velocity. Consequently, it is possible for the electron vorticity to cause
 68 fluctuations in the local magnetic field changing the topology, and therefore potentially
 69 impacting the local small-scale physics of magnetic reconnection. Ergun et al. (2019) pro-
 70 posed a model of electromagnetic drift waves propagating along the current sheet in the
 71 vicinity of a subsolar magnetopause EDR region. They show how these drift waves can
 72 also be represented as a series of electron vortices superimposed on a primary electron
 73 drift motion along the current sheet in a direction parallel to the X-line. The drift waves
 74 displace or corrugate the current sheet, which may also potentially displace the EDR.

75 Another important feature often associated with magnetic reconnection is the pres-
 76 ence of electrostatic solitary waves (ESWs). Graham et al. (2016) explored Cluster satel-
 77 lite observations of ESWs, consisting of solitary bipolar variations of the parallel elec-
 78 tric field in the vicinity of asymmetric reconnection exhausts at the dayside magnetopause.
 79 Due to the spin-plane limitation of the Cluster electric field instrument, without a ded-

80 icated electric field measurement along the spin axis, that ESW study was limited to very
 81 short time periods when burst mode electric field data were available in a fortuitous con-
 82 figuration relative to the magnetic field. Matsumoto et al. (2003) also reported a fortu-
 83 itous ESW encounter from Geotail satellite observations at the Earthward boundary of
 84 an asymmetric dayside magnetopause exhaust. The Geotail mission, as with the Clus-
 85 ter and THEMIS missions, suffers from the absence of a high-quality axial electric field
 86 measurement, with limited parallel electric field measurements to address the question
 87 of where ESWs occur relative to reconnection exhaust and separatrix regions, and whether
 88 the ESWs correspond to electron phase space holes. However, the MMS mission, with
 89 its complete set of electric field vector observations, is able to resolve small-scale exhaust
 90 structure and address the instabilities that generate ESWs to determine their impact on
 91 asymmetric reconnection. ESWs have the potential to be very important in asymmet-
 92 ric reconnection, since they may be able to dissipate strong currents and couple differ-
 93 ent electron populations, as suggested by Graham et al. (2016). Past observations and
 94 numerical simulations suggest that ESWs develop in the separatrix regions of magnetic
 95 reconnection exhausts. (Drake et al., 2003; Cattell et al., 2003; Fujimoto & Machida, 2006;
 96 Goldman et al., 2008; Graham et al., 2015) It has also been suggested that ESWs form
 97 near the EDR due to bi-streaming electron populations (Goldman et al., 2008; Swisdak
 98 et al., 2018), and they could potentially access the asymmetric exhaust, since the X-line
 99 is separated from the inflow stagnation point. Huang et al. (2014) explored PIC simu-
 100 lations of anti-parallel reconnection to suggest that electron phase space holes can form
 101 within the density cavity of the separatrix regions due to a two-stream electron beam
 102 instability between electrons flowing toward and away from the X-line. Goldman *et al.*
 103 (2014) employed a symmetric PIC simulation for weak guide magnetic field conditions
 104 to suggest that electron phase space holes may form along all four separatrix regions.
 105 Lapenta et al. (2011) performed PIC simulations of symmetric magnetic reconnection
 106 with different guide field strengths and they showed that a strong electron flow devel-
 107 ops along the separatrices in all cases, which is sufficient to lead to the onset of stream-
 108 ing instabilities and to form bipolar parallel electric field signatures. They also showed
 109 that the predominant strength of the inward flow along one pair of separatrices tends
 110 to increase the prevalence of bipolar signatures along those flow channels without sup-
 111 pressing them altogether along the other two separatrices. Asymmetric reconnection sim-
 112 ulations reported by Cazzola et al. (2015) (e.g., Fig. 14) and by Chang et al. (2021) (e.g.,

113 Fig. 8) also confirm the formation of bipolar fields (electron phase space holes) along the
 114 magnetosphere-side separatrix.

115 In this paper, we investigate a possible relation between electron vorticity and elec-
 116 tron phase space holes in a case of asymmetric magnetopause reconnection as recorded
 117 by the four MMS satellites in a high-quality tetrahedron formation near the subsolar re-
 118 gion. In this set of observations, a large magnitude of the obtained electron vorticity co-
 119 incides with the largest phase space hole magnitudes present along the earthward-side,
 120 magnetospheric separatrix of a northward exhaust. In section 2, we present an overview
 121 of the subsolar exhaust as observed by MMS to the north of a dayside magnetopause X-
 122 line and the signatures which are consistent with a magnetospheric separatrix earthward
 123 of a northward reconnection exhaust. In section 3, we explore the size and the speed of
 124 the electron vorticity structures using a timing analysis method of the four satellite ob-
 125 servations and we compare and contrast the observations with reported PIC simulation
 126 predictions. In section 4, we show that large amplitude phase space holes are present in
 127 the regions of vorticity enhancements and discuss their size and speed. Section 5 includes
 128 a summary and conclusion of this study into kinetic-scale separatrix structures.

129 2 MMS Event Overview

130 The four MMS satellites have repeatedly crossed the magnetopause boundary to
 131 capture the dynamics of dayside asymmetric magnetic reconnection.(Burch et al., 2015)
 132 The MMS satellites carry identical instruments capable of gathering high time resolu-
 133 tion particle and fields data while operating in burst mode.(Baker et al., 2016; Ergun
 134 et al., 2016; Kawa & Lin, 2016; Contel et al., 2016; Lindqvist et al., 2016; Pollock et al.,
 135 2016; Russell et al., 2016; Torbert et al., 2016)

136 The MMS satellites traversed the magnetopause from the magnetosheath into the
 137 magnetosphere at $\sim 10:09:45$ UT on 29 January 2017. The four satellites were in a tetra-
 138 hedron formation at this time with a separation distance of ~ 7 km at $[10.2, -2.9, 1.8] R_E$
 139 in GSM coordinates. Figure 1 shows an overview of the crossing as recorded by the MMS1
 140 satellite with all instruments operating in a burst mode. Panels a and b show ion and
 141 electron energy fluxes measured by the fast plasma investigation (FPI) instrument in 150
 142 ms cadence for ions and 30 ms cadence for electrons. The third and fourth panels (c and
 143 d), respectively, show the three components of the magnetic field and the electric field

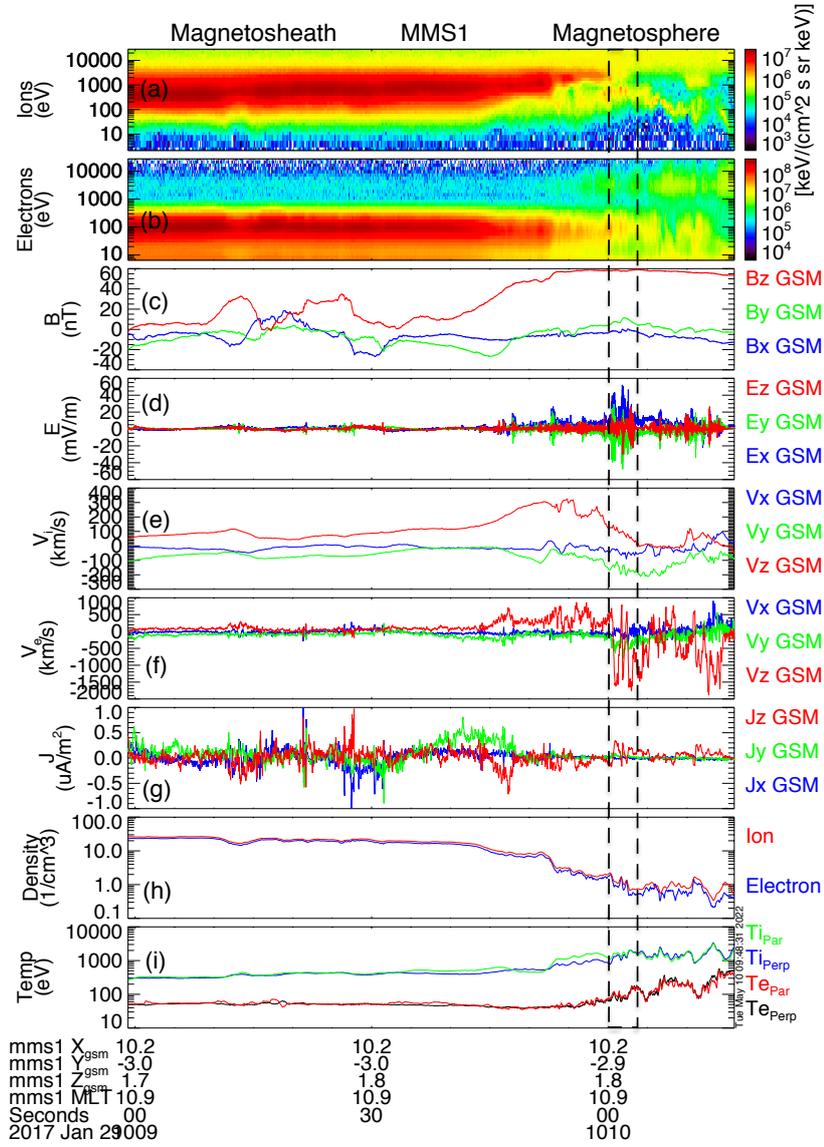


Figure 1. Overview of the magnetopause crossing using MMS1 data. (a) Ion energy flux (b) Electron energy flux (c) Magnetic field vector (d) Electric field (e) Ion velocity (f) Electron velocity (g) Current density from FPI instrument (h) Ion (red line) and electron density (blue line) (i) Ion and electron temperatures perpendicular and parallel to the background magnetic field.

144 in the GSM coordinate system. The magnetic field data are measured by the Fluxgate
 145 Magnetometer (FGM) instrument with 128 samples/second and electric field data are
 146 provided by the Electric field Double Probes (EDP) instruments with 8192 samples/second.
 147 Panels e and f show the three components of the ion and electron velocities in the GSM
 148 coordinate system as measured by the FPI instruments. The current density calculated
 149 using FPI instrument is shown in panel g. The ion and electron densities, and the tem-
 150 peratures perpendicular and parallel to the magnetic field, are displayed in panels h and
 151 i. The plasma densities decreased and the temperatures increased as the satellites moved
 152 from the magnetosheath into the magnetosphere as expected across the magnetopause.

153 At about 10:09:44 UT, the V_{iZ} component of the ion flow begins to increase toward
 154 a maximum value close to 300 km/s in a northward Z direction, at nearly the same time
 155 as the Z component of the magnetic field increases toward a maximum of about 60 nT.
 156 This $V_{iZ} \sim 300$ km/s jet speed roughly corresponds to half of the ion Alfvén speed ($v_A \sim$
 157 585 km/s), which is calculated from the average of B_Z and the average ion density in
 158 the interval of an increasing ion flow. This time period reflects an MMS crossing of a north-
 159 ward magnetic reconnection jet on the earthward side of the dayside magnetopause in
 160 agreement with a positive V_{iZ} and a positive B_Z .

161 A measured ~ 600 km/s electron flow, which is almost equal the estimated Alfvén
 162 speed, is also recorded in the same northward direction as the ion exhaust. As the ion
 163 flow starts to decrease, we observe several enhancements of the sunward X component
 164 of the electric field, with some periods showing E_X above ~ 40 mV/m. This positive elec-
 165 tric field, which is aligned with the magnetopause normal direction, is consistent with
 166 the expectations of a separatrix layer on the earthward side of an actively reconnecting
 167 magnetopause current sheet. (Swisdak et al., 2018) This approximately normal electric
 168 field acts to decelerate the ion inflow from the magnetosheath and it ensures an over-
 169 all charge neutrality with the much less dense population of magnetospheric electrons.
 170 A comparison of the measured components of the electric field with the components of
 171 the $V_e \times B$ electric field (not shown) suggests that the electrons are frozen-in to the mag-
 172 netic field. In this same region of enhanced sunward E_X , we also observe significant re-
 173 versals in V_{eZ} from northward ~ 600 km/s flows in the reconnection exhaust to south-
 174 ward flows toward the X line with amplitudes as large as 2000 km/s. These high veloc-
 175 ity electron flows are associated with an electron population with energies below 1000
 176 eV.

177 MMS also observed large amplitude parallel electric field fluctuations in this region
 178 that consists of several V_{eZ} reversals. The presence of multiple channels of fast ~ 1500
 179 km/s southward flows suggests that the MMS1 satellite may have traversed a system of
 180 filamentary currents at the exhaust boundary in the magnetospheric separatrix region.

181 **3 Observation of Electron Vorticity**

182 We calculate electron vorticity or $\nabla \times V_e$ from the measurements of electron ve-
 183 locity and position information at the four MMS satellite. Figure 2 shows a subset of
 184 electron and magnetic field measurements from all four MMS satellites in the 4 s region
 185 of interest from 10:10:00 UT to 10:10:04 UT when the large-amplitude electric field struc-
 186 tures and high-speed electron flow reversals are observed near the Earthward boundary
 187 of the reconnection exhaust. A box outlined by the dashed lines in Figure 1 marks this
 188 4 s region of interest. All vector components are shown in a boundary normal LMN co-
 189 ordinate system that we obtained as follows. We first find a large-scale boundary nor-
 190 mal coordinate system of the magnetopause current sheet in GSM coordinates with $L_0 = [-$
 191 $0.047, 0.318, 0.947]$, $M_0 = [-0.241, -0.923, 0.298]$ and $N_0 = [0.969, -0.214, 0.120]$. Here,
 192 N_0 is the direction of the cross-product normal for the average external magnetic fields
 193 at 10:08:52.700 UT and 10:10:05.300 UT, M_0 is the normalized cross-product of N_0 with
 194 the direction of maximum variance of the magnetic field for this same time period, and
 195 $L_0 = M_0 \times N_0$. However, the normal component of the magnetic field of this LMN
 196 system is found to point in a positive direction across the northward exhaust region. This
 197 is inconsistent with a reconnection X line to the south of the MMS satellites, which is
 198 expected to connect a northward magnetospheric field with an Earthward directed nor-
 199 mal component across the magnetopause layer to the adjacent magnetic field of the mag-
 200 netosheath. The large-scale LMN system, therefore, needs a slight correction that we
 201 achieve on the basis of a multi-satellite timing analysis of the V_{eZ} GSM component for
 202 a sharp southward flow reversal that MMS1 observed to be centered at 10:10:01.800 UT
 203 in the magnetospheric separatrix layer. Taking MMS1 as the reference satellite, and shift-
 204 ing the 30 ms cadence V_{eZ} observations forward in time by 41 ms for MMS2, 74 ms for
 205 MMS3 and 29 ms for MMS4, we obtain a timing analysis boundary normal vector of the
 206 V_{eZ} signal between the four MMS satellites. This unit vector is used to define a revised
 207 $L = [-0.040, -0.227, 0.973]$ due to its alignment with a Z_{GSM} direction. We find a cor-
 208 rected normal direction $N = [0.966, -0.259, -0.020]$ as the normalized cross-product of

209 this L with M_0 , while $M = [-0.256, -0.939, -0.230]$, completes the orthogonal system as
 210 the cross-product of N and L . This LMN system is now optimized for a local analy-
 211 sis of the observed electron flow reversals of the magnetospheric separatrix from a small
 212 8.5 degrees rotation of the normal vector. This correction of the direction of the normal
 213 vector also results in a negative B_N component of the magnetic field as expected across
 214 a northward reconnection exhaust region near the subsolar magnetopause.

215 Figure 2a displays a subset of measurements at all four MMS satellites with pan-
 216 els 1 and 2 showing the L and M components of the electron velocity, panels 3 and 4 show-
 217 ing the M and N components of the magnetic field, and panel 5 showing the electron
 218 density. Panel 6 presents the derived electron vorticity in LMN coordinates or $\nabla \times V_e$,
 219 which is obtained from the measurements of electron velocity and position information
 220 at the four MMS satellites. The vorticity corresponds to a rotation or a shear in the ve-
 221 locity field. Figure 2b displays the MMS satellite separations with a high tetrahedron
 222 quality factor (TQF=0.909) of the formation at 10:10:02 UT when MMS recorded the
 223 high-speed electron flow reversals and the corresponding electron vorticity displayed in
 224 Figure 2a. Figure 2c depicts a schematic of the asymmetric reconnection exhaust region
 225 with a red line showing the inferred, and highly localized, MMS path across the mag-
 226 netospheric separatrix on the basis of a multi-satellite timing analysis of a sharp south-
 227 ward V_{eZ} flow reversal at 2017-01-29/10:10:01.800 UT.

228 The electron vorticity vector is clearly enhanced in the region of electron flow re-
 229 versals at 10:10:01.5-10:10:03 UT. It is predominantly aligned with a positive M-direction
 230 and to a lesser extent with a negative N-direction as indicated between the two red dashed,
 231 vertical lines. In other words, it is perpendicular to the background mostly northward
 232 ($B_L > 0$) magnetic field. The large vorticity region coincides with changes in the B_M
 233 (guide field) and the B_N components of the magnetic field as also illustrated in panels
 234 3 and 4. The B_M component of the magnetic field (the reconnection guide field) is about
 235 -18 nT which is about 30% of the B_L or reconnecting magnetic field. The high shear in
 236 the L direction of the electron velocity drives the electron vorticity in the M and N di-
 237 rections.

238 In order to appreciate the connection between the electron flow reversals and the
 239 deduced electron vorticity, we display the V_e vector field components of the NL -plane
 240 as measured along the MMS1 and MMS3 satellite trajectories in Figure 3. Here, the tem-

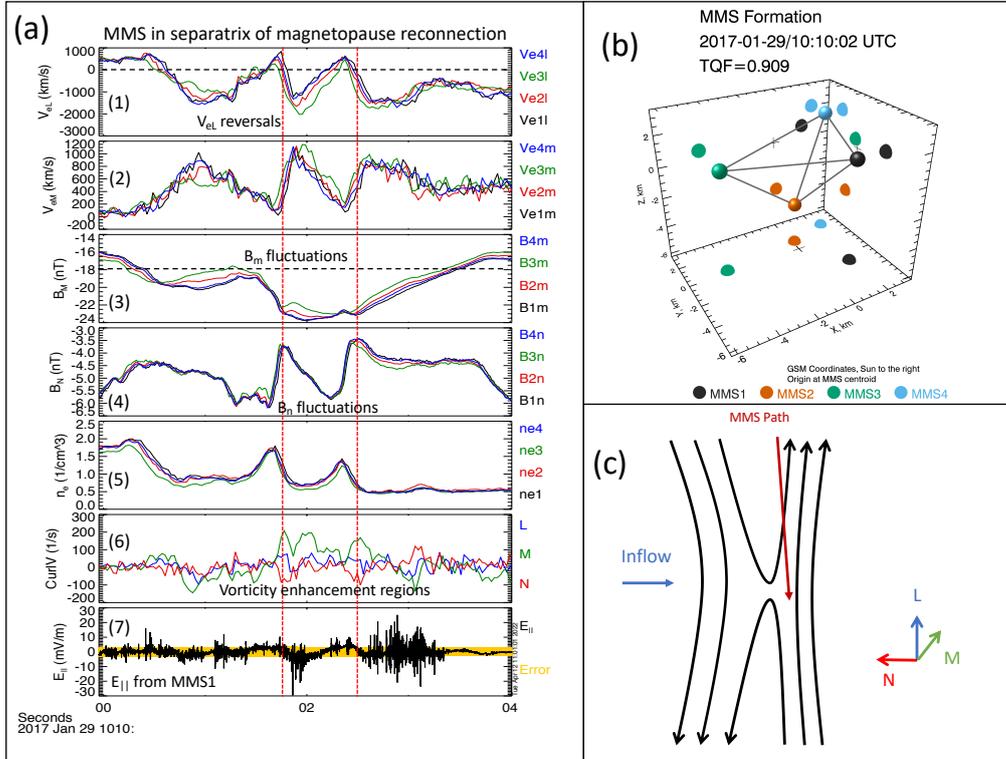


Figure 2. (a) Region of interest at 10:10:00-10:10:04 UT with a large electron vorticity associated with the magnetospheric separatrix region of dayside magnetopause reconnection. (1) L component of electron velocity (V_{eL}), (2) V_{eM} component, (3) B_M component of magnetic field, (4) B_N component of the magnetic field, (5) electron density for all four satellites, (6) multi-satellite electron vorticity calculation in LMN coordinates, and (7) parallel electric field from MMS1, (b) MMS formation at the 10:10:02 UT time of crossing, and (c) the sketch of MMS path across the local section of the magnetospheric separatrix layer. The two red vertical dashed lines at 10:10:01.680 UT and 10:10:02.490 UT mark the V_{eL} flow reversal of two electron vorticity structures.

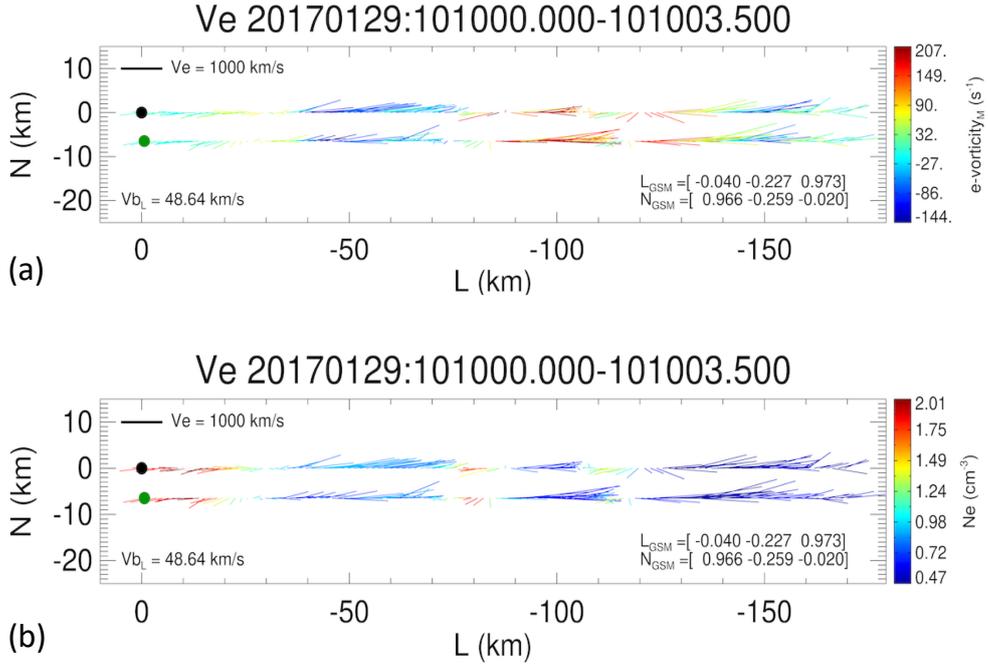


Figure 3. Vector analysis of electron velocity for MMS1 and MMS3. (a) The electron velocity vectors are displayed in the NL -plane with the M component of electron vorticity shown in color, (b) electron velocity shown in the NL -plane with color showing the measured electron density. We identify the center of two V_e rotations corresponding to the two strong electron vorticity events with the first rotation (vorticity) centered at $L \sim -85$ km for MMS3 and $L \sim -90$ km for MMS1, and the second rotation centered at $L \sim -115$ km for MMS3 and $L \sim -120$ km for MMS1.

241 poral information as recorded by the two satellites for a 3.5 s period from 10:10:00 UT
 242 are converted to spatial L -distances using the northward-pointing L -vector of the tim-
 243 ing method and the corresponding boundary speed $v_L \approx 48.6$ km/s by which the elec-
 244 tron velocity field and the associated electron vorticity structures are observed to prop-
 245 agate in a northward L -direction away from the X line. The northward electron vortic-
 246 ity propagation is reflected, therefore, in a southward or $-L$ MMS trajectory direction
 247 of this figure. The origin at $(N, L) = (0, 0)$ is taken as the location of MMS1 (black dot)
 248 at 10:10:00 UT with the location of MMS3 (green dot) shown with its NL -plane separa-
 249 tion, $\Delta L = -0.64$ km and $\Delta N = -6.47$ km, from MMS1 at this same time. The MMS2
 250 and MMS4 satellite trajectories are excluded here due to their close N -separation from
 251 MMS1, with MMS2 at $\Delta N = -2.43$ km and MMS4 at $\Delta N = -0.62$ km, such that their V_e
 252 vectors would directly overlap with the vector field at MMS1. The figure shows how the
 253 electron velocity V_e vectors are changing direction and amplitude as the two MMS satel-
 254 lites fly through this separatrix region. In panel a, the color represents the out-of-plane
 255 M -component of the electron vorticity and in panel b, color displays the observed elec-
 256 tron density along the two MMS trajectories.

257 Based on the V_e vector plots of Figure 3, we identify the center of two counter-clockwise
 258 V_e rotations coincident with the two strongest electron vorticity events. The center of
 259 one vorticity enhancement is present at $L \sim -85$ km for MMS3 and $L \sim -90$ km for
 260 MMS1, and a second vorticity structure is centered near $L \sim -115$ km at MMS3 and
 261 $L \sim -120$ km at MMS1. The electron velocity field of Figure 3 illustrates how the orig-
 262 inal V_e measurements are reproduced in the deduced electron vorticity. A comparison
 263 of the two panels of Figure 3 also suggests a presence of two density enhancement fea-
 264 tures (see panel b) just ahead of the northward-propagating vorticity structures at $|L| =$
 265 80 km and $|L| = 110$ km. These density enhancements indicates that the MMS satel-
 266 lites are entering the exhaust region as the magnetopause boundary oscillates.

267 The two electron vorticity structures are kinetic-scale with Figure 3 suggesting an
 268 approximate 10-20 km cross-section. This dimension corresponds to $2-4 d_e$ where $d_e \sim$
 269 5 km is an electron inertial length for the average 1.1 cm^{-3} plasma density recorded at
 270 this time (see Figure 2a, panel 5). Given that $d_i = (m_p/m_e)^{1/2} d_e$, the electron vortices
 271 are only about $0.05 - 0.09 d_i$ in terms of the local ion inertial length. They are mov-
 272 ing away from the X-line and along the northward-directed reconnecting magnetic field
 273 lines with a 48.6 km/s propagation speed or $v_L \sim 0.1 v_A$. These particular electron vor-

274 ticity structures that MMS observed at the magnetospheric separatrix layer are some-
 275 what smaller in size and their propagation speed is somewhat slower as compared with
 276 the typical $\sim 0.3d_i$ size electron vortices that Pritchett and Mozer (2009) simulated to
 277 move at about $\sim 0.3v_A$ in a 2D PIC code.

278 The counter-clockwise rotation of the electron velocity present in the NL -plane of
 279 Figure 3 corresponds to a clockwise loop of current that induces a magnetic field deflec-
 280 tion in a negative M direction. This localized electron vorticity is the likely origin of the
 281 highly localized and rather weak ~ 2 nT magnitude increases of the measured guide mag-
 282 netic field, which are shown to be centered at the times of each of the two red, dashed
 283 vertical lines of Figure 2a (see panel 3) that mark the locations of the two V_{eL} flow re-
 284 versals of each electron vorticity structure. The two electron vorticity structures are not
 285 perfectly contained in the NL -plane with a single M -component of the vorticity vector
 286 as indicated by the presence of an additional N -component of the vorticity vector in panel
 287 6 of Figure 2a. This finite N -component of the electron vorticity vector is reflected in
 288 the enhanced $V_{eM} > 0$ velocity changes, which are centered at the same times as the
 289 two red, dashed vertical lines (see Figure 2a, panel 2) of the V_{eL} reversals. These $V_{eM} >$
 290 0 velocity changes induce a magnetic field fluctuation in the positive N -direction as re-
 291 flected by the observed ~ 2 nT magnitude decreases of the B_N component of the mag-
 292 netic field (see Figure 2a, panel 4), which are co-located with the V_{eL} reversals of the elec-
 293 tron vorticity. The guide magnetic field B_M also displays a large-scale ~ 6 nT change in
 294 a negative M -direction over a 1.5 s duration period from $\sim 10:10:01.5$ UT to $\sim 10:10:03$
 295 UT. This change of the guide-field, albeit coincident with the region of enhanced elec-
 296 tron vorticity, is more likely supported by a larger-scale current system along the L -direction
 297 that may consist of the $V_{eL} > 0$ before 10:10:00.5 UT and the sustained $V_{eL} < 0$ af-
 298 ter 10:10:03 UT.

299 **4 Electron Phase Space Holes**

300 The observations of large electron vorticity across the magnetospheric separatrix
 301 region coincide with a presence of large amplitude bipolar parallel electric field struc-
 302 tures, which are also known ESWs. (Graham et al., 2016; Matsumoto et al., 2003) These
 303 bipolar structures are associated with amplitudes as high as 25 mV/m in this dayside
 304 event and they signal the presence of electron phase space holes. An electron phase space
 305 hole is equivalent to a very localized region of surplus positive charge and a diverging

306 electric field. An electron phase space hole moving along the background magnetic field
 307 will therefore generate a short-duration bipolar signal of the parallel electric field (E_{\parallel}).
 308 Figure 4 displays the instances of a large number of highly localized bipolar E_{\parallel} ESW struc-
 309 tures as primarily recorded by the axial double-probes of the EDP instrument at 8192
 310 Hz on all four MMS satellites across the northward reconnection exhaust region and the
 311 adjacent vorticity region of the magnetospheric separatrix. In panels d to g, the instances
 312 of bipolar ESWs are color-coded at each satellite with a red plus symbol corresponding
 313 to a positive-then-negative bipolar E_{\parallel} and each black plus symbol likewise indicating a
 314 negative-then-positive bipolar E_{\parallel} . The red ESWs likely represent fast electron phase space
 315 holes moving *parallel* with the magnetic field, while the black ESWs represent electron
 316 phase space holes moving in the *anti-parallel* magnetic field direction. The magnetic field
 317 does not change direction inside the vorticity region of interest, where it is predominantly
 318 directed in a northward L direction as shown in panel a. This means that the assumed
 319 electron phase space holes are propagating in opposite directions of the magnetospheric
 320 separatrix region. In comparing the propagation direction of electron phase space holes
 321 and the direction of the electron velocity, it is clear that the electron holes are stream-
 322 ing in the same direction as the electron flows, whether along the northward $V_{eL} \sim 600$
 323 km/s exhaust and away from the X-line before MMS encountered the separatrix region,
 324 or in a southward direction of the fast $V_{eL} \sim 1,500$ km/s flow channels that we asso-
 325 ciate with regions of enhanced electron vorticity structures.

326 Established techniques for directly determining the velocity of the electron phase
 327 space holes in this time interval cannot be used. The holes are too slow to estimate their
 328 velocity from magnetic field perturbations (Andersson et al., 2009) and they are too fast
 329 to estimate a propagation velocity from a timing difference between the potential mea-
 330 sured on two opposite electric field probes. The later limitation is particularly stringent
 331 because \mathbf{B} is primarily directed along Z_{GSM} and the associated ~ 1 ms duration of the
 332 ESWs is too short to resolve a time delay of the corresponding potential signature for
 333 the relatively short 29.2 m separation between the two opposite axial probes.

334 We can nevertheless make a rough estimate of the electron phase space hole veloc-
 335 ity as follows: The temporal separation of the largest (negative-to-positive) bipolar E_{\parallel}
 336 field peaks during this period are ~ 0.25 ms (e.g., Fig. 5b). Given a Debye length of $\lambda_e \sim$
 337 100m based on the density of ~ 0.5 cm $^{-3}$ and a parallel electron temperature of ~ 100 eV
 338 representative of the times when the largest bipolar fields are observed, an electron phase

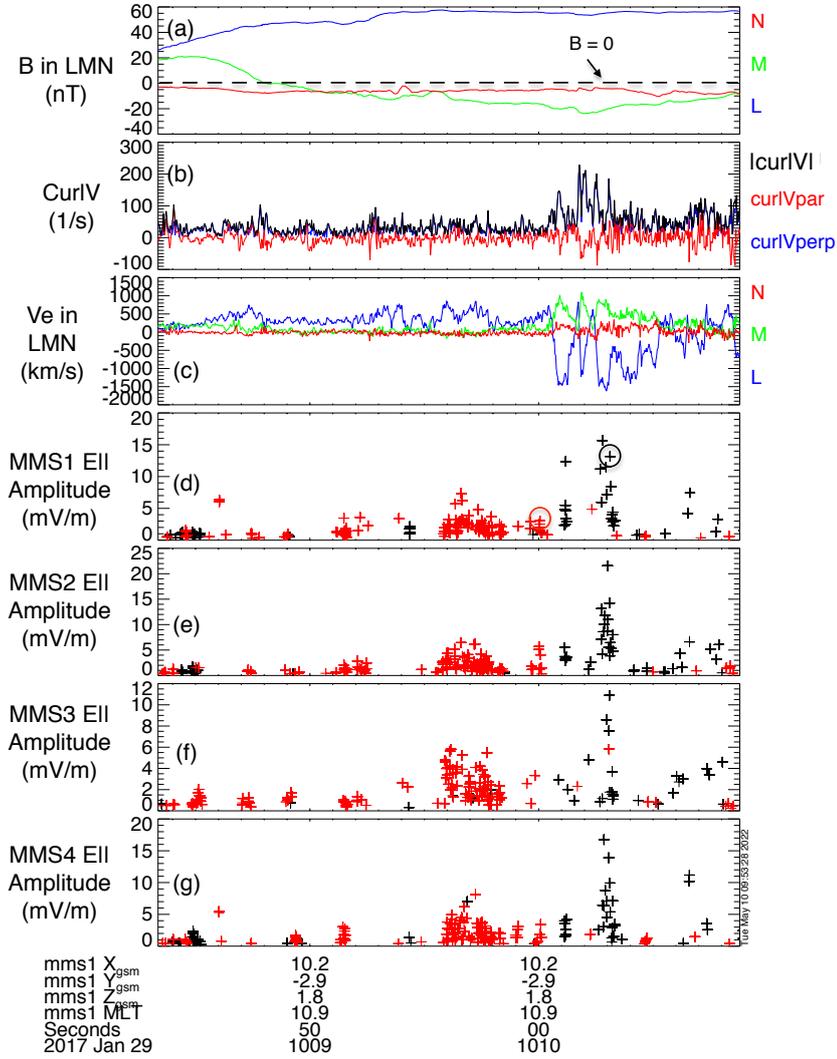


Figure 4. : Instances of observed phase space holes on all four MMS satellites. (a) Magnetic field, (b) Electron vorticity shown perpendicular (blue line), parallel (red line) to magnetic field and total (black line), (c) Electron velocity, (d) Phase space holes on MMS1, (e) Phase space holes on MMS2, (f) Phase space holes on MMS3 and (g) Phase space holes on MMS4. On panel d, marked phase space holes in circles are shown in figure 5.

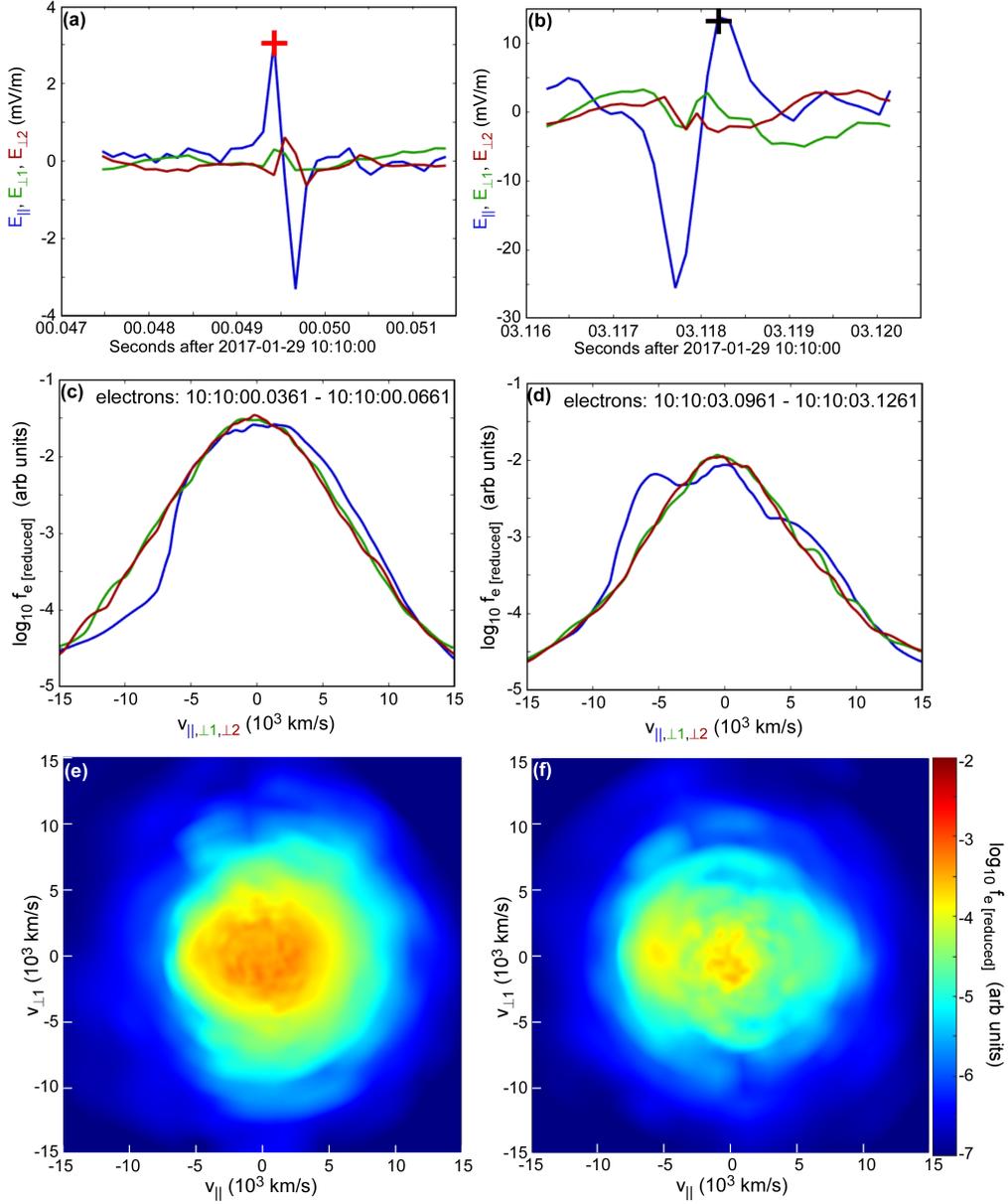


Figure 5. Panels (a) and (b) Electric field components in field-aligned coordinates for two bipolar E_{\parallel} events from MMS1 of opposite polarity (as indicated in panel d of Fig. 4). Each interval plotted covers ± 2 ms on either side of the detection time, which is marked by the respective red and black '+' signs. Panels (c) and (d) Reduced electron distributions along the three orthogonal field-aligned directions for the 30 ms FPI measurement interval containing the respective bipolar EH events in panels (a) and (b). (e) and (f) electron distributions in $v_{\parallel} - v_{\perp}$ plane at the interval of the EH events.

339 space hole with a typical spatial width of $\sim 2\text{--}10\lambda_e$ would have a velocity in the $-\mathbf{B}$ di-
 340 rection of $\sim 800\text{--}4000$ km/s. (Andersson et al., 2009; Hutchinson, 2017; Holmes et al., 2018)
 341 Such a velocity is consistent with the valley in the reduced parallel distribution in Fig. 5d.
 342 The weaker positive-to-negative bipolar E_{\parallel} in Fig. 5a is associated with a broad plateau
 343 in the simultaneous reduced distribution (Fig. 5c) in the $+v_{\parallel}$ direction, with no distinct
 344 valley. This is typical of times when weak bipolar fields of either polarity are observed
 345 during the time interval in figure 4. Note that the 30ms interval covered by each FPI
 346 distribution can contain multiple electron phase space holes—each of which has a du-
 347 ration of 1ms—and so the interpretation is complicated by temporal averaging. Figures
 348 5e and 5f show 2D electron distributions in the $v_{\parallel}\text{--}v_{\perp}$ plane for the interval of the elec-
 349 tron hole events. The left distribution (Fig. 5e) is characterized by a broad plateau sur-
 350 round by a hotter population. This distribution is shifted toward positive v_{\parallel} . However,
 351 there are no features that can be easily correlated with the observed bipolar field. The
 352 right distribution (Fig. 5f), on the other hand, appears to consist of a beam moving in
 353 the $-v_{\parallel}$ direction that is clearly separated from the core population. The beam-like com-
 354 ponent may be associated with a beam-plasma instability that saturates by producing
 355 electron holes. However, a beam-like feature can also result from passing electrons be-
 356 ing accelerated by the positive potential of an electron hole as it crosses the region. Since
 357 electron holes pass the detector over times short compared to the 30 ms interval mea-
 358 sured by FPI, the observed distribution is likely a mixture of an unmodified component
 359 and a component temporarily modified by the presence of electron holes.

360 5 Summary and Conclusions

361 In this study, we presented observations of electron vorticity across the magneto-
 362 spheric separatrix region of dayside magnetopause reconnection as recorded by the MMS
 363 satellites and their relation with electron phase space holes. The enhanced electron vor-
 364 ticity is directly associated with a region of high-speed electron flow reversals between
 365 ~ 600 km/s northward exhaust flows and several bursts of $\sim 1,500$ km/s southward elec-
 366 tron flows as identified in the data. The electron vorticity, which is dominant in the di-
 367 rections perpendicular to the reconnecting (L component) magnetic field, and the as-
 368 sociated current loops induces a ~ 2 nT magnetic field perturbation of both the guide
 369 magnetic field along the out-of-plane $-M$ -direction and the normal component of the
 370 magnetic field. The positive M -component of the electron vorticity vector, therefore, en-

371 hances the local magnitude of a negative guide field and the negative N -component of
 372 the vorticity vector reduces the magnitude of a negative normal magnetic field. Using
 373 a multi-satellite timing analysis of the V_{eZ} component of the electron velocity, we find
 374 that the vorticity structures are moving away from the X line along the northward L di-
 375 rection of the reconnecting magnetic field at $v_L \approx 49$ km/s or $v_L \sim 0.1v_A$. The esti-
 376 mated ~ 10 - 20 km size of two vorticity structures corresponds to a kinetic-scale dimen-
 377 sion of $2 - 4d_e$ or $0.05 - 0.09d_i$. These electron vorticity structures are roughly three
 378 times as small in size and propagation speed as compared with similar predictions of elec-
 379 tron vorticity structures from a 2D PIC numerical simulation by Pritchett and Mozer
 380 (2009) for a similar strength of a guide field. In comparison, Ergun et al. (2019) mod-
 381 eled a series of electron vorticity structures propagating along the current sheet in the
 382 vicinity of a subsolar magnetopause EDR region in a direction parallel to the X line or
 383 along the M direction, while our observations of the electron vorticity structures in the
 384 separatrix of magnetopause reconnection indicates they are propagating in the L direc-
 385 tion or along the reconnecting magnetic field.

386 The MMS satellites recorded large-amplitude electron phase space holes in the re-
 387 gion of electron vorticity structures at the exhaust boundary with the magnetospheric-
 388 side separatrix. The velocity shear of two counter-streaming electron beams that gen-
 389 erated this kinetic-scale electron vorticity likely also supported the formation of these
 390 electron phase space holes. MMS observed a northward propagation of moderate strength
 391 ~ 10 mV/m electron phase space holes along the ~ 600 km/s northward exhaust flow re-
 392 gion away from the X line. More intense ~ 25 mV/m phase space holes were subsequently
 393 observed to propagate toward the X-line and mostly within channels of fast $\sim 1,500$ km/s
 394 southward electron flows associated with the magnetospheric separatrix layer. Using re-
 395 duced parallel distributions of electrons, we showed that a typical electron phase space
 396 hole with a spatial width of ~ 2 - $10\lambda_e$ in a general agreement with the observed elec-
 397 tron density and electron temperature would have a velocity in the $-\mathbf{B}$ direction of ~ 800 -
 398 4000 km/s. Based on the PIC simulations performed by Cazzola et al. (2015) and Chang
 399 et al. (2021), phase space holes are dominant on magnetospheric side of reconnection sep-
 400 aratrix in asymmetric magnetic reconnection in agreement with these MMS observations.

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