

# A near-global climatology of oceanic coherent eddies

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## Key Points:

- Coherent eddies contain around 50% of the total surface ocean kinetic energy budget.
- Seasonal cycle of the number of coherent eddies and the coherent eddy amplitude reveals a 3-6 month lag to wind forcing.
- The seasonal lag between the number and the amplitude of coherent eddies suggests a role for the inverse cascade.

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**Abstract**

Ocean eddies influence regional and global climate through mixing and transport of heat and properties. One of the most recognizable and ubiquitous feature of oceanic eddies are coherent vortices with spatial scales of tens to hundreds of kilometers, frequently referred as “mesoscale eddies”. Coherent mesoscale eddies are known to transport properties across the ocean and to locally affect near-surface wind, cloud properties, and rainfall patterns. Although coherent eddies are ubiquitous, their climatology, seasonality, and long-term temporal evolution remains poorly understood. Here, we examine the kinetic energy contained by coherent eddies and present the seasonal, interannual and long-term variability using satellite observations between 1993 to 2019. A total of  $\sim 37$  million coherent eddies are detected in this analysis. Around 50% of the kinetic energy contained by ocean eddies corresponds to coherent eddies. Additionally, a strong seasonal cycle is observed, with a 3–6 months lag between the wind forcing and the response of the coherent eddy field. The seasonality of the number of coherent eddies and their amplitude reveals that the number of coherent eddies responds faster to the forcing ( $\sim 3$  months), than the coherent eddy amplitude (which lags by  $\sim 6$  months). This seasonal cycle is spatially variable, so we also analyze their climatology in key oceanic regions. Our analysis highlights the relative importance of the coherent eddy field in the ocean kinetic energy budget, implies a strong response of the eddy number and eddy amplitude to forcing at different time-scales, and showcases the seasonality, and multidecadal trends of coherent eddy properties.

**Plain language summary**

Coherent eddies are the most common feature of ocean variability observable from satellites. They are crucial in ocean dynamics as they can transport properties over long distances and interact with the atmosphere. Our study investigates the seasonal, interannual, and long-term changes in the abundance and intensity of coherent eddies, by automatically identifying individual eddies over the available satellite altimeter record. The seasonal cycle suggests a transition from numerous, smaller, and weaker coherent eddies, to fewer and larger, and stronger coherent eddies over the season. In addition, a long-term adjustment of the coherent eddy field is identified with possible links to long-term changes in the climate system.

## 1 Introduction

Mesoscale ocean variability with spatial scales of tens to hundreds of kilometers is comprised of processes such as vortices, waves, and jets (Ferrari & Wunsch, 2009; Fu et al., 2010). These mesoscale processes are highly energetic, and they play a crucial role in the transport of heat, salt, momentum, and other tracers through the ocean (Wunsch & Ferrari, 2004; Wyrтки et al., 1976; Gill et al., 1974). One of the most recognizable and abundant ocean processes observable from space are mesoscale vortices. Although mesoscale vortices are commonly referred to in the literature as “mesoscale eddies”, this term is also often used to describe the total mesoscale ocean variability (the time-varying component of the mesoscale flow), thus, to avoid ambiguity we will refer to mesoscale vortices as *coherent eddies*. Coherent eddies are abundant and energetic; they are essential to ocean dynamics as concluded by many previous studies (Hogg & Blundell, 2006; Siegel et al., 2011; Beron-Vera et al., 2013; Frenger et al., 2013, 2015; Pilo et al., 2015; Schubert et al., 2019; Patel et al., 2020).

Coherent eddies are quasi-circular geostrophic currents. According to their rotational direction and the sign of the Coriolis parameter, the sea surface height anomaly within a coherent eddy can have a negative or positive sea surface height anomaly (cold-core and warm-core coherent eddies, respectively). This characteristic sea surface height signature of coherent eddies has been utilized to identify and track coherent eddies from satellite altimetry (e.g., Chelton et al., 2007; Faghmous et al., 2015; Ashkezari et al., 2016; Martínez-Moreno et al., 2019; Cui et al., 2020). Automated identification algorithms of coherent eddies have revealed their ubiquity in the oceans, with a predominant influence at hotspots of eddy activity such as in boundary current extensions and the Antarctic Circumpolar Current. In these regions, it has been estimated that coherent eddies contribute around 40–50% of the net mesoscale kinetic energy (Chelton et al., 2011) and thus a significant fraction of the total kinetic energy (Ferrari & Wunsch, 2009). Although this estimate showcases the importance of the mesoscale coherent eddy field, the energy contained by coherent eddies was estimated by extracting the total geostrophic velocity within the radius of each detected coherent eddy; thus, it is possible that this estimate may contain energy from other processes. Here we extend on this past work by reconstructing the surface imprint of coherent eddies using a new eddy tracking algorithm and using the latest available satellite record.

76           There is broad consensus that mesoscale eddy kinetic energy has a pronounced sea-  
77           sonal variability (Qiu, 1999; Qiu & Chen, 2004; Kang & Curchitser, 2017; Uchida et al.,  
78           2017). Several hypotheses have been proposed to explain this seasonality including: sea-  
79           sonal variations of atmospheric forcing (Sasaki et al., 2014), seasonality of the mixed layer  
80           depth (Qiu et al., 2014; Callies et al., 2015), seasonality of the intensity of barotropic in-  
81           stability (Qiu & Chen, 2004), the variability of the baroclinic instability due to the sea-  
82           sonality of the vertical shear (Qiu, 1999), and a seasonal lag of the inverse energy cas-  
83           cade (i.e. energy is transported between scales, from small to large; Arbic et al., 2013)  
84           in combination with the presence of a front in the mixed layer, which can lead to a sea-  
85           sonal cycle of the baroclinic instability (Qiu et al., 2014). On one hand, processes such  
86           as barotropic and baroclinic instabilities control the seasonality of coherent eddies in the  
87           ocean. On the other hand, recent studies using observations and eddy-permitting climate  
88           models suggest slower adjustments of the global ocean that create long-term changes in  
89           the coherent eddy field. Such readjustments include a multidecadal increase in the ocean  
90           stratification resulting from temperature and salinity changes (Li et al., 2020), a hori-  
91           zontal readjustment of sea surface temperature gradients (Cane et al., 1997; Bouali et  
92           al., 2017; Ruela et al., 2020), and an intensification of the kinetic energy, eddy kinetic  
93           energy, and mesoscale eddy kinetic energy over the last 3 decades as a consequence of  
94           an increase in wind forcing (Hu et al., 2020; Wunsch, 2020; Martínez-Moreno et al., 2021).  
95           All of these seasonal factors and long-term readjustments directly influence the annual  
96           and decadal response of the coherent eddy field, however, the seasonality of the coher-  
97           ent component of the eddy kinetic energy, as well as the seasonal cycle and trends of the  
98           coherent eddy statistics, remain unknown.

99           Here we present a new global climatology of the coherent eddy kinetic energy by  
100           reconstructing the coherent eddy signature from satellite observations. Our study doc-  
101           uments the seasonal cycle of the coherent eddy kinetic energy, and the seasonal cycle and  
102           long-term trends of the coherent eddy properties over the satellite record. Moreover, we  
103           conduct more detailed analyses in regions where coherent eddies dominate the eddy ki-  
104           netic energy field. The rest of this paper is structured as follows: the data sources and  
105           methodology are described in Section 2. Then, we present the climatology, energy ra-  
106           tios, and global seasonality of the coherent eddy kinetic energy in Section 3. Section 4  
107           outlines the global climatology and seasonality of coherent eddy properties, followed by  
108           long-term changes of the coherent eddy properties (Section 5). Then we focus our at-

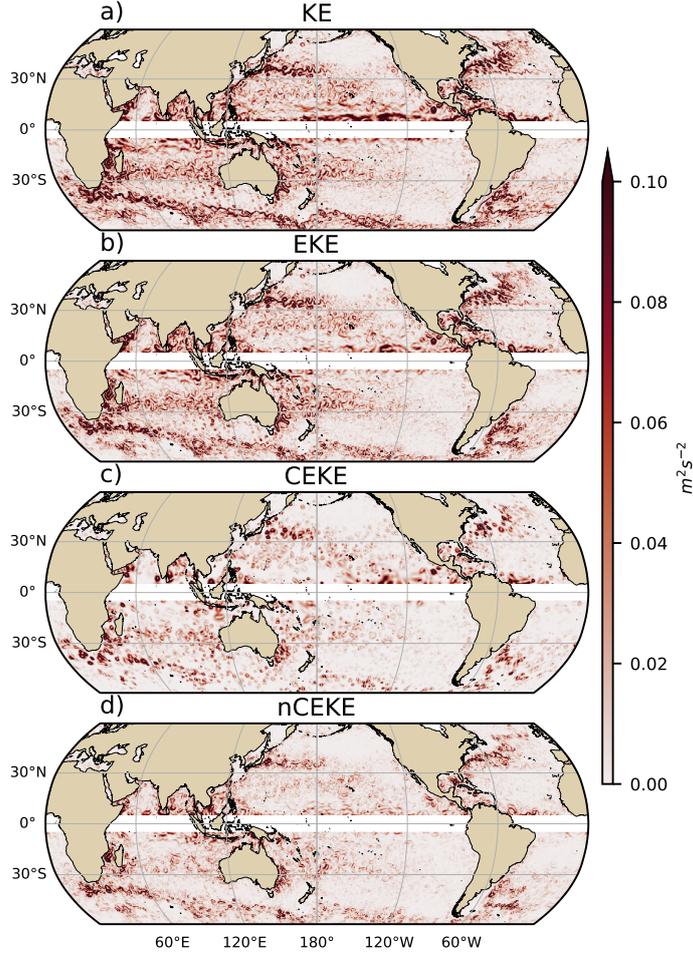
109 attention on the seasonal cycle and coherent eddy properties in regions dominated by co-  
110 herent eddies (Section 6). Finally, Section 7 summarizes the main results and discusses  
111 the implications of this study.

## 112 2 Methods

113 We use daily sea surface height (SSH) data made available by the Copernicus Ma-  
114 rine Environment Monitoring Service in near real time (CMEMS, 2017). This gridded  
115 product contains the sea surface height and geostrophic velocities with daily  $0.25^\circ$  res-  
116 olution from January 1993 to 2019. The daily geostrophic velocities allow us to compute  
117 the kinetic energy (KE) and eddy kinetic energy (EKE) over the satellite record. The  
118 main source of EKE is the time-varying wind (Ferrari & Wunsch, 2009); thus, we also  
119 compute the seasonal cycle of the wind magnitude from the JRA55 reanalysis (Japan  
120 Meteorological Agency, Japan, 2013) using wind velocities at 10m above the ocean's sur-  
121 face.

122 Over the same record, coherent eddy statistics from Martínez-Moreno et al. (2019),  
123 hereafter MM19, are analyzed and compared with those released by Chelton & Schlax  
124 (2013), hereafter CS13. Both datasets are gridded in a  $1^\circ$  resolution and are produced  
125 via automated eddy identification algorithms using closed contours of SSH. However, these  
126 datasets have important differences in the criteria they use to identify and record coher-  
127 ent eddies statistics. The major differences include: (i) MM19's algorithm requires an  
128 adjustment between a 2D Gaussian and the SSH anomaly (SSHa) surface within the iden-  
129 tified closed contour, while CS13's only uses the outermost closed contour of SSH; (ii)  
130 MM19's dataset reports the maximum SSHa within the identified coherent eddy, while  
131 CS13's algorithm reports the maximum SSH value minus the discrete level in which the  
132 coherent eddy was identified; and (iii) MM19's dataset includes all detected coherent ed-  
133 dies, while CS13's dataset excludes coherent eddies with lifetimes shorter than four weeks  
134 and coherent eddy amplitudes smaller than 1cm. Moreover, MM19's algorithm allows  
135 the reconstruction of the coherent eddy field under the assumption that coherent eddies  
136 have a 2D Gaussian imprint in the sea surface height. This Gaussian reconstruction of  
137 the coherent eddy field then allows us to estimate the coherent geostrophic eddy veloc-  
138 ities and thus the kinetic energy contained only by coherent eddies.





155 **Figure 1.** Snapshot of surface kinetic energy ( $\overline{KE}$ ), surface eddy kinetic energy ( $\overline{EKE}$ ),  
 156 surface coherent eddy kinetic energy ( $\overline{CEKE}$ ), and surface non-coherent eddy kinetic energy  
 157 ( $\overline{nCEKE}$ ) for the 1st of January 2017.

167 Note that the  $cEddy_{amp}^+$  and  $cEddy_{amp}^-$  are sign definite, thus the difference will always  
 168 be positive, whereas the gridded averaged  $cEddy_{amp}$  can be negative or positive noting  
 169 the dominant polarity of coherent eddies in the region, and the absolute value of  $cEddy_{amp}$   
 170 is denoted by  $cEddy_{|amp|}$ . We analyze the climatology and trends of the above eddy statis-  
 171 tics over the available satellite record, namely between 1993 and 2019. We exclude the  
 172 equatorial region ( $10^\circ S - 10^\circ N$ ) and regions poleward of  $60^\circ$ , because the geostrophic ap-  
 173 proximation is invalid near the Equator and the satellite spatial coverage at high-latitudes  
 174 is unable to resolve the coherent eddy scales polewards of  $60^\circ$ . Note that the climatol-  
 175 ogy of  $cEddy_n$  is computed by adding all the identified eddies over the record, while all

176 other climatological statistics are computed as the time-average over the record. Sea-  
 177 sonal climatologies are calculated for the monthly average of each coherent eddy statis-  
 178 tic, while hemispheric time-series are filtered with a running average of 90 days. Trends  
 179 of  $cEddy_n$  and  $|cEddy_{amp}|$  are calculated by coarsening the dataset to a  $5^\circ$  grid, and then  
 180 linear trends are computed for each grid point. The statistical significance of trends is  
 181 assessed by a modified Mann-Kendall test above the 95% confidence level (Yue & Wang,  
 182 2004).

183 Time averages are denoted by  $\overline{\quad}$ , while area-weighted averages are denoted using  
 184  $\langle \quad \rangle$ , where the area-weighted average of a function  $f$  is:

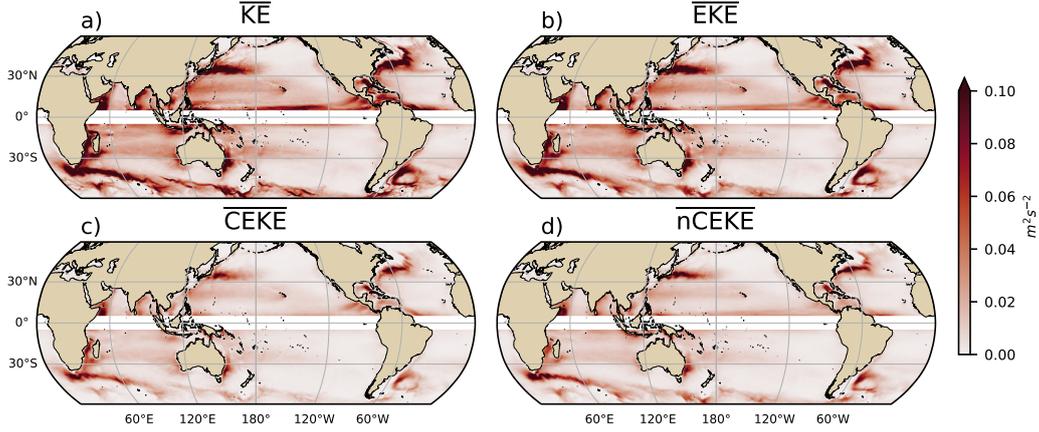
$$\langle f \rangle = \frac{\int f \xi dx dy}{\int \xi dx dy}, \quad (3)$$

185 where  $\xi$  is a mask that is set to zero in grid cells where no coherent eddies were iden-  
 186 tified and one elsewhere.

### 187 **3 Global Coherent Eddy Energetics**

188 The kinetic energy decomposition estimated from sea surface height measured by  
 189 satellite altimeters averaged from 1993-2019 is shown in Figure 2. These maps show that  
 190 many regions of the global ocean are highly energetic in mean KE ( $\overline{KE}$ ), mean EKE ( $\overline{EKE}$ ),  
 191 mean coherent eddy kinetic energy ( $\overline{CEKE}$ ) and mean non-coherent eddy kinetic energy  
 192 ( $\overline{nCEKE}$ ). The spatial pattern highlights well-known regions of the ocean where mesoscale  
 193 processes are abundant, such as the western boundary current (WBC) extensions and  
 194 the Antarctic Circumpolar Current. The spatial distribution of the energy contained by  
 195 the reconstructed mesoscale coherent eddies and non-coherent components are similar  
 196 (Figures 2c,d). However, there are some regions where coherent eddies dominate over  
 197 non-coherent, and vice-versa. Overall, this decomposition suggests that boundary cur-  
 198 rent extensions and other energetic regions of the ocean contain both coherent and non-  
 199 coherent components of the kinetic energy.

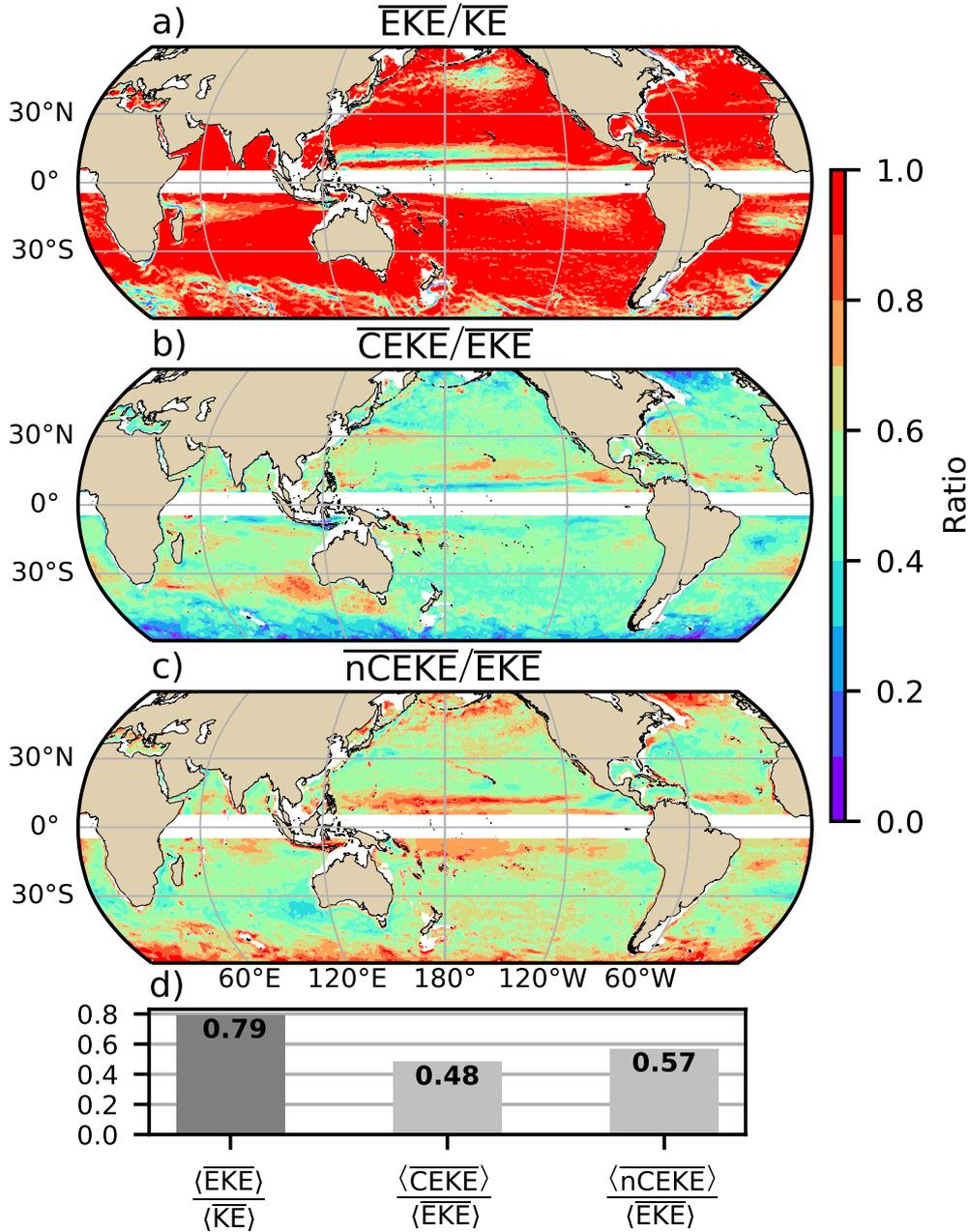
212 Eddy kinetic energy is known to be more than an order of magnitude greater than  
 213 kinetic energy of the mean flow (MKE; Gill et al., 1974); this result is clearly shown in  
 214 Figure 3a, which indicates that  $\overline{EKE}$  is responsible for almost all the  $\overline{KE}$  across the ocean,  
 215 except for regions with persistent currents over time. Such regions are located in the mean  
 216 boundary extension locations, the equatorial Pacific currents and regions in the Antarc-  
 217 tic Circumpolar Current, where the  $\overline{EKE}$  explains around 40% of the  $\overline{KE}$ . In a previ-



200 **Figure 2.** a) Mean surface kinetic energy ( $\overline{\text{KE}}$ ); b) surface eddy kinetic energy ( $\overline{\text{EKE}}$ ); c)  
 201 surface coherent eddy kinetic energy ( $\overline{\text{CEKE}}$ ), and d) surface non-coherent eddy kinetic energy  
 202 ( $\overline{\text{nCEKE}}$ ) averaged between 1993-2018.

218 our study, Chelton et al. (2011) estimated that the EKE within coherent eddies with life-  
 219 times greater than 4 weeks contain between 40-60% of the  $\overline{\text{EKE}}$ . Our method to recon-  
 220 struct the coherent eddy signature (Figure 3b) further corroborates that the coherent  
 221 eddy component ( $\langle \overline{\text{CEKE}} \rangle$ ) has  $\sim 48\%$  of the  $\langle \overline{\text{KE}} \rangle$  (Figure 3d). Furthermore, global area  
 222 averages of the ratios show that  $\langle \overline{\text{EKE}} \rangle$  explains  $\sim 78\%$  of the ocean  $\langle \overline{\text{KE}} \rangle$  field, while  
 223 non coherent eddy features contain  $\sim 57\%$  percent of the  $\langle \overline{\text{EKE}} \rangle$ . Note that the globally  
 224 averaged coherent and non coherent components do not add to 100% as the cross terms  
 225 ( $\mathcal{O}_c^2$ ) are non-zero. The spatial pattern reveals a dominance of the  $\overline{\text{CEKE}}$  equatorward  
 226 from the boundary current extensions and in areas with large coherent eddy contribu-  
 227 tions of around 80% of the region's eddy kinetic energy, such as south of Australia, in  
 228 the Tehuantepec Gulf, and in the tropical Atlantic. An evident signal is a reduction of  
 229 the energy contained by coherent eddies at high latitudes and an increase in the energy  
 230 explained by non-coherent eddies; this signal could be a consequence of the inability of  
 231 the  $0.25^\circ$  satellite resolution ( $\sim 13$  km at  $60^\circ$  latitude) to resolve coherent eddies with  
 232 scales smaller than  $\sim 10$  km (first baroclinic Rossby radius at  $60^\circ$ ; Chelton et al., 1998).

233 Figure 4 shows the seasonal cycle of the area-weighted EKE and CEKE for the North-  
 234 ern Hemisphere ( $\langle \text{EKE} \rangle_{NH}$  and  $\langle \text{CEKE} \rangle_{NH}$ ;  $10^\circ\text{N} - 60^\circ\text{N}$ ) and Southern Hemisphere  
 235 ( $\langle \text{EKE} \rangle_{SH}$  and  $\langle \text{CEKE} \rangle_{SH}$ ;  $60^\circ\text{S} - 10^\circ\text{S}$ ). In both hemispheres, the  $\langle \text{EKE} \rangle$  and  $\langle \text{CEKE} \rangle$   
 236 peak during summer. In the Northern Hemisphere, the largest  $\langle \text{EKE} \rangle_{NH}$  and  $\langle \text{CEKE} \rangle_{NH}$



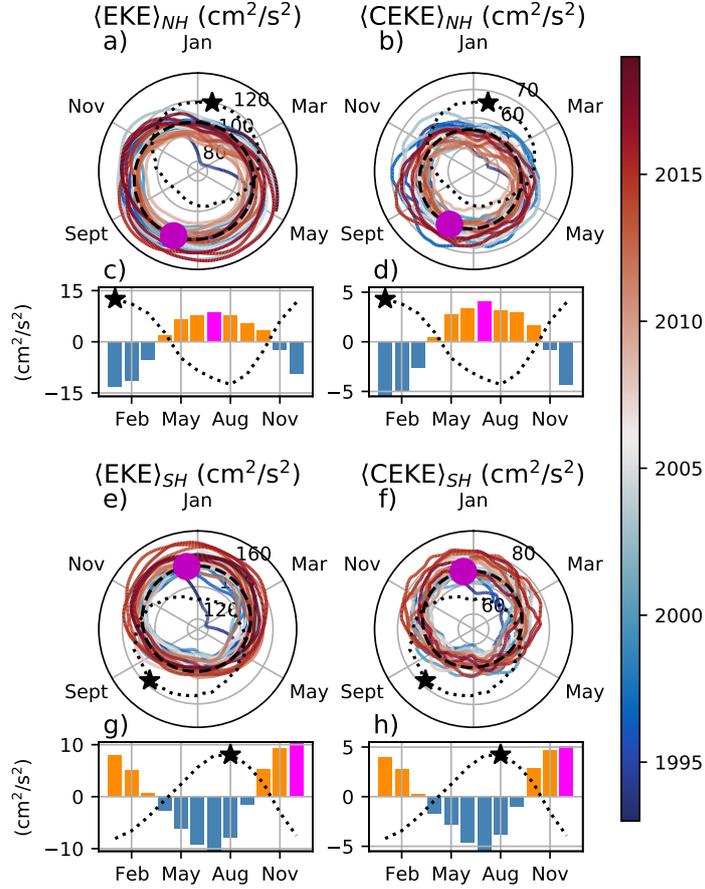
203 **Figure 3.** Ratios of the kinetic energy components. a) Map of the proportion of mean eddy  
 204 kinetic energy ( $\overline{EKE}$ ) versus mean kinetic energy ( $\overline{KE}$ ); b) Map of the fraction of mean coherent  
 205 eddy kinetic energy ( $\overline{CEKE}$ ) versus mean eddy kinetic energy ( $\overline{EKE}$ ); c) Map of the fraction of  
 206 mean non-coherent eddy kinetic energy ( $\overline{nCEKE}$ ) versus mean eddy kinetic energy ( $\overline{EKE}$ ); d)  
 207 Global time and area averaged (represented by  $\langle \rangle$ ) fraction of mean eddy kinetic energy ( $\langle \overline{EKE} \rangle$ )  
 208 versus the global mean kinetic energy ( $\langle \overline{KE} \rangle$ ), area averaged fraction of mean coherent eddy  
 209 kinetic energy ( $\langle \overline{CEKE} \rangle$ ) and mean non coherent eddy kinetic energy ( $\langle \overline{nCEKE} \rangle$ ) versus global  
 210 mean eddy kinetic energy ( $\langle \overline{EKE} \rangle$ ). Regions where the depth of the ocean is shallower than  
 211 1000m are removed from the ratio estimation.

237 occurs in July,  $\sim 6$  months after the maximum winds in January (purple bar and black  
 238 star in Figure 4c and d). Meanwhile, the Southern Ocean  $\langle \text{EKE} \rangle_{SH}$  and  $\langle \text{CEKE} \rangle_{SH}$  sea-  
 239 sonal maxima arises during December,  $\sim 4$  months after the maximum winds in August  
 240 (purple bar and back star in Figure 4g, and h). This lag between winds and the eddy  
 241 and coherent eddy energy components is further discussed in Section 4.

242 The cyclic plots in Figure 4 show the temporal evolution of  $\langle \text{EKE} \rangle$  and  $\langle \text{CEKE} \rangle$ .  
 243 Note that high frequency variability can be observed in the  $\langle \text{CEKE} \rangle$  field with tempo-  
 244 ral scales of a few months, this variability could be attributed to regional dynamics av-  
 245 eraged over the hemisphere (boundary currents, ocean gyres, etc.), as well as errors within  
 246 the coherent eddy reconstruction. Additionally, concentric changes in the cyclic plots high-  
 247 light long-term changes over the record. For example, the Northern Hemisphere winters  
 248 during early years of the record (blue) had a more energetic coherent eddy field, which  
 249 has transitioned to weaker coherent energy content since 2010 (red), in other words, the  
 250 intensity of the  $\langle \text{CEKE} \rangle_{NH}$  field has decreased. A larger long-term change can be ob-  
 251 served in the Southern Hemisphere, where concentric growth over time in  $\langle \text{EKE} \rangle_{SH}$  and  
 252  $\langle \text{CEKE} \rangle_{SH}$  supports the previously observed strengthening of the eddy field in the South-  
 253 ern Ocean (Hogg et al., 2015; Martínez-Moreno et al., 2019; Martínez-Moreno et al., 2021).

## 263 4 Global Coherent Eddy Statistics

264 Coherent eddy kinetic energy allows us to quantify and study the energy of the eddy  
 265 field, but the coherent eddy properties computed by automated coherent eddy identi-  
 266 fication algorithms allow us to further investigate the contribution and temporal changes  
 267 of their abundance (i.e. the number of eddies) and their intensity (both their amplitude  
 268 and diameter). Figure 5 shows gridded estimates of the number of eddies and the eddy  
 269 amplitude. In this analysis, we contrast our MM19 eddy count with that of CS13 (Chel-  
 270 ton et al., 2007; Figure 5a-b). Although the number of identified eddies is larger in MM19,  
 271 possibly due to the lifespan filter implemented by CS13, both datasets reveal consistent  
 272 spatial patterns. For example, both datasets show an important meridional variation in  
 273 the abundance of eddies, with high numbers of eddies in mid-latitudes and fewer eddies  
 274 in the tropics and at high-latitudes ( $\sim 60^\circ$ ). Additionally, there is a tendency at mid-latitudes  
 275 ( $30^\circ$ ) for higher numbers of eddies in the eastern side of ocean basins (e.g. the East North  
 276 Pacific, East North Atlantic, East South Pacific, and East South Atlantic). Another in-  
 277 teresting pattern emerges in both eddy count datasets, where small scale structures ap-



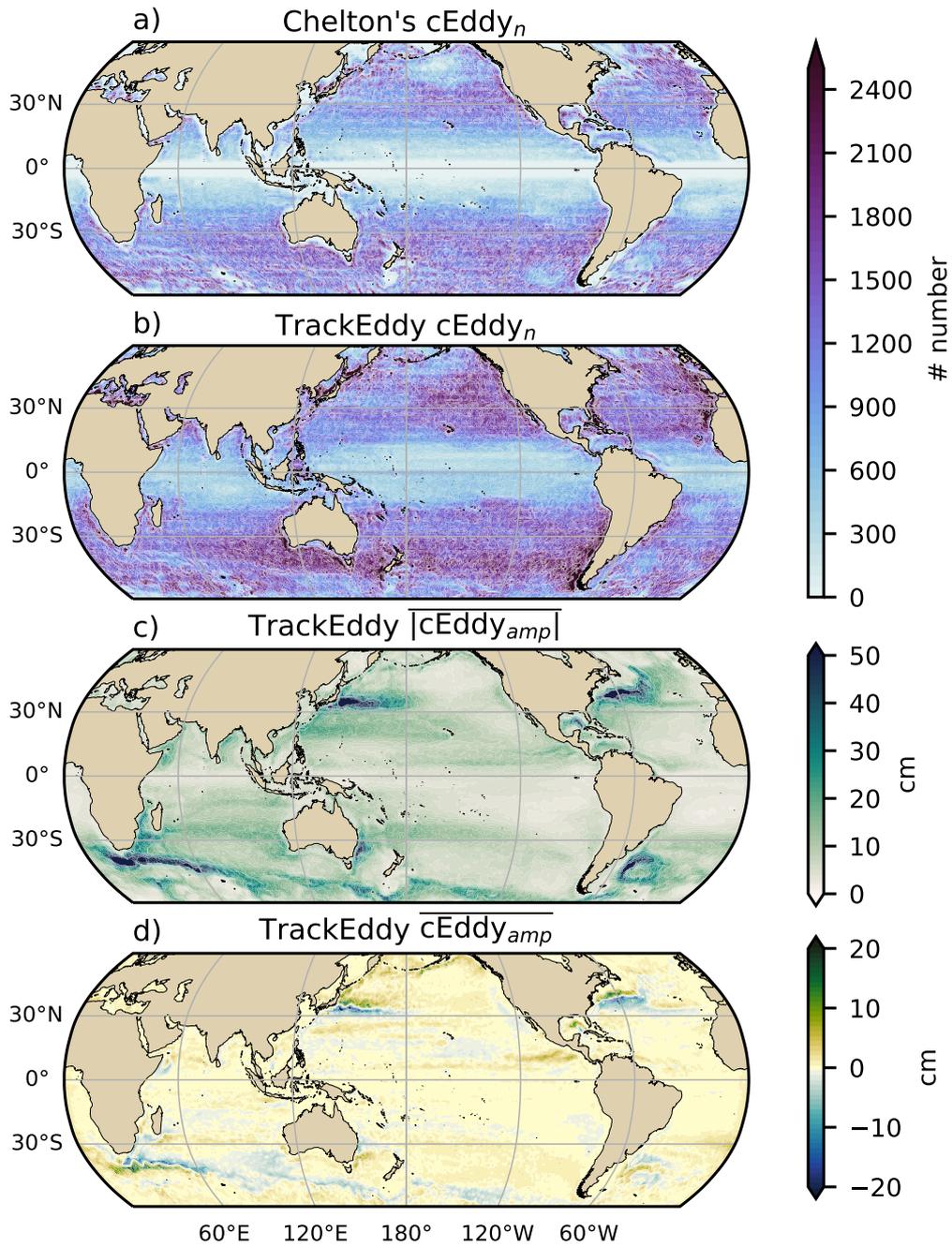
254 **Figure 4.** Seasonality of the area-weighted eddy kinetic energy ( $\langle \text{EKE} \rangle$ ) and coherent eddy  
 255 kinetic energy ( $\langle \text{CEKE} \rangle$ ). Panels a) and b) show the time-series of the Northern Hemisphere,  
 256 while panels e) and f) correspond to the Southern Hemisphere. Panels c) and d) show the sea-  
 257 sonal cycle of the  $\langle \text{EKE} \rangle_{NH}$  and  $\langle \text{CEKE} \rangle_{NH}$  in the Northern Hemisphere, and panels g) and h)  
 258 show the Southern Hemisphere ( $\langle \text{EKE} \rangle_{SH}$  and  $\langle \text{CEKE} \rangle_{SH}$ ). Dashed lines correspond to the sea-  
 259 sonal cycle of the fields and dotted lines show the seasonal cycle of the wind magnitude smoothed  
 260 over 120 days (moving average). The black stars and magenta markers (circle and bar) show  
 261 the maximum of the seasonal cycle for the kinetic energy components and the wind magnitude,  
 262 respectively. In the cyclic plots, line colors shows the year.

278 pear in the eddy count field. These small structures highlight preferred coherent eddy  
 279 paths observable in boundary current extensions and over regions of the Southern Ocean.  
 280 These structures and paths of coherent eddies could be associated with topographic fea-  
 281 tures, with overall consistency between the eddy count patterns using the two different  
 282 eddy identification methods.

288 Regions with large counts of eddies have, in general, small absolute amplitudes (Fig-  
 289 ure 5c), for example, the eastern side of mid-latitude ocean basins. The ocean gyre in-  
 290 teriors have a larger absolute amplitude and finally regions such as the boundary cur-  
 291 rent extensions and the Antarctic Circumpolar Current have the largest coherent eddy  
 292 absolute amplitudes, as also shown by Chelton et al. (2011). Eddy amplitude highlights  
 293 regions dominated by a given coherent eddy polarity, for example, boundary extensions  
 294 have a preferred sign (Figure 5 d); namely, positive amplitude polewards of the bound-  
 295 ary current extension mean location, and negative amplitude equatorwards. This sign  
 296 preference is consistent with the preferential way that coherent eddies are shed from bound-  
 297 ary current extensions; with warm core eddies (positive) polewards of the boundary cur-  
 298 rent extension, and equatorward for cold core eddies (negative) (Chelton et al., 2007, 2011;  
 299 Kang & Curchitser, 2013). These global statistics reveal the absolute coherent eddy am-  
 300 plitude as a proxy for the CEKE with similar spatial patterns (Figure 2 & Figure 5c)  
 301 and showcases that in regions where  $\overline{\text{CEKE}}$  has a large proportion of  $\overline{\text{EKE}}$  (Figure 3),  
 302 the absolute coherent eddy amplitude is also large.

303 To further understand the seasonal cycle of  $\langle \text{CEKE} \rangle$ , we compute the climatology  
 304 of coherent eddy properties in each hemisphere (Figure 6). The seasonality of the num-  
 305 ber of eddies in the Northern Hemisphere peaks in April (Figure 6a, c), while the South-  
 306 ern Hemisphere maximum number of eddies occurs during October (Figure 6e, g). Mean-  
 307 while, the seasonality of the eddy amplitude ( $\langle |c\text{Eddy}_{amp}| \rangle$ ) peaks in August and Jan-  
 308 uary for the Northern and Southern Hemispheres respectively (Figure 6b, d, f, and h).  
 309 As expected, the seasonality of  $\langle |c\text{Eddy}_{amp}| \rangle$ , equivalent to the intensity of the coher-  
 310 ent eddies, is consistent with the seasonal cycle of  $\langle \text{CEKE} \rangle$ .

311 A key feature of Figure 6 is a distinct lag of  $\sim 3$  months between the winds and eddy  
 312 count, while the eddy amplitude maximum occurs  $\sim 6$  months after the seasonal max-  
 313 ima in winds. We suggest that the eddy number increases earlier in the year and, through  
 314 eddy-eddy interactions (merging of coherent eddies), the coherent eddy amplitude in-

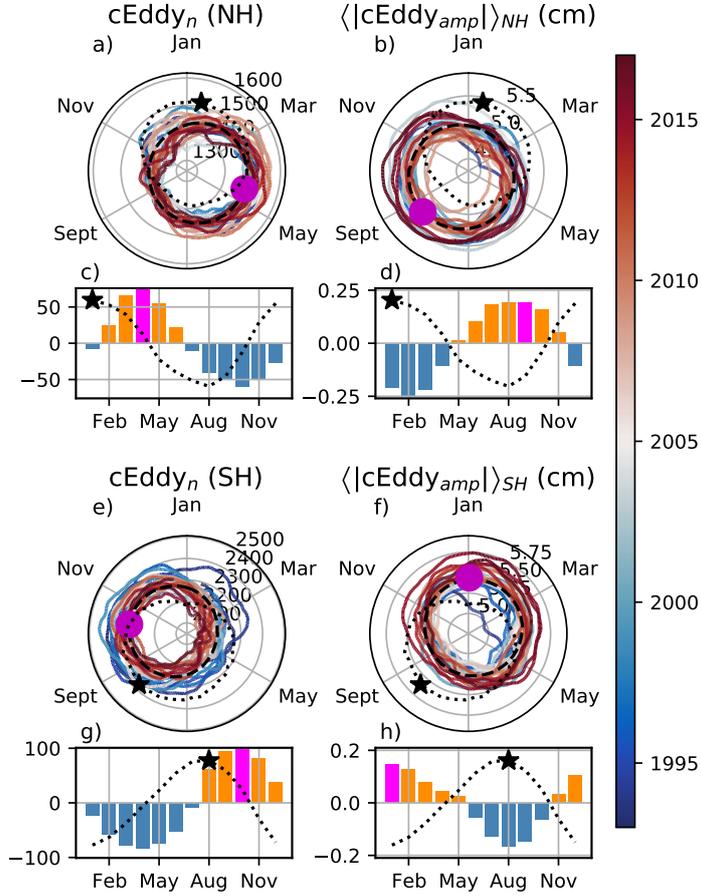


283 **Figure 5.** Averaged coherent eddy statistics. a) Climatology of the number of coherent eddies  
 284 ( $cEddy_n$ ) identified by Chelton et al. (2007); b) Climatology of the number of coherent eddies  
 285 ( $cEddy_n$ ) identified by Martínez-Moreno et al. (2019); c) Climatology of the mean absolute co-  
 286 herent eddy amplitude ( $cEddy_{amp}$ ), and d) Climatology of the mean coherent eddy amplitude  
 287 ( $\overline{cEddy_{amp}}$ ).

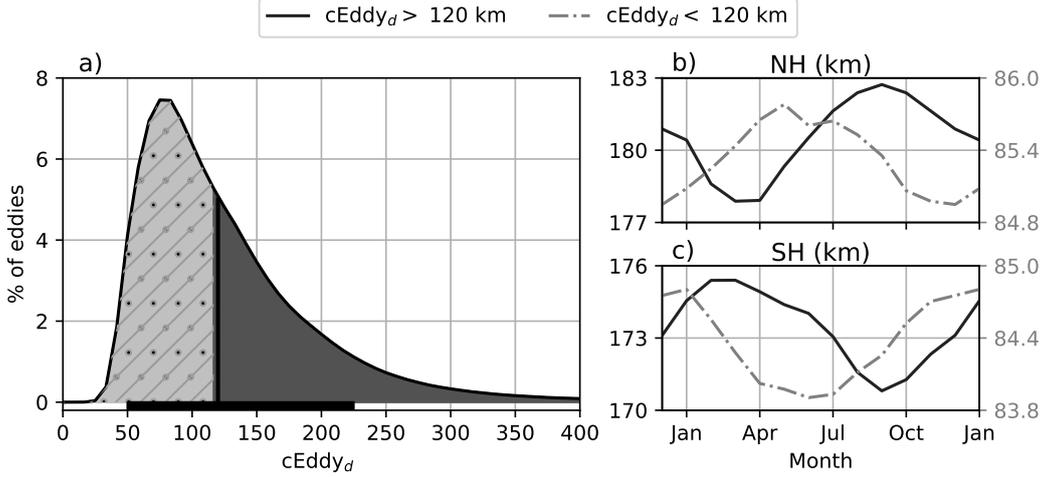
315 creases  $\sim 3$  months after. This seasonal lag and summer maxima is consistent with pre-  
 316 vious studies which suggest that a time-lag of the inverse cascade (Sasaki et al., 2014;  
 317 Qiu et al., 2014) is responsible for the EKE seasonal cycle, where winter has the high-  
 318 est energy at the smallest scales (non-resolvable with satellite observations), spring and  
 319 autumn have the highest and lowest energy at scales of 50-100 km, and summertime has  
 320 the highest energy at the largest scales ( $> 100$  km; Uchida et al., 2017). Thus, the max-  
 321 imum of  $\langle \text{EKE} \rangle$ ,  $\langle \text{CEKE} \rangle$ , and  $\langle |\text{cEddy}_{amp}| \rangle$  located during summertime suggests that  
 322 the seasonality of eddies and coherent eddies could be dominated by scales larger than  
 323 100 km.

324 This result can be further explored by looking at the seasonal evolution of the eddy  
 325 diameter ( $\text{cEddy}_d$ ). Note that 90% of identified coherent eddies have diameters between  
 326 50 to 220 km (Figure 7a). We partition eddies into large-scale coherent eddies (diam-  
 327 eter  $> 120$  km) and small-scale coherent eddies (diameter  $< 120$  km; Figure 7a). In the  
 328 Northern Hemisphere, small-scale eddies have a seasonal peak in diameter during May,  
 329 while large-scale eddies have the greatest diameter in September (Figure 7b). Meanwhile,  
 330 in the Southern Hemisphere, the small-scale coherent eddies exhibit maximum diame-  
 331 ter in December, while the diameter of large-scale coherent eddies peaks in February (Fig-  
 332 ure 7 c). This result suggests that wind driven baroclinic instabilities generate small co-  
 333 herent eddies early in the season, which then merge and grow to become larger in diam-  
 334 eter and amplitude, and thus, more energetic. This process is likely associated with the  
 335 inverse energy cascade, and suggests that this mechanism not only drives EKE season-  
 336 ality, but also may be responsible for the seasonal cycle of coherent eddies.

352 Long-term changes can be observed in Figure 6a,b, e, and f where growing/shrinking  
 353 concentric circles over time denote an increase/decrease trend of the field. This trend  
 354 is particularly evident in the Southern Hemisphere, where the number of eddies has de-  
 355 creased, while the eddy amplitude has increased. This result is consistent with the ob-  
 356 served trends in EKE and mesoscale EKE in the Southern Ocean (Hogg et al., 2015; Martínez-  
 357 Moreno et al., 2019). The coherent eddy amplitude from positive coherent eddies and  
 358 negative coherent eddies show similar seasonal cycles to the absolute eddy amplitude.  
 359 The Northern Hemisphere decrease in absolute eddy amplitude is driven by a decrease  
 360 of the amplitude of negative coherent eddies in the Northern Hemisphere. Meanwhile  
 361 in the Southern Ocean, the increase in absolute eddy amplitude is corroborated by a strength-  
 362 ening of both coherent eddy polarities since the early 90s.



337 **Figure 6.** Seasonality of the count of number of eddies ( $cEddy_n$ ) and the area-weighted polar-  
 338 ity independent coherent eddy amplitude ( $\langle |cEddy_{amp}| \rangle$ ); Panels a and b show the time-series of  
 339 the Northern Hemisphere, while panels e and f correspond to the Southern Hemisphere. Panels  
 340 c and d show the seasonal cycle of  $cEddy_n$  and  $\langle |cEddy_{amp}| \rangle_{NH}$  in the Northern Hemisphere,  
 341 and panels g and h show the Southern Hemisphere,  $cEddy_n$  and  $\langle |cEddy_{amp}| \rangle_{SH}$ . Dashed lines  
 342 correspond to the seasonal cycle of the fields and dotted lines show the seasonal cycle of the wind  
 343 magnitude, smoothed over 120 days (moving average). The black stars and magenta markers  
 344 (circle and bar) indicate the maximum of the seasonal cycle for the eddy property, and the wind  
 345 magnitude, respectively. In the cyclic plots, line colors show the year from 1993-2019.

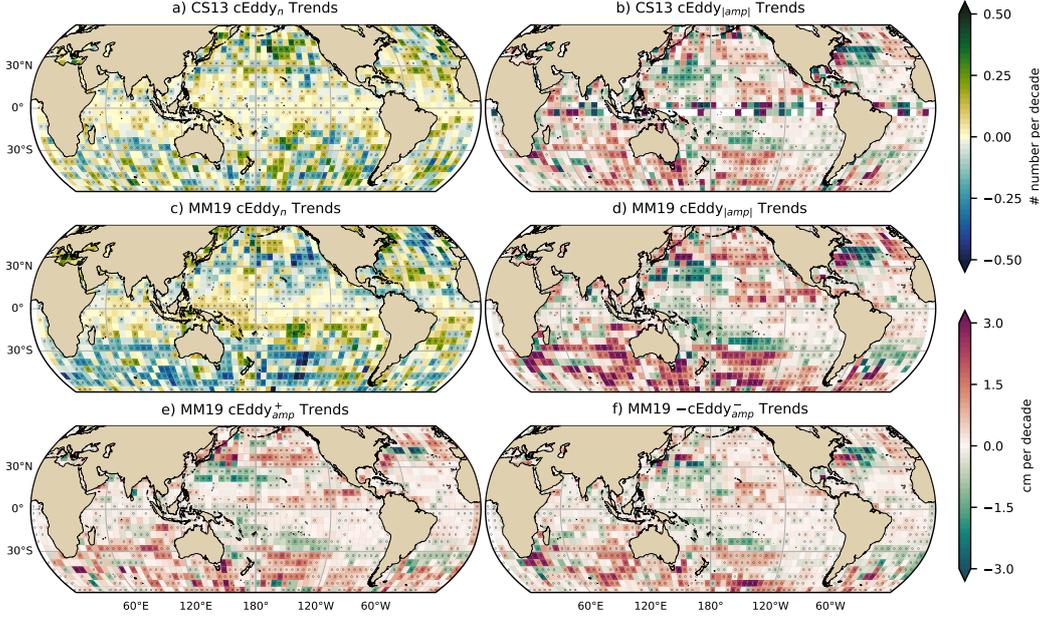


346 **Figure 7.** Distribution of the identified eddy diameter ( $cEddy_d$ ; km) and hemispherical  
 347 seasonality of the coherent eddy diameter. a) Distribution in percentage of identified eddy am-  
 348 plitude, solid bar below distribution represents 90% of the identified eddies. Seasonal cycle of  
 349 the eddy diameter for the b) Northern Hemisphere and c) Southern Hemisphere. Dark solid line  
 350 and area corresponds to coherent eddies with diameters larger than 120 km, while light gray  
 351 dash-dotted line and area shows coherent eddies with diameters smaller than 120 km.

## 363 5 Trends

364 The results presented in Figures 4 and 6 suggest a long-term readjustment of the  
 365 coherent eddy field. The long-term trends of the number of coherent eddies, absolute co-  
 366 herent eddy amplitude, and coherent eddy amplitude polarities are further explored in  
 367 Figure 8 contrasting the MM19 and CS13 methods. Both MM19 and CS13 datasets show  
 368 consistent spatial patterns in the trends and significance of the number of coherent ed-  
 369 dies and the absolute coherent eddy amplitude. Several regions in the ocean, such as the  
 370 Southern Ocean, North Atlantic and North Pacific, show a decrease in the number of ed-  
 371 dies. Those same regions also have a clear increase in the absolute coherent eddy am-  
 372 plitude. These trends are similar to those observed in mesoscale eddy kinetic energy (Martínez-  
 373 Moreno et al., 2021) and provide additional evidence of a readjustment of the mesoscale  
 374 eddy field over the last 3 decades.

375 The observed trends of  $cEddy_{|amp|}$  in several oceanic regions have the same scale  
 376 as sea level rise ( $\sim 3\text{cm}$  per decade). By analyzing the positive and negative coherent eddy  
 377 amplitude, we filter out the observed trends that come from a net increase in sea level.



385 **Figure 8.** Trends of coherent eddy statistics. a) and b) Trends of the number of identified  
 386 coherent eddies from satellite observations identified using the TrackEddy scheme of MM19,  
 387 and those reported in CS13’s dataset. c) and d) Trends of the absolute value of identified coher-  
 388 ent eddy amplitude ( $cEddy_{|amp|}$ ) from satellite observations identified using TrackEddy (after  
 389 MM19), and those reported by CS13. e) and f) Trends of the eddy amplitude polarity using  
 390 TrackEddy ( $cEddy_{amp}^+$  and  $cEddy_{amp}^-$ ). Gray stippling shows regions that are statistically signifi-  
 391 cant above the 95% confidence level.

378 In fact, each coherent eddy polarity has intensified in the Southern Ocean and North East  
 379 Pacific and Atlantic. In other words, the amplitude of each polarity has increased over  
 380 time, and thus this strengthening is an intrinsic response of the coherent eddy field. Note  
 381 that the negative coherent eddy amplitude dominates the global  $|cEddy_{amp}|$  trends (Fig-  
 382 ure 8e, f). However, different trend patterns can be observed in both positive and neg-  
 383 ative coherent eddy amplitudes in the North Atlantic and North Pacific, where the neg-  
 384 ative coherent eddy amplitude in the Western Boundary Currents appears to decrease.

## 392 6 Regional Climatology

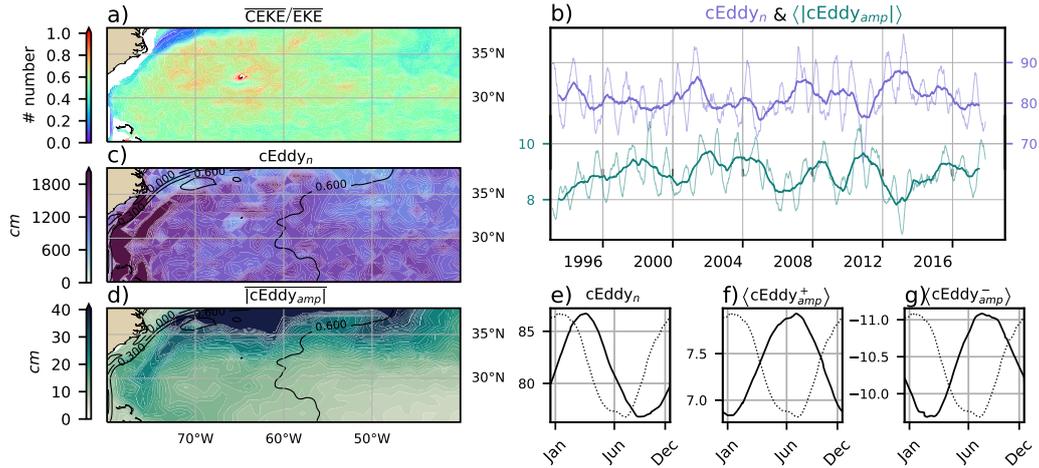
393 For regions with relatively large proportions of CEKE located at WBC extensions  
 394 and eastern boundary currents, we investigate the seasonal and long-term variability of  
 395 the coherent eddy properties. The most energetic WBC include the Gulf Stream, the Kuroshio

396 Current, and the Agulhas Current (Figures 9, 10, and 11). Coherent eddy generation in  
 397 boundary current extensions occurs through baroclinic and barotropic instabilities of the  
 398 mean current, thus all these regions share similar generation dynamics. In all these re-  
 399 gions without exception; (i) CEKE contains 50-80% of the EKE in regions equatorward  
 400 from the mean WBC extensions, (ii) the number of eddies is consistently small over the  
 401 mean WBC extensions, and (iii) the eddy amplitude is larger over the mean WBC ex-  
 402 tensions.

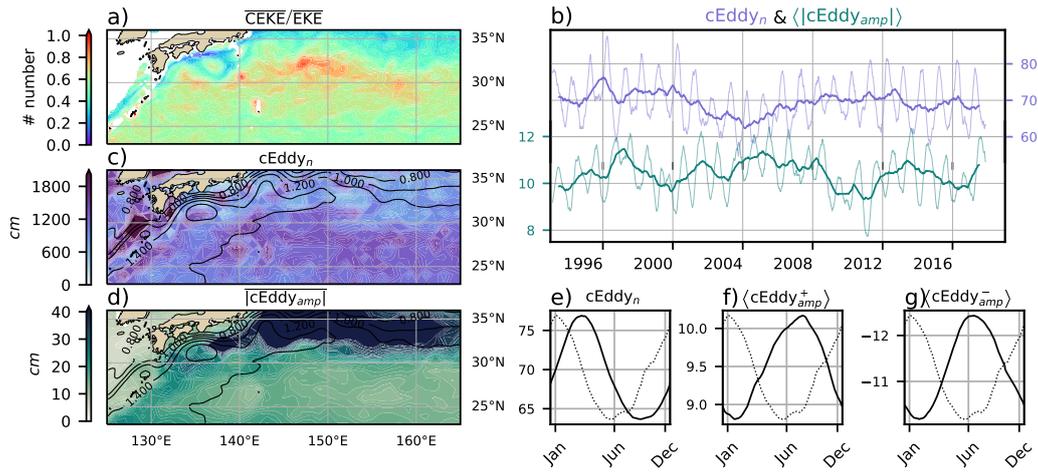
403 In the Gulf Stream, the energy ratio between CEKE and EKE is  $\sim 56\%$  (Figure 9).  
 404 The highest energy ratio occurs in regions with numerous eddies, colocated with regions  
 405 where the largest  $|\text{cEddy}_{amp}|$  gradients occur. The time series of  $\text{cEddy}_n$  and  $\langle |\text{cEddy}_{amp}| \rangle$   
 406 are anti-correlated ( $-0.52$ ), and they display interannual and seasonal variability. Although  
 407 Chaudhuri et al. (2009) observed that a positive phase of the North Atlantic Oscillation  
 408 (NAO) exhibits higher EKE, due to an increase in baroclinic instability, thus suggest-  
 409 ing more coherent eddies, we do not find a correlation between the  $\text{cEddy}_n$  or the  $\langle |\text{cEddy}_{amp}| \rangle$   
 410 in the Gulf Stream and the NAO index. Similar to the signal observed in the hemispheric  
 411 analysis, the eddy count seasonal cycle follows the wind maximum lagging by  $\sim 3$  months,  
 412 while the amplitude of the coherent eddies lags by  $\sim 6$  months.

421 The variability of the  $\text{cEddy}_n$  and  $\langle |\text{cEddy}_{amp}| \rangle$  in the Kuroshio Current are weakly  
 422 anti-correlated ( $-0.41$ ; Figure 10). However, on average 56% of the energy in the region  
 423 corresponds to CEKE. As observed in the Gulf Stream, there is an important seasonal  
 424 cycle in the boundary extension, where the eddy count seasonal cycle occurs in March,  
 425 lagging the wind maximum by  $\sim 3$  months (January). Meanwhile, the amplitude of the  
 426 coherent eddies lags the wind maximum by  $\sim 6$  months (June).

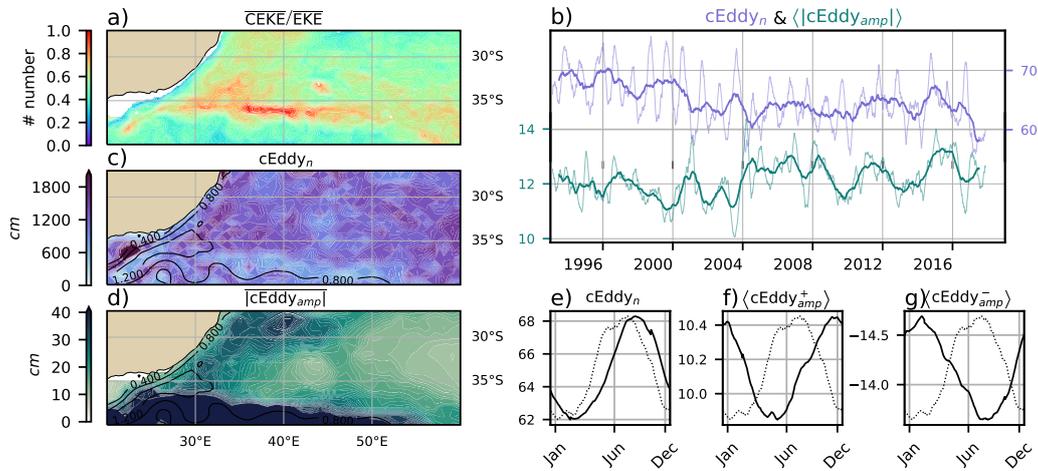
433 In the Southern Hemisphere the strongest boundary current, the Agulhas Current,  
 434 shows similar behavior to its counterparts in the Northern Hemisphere (Figure 11). On  
 435 average, coherent eddies in the Agulhas Current contain  $\sim 56\%$  of the energy, meanwhile  
 436 the  $\text{cEddy}_n$  seasonal peak occurs in August, while the  $\langle |\text{cEddy}_{amp}| \rangle$  peak occurs in January-  
 437 February. The seasonal lag between the winds, eddy count, and eddy amplitude in each  
 438 of the WBC extensions is interpreted as being analogous to the lagged response of co-  
 439 herent eddy properties (Figure 6) due to eddy-eddy interactions, consistent with the in-  
 440 verse cascade of energy.



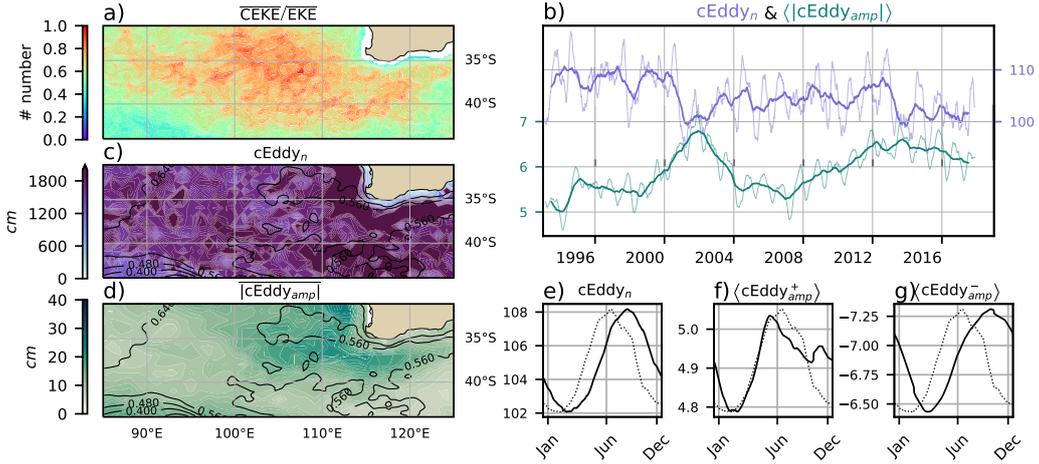
413 **Figure 9.** Climatology of the eddy field and coherent eddy field in the Gulf Stream. a) Ratio  
 414 of mean coherent eddy kinetic energy ( $\overline{CEKE}$ ) versus mean eddy kinetic energy ( $\overline{EKE}$ ); b) Thick  
 415 lines show the running average over 2 years and thin lines show the running average over 90 days  
 416 of the coherent eddy number sum and the average coherent eddy amplitude; c) Map of the number of  
 417 eddies; d) Map of the average coherent eddy amplitude; e) Seasonal cycle of the number  
 418 of eddies ( $\langle cEddy_n \rangle$ ); f) Seasonal cycle of the positive coherent eddy amplitude ( $\langle cEddy_{amp}^+ \rangle$ ),  
 419 and g) Seasonal cycle of the negative coherent eddy amplitude ( $\langle cEddy_{amp}^- \rangle$ ). Contours in maps  
 420 correspond to mean sea surface height (m).



427 **Figure 10.** As in Figure 9, only showing the climatology of the eddy field and coherent eddy  
 428 field in the Kuroshio extension. a) Ratio of mean coherent eddy kinetic energy ( $\overline{\text{CEKE}}$ ) versus  
 429 mean eddy kinetic energy ( $\overline{\text{EKE}}$ ); b) Time-series of the coherent eddy number and the average  
 430 coherent eddy amplitude; c) Map of the number of eddies; d) Map of the average coherent eddy  
 431 amplitude; Seasonal cycle of the e) number of eddies; f) positive coherent eddy amplitude, and g)  
 432 negative coherent eddy amplitude.



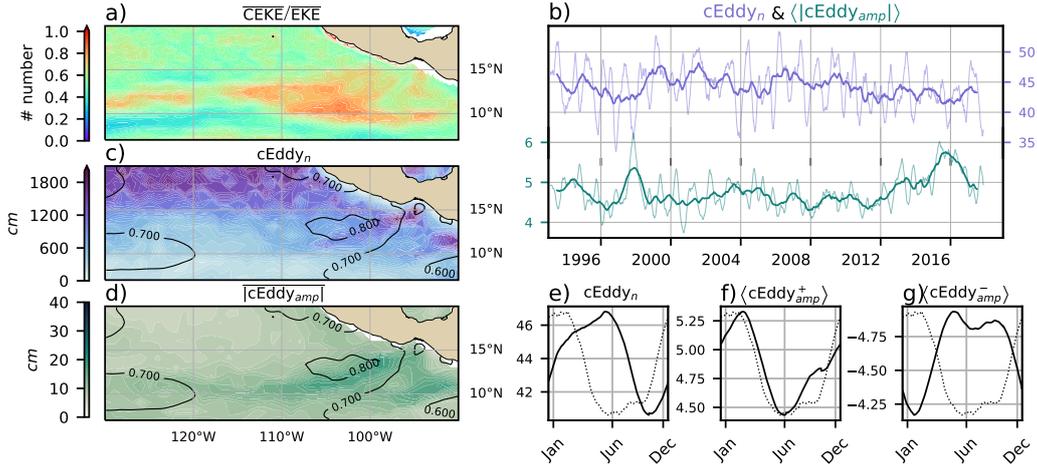
441 **Figure 11.** As in Figure 9, only showing the climatology of the eddy field and coherent eddy  
 442 field in the Agulhas Current. a) Ratio of mean coherent eddy kinetic energy ( $\overline{\text{CEKE}}$ ) versus  
 443 mean eddy kinetic energy ( $\overline{\text{EKE}}$ ); b) Time-series of the coherent eddy number and the average  
 444 coherent eddy amplitude; c) Map of the number of eddies; d) Map of the average coherent eddy  
 445 amplitude; Seasonal cycle of the e) number of eddies; f) positive coherent eddy amplitude, and g)  
 446 negative coherent eddy amplitude.



457 **Figure 12.** As in Figure 9, only showing the climatology of the eddy field and coherent eddy  
 458 field in the Leeuwin Current. a) Ratio of mean coherent eddy kinetic energy ( $\overline{\text{CEKE}}$ ) versus  
 459 mean eddy kinetic energy ( $\overline{\text{EKE}}$ ); b) Time-series of the coherent eddy number and the average  
 460 coherent eddy amplitude; c) Map of the number of eddies; d) Map of the average coherent eddy  
 461 amplitude; Seasonal cycle of the e) number of eddies; f) positive coherent eddy amplitude, and g)  
 462 negative coherent eddy amplitude.

447 Coherent eddies dominate the EKE field in other regions such as the Leeuwin Cur-  
 448 rent (Figure 12), where 65% of the energy is contained by coherent eddies. The Leeuwin  
 449 region is not characterized by having a large EKE, however, a considerable abundance  
 450 of eddies and large eddy amplitudes are observed in the region. The time-series reveal  
 451 a significant increase in the  $\langle |c\text{Eddy}_{amp}| \rangle$ , while the  $c\text{Eddy}_n$  has decreased over the last  
 452 3 decades. The seasonal cycle shows that the  $c\text{Eddy}_n$  peak occurs in August, 3 months  
 453 after the maximum winds (June). Meanwhile, the  $\langle c\text{Eddy}_{amp}^+ \rangle$  responds in synchrony  
 454 to the winds, and the  $\langle c\text{Eddy}_{amp}^- \rangle$  is in phase with the seasonal cycle of the eddy num-  
 455 ber ( $c\text{Eddy}_n$ ). Hence, this region contrasts the behavior of WBC extensions, and show-  
 456 cases the spatial variability of the seasonal cycle of coherent eddies.

463 Another region with important contributions to the coherent eddy field is the East  
 464 Tropical Pacific (Tehuantepec region; Figure 13), where coherent eddies contain  $\sim 58\%$   
 465 of the energy. In fact, coherent eddy generation in this region is modulated by winds and  
 466 coastally trapped waves which produce a strong horizontal and vertical shear (baroclinic  
 467 and barotropic instabilities; Zamudio et al., 2006). Furthermore, the equatorial gener-



476 **Figure 13.** As in Figure 9, only showing the climatology of the eddy field and coherent eddy  
 477 field in the East Tropical Pacific. a) Ratio of mean coherent eddy kinetic energy ( $\overline{CEKE/EKE}$ ) versus  
 478 mean eddy kinetic energy ( $\overline{EKE}$ ); b) Time-series of the coherent eddy number and the average  
 479 coherent eddy amplitude; c) Map of the number of eddies; d) Map of the average coherent eddy  
 480 amplitude; Seasonal cycle of the e) number of eddies; f) positive coherent eddy amplitude, and g)  $\langle cEddy_{amp}^- \rangle$   
 481 negative coherent eddy amplitude.

468 ated waves propagating along the coast have an important interannual variability ob-  
 469 servable in the  $\langle |cEddy_{amp}| \rangle$  time-series, where El Niño events are notable during 1997  
 470 and 2015 (Figure 13b). The seasonal cycle of  $cEddy_n$ ,  $\langle cEddy_{amp}^+ \rangle$ , and  $\langle cEddy_{amp}^- \rangle$  sup-  
 471 port the idea of a coherent eddy response to two different coherent eddy generation mech-  
 472 anisms; the number of eddies lags by  $\sim 3$  months from the winds, while the  $\langle cEddy_{amp}^+ \rangle$   
 473 is in phase with the winds and the time of maximum trapped wave activity (winter; Za-  
 474 mudio et al., 2006), while the  $\langle cEddy_{amp}^- \rangle$  could be a consequence of eddy-eddy inter-  
 475 actions.

## 482 7 Discussion and Conclusions

483 We have investigated the contribution of coherent eddies to the total kinetic en-  
 484 ergy field using available satellite observations. We found that around half of the EKE  
 485 is explained by coherent eddies. This half is concentrated in eddy-rich regions where a  
 486 recent multi-decadal intensification of the eddy field has been observed (Martínez-Moreno  
 487 et al., 2021). The energy contained by eddies is larger than the previous estimate of 40%

488 by Chelton et al. (2011). Although there are differences in the identification criteria of  
489 both eddy identification methods, the main cause of the difference is likely to be the lifes-  
490 pan and amplitude filters. These filters are widely used to track individual eddies in space  
491 and time, however, interactions between eddies in energetic regions may obscure the abun-  
492 dance and influence of short-lived coherent eddies. Filters are not used in this study, and  
493 indeed a lack of filters could facilitate an over-estimation of the the energy contained by  
494 coherent eddies, when mis-identifying or mis-fitting a coherent eddy.

495 It should also be noted that regions with first baroclinic Rossby radius of defor-  
496 mation smaller than 10km cannot be resolved by satellite observations. Thus, the en-  
497 ergy contained by coherent eddies around latitudes of  $60^\circ$  and those near the shore are  
498 missed from this estimate, and their role in the seasonal cycle and local dynamics remains  
499 unknown . New satellite altimeter missions (e.g. Surface Water and Ocean Topography;  
500 SWOT) may allow estimates of the energy contained by mesoscale coherent eddies out-  
501 side the subtropical regions and over the continental slope.

502 Hemisphere-wide variability indicates a strong seasonal cycle of the EKE, CEKE,  
503 and eddy properties. The seasonal cycle of the CEKE in each hemisphere occurs as a  
504 consequence of numerous small coherent eddies interacting with each other (eddy-eddy  
505 interactions) and resulting in stronger, larger and more energetic (but fewer) coherent  
506 eddies during summer, after a few months of the yearly coherent eddy number maxima.  
507 This result reveals eddy-eddy interactions and thus the transfer of energy from smaller  
508 coherent eddies to larger coherent eddies could explain the observed seasonal cycle of CEKE  
509 and coherent eddies properties.

510 Coherent eddy properties reveal a non-uniform long-term readjustment of the mesoscale  
511 eddy field. Overall, the eddy number has decreased globally at a significant rate of  $\sim 35$   
512 eddies per decade from  $\sim 4000$  eddies identified globally on average each day. Despite the  
513 small changes in the total eddy numbers, large proportions of the ocean show a major  
514 strengthening of the mesoscale coherent eddy amplitude at rates greater than  $\sim 1$  cm per  
515 decade. This strengthening of the coherent eddy amplitude is attributed to an intensi-  
516 fication of each coherent eddy polarity, rather than a readjustment of the coherent eddy  
517 field to sea level rise. In other words, the coherent eddy amplitude intensification is oc-  
518 ccurring in both coherent eddy polarities and explains a proportion of the previously ob-  
519 served readjustments in the eddy field to long-term changes in the ocean forcing (Hu et

520 al., 2020; Wunsch, 2020; Martínez-Moreno et al., 2021). This long-term readjustment re-  
521 veals an intensification of the coherent eddy field, possibly due to long-term readjust-  
522 ments in the ocean baroclinic and barotropic instabilities, as well as the strength of the  
523 winds.

524 The reconstruction of the coherent eddies and their statistics has revealed regions  
525 with important coherent eddy contributions and a distinct seasonal evolution of the co-  
526 herent eddies. Western boundary current (WBC) extensions generate eddies through the  
527 instability of the main currents and the seasonal cycle of coherent eddies, CEKE, and  
528 thus EKE could be associated with an inverse energy cascade observable through lagged  
529 seasonal cycles in the coherent eddy statistics. In addition, the amplitude of the seasonal  
530 cycle in WBC extensions is two times larger than any other region, thus the seasonal-  
531 ity of the coherent eddies in WBC extensions dominates the hemispheric seasonal cy-  
532 cle. Furthermore, the seasonal lag of the inverse energy cascade is coupled with the pres-  
533 ence of fronts (Qiu et al., 2014), such as the case for WBC extensions, and our results  
534 are consistent with the notion of baroclinic instability generating eddies and, via eddy-  
535 eddy interactions, a lagged inverse energy cascade.

536 The use of satellite observations in this study limits our ability to quantify the im-  
537 portance of the inverse energy cascade seasonality in the control of the coherent eddy  
538 seasonal cycle. As mentioned above, there is robust evidence of an increase in eddy-eddy  
539 interactions, however we cannot discard important contributions from other processes  
540 such as the seasonal cycle of forcing, stratification, and instabilities, which are crucial  
541 in the generation of coherent eddies. Although this study can provide a descriptive re-  
542 sponse of the coherent eddy field, further work is needed to assess the role of eddy-eddy  
543 interactions in our changing climate, ocean dynamics, and biogeochemical processes. Fur-  
544 thermore, the SWOT mission could allow us to advance our understanding of eddy-eddy  
545 interactions and the seasonal cycle of scales smaller than mesoscale, which may provide  
546 further evidence of the inverse energy cascade driving the coherent eddy seasonality. Cur-  
547 rent generation climate models have just started to resolve mesoscale dynamics, thus,  
548 the presented estimate of energy in coherent eddies from satellite observations could be  
549 used as a benchmark that facilitates the evaluation of such models, and to quantify the  
550 energy contained by mesoscale and more specifically coherent eddies in future climate  
551 projections.

## Acknowledgments

The Chelton & Schlax (2013) dataset was produced by SSALTO/DUACS and distributed by AVISO+ (<https://www.avis0.altimetry.fr/>) with support from CNES, developed and validated in collaboration with E.Mason at IMEDEA. Global coherent eddy reconstruction, coherent and non-coherent eddy kinetic energy datasets, in addition to gridded coherent eddy tracking datasets are publicly available at (<https://doi.org/10.5281/zenodo.4646429>). All analyses and figures in this manuscript are reproducible via Jupyter notebooks and instructions can be found in the Github repository `CEKE.climatology` (<https://github.com/josuemtzmo/CEKE.climatology>). Trends used the Python Package `xarrayMannKendall` (<https://doi.org/10.5281/zenodo.4458776>). J.M.-M. was supported by the Consejo Nacional de Ciencia y Tecnología (CONACYT), Mexico funding. M.H.E. is supported by the Centre for Southern Hemisphere Oceans Research (CSHOR), a joint research centre between Qingdao National Laboratory for Marine Science and Technology (QNLN), Commonwealth Scientific and Industrial Research Organisation (CSIRO), University of New South Wales (UNSW), and the University of Tasmania (UTAS). Analyses were undertaken on the National Computational Infrastructure in Canberra, Australia, which is supported by the Australian Commonwealth Government.

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