

1 **Venusian Habitable Climate Scenarios: Modeling Venus**
2 **through time and applications to slowly rotating**
3 **Venus-Like Exoplanets**

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9 **Key Points:**

- 10 • Venus could have had habitable conditions for nearly 3 billion years.
11 • Surface liquid water is required for any habitable scenario.
12 • Solar insolation through time is not a crucial factor if a carbonate-silicate cycle
13 is in action.

Abstract

One popular view of Venus’ climate history describes a world that has spent much of its life with surface liquid water, plate tectonics, and a stable temperate climate. This leads to the conclusion that Venus’ present state is the result of a runaway greenhouse from a gradually warming sun since its formation. Part of the basis for this optimistic scenario is the high deuterium to hydrogen ratio from the Pioneer Venus mission that was interpreted to imply Venus had a shallow ocean’s worth of water throughout much of its history. Another view is that Venus had a long lived (100 million year) primordial magma ocean with a CO₂ and steam atmosphere because of its closer proximity to the Sun. This would differ from Earth’s evolution since Earth’s magma ocean lifetime was probably around 1 million years. Venus’ long lived steam atmosphere would allow sufficient time to dissociate most of the water vapor, allow significant hydrogen escape and oxidize the magma ocean. Maybe the high deuterium to hydrogen ratio is from this period. However, a dearth of observational data due to a lack of missions to Venus makes it impossible to choose between these two scenarios. Using a general circulation model we demonstrate the viability of the optimistic scenario using the few observational constraints available. The results have implications for what astronomers term “the habitable zone,” and if Venus-like exoplanets exist with clement conditions akin to modern Earth we propose to place them in what we term the “optimistic Venus zone.”

Plain Language Summary

We have little data on our neighbor Venus to help us understand its climate history. Yet Earth and Venus are sister worlds: they initially formed close to one another, and have nearly the same mass and radius. Despite the differences in their current atmospheres and surface temperatures, they likely have similar bulk compositions, making comparison between them extremely valuable for illuminating their distinct climate histories. We analyze our present data on Venus with knowledge about Earth’s climate history to make a number of exciting claims. Evaluating several snapshots in time over the past 4+ billion years, we show that Venus could have sustained liquid water and moderate temperatures for most of this period. Cloud feedbacks from a slowly rotating world with surface liquid water reservoirs were the keys to keeping the planet clement. Contrast this with its current surface temperature of 450 degrees and an atmosphere dominated by Carbon Dioxide and Nitrogen. Our results demonstrate that it was not the gradual warming of the sun over the eons that contributed to Venus present hot-house state. Rather we speculate that Venus’ global resurfacing 100s of millions of years ago played a key role in ending the clement period in its history.

1 Introduction

The case for the early habitability of Venus has been made by a number of authors in recent years (e.g., Grinspoon & Bullock, 2007; Way et al., 2016). If so, then if habitability was widespread and persisted over geological timescales (e.g. 10⁹ years), it is possible that organisms were capable of filling a large variety of environmental niches as occurred on Earth via evolutionary processes. This has led to speculation about possible remaining life-friendly niches (e.g., Morowitz & Sagan, 1967; D. H. Grinspoon, 1997; Cockell, 1999; Schulze-Makuch & Irwin, 2002; Schulze-Makuch et al., 2004; Dartnell et al., 2015; Limaye et al., 2018). These possibilities alone provide sufficient motivation to understand whether early Venus was habitable or not. However, Venus is also interesting from the perspective of the large number of rocky exoplanets discovered to date inside the traditional inner edge of the habitable zone but far enough from their host stars to maintain an atmosphere (Kane et al., 2019). If some of these planets in close proximity to their host stars have long periods of habitability it may overturn traditional notions of the habitable zone (e.g., Kopparapu et al., 2013) and influence target selection

64 for characterization of the atmospheres of these planets. Hence, what appears to be a
 65 modern Venus-like world close to its parent star might host surface liquid water. We re-
 66 fer to such habitable worlds as residing in an “Optimistic Venus Zone,” a subset of the
 67 planets in the “Venus Zone” described in Kane et al. (2014).

68 There are three primary requirements for the early habitability of Venus. The first
 69 is that temperatures were low enough that liquid water was capable of condensing on
 70 the surface of Venus. The second is that Venus had a sufficient inventory of water to cre-
 71 ate the conditions believed necessary for the rise of life on terrestrial worlds (e.g., Brack
 72 et al., 2010). Third, volatile cycling and a geologically active surface must exist to regu-
 73 late the climate as possibly supported by the work of Noack et al. (2012). For surficial
 74 water reservoirs most research has focused on the poorly constrained measurement of
 75 the D/H ratio of $\sim 150 \pm 30$ (times Earth’s value) by the Pioneer Venus Mission¹ (Donahue
 76 et al., 1982, 1997). A number of other ground based measurements have been made (De Bergh
 77 et al., 1991; Marcq, 2006; Bjoraker et al., 1992; Fedorova et al., 2008) that have gener-
 78 ally revealed values of D/H greater than 100. For a recent review of D/H measurements
 79 of Venus see Section 6.3 of Marcq et al. (2017). We review the work of authors who have
 80 considered Venus’ climate evolution below.

81 The first simple gray radiative transfer calculations of Venus’ climate history sug-
 82 gested an early runaway greenhouse effect (e.g., Ingersoll, 1969; Rasool & de Bergh, 1970).
 83 Any water would have remained in vapor form throughout Venus’ early history until its
 84 loss to space via photodissociation processes (Goody & Walker, 1972; Walker, 1975). CO₂
 85 would continue to be outgassed and would accumulate (offset by atmospheric loss over
 86 the last 4.5Gyr) to the values we see today.

87 Pollack (1971) used a similar model but with non-gray radiative transfer. This was
 88 probably the first work to show that Venus could have hosted liquid water on its sur-
 89 face in its early history, but they also demonstrated that a runaway greenhouse was pos-
 90 sible as well. The difference between the two outcomes was *mostly* due to the spec-
 91 ified cloud cover and water vapor content of the atmosphere. With 50% cloud fraction
 92 the result was usually a runaway greenhouse. For an N₂ dominated atmosphere with a
 93 modern Earth water vapor profile, 100% cloud cover, and a 30% less luminous sun (akin
 94 to ~ 4.5 Ga) he achieved surface temperatures ~ 300 K. These purely radiative models did
 95 not include convection, which would limit the water vapor mixing ratio by precipitation
 96 and limit the lapse rate of temperature to the moist adiabatic value (Rampino & Caldeira,
 97 1994) and lessen the strength of the greenhouse effect to some degree. Likewise, they did
 98 not include an atmospheric general circulation that would determine the cloud fraction
 99 self-consistently. Donahue and Pollack (1983) replicated some aspects of the Pollack (1971)
 100 work and again showed that clement conditions were possible in Venus’ early history.

101 Kasting et al. (1984) used a 1-D radiative-convective model to demonstrate that
 102 a planet with an insolation $S_0X=1.45$ times that of modern Earth ($173 \text{ W m}^{-2} \sim 3.8$ Gya
 103 at Venus’ orbit) would have temperatures $\sim 100^\circ \text{ C}$ for a 2 bar N₂ H₂O atmosphere. This
 104 atmosphere contained a wet stratosphere, and thus photodissociation of H₂O and sub-
 105 sequent escape of hydrogen (e.g., Goody & Walker, 1972) could explain the long-term
 106 loss of Venus’ primordial ocean and its present dry conditions. In a subsequent study
 107 Kasting (1988) claimed that clouds would cool the atmosphere enough to keep it in a
 108 moist greenhouse state, rather than the runaway state. At the same time the Kasting
 109 (1988) surface temperature for his “Early Venus” (Figure 7) cloud-free model was ~ 500 K
 110 and hovered just at the margin of the moist vs. runaway greenhouse states. Thus it was
 111 clear early on that maintaining liquid water on ancient Venus required high fractional
 112 cloud cover, but whether physical mechanisms exist to produce these cloudy conditions
 113 was not addressed.

¹ https://nssdc.gsfc.nasa.gov/planetary/pioneer_venus.html

114 Grinspoon and Bullock, in a number of abstracts (e.g., D. H. Grinspoon & Bul-
 115 lock, 2003; D. Grinspoon, 2008), described a Venus climate evolution scenario consistent
 116 with that described later in this paper. This work was never published, though, so the
 117 details of their calculations and the justifications for their conjectures are not documented.

118 The first three-dimensional (3-D) General Circulation Model (GCM) exploration
 119 of issues relevant to ancient Venus was that of Abe et al. (2011). Their study focused
 120 on land planets (analogous to the planet Arrakis of *Dune*) with no permanent surface
 121 water reservoirs, only limited fixed amounts of ground water. In some scenarios they found
 122 that “in principle Venus could have been a habitable land planet as recently as 1 billion
 123 years ago.” Their conclusions resulted from the limited water vapor in the atmosphere
 124 caused by the modest subsurface reservoir, its limited greenhouse effect, and the albedo
 125 of the planet. Thereafter Leconte et al. (2013) used a 3-D GCM to look at climate sce-
 126 narios for Gl581 c and HD 85512 b. HD 85512 b orbits a K-dwarf with a synchronous
 127 rotation period of ~ 58 days and receives approximately 1.86 times the insolation of present
 128 day Earth. This is slightly less than present day Venus’ insolation (1.9), but its host star
 129 is a K-dwarf rather than a G-dwarf and hence the peak of its Planck blackbody spec-
 130 trum is at longer wavelengths. They state, “if not for its thick atmosphere, Venus’ cli-
 131 mate would be very close to the one of HD 85512 b.” They modeled a dry planet akin
 132 to that of Abe et al. (2011) using an N_2 dominated atmosphere with 376ppm CO_2 , and
 133 no water vapor. The model produced large temperate regions for a number of different
 134 atmospheric pressures (See Leconte et al. (2013) Figures 1, 2) for their land planet setup,
 135 but they also found similar behavior in their ‘collapse’ scenario with different atmospheric
 136 water vapor profile amounts.

137 A rationale for high albedo cloud cover on ancient Venus was first presented by Yang
 138 et al. (2014). Yang et al. (2014) was a large parameter study looking at the inner edge
 139 of the habitable zone around solar type stars using a modern GCM² with comprehen-
 140 sive atmospheric physics but a thermodynamic ocean. In essence they stepped the side-
 141 real rotation rate of an Earth-like world from 1 to 256 sidereal days. At the same time
 142 they increased the insolation from that of modern Earth to as much as 2.6 times the mod-
 143 ern Earth for their most slowly rotating world of 256 sidereal days period. One of their
 144 key conclusions was that the slowest rotators would have had a day-night general cir-
 145 culation that would generate an optically thick contiguous cloud bank in the substellar
 146 region. This cloud deck would greatly increase the planetary albedo, keeping the sur-
 147 face temperature moderate even for quite high values of insolation. These results were
 148 later confirmed with a completely different 3-D GCM with a fully-coupled dynamic ocean
 149 (Way et al., 2018). Yang et al. (2014) also included a simulation with modern Venus or-
 150 bital parameters, spin rate, insolation while using modern Earth topography and land/ocean
 151 mask, but again with a thermodynamic ocean. These studies provide a possible ration-
 152 ale for the cloud cover needed to produce temperate surface conditions as first postu-
 153 lated by Pollack (1971). Finally Way et al. (2016, hereafter Paper I) took things one step
 154 further by exploring different topographies, insulations and rotation rates to put tighter
 155 constraints on possible habitable conditions for ancient Venus.

156 In this paper, we extend those parameter studies to consider a wider variety of plan-
 157 ets and attempt to justify our modeling assumptions in light of possible scenarios for Venus’
 158 evolution. Sections 2–5 review existing observational constraints and hypotheses about
 159 the composition, thickness, and evolution of Venus’ atmosphere and water history; its
 160 surface and interior; and its rotation and obliquity. In Section 6 we draw upon this in-
 161 formation to inform a series of 3-D global climate model simulations to illustrate possi-
 162 ble scenarios for an early habitable Venus that transitioned to its current inhospitable
 163 state relatively late in its history. We discuss the implications of our results for the de-
 164 sign of future missions to Venus and for the potential habitability of exoplanets inside

² The National Center for Atmospheric Research (NCAR) Community Atmosphere Model (CAM).

165 the inner edge of the traditional “habitable zone” in Section 7. Finally, recently published
 166 complimentary work by Weller and Kiefer (2019) supports many of our conclusions.

167 2 Atmospheric Composition and Pressure

168 Assuming that Venus and Earth formed from the same parts of the protoplanetary
 169 disk and thus with similar compositions (e.g., Raymond et al., 2004), we are guided by
 170 the history of Earth whose early atmosphere was likely CO₂-rich and possibly cool, but
 171 not frozen due to the faint young sun for the late Hadean and early Archean (4.2-4.0 Ga)
 172 (e.g., Owen et al., 1979; Kasting, 1993; K. Zahnle & Sleep, 2002; J. W. Valley et al., 2002;
 173 K. J. Zahnle, 2006; K. Zahnle et al., 2010; Kunze et al., 2014; Catling & Kasting, 2017;
 174 Krissansen-Totton et al., 2018; Mello & Friaa, 2019; S. J. Mojzsis et al., 2019)³. The pic-
 175 ture of a “Cool Early Earth” promoted in the early 2000s by J. W. Valley et al. (2002);
 176 J. Valley (2005) and with more recent zircon data (Valley et al., 2014) may also be ap-
 177 plicable to Venus’ early evolutionary history if Venus survived its magma ocean phase
 178 with some liquid surface water. Chassefière et al. (2012) was probably the first to make
 179 the comparison of Hadean Earth with Venus, calling his a “cool early Venus” hypoth-
 180 esis. In some ways, the prospects for this cool early Earth and Venus hypothesis have
 181 improved with increasing evidence that the impact of the late heavy bombardment (LHB)
 182 on Earth’s accretional history may have been overestimated. The intensity of the LHB
 183 was imagined so intense that the surface temperatures of Earth would return to those
 184 before the cool early Earth period, perhaps evaporating the oceans and/or raising sur-
 185 face temperatures to values exceeding 100°C. Recent work by S. J. Mojzsis et al. (2019)
 186 has demonstrated that the LHB was either a very minor event or never took place and
 187 that surface temperatures continued their downward trend from the cool early Earth pe-
 188 riod (see S. J. Mojzsis et al. (2019) Figure 1). Figure 4 of S. J. Mojzsis et al. (2019) sug-
 189 gests that Venus accreted less than 0.01 wt% between 4.3 and 4.1Ga, about the same
 190 as Earth. As S. J. Mojzsis et al. (2019) notes, “Results show that an abating impact flux
 191 from late accretion is inadequate to sterilize the surface zone.” Hence, the work of S. J. Mo-
 192 jzsis et al. (2019) demonstrates that the flux from the LHB was originally overestimated
 193 and that the cool early Earth continued through the originally proposed time period of
 194 the LHB from ~ 3.8–4.1Ga. Hence we will begin our most ancient Venus simulations with
 195 CO₂ dominated atmospheres at ~4.2Ga (during the J. W. Valley et al. (2002) cool early
 196 Earth period) that evolves over time to N₂ dominated atmospheres. If the LHB were a
 197 real event we would want to begin our simulations at 3.8Ga, rather than 4.2Ga, as oth-
 198 ers have speculated previously (e.g., Rampino & Caldeira, 1994; Lammer et al., 2018).
 199 Regardless, there are still many unanswered questions regarding the early post-MO his-
 200 tory of the atmospheres of Venus, Earth and Mars. Even though we have far more in-
 201 formation to discern Earth’s early post-MO atmosphere it remains a complicated story
 202 yet to be fully resolved (e.g., Hirschmann, 2012; Nikolaou et al., 2019) and may depend
 203 on atmospheric pressure more than previously assumed (e.g., Gaillard & Scaillet, 2014;
 204 Bower, Dan J. et al., 2019).

205 The carbonate-silicate cycle (e.g., Walker et al., 1981; Stewart et al., 2019) is the
 206 key to keeping most of Earth’s CO₂ locked up in rocks for much of its history rather than
 207 in the atmosphere as on present day Venus. The carbonate-silicate cycle on Earth func-
 208 tions via subductive-type plate tectonics, the presence of a hydrosphere and continen-
 209 tal crust. For Earth there are several lines of evidence to suggest these may go back to
 210 the Hadean (e.g., S. Mojzsis et al., 2001; Hopkins et al., 2008; Harrison, 2009; Korenaga,
 211 2013; Harrison et al., 2017; O’Neill et al., 2017; Rozel et al., 2017; Kusky et al., 2018;
 212 Korenaga, 2018; Maruyama et al., 2018, and references therein). Dehant et al. (2019)
 213 reviews the literature for a later beginning of plate tectonics on Earth (Section 3.2). The
 214 requirements for how plate tectonics begins are still not fully understood and hence re-

³ For Catling and Kasting (2017) see section 11.4

215 mains an active area of research (e.g., Lenardic et al., 2019). If Venus had a similar early
 216 atmospheric and interior evolution to that of Earth then early volatile cycling via some
 217 form of plate tectonics is a viable hypothesis⁴ Hence in this scenario a carbonate silicate
 218 cycle is hypothesized for Venus after the magma ocean phase and before its resurfacing
 219 period to keep CO₂ largely sequestered in crustal carbonates as on modern Earth.

220 We are motivated to choose 1 bar atmospheres for our epochs of interest based on
 221 geological atmospheric pressure proxies for Earth that imply an atmospheric density of
 222 $\sim 0.25\text{--}1.1$ bar for most of the past 4 billion years (Nishizawa et al., 2007; Goldblatt et
 223 al., 2009; D. Som S. Catling et al., 2012; S. Som et al., 2016; Marty et al., 2013; Avicé
 224 et al., 2018). However, some models of Nitrogen cycling imply that the atmospheric pres-
 225 sure could have been higher in the Archean (Johnson & Goldblatt, 2015; Mallik et al.,
 226 2018). Differences in N₂ of factors of a few have relatively small net impacts on climate,
 227 primarily due to small decreases/increases in the pressure broadening of CO₂ lines and
 228 partly offsetting decreases/increases in Rayleigh scattering. Given the likely similar geo-
 229 chemistry and volatile histories of Venus and Earth (e.g., Ringwood & Anderson, 1977;
 230 Lécuyer et al., 2000; Treiman, 2009; Chassefière et al., 2012; Rubie et al., 2015), a sim-
 231 ilar evolution in their early composition and pressure histories is plausible. Although whether
 232 they did indeed start out with similar volatile inventories is still an open research ques-
 233 tion (e.g. Horner et al., 2009).

234 Additional work demonstrates that the present day CO₂ and N₂ inventories of Earth
 235 and Venus may be similar (Donahue & Pollack, 1983; Goldblatt et al., 2009) if Venus
 236 is more degassed than Earth as Donahue and Pollack (1983) speculated. Early work by
 237 Rasool and de Bergh (1970); Kasting (1988) estimated nearly as much CO₂ was locked
 238 up in carbonate rocks on Earth as exist in the atmosphere of Venus. Since there is no
 239 carbonate-silicate cycle active on Venus today most of the CO₂ that would otherwise be
 240 locked up in the interior is in the atmosphere, having degassed over the past several hun-
 241 dred million years (at least). As well, nearly 4 times as much N₂ is found in Venus’ present
 242 day atmosphere as in Earth’s when scaled by planetary mass. Earth’s internal N₂ bud-
 243 get is estimated at $\sim 7 \pm 4$ times that of the atmospheric mass (Johnson & Goldblatt,
 244 2015).

245 However, Argon (Ar) modeling studies (Kaula, 1999; ORourke & Korenaga, 2015)
 246 imply that Venus is less degassed than Earth, although see Donahue and Pollack (1983)
 247 and Halliday (2013) for alternative explanations for the current ⁴⁰Ar abundance differ-
 248 ence between Earth & Venus. Watson et al. (2007) claims that ⁴⁰Ar is not a reliable in-
 249 dicator of degassing, although “the interpretation of their data is controversial” accord-
 250 ing to ORourke and Korenaga (2015). Halliday (2013) also mentions that Venus’ radio-
 251 genic ⁴⁰Ar implies Venus is less degassed than Earth, but he notes that Venus’ atmospheric
 252 non-radiogenic ³⁶Ar has concentrations roughly two orders of magnitude greater than
 253 Earth’s and that “this is hard to explain unless it (Venus) is heavily degassed with a larger
 254 inventory of primordial volatiles.” In further support of a degassed Venus Halliday (2013)
 255 notes that the similar Earth and Venus budgets of C and N support a degassed Venus
 256 while helping to explain the high ²⁰Ne and ³⁶Ar budgets of Venus. One of the largest
 257 stumbling blocks in deciding whether the ⁴⁰Ar modeling inference is correct is a proper
 258 measurement of K/U on Venus, which is presently highly unconstrained. For this rea-
 259 son modeling studies generally use Earth values. Namiki and Solomon (1998) use He to
 260 confirm the ⁴⁰Ar estimates, but they require a Venus in-situ mission to make the appro-
 261 priate Th and U abundance measurements to characterize the geochemistry necessary
 262 for He production. One also needs to consider He escape, an under-explored area of re-

⁴ In fact recent work demonstrates convection regimes like that of plate tectonics in Venus’ recent his-
 tory (Gillmann & Tackley, 2014; Davaille et al., 2017). Dehant et al. (2019) reviews plate tectonic mecha-
 nisms that may not require as much water as previously believed, which may be relevant if Venus retained
 some water but less than Earth.

263 search given that it depends not only on the ability of He to remain charged or neutral,
 264 but also whether Venus had a past magnetic field and how present day measurements
 265 can distinguish in situ He in the upper atmosphere from solar wind deposition.

266 Another limiting factor in comparing Earth and Venus is the lack of good constraints
 267 on the bulk water content of Earth. Some papers refer only to surface water amounts
 268 of 0.023 wt% of the Earth’s total mass, but there are compelling arguments for possi-
 269 bly deep reservoirs within the lower mantle or core (e.g., Raymond et al., 2006; Schmandt
 270 et al., 2014; Genda, 2016; Ikoma et al., 2018). For Venus, interior water amount estimates
 271 are mostly speculation and are restricted to planetary formation studies (e.g., Ikoma et
 272 al., 2018).

273 **3 Venus’ Early Evolution & Evidence for Water**

274 Venus & Earth likely received similar initial water inventories during their forma-
 275 tion histories, as has been shown in a number of works (e.g., Raymond et al., 2006). It
 276 is also tied to their composition in general, as discussed above in Section 2.

277 The Pioneer Venus mass spectrometer measured a very high D/H ratio of 150 ± 30
 278 times that of terrestrial water (Donahue et al., 1997) in a trapped droplet of sulfuric acid.
 279 It is the only such published in-situ measurement. Other non in-situ measurements have
 280 been made as noted in Section 1, and work continues apace (e.g. Tsang et al., 2017). The
 281 original Venus D/H discovery paper by Donahue et al. (1982) was titled “Venus was Wet:
 282 A measurement of the Ratio of Deuterium to Hydrogen.” This was a tantalizing prospect,
 283 but limited by our knowledge of atmospheric escape processes (Donahue & Pollack, 1983;
 284 Donahue & Russell, 1997; Donahue, 1999) and the D/H of delivered materials over the
 285 aeons.

286 It is possible that the D/H ratio is not an indicator of large amounts of water in
 287 Venus’ ancient history. D. H. Grinspoon (1993) pointed out that a short residence time
 288 for water in the present atmosphere of Venus works against the primordial ocean hypoth-
 289 esis. D. H. Grinspoon (1993) also noted that updated theoretical calculations at that time
 290 that implied higher deuterium escape efficiency put constraints on the D/H source wa-
 291 ter of 10-15. That would rule out source material such as meteorites, comets and dust
 292 particles with high D/H ratios (e.g., Irvine et al., 2000; Charnley & Rodgers, 2009). As
 293 well, measurements of D/H and Xenon isotopes in Comet 67P Altwegg et al. (2015); Marty
 294 et al. (2017) imply that Earth’s ocean has a much lower contribution from cometary ob-
 295 jects than previously thought. This would also likely rule out a large cometary contri-
 296 bution to the high D/H ratio measured on Venus.

297 A number of authors have tried to model changes in the D/H fractionation over
 298 time (Gurwell & Yung, 1993; Gurwell, 1995; Hartle et al., 1996) to put some constraints
 299 on when the water was lost. Given the lack of data from Venus it is equally difficult to
 300 constrain or move these models forward.

301 The possibility that the high D/H ratio implies long-lived surface water is also lim-
 302 ited by in-situ measurements. There is some circumstantial evidence of surface water from
 303 surface emissivity observations from the Galileo NIMS instrument (Hashimoto et al., 2008)
 304 and the Venus Express VIRTIS instrument (Mueller et al., 2008). These observations
 305 may imply that the highland “tessera” regions are mostly composed of felsic rocks, and
 306 if they are indeed granitic they would have required surface water to form (Campbell
 307 & Taylor, 1983). M. S. Gilmore et al. (2015); M. Gilmore et al. (2017) find that at least
 308 one tessera region observed with VIRTIS (Alpha Regio) appears to be more felsic than
 309 surrounding plains. This also suggests that these possibly older stratigraphic units (Ivanov
 310 & Basilevsky, 1993; M. S. Gilmore et al., 1997) are granitic crustal remnants, but recent
 311 work by Wroblewski et al. (2019) shows that parts of the Ovda Regio highland tessera

312 are not in fact of granitic origin. However, it is not yet possible to generalize the work
 313 of Wroblewski et al. (2019) to the entirety of tessera.

314 Nikolayeva (1990) and Shellnutt (2019) analyzed surface rock measurements from
 315 Venera 8. As Shellnutt (2019) explains “it is possible that the Venera 8 probe encoun-
 316 tered a fragment of crust that resembles a terrestrial greenstone belt.” Zolotov et al. (1997)
 317 and Johnson and Fegley (2000) have also demonstrated that signatures of water in hy-
 318 drous minerals may persist on the surface of Venus for long periods even after the sur-
 319 face morphology has changed. This motivates an in-situ mission to Venus to search for
 320 such materials, if they exist. Watson et al. (2007) tried to demonstrate that the ^{40}Ar in
 321 Earth’s atmosphere is related to the hydration of the oceanic lithosphere consisting of
 322 relatively Ar-rich olivine and orthopyroxene. If the results from Watson et al. (2007) are
 323 correct, and there is skepticism (Ballentine, 2007), this would lead one to believe that
 324 the ^{40}Ar in Venus’ atmosphere today implies that water oceans could have persisted for
 325 some time.

326 An outstanding unsolved and understudied problem is what happened in the epoch
 327 of Venus’ magma ocean (MO) as it cooled, as this may greatly affect the long-term wa-
 328 ter inventory of the planet. The timescale of the MO crystallization could be of order
 329 a few million years (My) as for Earth (e.g., Katyal et al., 2019; Nikolaou et al., 2019)
 330 or greater than 100 My (Hamano et al., 2013; Lebrun et al., 2013). The longevity of the
 331 MO and associated hot steam and CO_2 atmosphere is vital to understanding the volatile
 332 history of Venus. If the MO and steam atmosphere persist too long then much of the pri-
 333 mordial water inventory of Venus could have been lost in its very early history from a
 334 stronger solar wind (Chassefière, 1997; Lichtenegger et al., 2016). It is possible that the
 335 high D/H ratio we see today (Donahue et al., 1982, 1997) is a relic of that time. If the
 336 MO cooled quickly, then there was an opportunity to build up a surface ocean and at-
 337 mosphere as is believed to have happened in Earth’s early history. The question is whether
 338 Venus’ surface conditions as a result of its closer proximity to the Sun would prevent the
 339 condensation of water on its surface or not. The answer is more complicated than it may
 340 seem since water can condense under hot high pressure multi-bar atmospheres. Matsui
 341 and Abe (1986) allow for temperatures up to 600K, while later work by Liu (2004) al-
 342 low temperatures approaching 720K.

343 Additionally the answer may reside in the planet’s rotation history, what role clouds
 344 played, and the outgassing rates of H_2O and CO_2 . As we will show in Section 5 it is pos-
 345 sible for Venus to reach a tidally locked state in less than a few hundred My using con-
 346 stant phase lag dissipation theory, suggesting that the planet’s rotation rate could have
 347 been slow early on. This affects the clouds and atmospheric dynamics. As shown in pre-
 348 vious work (Way et al., 2016, 2018) as long as a planet is in the slowly rotating regime
 349 (length of day greater than ~ 16 Earth sidereal days) its climate dynamics work to al-
 350 low liquid water to persist on the surface for insulations up to ~ 2.6 times that of present
 351 day Earth.

352 The timing of the MO termination is critical in more than one way. If the steam
 353 and CO_2 atmosphere cooled sufficiently for MO crystallization to occur by the time of
 354 the Late Veneer then even if Venus lost most/all of its primordial H_2O through escape
 355 processes (Gillmann et al., 2009; Hamano et al., 2013; Lichtenegger et al., 2016) there
 356 may have been a second chance to obtain a surface ocean, albeit a shallow one. Recent
 357 work by Greenwood et al. (2018) implies that Earth may have received as much as 30%
 358 of its H_2O inventory in post-accretion impact delivery, consistent with research that shows
 359 that the entire H_2O budget cannot come from the late veneer (Morbidelli & Wood, 2015).
 360 Halliday (2013) concludes that if veneers were common they should be proportional to
 361 planetary mass, and hence Venus would have received a percentage of late veneer H_2O
 362 similar to that of Earth. If Venus was left dry after a long-lived magma ocean phase (Hamano
 363 et al., 2013), then this amount of H_2O veneer also fits within the error bounds of Venus’
 364 measured D/H ratio (Donahue et al., 1982, 1997). It should be noted that the work of

Greenwood et al. (2018) can also fit within the Ruthenium studies of Fischer-Gödde and Kleine (2017). For a contrary point of view see Gillmann et al. (2019), who claim that most of the late veneer impactors would have been Enstatite/ordinary chondrites which are water-poor (also pointed out in Fischer-Gödde and Kleine (2017)), as opposed to water-rich carbonaceous chondrites that would have been a mere 0-2% of the total chondrite delivery.

At the same time, work by Gillmann et al. (2009), Morbidelli et al. (2000), and Raymond et al. (2006) shows that Venus' initial water inventory at formation could be as much as two terrestrial ocean's worth while large planetary embryos could deliver much more within 200 *My* of formation. If true, and if the magma ocean lifetime on Venus was shorter rather than longer, then our estimates of the water content on Venus from Pioneer Venus D/H ratios (Donahue et al., 1982, 1997) should be more toward the higher end, $\sim 16\%$ of a present day Earth's ocean (Donahue et al., 1997).

Another hypothesis (e.g., Rampino & Caldeira, 1994) states that because of its proximity to the Sun, Venus could never condense water on its surface and hence its surface temperature has always been 300K or higher (see Figure 1 in Rampino & Caldeira, 1994) and that most of this water was lost by photodissociation (Goody & Walker, 1972). At the same time the lack of water prevents silicate rock weathering (on Earth this removes CO₂ from the atmosphere), hence the CO₂ builds up in the atmosphere driving temperatures ever higher due to the greenhouse effect as seen today.

4 Surface History, Impactors and Climate Evolution

Understanding the surface history of Venus is crucial to constraining any theory of its long-term climate evolution. Smrekar et al. (2018) reviews the literature on Venus' internal structure and dynamics. In this section we mostly focus on implications for the surface features we see today and how those might be consistent with a hypothesis for the long-term habitability of Venus and a transition to a more recent (~ 1 Gyr) hothouse state.

Up to 80% of the Venus surface has volcanic plains and tectonic structures emplaced over a relatively short geological interval as determined from crater counts (Ivanov & Head, 2013, 2015). The cratering record seen in the plains regions imply surface ages ranging, for example, from ~ 180 Ma (Bottke et al., 2016), to ~ 300 Ma (Strom et al., 1994) to ~ 750 Ma (McKinnon et al., 1997). The relative youth of most of Venus' surface may be the result of a large scale lithospheric overturn event known as the Global Resurfacing Event (GRE), or it may be due to the latest GRE in a long sequence of episodic resurfacing events (e.g., Turcotte, 1993; Strom et al., 1994). For example, Kaula (1999) constructed a simple model with outgassing events staggered at time periods of 4.1, 3.8, 3.5, 3.1, 2.6, 2.1, 1.5 and 0.7 Ga constrained by ⁴⁰Ar measurements. The other hypothesis for the young surface of Venus is from continuous volcanic resurfacing (e.g. Basilevsky et al., 1997; Bjonnes et al., 2012; King, 2018).

The highland tesserae may be one of the keys to understanding this history. They are of particular interest because they may contain information about past crustal differentiation and other processes prior to the loss of any surface water. Some crater age estimates from the Magellan Mission imply that the tesserae are $\sim 40\%$ older than the plains (Ivanov & Basilevsky, 1993; M. S. Gilmore et al., 1997). However, Strom et al. (1994) did not agree with this conclusion. Additional work by V. Hansen and López (2010) points to the possibility that the Ribbon Tessera Terrain are older than the surface units identified with the GRE. Later analysis by Ivanov and Head (2013) implied that tesserae are the oldest stratigraphic unit and that they were created near the beginning of Venus' surface observable history during the "tectonically dominated regime."

414 However, *how* the large basaltic plains were emplaced remains controversial. A num-
 415 ber of authors (e.g., Herrick, 1994; Strom et al., 1994; Basilevsky & Head, 1996) postu-
 416 lated a nearly global ($\sim 80\%$) geologically instantaneous (10-100My) thick ($>1\text{km}$) de-
 417 position of basaltic material from volcanic type outflows (GIBVO) that would have buried
 418 older craters we cannot observe today (akin to the GRE mentioned above). The outflow
 419 depth requirements are determined by the size of the largest impact craters that would
 420 have to be completely covered. However, as Ivanov and Head (2013) point out it is pos-
 421 sible that the cratering record previous to GIBVO could have also been erased in some
 422 manner. The GIBVO model was later augmented and became known as the global stratig-
 423 raphy hypothesis (e.g., Basilevsky & Head, 1996; Basilevsky et al., 1997; Basilevsky &
 424 Head, 1998; J. W. Head & Basilevsky, 1998). Yet another hypothesis to explain the Venus
 425 surface record was initially put forward by Phillips et al. (1992) and is termed the Equi-
 426 librium Resurfacing Model (ERM). In this model the number of craters observed on Venus
 427 today is the result of an equilibrium between constant crater formation (via impacts) and
 428 the removal of such craters via on-going tectonic or volcanic methods. Strom et al. (1994)
 429 did Monte Carlo calculations to demonstrate why the ERM was not feasible, and that
 430 the GIBVO was a better fit to their data. More recent Monte Carlo calculations by Bjonnes
 431 et al. (2012) show that the ERM is able to fit the observations.

432 V. Hansen and Young (2007) strove to demonstrate why none of these hypothe-
 433 ses fit all available observational constraints. V. Hansen and Young (2007) then proposed
 434 what they termed the Spatially Isolated Time-Transgressive Equilibrium Resurfacing (SPIT-
 435 TER) hypothesis to explain more of the observational constraints. It is not clear that
 436 the Venus geological community has settled on any of these hypotheses. Perhaps one of
 437 the largest problems with the global lava hypothesis is the timescale, volume and depth
 438 of the basaltic flows required, none of which have been observed on any present or pre-
 439 viously active volcanic body in the solar system (including in Earth's past). The largest
 440 known outflow to date in Earth's history is the mid-Cretaceous Superplume (Larson, 1991),
 441 which is small by comparison to those envisioned to describe Venus' resurfacing. At the
 442 same time the superplume hypothesis for Venus is compelling as large amounts of CO_2
 443 could have been released at the same time as the plume event (Caldeira & Rampino, 1991).
 444 Large overturn events have been proposed as an explanation for Venus' present surface
 445 state, but in such a scenario it is possible to sequester large amounts of CO_2 in fresh flood
 446 basalt outflows due to enhanced planetary weatherability (e.g., Goddard et al., 2003; Cox
 447 et al., 2016). Large Igneous Provinces (LIPs), on the other hand, can release copious amounts
 448 of CO_2 sequestered in some sedimentary materials (e.g., Ganino & Arndt, 2009) while
 449 avoiding the sequestration issues of a large overturn event. LIPs have been proposed as
 450 an explanation for Venus' present day state as we will discuss below.

451 Previous simulations by Way et al. (2016) showed that Venus could have had tem-
 452 perate conditions for nearly 2 billion years providing it had a shallow ocean, slow rota-
 453 tion rate, and modern orbital elements. Venus might even have experienced more sta-
 454 ble conditions than Earth in its early history since studies by Correia and Laskar (2001)
 455 and J. W. Barnes et al. (2016) have shown that low obliquity states (like that of mod-
 456 ern Venus) may be stable over billions of years and we know that the much shorter Mil-
 457 ankovich cycles have had a strong influence on Earth's climate through time. Deitrick
 458 et al. (2018) reviews the influence of such cycles on the climate of Earth and possible in-
 459 fluences on exoplanets. In addition, Weller et al. (2018) has also shown from geological
 460 models that early Venus could have avoided glaciations more easily than early Earth,
 461 which experienced several partial or total snowball periods in its history. If long-term
 462 stable surface conditions are a requirement for life, Venus might have been more stable
 463 and allowed primitive life to fill more ecological niches more quickly than on Earth. This
 464 gives rise to the possibility that life may still exist in Venus' upper atmosphere (Limaye
 465 et al., 2018).

466 Ernst et al. (2017) speculate that “On Venus, voluminous LIP volcanism produced
 467 high levels of CO₂ that led to run-away greenhouse effect, and high levels of SO₂ that
 468 caused acid rain,” but with little supporting evidence. Bullock and Grinspoon (2001)
 469 present a similar hypothesis that involves outgassing of SO₂ and H₂O that eventually
 470 drive the planet, over 100s of *My*, into a runaway greenhouse state, but do not mention
 471 CO₂. If Venus had LIP volcanism then CO₂ as well as SO₂ can be outgassed if trapped
 472 in sediments in the crust as is seen on Earth (e.g., J. W. Head III & Coffin, 1997; V. L. Hansen,
 473 2007; Ernst et al., 2017; Ernst & Youbi, 2017). Hence if Venus had an earlier epoch of
 474 liquid water habitability then it is logical to assume that CO₂ would have been trapped
 475 in the crust of the planet in the same way it is trapped on Earth today.

476 However, as noted in Macdonald and Wordsworth (2017) when the surface tem-
 477 perature is warmer ($T > 300\text{K}$, see their Figure 2) more water vapor is injected into the
 478 stratosphere, which stabilizes the lapse rate. Such warm climates (as seen in the Venus
 479 models herein) would prevent the largest plumes from injecting SO₂ into the stratosphere,
 480 allowing CO₂ warming without offsetting cooling by H₂SO₄ aerosols.

481 Another well known mechanism to get Venus from a cool clement state to its present
 482 day hot and dry state was proposed by Taylor and Grinspoon (2009) who speculated that
 483 water loss via upper atmospheric dissociation and then hydrogen escape would have even-
 484 tually made the planet dry. Then “With the loss of water, the removal mechanism for
 485 CO₂ would be eliminated, and carbonate rocks on the surface would presumably even-
 486 tually be subducted and lost to thermal decomposition, with the CO₂ being irreversibly
 487 returned to the atmosphere through outgassing.” This model fits in with more recent re-
 488 search by Wordsworth (2016a) who states that the oxygen left over would eventually find
 489 its way to oxidize the mantle and change its redox state, allowing for enhanced nitro-
 490 gen outgassing which is compatible with the nearly 3 bars of N₂ we see in Venus’ atmo-
 491 sphere today. However, an alternative hypothesis is proposed by Gillmann et al. (2009),
 492 who suggest that the oxidation of the mantle occurred in the first 100 *My* of Venus’ his-
 493 tory. They assume the surface was never cool enough to allow liquid water to condense.
 494 The water would again be photodissociated and the hydrogen would have been lost to
 495 space as described above. The leftover oxygen would have dissolved in the magma ocean.

496 Genda and Abe (2005) have proposed that the lack of water on Venus and in Venus’
 497 protoplanetary impactors in its early history (in contrast to that of Earth and its wa-
 498 ter rich impactors) would explain differences in most of the noble gas abundances be-
 499 tween Venus and Earth because oceanic protoplanets would enhance atmospheric loss,
 500 implying that Venus’ original noble gas abundant proto-atmosphere survived to present
 501 day on Venus, unlike that of Earth. A lack of water being detrimental to subductive plate
 502 tectonics (see Section 8). Sakuraba et al. (2019) have also attempted to get the presently
 503 observed nitrogen and noble gas abundances via impact degassing and atmospheric ero-
 504 sion (also see work by, Pham et al., 2011), but unlike Genda and Abe (2005) they be-
 505 lieve late accretion may have further influenced the atmosphere of Venus.

506 More recent work by Gillmann et al. (2016) show that large impactors (400-800km
 507 in diameter) can cause atmospheric erosion and escape and deposit energy in the crust
 508 and mantle. They believe the latter can cause a thermal anomaly in the crust and man-
 509 tle triggering large scale volcanic events at the impact region and the antipode. This in
 510 turn may deplete the upper mantle of volatiles and lead to water loss in the early atmo-
 511 sphere, or conversely provide a volatile heavy atmosphere with extreme temperatures
 512 for billions of years. In a sense this is similar to a theory by Davies (2008) who propose
 513 a mega-collision (akin to that of the Earth’s moon-forming impact) to dry out the in-
 514 terior of the planet. But thus far no large Venus impactor simulations have been utilized
 515 to examine such a scenario, as has been done for Earth’s moon-forming collision (e.g.,
 516 Canup, 2004).

517 To summarize, a number of mechanisms exist by which early Venus could have con-
 518 densed liquid water on its surface. The key ingredient is that it must have been cool enough
 519 for long enough in its early history. As shown by Yang et al. (2014); Way et al. (2016,
 520 2018), the rotation rate of a planet greatly affects its climate dynamics. Specifically, for
 521 very slow rotation a large contiguous water cloud forms at the substellar point, increas-
 522 ing the Bond albedo markedly and keeping surface temperatures moderate for insola-
 523 tion values up to nearly three times that of modern Earth’s 1361 W m^{-2} . In Paper I Way
 524 et al. (2016) we demonstrated that early Venus could have had consistently habitable
 525 conditions throughout its early history if it began with sufficiently slow rotation. In the
 526 next section we review what is understood about the possible evolution of Venus’ spin-
 527 orbit state.

528 In our scenario, early Venus’ has the earliest consistent liquid water habitability
 529 in the solar system followed by Earth and then Mars. This is a broader statement of the
 530 Faint Young Sun Paradox (FYSP), the challenge of explaining how early Earth, not to
 531 mention Mars, could have been warm and wet early in their histories when the Sun was
 532 25-30% dimmer than today (e.g., Feulner, 2012). There is still debate in the ancient Earth
 533 GCM community about the actual composition and thus temperature of early Earth’s
 534 atmosphere given observational proxies for CO_2 that span orders of magnitude, though
 535 models suggest that the range encompasses several viable scenarios (e.g., Charnay et al.,
 536 2013; Wolf & Toon, 2013; Kunze et al., 2014; Le Hir et al., 2014; Charnay et al., 2017;
 537 Krissansen-Totton et al., 2018). These GCM studies and most proxies (e.g., Spencer, 2019)
 538 are from the Archean rather than the late Hadean, but there is some evidence that hab-
 539 itable surface conditions existed well back into the Hadean (e.g., Harrison, 2009; Arndt
 540 & Nisbet, 2012).

541 It is interesting to note that recent atmospheric pressure proxies from the late Archean
 542 imply an atmospheric pressure less than half that of today (D. Som S. Catling et al., 2012;
 543 S. Som et al., 2016). Atmospheres thinner than modern Earth’s are less likely to avoid
 544 snowball conditions, yet the literature above notes that there is geological evidence that
 545 Earth was not in a snowball state during much of the late Archean that the pressure prox-
 546 ies correspond to. Regardless, for this reason we feel it is necessary to explore the pos-
 547 sibilities of lower atmospheric surface pressures in Venus’ climatic history as described
 548 for Simulations 26-30 in Section 6.

549 The FYSP for Mars remains difficult to resolve (e.g., Wordsworth, 2016b). Long-
 550 standing solutions involving large amounts of atmospheric CO_2 are inconsistent with un-
 551 observed carbonate deposits expected from such CO_2 dominated atmospheres (Shaw,
 552 2018) and are insufficient in isolation to produce above-freezing conditions. Other so-
 553 lutions to Mars’ FYSP exist (e.g., Ramirez & Craddock, 2018), although there appears
 554 to be little consensus in the community. In general 3-D GCMs have trouble consistently
 555 sustaining large-area liquid water conditions over millions of years (e.g., Goldblatt et al.,
 556 2009; Kasting, 2010; Kienert et al., 2012; Feulner, 2012; Haqq-Misra et al., 2008) with-
 557 out snowball type conditions.

588 **5 Rotation and Obliquity evolution**

589 To the best of our abilities we would like to constrain the obliquity and rotational
 590 history of Venus to better constrain these important inputs for climate models. This is
 591 limited by the absence of any direct information about Venus’ initial rotation and obliqu-
 592 ity and the fact that impacts likely play a significant role in the early rotational his-
 593 tory of the terrestrial planets (e.g., Lissauer & Kary, 1991; Dones & Tremaine, 1993).
 594 On Earth a variety of means exist to obtain some constraints using dynamical model-
 595 ing combined with geological data when available (e.g., Hays et al., 1976; Park & Her-
 596 bert, 1987; Imbrie et al., 1992; Matthews et al., 1997; Petit et al., 1999; Pälike & Shack-
 597 leton, 2000; Pälike et al., 2004; Olsen et al., 2019) and there has been modest success

568 doing the same for Mars (e.g., Cutts & Lewis, 1982; Laskar et al., 2002, 2004; Byrne, 2009;
 569 Dickson et al., 2015; Bierson et al., 2016). For Earth, an additional constraint is provided
 570 by the Moon, which has predictably affected the evolution of Earth’s rotation and damped
 571 obliquity excursions over its history (K. Zahnle & Walker, 1987; Lissauer & Chambers,
 572 2011). However, until and unless geological observables become available to constrain
 573 dynamical models, only plausible scenarios for the rotational and obliquity history of Venus
 574 can be defined.

575 Hoolst (2015) summarizes much of the literature on the rotational evolution of Venus
 576 throughout its history. We summarize some of the work on this subject below and add
 577 some additional estimates. First we look at the history of studies of the possible spin evo-
 578 lution of Venus.

579 In the 1960-70s several authors investigated the possibility that Venus’ rotation pe-
 580 riod was correlated with its synodic period (Goldreich & Peale, 1966; Gold & Soter, 1969,
 581 1979). Goldreich and Peale (1966) states, “the presence of the Earth may have stabilized
 582 the sidereal rotation period of Venus at the value of 243.16 days retrograde.” An equi-
 583 librium between the atmospheric and body tide of Venus was first proposed by Gold and
 584 Soter (1969) to explain Venus’ non-synchronous rotation period, based on the incorrect
 585 belief at that time that Venus always showed the same face at each inferior conjunction
 586 with Earth as proposed by Goldreich and Peale (1966).

587 The first work to analytically look at Venus’ rotation rate and the role of atmo-
 588 spheric tides was by Ingersoll and Dobrovolskis (1978) who extended the earlier work of
 589 Lord Kelvin Thomson (1882), Chapman and Lindzen (1970) and Munk and MacDon-
 590 ald (1960). They mention that “Venus probably originated with a retrograde rotation
 591 in order to have evolved to the current retrograde state.” In the 1980s this work was fur-
 592 ther extended in a series of papers (Dobrovolskis & Ingersoll, 1980; Dobrovolskis, 1980,
 593 1983). It was clear that Venus’ rotation rate was probably determined by a balance be-
 594 tween the solid body tidal dissipation and the thermal tides of its thick atmosphere with
 595 the sun. Core-mantle friction (CMF) can also play an important role in slowing the spin
 596 rate of Venus, as first explored by Goldreich and Peale (1970). Goldreich and Peale (1970)
 597 were also the first to demonstrate that core-mantle viscous coupling can drive the obliq-
 598 uity to 0° when less than 90° and to 180° if it is greater than 90° over time.

599 This remained the state of understanding of Venus’ rotational history until the early
 600 2000s when the long-term evolution of its spin state of Venus was investigated in a se-
 601 ries of papers by Correia and Laskar (2001); Correia et al. (2003); Correia and Laskar
 602 (2003), who suggested that Venus may have rotated faster in the past, and possibly pro-
 603 grade. It also became clear that at faster spin rates CMF plays an important role in slow-
 604 ing the rotation of the planet, but less so at slower spin rates.

605 Once a planet is spinning more slowly CMF may play an important role in obliq-
 606 uity variations (e.g., Correia et al., 2003). Correia and Laskar (2001) explored a num-
 607 ber of stable obliquity and spin states of Venus while more recent work by J. W. Barnes
 608 et al. (2016) has investigated how stable the obliquity of Venus might be though time.

609 The work on the thermal tides of Venus had led researchers to assume that its ef-
 610 fects would be minor (as it is for Earth) for atmospheres of modest density (e.g. 1 bar).
 611 However, more recent work by Leconte et al. (2015) has demonstrated that thermal tides
 612 arising from even 1 bar atmospheres can be significant depending on the distance to the
 613 host star and the host star’s mass. Leconte et al. (2015) show that even if modern Venus
 614 had a 1 bar atmosphere the tidal torques would still be quite significant.

615 R. Barnes (2017) used an equilibrium tide model with a constant phase lag (CPL)
 616 to find that Earth could have ended up tidally locked today (after 4.5Gyr) had it started
 617 with a rotation rate of 3 Earth Days or longer (the latter more likely if Earth had no satel-

618 lite). We have applied the same Equilibrium Tide Model (EqTide)⁵ from R. Barnes (2017)
 619 to Venus to explore how long it would take Venus to reach a tidally locked state only from
 620 solid body tides. As shown in Figure 1 using CPL theory we find that Venus could have
 621 been tidally locked within 684Myr if it started with a prograde rotation period of 3 Earth
 622 days and zero obliquity. Unfortunately the EqTide model we utilize does not support
 623 retrograde spin states, but we expect the differences to be minor. We will continue to
 624 explore these issues in a future work using the simulator vplanet (R. Barnes et al., 2019)
 625 once this functionality is added. Figure 1 gives further examples for CPL and Constant
 626 Time Lag (CTL) theory results using EqTide. For input parameters we assume that the
 627 tidal dissipation factor $Q=12$ and Love number of degree 2 $k_2=0.3$. These are the same
 628 numbers used for the modern Earth in R. Barnes (2017). Recent work by Henning and
 629 Hurford (2014) demonstrates that our choice for Q may not be unreasonable for Venus.
 630 Henning and Hurford (2014) give estimates of Q for Earth-like planets (see their Fig 15,
 631 top-center-row plot) with orbital periods from 0 to 200 days. Venus' 224 d period is slightly
 632 outside the range they explore (but can be anticipated from the trend visible in their fig-
 633 ure). Our assumption of $Q=12$ is not far off the Henning and Hurford (2014) 'Warm Earth
 634 2' estimate in their Fig 15. As an aside, Q and k_2 are poorly constrained for present-day
 635 Venus. We have even fewer constraints on these values for an ancient Venus, but per-
 636 haps those values would be more Earth-like than present day Venus. For example, present
 637 day Venus' time lag may not be the same as Earth's because of higher internal temper-
 638 atures (MacDonald, 1962; Henning & Hurford, 2014). Historically Goldreich and Soter
 639 (1966) estimated that $Q<17$ for Venus, Lago and Cazenave (1979) had values up to $Q\sim 40$
 640 while Leconte (2018) estimate $Q\sim 100$.

641 More recently Venus' tidal love number was estimated by Konopliv and Yoder (1996)
 642 using Magellan and Pioneer Venus Orbiter data to be $k_2=0.295\pm 0.066$ implying the core
 643 is liquid (Yoder, 1997). Work by Zhang (1992); Xia and Xiao (2002) have estimated $k_2=0.18$
 644 ~ 0.26 . A smaller value ($k_2=0.17$) would imply a solidified iron core which is not con-
 645 sistent with Konopliv and Yoder (1996). Modeling work by Dumoulin et al. (2017) are
 646 consistent with the work of Konopliv and Yoder (1996) as well as our own modeling choices
 647 (discussed above) of $Q=12$ and Love number of degree 2 $k_2=0.299$ (see Table 3 in Dumoulin
 648 et al. (2017)). Regardless, if one uses higher values of Q and/or lower values of k_2 for
 649 ancient Venus it is sufficient to say that equilibrium tide theory predicts that the CPL
 650 and CTL for Venus estimates for tidal locking will be longer than those presented in Fig-
 651 ure 1. The values in Figure 1 then represent *lower limits* to tidal locking for a given start-
 652 ing rotational period. As a caveat there is a debate in the dynamics community about
 653 the appropriateness of the CPL and CTL approaches (Efroimsky & Williams, 2009; Efroim-
 654 sky & Makarov, 2013; Touma & Wisdom, 1994; Greenberg, 2009). so these tidal lock-
 655 ing timescales should be viewed with some caution in the context of the CPL and CTL
 656 models used herein. Of course we do not take into account magnetic braking to see how
 657 the Sun's natural spin-down might affect the tidal evolution of Venus, nor do we assume
 658 that Venus' orbital characteristics would have changed over the timescale of our calcu-
 659 lations, the latter being one of the criticism when applying CPL/CTL to evolving sys-
 660 tems (Efroimsky & Williams, 2009; Efroimsky & Makarov, 2013).

661 Recent work (Green et al., 2019) has investigated the influence of a hypothetical
 662 shallow ocean on Venus (water equivalent layers of ~ 330 meters deep and 830m) using
 663 present day topography and a range of initial rotation periods. The most dissipative sce-
 664 nario predicts a slow down of 72 days per million years. The latter result may be sur-
 665 prising until one recalls that tidal dissipation in Earth's oceans is larger than that of the
 666 Earth's solid body tides (e.g., Munk & MacDonald, 1960).

667 It is likely that Venus was initially a prograde spinning body (like the other 3 ter-
 668 restrial planets in our solar system). From that starting point we find the following per-

⁵ <https://github.com/RoryBarnes/EqTide>

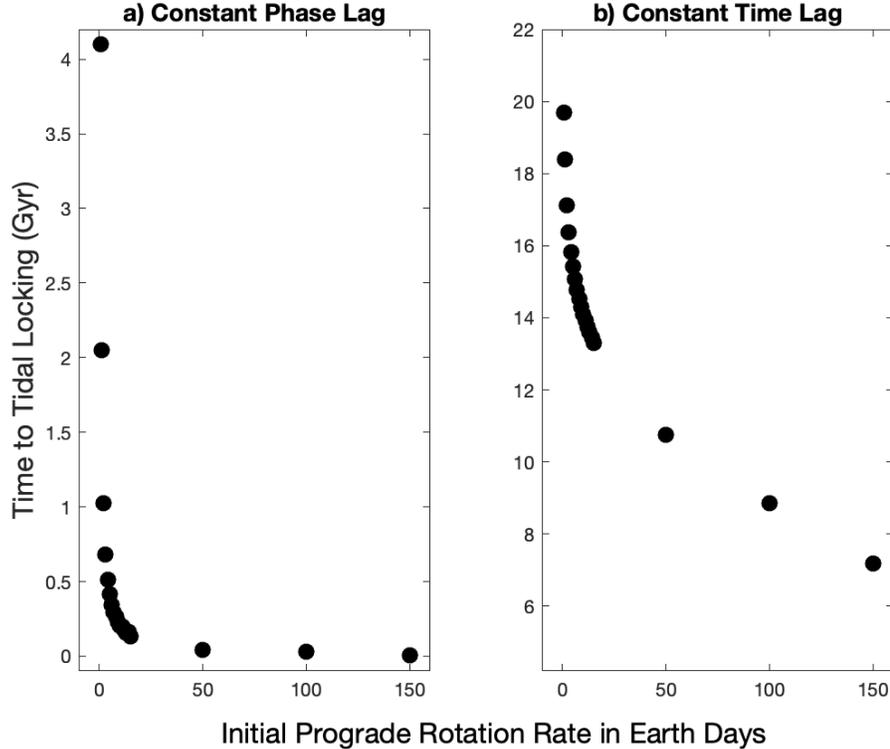


Figure 1. Venus tidal locking timescales using equilibrium tide constant phase (a) and time (b) lag theory. Note that the y-axes have different limits.

669 haps the most compelling answer to Venus’ present day spin state. As shown above there
 670 are models that can drive the planet toward a tidally locked state rather quickly. Once
 671 spinning more slowly, but still prograde, core-mantle friction could cause its spin axis
 672 to orient 180° . Atmospheric tides would prevent the planet from being tidally locked and
 673 that is the state the planet has been in since that time as shown in some of the work of
 674 Correia and Laskar (2001).

675 Yet it has long been speculated that Venus’ current retrograde rotation state is the
 676 result of a large impactor early in its history (McCord, 1968; Singer, 1970; French & Singer,
 677 1971; Counselman, 1973; Burns, 1973; Ward & Reid, 1973; Harris, 1978; Alemi & Steven-
 678 son, 2006; Davies, 2008). The large impactor hypothesis may also explain a possibly very
 679 dry Venus interior, a lack of oxygen in the atmosphere and reconcile the ^{40}Ar results that
 680 imply it is less degassed than Earth (Davies, 2008). However, as pointed out in Ward
 681 and Reid (1973) it is possible for a small impactor (less than 1% of the Moon’s mass)
 682 to drive Venus retrograde if the planet has already spun down considerably due to tidal
 683 dissipation with the Sun (Figure 1).

684 Unfortunately, there is little hope that we will ever truly know the rotation rate
 685 of Venus through time without a way to either measure its “Geological Orrery” as on
 686 Earth (Olsen et al., 2019) and tie that to dynamical models, or find evidence that an im-
 687 pactor played a role in its rotational and hence geochemical evolution as seen in inves-
 688 tigations on Earth related to late accretion and its effect on different isotope abundances
 689 (e.g., Varas-Reus et al., 2019). In fact Brassier et al. (2016) and S. J. Mojzsis et al. (2019)
 690 prefer the hypothesis that the Earth’s Late Veneer was mainly delivered by a single Pluto-
 691 or Ceres-sized impactor. Hence if a larger object was involved in the late evolution of

692 Venus’ spin or obliquity it may be possible to detect its geochemical fingerprints in a fu-
693 ture in-situ mission.

694 For *most* of our climate simulations we assume, rightly or wrongly as discussed above,
695 that Venus has had the same retrograde rotation and its present day obliquity for the
696 past 4.2×10^9 years. We have included a few faster rotation rates that approach the ‘fast
697 rotators’ described in Yang et al. (2014); Way et al. (2018), meaning sidereal day lengths
698 of 16 and 64 times that of present day Earth.

699 **6 Methods**

700 All our simulations use ROCKE-3D (Way et al., 2017) a three-dimensional (3-D)
701 General Circulation Model (GCM) developed at the NASA Goddard Institute for Space
702 Studies (GISS). Most simulations use modern Venus’ current orbital parameters, slow
703 retrograde rotation period (-243 Earth sidereal days in length) and orbital period (224
704 sidereal days). In Paper I (plotted herein with ID = B) we looked at a faster rotation
705 period (16 x modern Earth’s sidereal day length) to see how the planet’s early climate
706 might have responded, and we also look at 16 & 64 day retrograde rotation periods in
707 Venus’ early history in this work. Our focus is on changing insolation, topography, land/sea
708 mask, surface water availability and atmospheric constituents. We motivate our choices
709 below.

710 All simulations discussed below are outlined in Table 1. Our focus is on the reten-
711 tion and stability of surface liquid water on Venus over time. Therefore we simulate four
712 types of planets with surfaces that differ in the amount of water they contain and how
713 that water is allowed to interact with the atmosphere. ROCKE-3D allows for 3 types
714 of surface water: Soil moisture at and beneath the surface with no standing bodies of
715 water; “dynamic” lakes whose depth and area vary with time and that can appear or dis-
716 appear as the competition between precipitation and evaporation dictates; and deeper
717 oceans with permanent boundaries and an effectively infinite source of water for the at-
718 mosphere. All planets with oceans are fully dynamic. For more details on the capabil-
719 ities of such oceans see Way et al. (2017).

- 720 • Arid Venus: This planet has modern Venus topography, but only contains 20cm
721 of water in the subsurface soil layers, soil consisting of 100% sand, and no surface
722 standing water at the start of the simulation. The atmosphere is initialized with
723 zero water vapor and an isothermal temperature profile at 300K. This initial con-
724 dition is similar to that of Kodama et al. (2019); Abe et al. (2011) who attempt
725 to limit the amount of water vapor in the atmosphere (a strong greenhouse gas)
726 and subsequently push the inner edge of the habitable zone farther inward. How-
727 ever, Kodama et al. (2019); Abe et al. (2011) use modern Earth’s rotation rate
728 for all their experiments.
- 729 • 10m-Venus: Uses modern Venus topography and places a 10 meter liquid water-
730 equivalent layer in the lowest lying topographic areas. These are treated by the
731 model as lakes, which have no circulation. The soil is a 50/50 sand/clay mix as
732 used in Yang et al. (2014); Way et al. (2016, 2018).
- 733 • 310m-Venus: Similar to 10m-Venus, except with a 310 meter water equivalent layer
734 again spread in the lowest lying regions. This is the same topography used in Way
735 et al. (2016) simulations A,B and D.
- 736 • 158m-Aqua: This is a simple aquaplanet configuration that is commonly used in
737 the exoplanet community. It uses a fixed 158 meter deep ocean, which corresponds
738 to the bottom of the fifth layer of the ROCKE-3D ocean model. It is a bit shal-
739 lower than the mean depth of the 310m-Venus ocean, and therefore comes into equi-
740 librium a bit faster while still having a similar heat capacity, while including hor-
741 izontal heat transport as well as wind-driven and thermohaline overturning cir-
742 culations.

- 743 • 310m-Earth: Similar to 310m-Venus, but using a modern Earth-like land/sea mask
744 with a 310m deep bathtub dynamic ocean (i.e., every ocean grid cell is of a fixed
745 depth of 310m). We call this an Earth-like land/sea mask since it is not exactly
746 modern Earth, but has some modest changes as shown in Way et al. (2018) Fig-
747 ure 8.

748 The five planets above are then given four types of atmospheres and four differ-
749 ent insulations as described below:

- 750 • Simulations 1-5: These have a 10 bar 100% CO₂ atmosphere using a solar spec-
751 trum and insolation from 4.2Ga from the work of Claire et al. (2012). CO₂ was
752 probably the dominant gas in Earth's early atmospheric evolution (e.g., Kasting,
753 1993). We pick atmospheric pressures of 1 bar (see next bullet point) and 10 bar
754 to cover the Kasting (1993) ranges (see their Fig 2). Although our 10 bar results
755 are not in equilibrium and the temperatures attained are beyond the upper lim-
756 its of our radiation tables, the temperatures reached would still allow for water
757 to condense on the surface as shown in Matsui and Abe (1986); Kasting and Ack-
758 erman (1986). They are also not inconsistent with the 1-D results presented in Kasting
759 and Ackerman (1986). At the very least they provide a future guide. The 10 bar
760 simulations use a modern Venus rotation rate and obliquity.
- 761 • Simulations 6-10: Similar to Simulations 1-5, but these use a 1 bar 97% CO₂ and
762 3% N₂ atmosphere at 4.2Ga.
- 763 • Simulations 11-15: As in Simulations 6-10 but with a rotation period of -16 side-
764 real Earth days to place the planet on the edge of the fast rotator regime as de-
765 scribed in Yang et al. (2014); Way et al. (2018). This allows us to explore the pos-
766 sibility that the planet was rotating more quickly in its early history than today.
767 The choice of a retrograde rotation rate was chosen to be consistent with the present
768 day retrograde rotation, but unpublished simulations with prograde rotation rates
769 with these values produce very similar temperatures. Note that the work of Correia
770 and Laskar (2001, 2003) indicate that prograde rotation rates of 16 days for Venus
771 put its spin axis (obliquity) in a possibly chaotic regime. However, other work by
772 J. W. Barnes et al. (2016) indicate that low obliquity retrograde rotation rates gen-
773 erally have more stable spin axes. Even if there are spin axis variations on geo-
774 logical timescales, it is not possible for us to model those here given that ROCKE-
775 3D simulations are limited to < 10,000 years in length.
- 776 • Simulations 16-20: As in Simulations 6-10 but with a rotation period of -64 side-
777 real Earth days. This allow us to explore the possibility that the planet was ro-
778 tating somewhat more quickly in its early history but still in the slowly rotating
779 dynamical regime. Again, prograde rotation rates were also used in unpublished
780 results and have similar global surface temperature values.
- 781 • Simulations 21-25: These simulations use an atmospheric composition and pres-
782 sure very similar to modern Earth, namely an N₂-dominated atmosphere with 400ppmv
783 CO₂ and 1ppmv CH₄ with a 1013mb surface pressure. They also use a solar spec-
784 trum and insolation at 2.9Ga from Claire et al. (2012). The rotation rate is the
785 same as modern Venus.
- 786 • Simulations 26-30: Similar to Simulations 21-25, but with a lower atmospheric sur-
787 face pressure of 250mb. This is again in the interest of comparative climatology
788 since the Archean atmospheric pressure proxy work of D. Som S. Catling et al.
789 (2012); S. Som et al. (2016) suggests that Earth may have had a surface pressure
790 similar to 250mb at this time.
- 791 • Simulations 31-35: Similar to Simulations 21-25, but now using a solar spectrum
792 and insolation from 0.715Ga from the work of Claire et al. (2012).
- 793 • Simulations 36-40: Again, similar to Simulations 21-25, but now using a modern
794 solar spectrum and insolation.

- 795 • Simulations 41-45: Similar to Simulations 21-25, but now using a modern solar
 796 spectrum, but with insolation set to 1.26 times Venus' present day insolation (2.4
 797 times modern day Earth's insolation) to test the boundaries of the inner edge of
 798 the habitable zone as in Way et al. (2018).

Table 1. Experiments

ID	Topography ^a	Epoch Ga	Insolation ^b S0X/W m ⁻²	P ^c bar	Spin days	N ₂ ppmv	CO ₂ ppmv	CH ₄ ppmv	Soil Type ^d
01	Arid-Venus	4.2	1.396/1913.6	10	-243	0	1000000	0	S
02	10m-Venus	"	"	"	"	"	"	"	S/C
03	310m-Venus	"	"	"	"	"	"	"	S/C
04	158m-Aqua	"	"	"	"	"	"	"	-
05	310m-Earth	"	"	"	"	"	"	"	S/C
06	Arid-Venus	4.2	1.396/1913.6	1	-16	43000	970000	0	S
07	10m-Venus	"	"	"	"	"	"	"	S/C
08	310m-Venus	"	"	"	"	"	"	"	S/C
09	158m-Aqua	"	"	"	"	"	"	"	-
10	310m-Earth	"	"	"	"	"	"	"	S/C
11	Arid-Venus	4.2	1.396/1913.6	1	-64	43000	970000	0	S
12	10m-Venus	"	"	"	"	"	"	"	S/C
13	310m-Venus	"	"	"	"	"	"	"	S/C
14	158m-Aqua	"	"	"	"	"	"	"	-
15	310m-Earth	"	"	"	"	"	"	"	S/C
16	Arid-Venus	4.2	1.396/1913.6	1	-243	43000	970000	0	S
17	10m-Venus	"	"	"	"	"	"	"	S/C
18	310m-Venus	"	"	"	"	"	"	"	S/C
19	158m-Aqua	"	"	"	"	"	"	"	-
20	310m-Earth	"	"	"	"	"	"	"	S/C
21	Arid-Venus	2.9	1.47/2001.0	1	-243	1012599	400	1	S
22	10m-Venus	"	"	"	"	"	"	"	S/C
23	310m-Venus	"	"	"	"	"	"	"	S/C
24	158m-Aqua	"	"	"	"	"	"	"	-
25	310m-Earth	"	"	"	"	"	"	"	S/C
26	Arid-Venus	2.9	1.47/2001.0	0.25	-243	1012599	400	1	S
27	10m-Venus	"	"	"	"	"	"	"	S/C
28	310m-Venus	"	"	"	"	"	"	"	S/C
29	158m-Aqua	"	"	"	"	"	"	"	-
30	310m-Earth	"	"	"	"	"	"	"	S/C
31	Arid-Venus	0.715	1.71/2358.9	1	-243	1012599	400	1	S
32	10m-Venus	"	"	"	"	"	"	"	S/C
33	310m-Venus	"	"	"	"	"	"	"	S/C
34	158m-Aqua	"	"	"	"	"	"	"	-
35	310m-Earth	"	"	"	"	"	"	"	S/C
36	Arid-Venus	0.0	1.9/2601.0	1	-243	1012599	400	1	S
37	10m-Venus	"	"	"	"	"	"	"	S/C
38	310m-Venus	"	"	"	"	"	"	"	S/C
39	158m-Aqua	"	"	"	"	"	"	"	-
40	310m-Earth	"	"	"	"	"	"	"	S/C
41	Arid-Venus	Future	2.4/3266.0	1	-243	1012599	400	1	S
42	10m-Venus	"	"	"	"	"	"	"	S/C
43	310m-Venus	"	"	"	"	"	"	"	S/C
44	158m-Aqua	"	"	"	"	"	"	"	-
45	310m-Earth	"	"	"	"	"	"	"	S/C
D ^d	310m-Venus	2.9	1.47/2001.0	1	-16	1012599	400	1	S/C

^aTopography: Arid-Venus=Only Ground Water, no surficial reservoirs, 20cm water in soil, with modern Venus topography; 10m-Venus=10m Water Equivalent Layer (WEL) spread in lowest elevations as lakes with modern Venus topography; 310m-Venus=310m deep ocean with modern Venus topography; 310m-Earth=Modern Earth-like topography with 310m deep ocean; 158m-Aqua=158m deep aquaplanet.

^bInsolation: S0X = multiple of amount that Earth receives today in insolation (S0=1361 W m⁻¹).

^c Pressure in bar. ^d S=100% Sand, S/C=50/50% Sand/Clay, - = Not Applicable, 100% ocean.

^d Simulation D from Paper 1 (Way et al., 2016). Most similar to ID 8.

799 Most simulations except 158m-Aqua use a fixed ground albedo of 0.2 (thermal con-
 800 ductivity = $0.26 \text{ W m}^{-1} \text{ K}^1$) and 50/50 mix of sand/clay soil following the work of Yang
 801 et al. (2014); Way et al. (2016, 2018). The Arid-Venus simulation uses the same albedo
 802 (0.2), but utilizes a 100% sand soil, rather than the sand/clay mix in other simulations.
 803 The advantage of using sand is that it more quickly loses and absorbs water. This al-
 804 lows the ground hydrology to come into balance more quickly than other soil types. This
 805 is because in the Arid-Venus simulations we are focused on water availability to/from
 806 the atmosphere from/to the soil and hence the amount of total water vapor acting as
 807 a greenhouse gas in the atmosphere.

808 7 Results and Discussion

809 Simulations 1-5, all with a 10 bar pure CO_2 atmosphere and 4.2 Ga insolation, are
 810 uniformly uninhabitable regardless of the surface water reservoir and topography (see
 811 Table 2 in the Supplementary Information). The driest planet (Arid Venus) does reach
 812 equilibrium, but with a surface temperature of 262°C , well above the accuracy limits of
 813 the radiation parameterization used by ROCKE-3D. The other four planets are also well
 814 above 100°C at the point at which they were terminated and are not converging to equi-
 815 librium. Given the greater water reservoirs in these simulations, they are likely to be ap-
 816 proaching a runaway greenhouse state.

817 In Figure 2 we show several different possible evolutionary scenarios for Venus de-
 818 rived from the other experiments in Table 1. In all such scenarios we assume that Venus
 819 had surface liquid water in varying amounts at model start, as described in Section 6.
 820 The colors in this figure differentiate groups of simulations with different insolation, ro-
 821 tation, surface pressure, and/or atmospheric compositions, while the numbers 1-5 and
 822 corresponding symbols for each color delineate the range of climates obtained for differ-
 823 ent surface water reservoir and topography assumptions.

824 7.1 4.2Ga

825 Since there are major uncertainties about what Venus' initial rotation rate was (See
 826 Section 5) we explored early post-magma ocean scenarios at 4.2Ga with three different
 827 retrograde initial rotation periods in the left hand part of Figures 2,3,4: -16 days (gray;
 828 experiments 6-10), -64 days (magenta; experiments 11-15) and -243 days (red; exper-
 829 iments 16-20). Each assumes a 1 bar CO_2 -dominated atmosphere. As one would expect
 830 from the studies of Yang et al. (2014); Way et al. (2018) the faster spin rate simulations
 831 generally have higher temperatures because of the cloud processes discussed in those pa-
 832 pers, but almost all of them reach equilibrium at a habitable global mean surface tem-
 833 perature. However, the clouds also differ to some degree because of water availability.
 834 Contrary to the work of Abe et al. (2011); Kodama et al. (2019) the Arid-Venus cases
 835 all have higher surface temperatures than their counterparts. This is because those pre-
 836 vious works used modern Earth's rotation rate, whereas the cloud processes on these slower
 837 rotating worlds better regulate the climate, more so the more water that is available for
 838 cloud formation. This analysis is backed up by Figure 3 where we plot the shortwave cloud
 839 radiative forcing (SWCRF). The Arid-Venus simulations have the smallest (in magni-
 840 tude) values, because a drier planet has less reflective clouds with less condensed water.
 841 In Figure 4 we show the percentage of high level clouds (PCLDH), the dominant of the
 842 three cloud types (high, medium, low) in Table 2. Here the distinction between the Arid-
 843 Venus simulations and the others is not consistent across the different rotation periods,
 844 suggesting that in some cases middle and/or low level clouds make important contribu-
 845 tions to SWCRF.

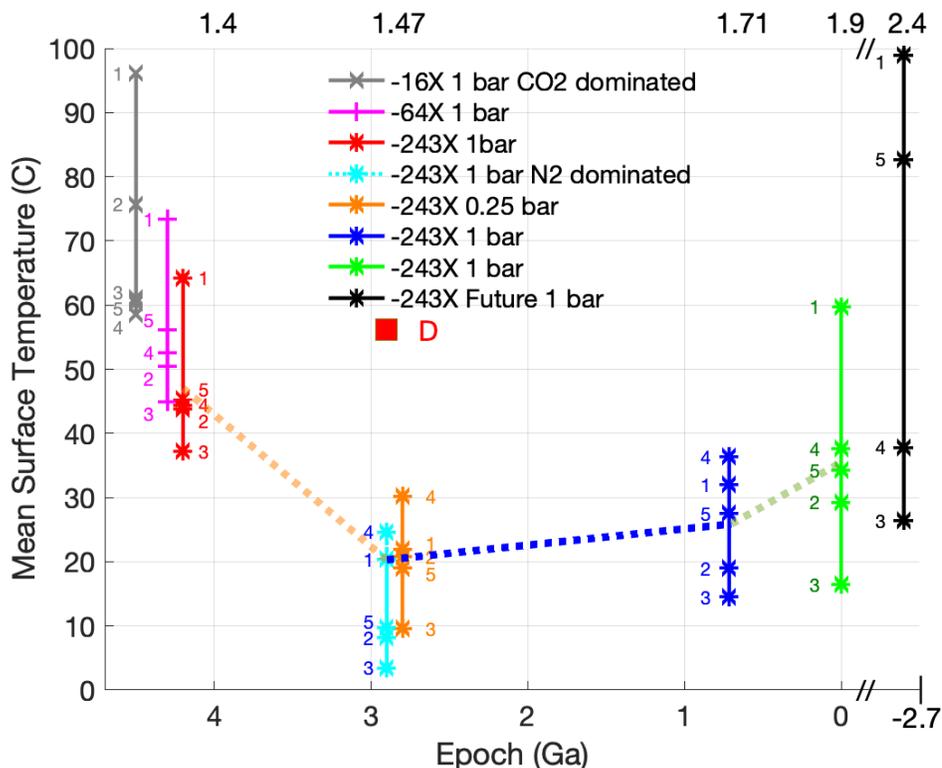


Figure 2. Atmospheric evolutionary scenarios for Venus. Top x-axis is insolation relative to modern Earth ($1.4 = 1.4 \times 1361 \text{ W m}^{-2}$). Note that the gray, magenta and red data in the left-most section of this plot are all for 4.2 Ga simulations. They are separated purely for visual effect and do not reflect differences in epoch or insolation. The same is true for the cyan, and orange data, all of which correspond to 2.9 Ga. The numbers oriented vertically along each set of simulations correspond to the different water reservoir/topography types: 1=Arid-Venus, 2=10m-Venus, 3=310m-Venus 4=158m-Aqua, 5=310m-Earth. See Table 2.

7.2 2.9Ga

Here we plot two different sets of simulations for N_2 -dominated atmospheres: 1 bar (cyan; simulations 21-25) and 250mb (orange; simulations 26-30). In effect these portray representative possible atmospheres for an ancient Venus with liquid water that has evolved from an early CO_2 -dominated atmosphere to a more Earth-like composition via the carbonate-silicate cycle feedback that is believed to regulate CO_2 on planets with liquid water. In both cases the Aqua-158m simulations have the highest mean surface temperatures with the Arid-Venus a close second, but all 10 simulations have moderate surface temperatures fairly similar to modern Earth. However as for simulations 6-20 the shortwave cloud radiative forcing is again the smallest for the Arid-Venus simulations (Figure 3) while also having less high cloud in (Figure 4) than the simulations with more surface water. Unsurprisingly, the thin 250mb atmospheres (simulations 26-30) have cooler surface temperatures in Figure 2. Simulations 27 & 28 have lower mean surface temperatures than modern Earth. The surface temperature field for simulation 28 is plotted in Figure 5 for reference. It exhibits fairly uniformly warm oceans, a signature of slowly rotating planets. Continental temperatures are cooler, slightly below freezing on average, due to night-

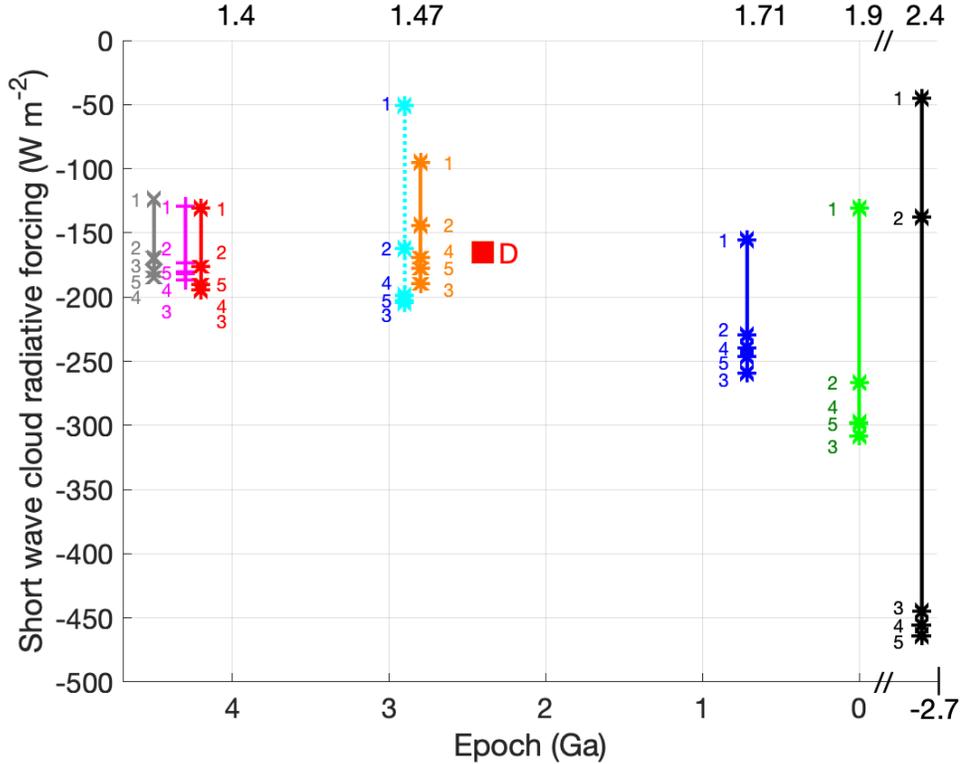


Figure 3. Shortwave cloud radiative forcing for the same evolutionary scenarios, defined as the difference between the solar radiation actually absorbed by the planet and how much would be absorbed if clouds were transparent. This is an estimate of how effective the clouds are at shielding the planet from the star’s intense radiation.

862 time cooling that offsets daytime warming. We also plot Simulation D from Paper I⁶ which
 863 is an N₂-dominated atmosphere, but is otherwise similar to ID 8 (a CO₂ dominated at-
 864 mosphere) in Tables 1 and 2. It has a lower mean surface temperature than ID 8 as ex-
 865 pected, but is significantly higher than the other simulations with larger rotation peri-
 866 ods (ID=21–30) at 2.9Ga. It has less short wave cloud radiative forcing (Figure 3) and
 867 lower percentage of high level clouds (Figure 4) compared to the other 310m-Venus sim-
 868 ulations at 2.9Ga. This is expected given its faster rotation period, stronger Coriolis force,
 869 and less contiguous clouds at the substellar point as discussed in Paper I.

870 7.3 0.715Ga

871 This epoch captures a possible final habitable phase on Venus, if the thick CO₂ at-
 872 mosphere we see today was created by volcanic emissions during the global resurfacing
 873 event(s). The spread in surface temperatures between simulations remains about the same
 874 as in previous epochs, but for all surface types the temperature is warmer than at 2.9
 875 Ga because of the brighter Sun. Again, the Aqua-158m is has the highest surface tem-
 876 peratures with the Arid-Venus close behind as in the 2.9Ga epoch. The SWCRF is some-
 877 what larger in general, due to both the stronger insolation and slightly reduced high cloud,

⁶ The other 3 simulations from Paper I have similar values to their corresponding simulations herein.

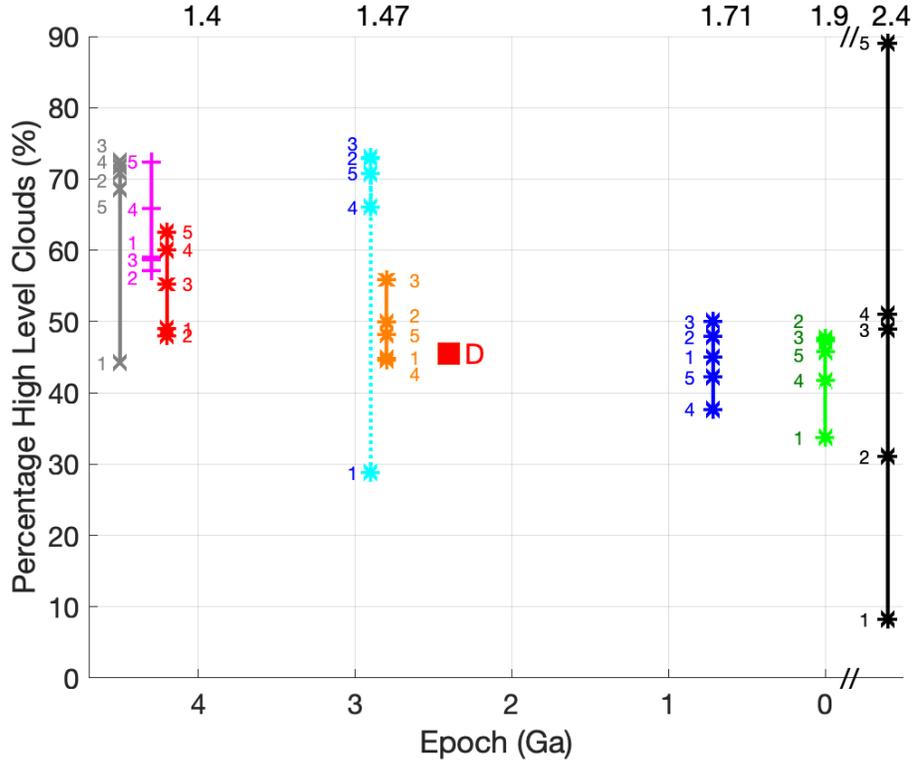


Figure 4. Percentage areal coverage of high level clouds (PCLDH). In Way et al. (2018) it was shown that PCLDH plays a key role in the cloud albedo feedback for slow rotators.

878 but again the Arid-Venus has the smaller effect. The spread in high level clouds has shrunk
 879 considerably from the 1 bar simulations at 2.9Ga.

880 **7.4 Present Day**

881 This suite of simulations at Venus’ present day insolation are designed to demon-
 882 strate that even under today’s Sun the slow rotation cloud feedback effect would have
 883 remained strong as Venus’ atmospheric pressure and composition remain unchanged. This
 884 points to the idea that it was not an increase in insolation that drastically changed Venus’
 885 clement climate of earlier epochs, but rather something else, which we speculate to be
 886 multiple/simultaneous large igneous provinces. The Arid-Venus simulations again have
 887 the highest temperatures and corresponding smallest SWCRF. This is more along lines
 888 of what we saw with the simulations at 4.2Ga, with climate forcing by a stronger Sun
 889 replacing climate forcing by a thicker greenhouse gas atmosphere as the primary reason
 890 for a warm climate.

891 **7.5 Future**

892 Our last set of simulations at insolation values 2.4 times that of present day Earth
 893 are meant to show how long a temperate Venus-like world could have remained habit-
 894 able for a given surface type. Our 10m-Venus simulation is not in equilibrium and so it
 895 is not plotted. The Arid-Venus simulation is at nearly 100° Celsius. It appears to be ap-
 896 proaching radiative equilibrium, but the simulation crashed after 20 years so it is dif-
 897 ficult to be certain. At this point the cloud/albedo feedback for the Arid-Venus case has

Simulation 28: 310m-Venus : 2.9Ga : -243X : 250mb : N2 dominated

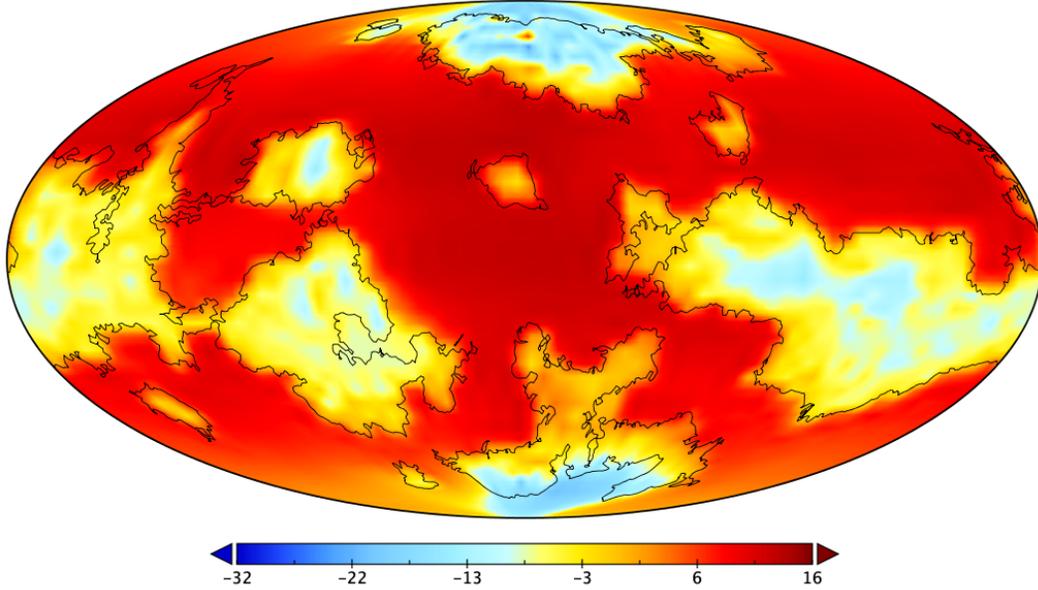


Figure 5. Simulation 28: mean surface temperature over 1/6th of a diurnal cycle.

898 decreased to Earth-like values (the SWCRF is a mere -50 W m^{-2}). This simulation has
 899 the lowest value of PCLDH, which makes it hard to counter the increased insolation at
 900 this time in order to keep mean surface temperatures below the boiling point of water.

901 7.6 General Trends

902 A few relatively consistent trends are apparent from our simulations. First, the Arid-
 903 Venus simulations tend to have the highest surface temperatures, smallest values of SWCRF
 904 and lowest percentages of PCLDH. In many cases the 10m-Venus simulations are next,
 905 but not always. This may point to the fact that even with 10m of available water the
 906 cloud/albedo feedback is generally effective at shielding this slowly rotating world from
 907 the intense solar radiation at all epochs when considering what modern Earth receives.
 908 The simulations with generally higher water availability and similar percentages of land-
 909 to-sea (310m-Venus and 310m-Earth) tend to cluster together in Figures 2,3,4. The 158m-
 910 Aqua simulations seem to float in between, perhaps because of the lower surface albedo
 911 in combination with the cloud albedo feedback.

912 Even without a transition to a runaway greenhouse, prior water loss due to the on-
 913 set of a “moist greenhouse” state may have been important to Venus’ evolution. A num-
 914 ber of simulations in Table 2 contain stratospheric water concentrations (Q_{top}) greater
 915 than $3 \times 10^{-3} \text{ kg kg}^{-1}$, the traditional Kasting et al. (1993) limit for onset of the moist
 916 greenhouse. However, recent work by Chen et al. (2019) for M-star planets has demon-
 917 strated that previous work may have overestimated water loss rates. Hence we should
 918 exercise more caution in using $3 \times 10^{-3} \text{ kg kg}^{-1}$ as a hard value for the moist greenhouse
 919 until similar models are applied to G-star planets. Column Q_{surf} is a check on the amount
 920 of water vapor at the surface of the model. ROCKE-3D runs with a fixed molecular mass
 921 at model start and ignores the spatially/temporally variable mass of water in calculat-
 922 ing pressure gradients, so it is important to keep track of whether water becomes a non-

negligible fraction of atmospheric mass (e.g., 20% of the total or more) as the dynamics in the model will begin to be outside an acceptable range. Only in simulation 45 does this value go over the 20% limit.

Figure 6 shows calculations for what exoplanet astronomers might find for a population of “exo-Venuses,” some of them habitable and some not, in future observations. We use an ensemble of ROCKE-3D simulations of a variety of rocky planet types from which predictors for Bond albedo and surface temperature have been derived using insolation and star temperature as inputs (Del Genio et al., 2019, hereafter DG19). Figure 6 applies the predictor to our Venus evolutionary scenarios to determine the predictability of albedo and surface temperature. In general Bond albedo increases with insolation in the Venus simulations (upper panel), the exceptions being 4 of the Arid-Venus cases with limited surface water, fewer clouds and thus lower albedos than our other simulations. The DG19 predictor (middle panel) works well for all but these 4 cases, since it predicts a high cloud-controlled albedo for the wetter planets that have more and/or thicker clouds than others. For surface temperature (bottom panel), the predictor tends to underestimate the actual temperatures by roughly 20° or less in most cases, but by up to 50° for the hotter, drier, marginally habitable Arid-Venus cases.

Figure 7 shows (left panels) the vertically integrated soil moisture and (right panels) lake fraction for two of our Arid-Venus simulations at 2.9 Ga with Earth-like atmospheres: Experiment 21 (1 bar, top panels) and Experiment 26 (250 mb, bottom panels). The Arid-Venus cases are of particular interest because they are initialized with a spatially uniform subsurface soil water reservoir and no standing water bodies, and they then equilibrate to a heterogeneous distribution of surface and subsurface water depending on the climate and circulation and thus the local precipitation-evaporation competition. The 1 bar atmosphere is typical of the behavior of most of the experiments in Table 1: In equilibrium, subsurface water collects primarily in the equatorial region where rising motion and precipitation is prevalent during the daytime, and especially in highland regions such as Aphrodite Terra. Lakes (which are not present in the initial condition) also form preferentially over the equatorial highlands. The 250 mbar simulation is an outlier, with subsurface water and lakes arising primarily at high latitudes. This is reminiscent of what is observed for the methane-ethane lakes on Titan, which Mitchell (2008) is able to reproduce in a GCM when a limited subsurface methane reservoir is assumed.

Figure 8 shows the same quantities for the analogous 10-m Venus simulations (Experiments 22 and 27). These differ from the Arid-Venus cases not only because the planet contains more water, much of it in surface lakes rather than subsurface soil moisture, but also because the lakes fill the lowlands at the start of each simulation rather than being distributed uniformly across the planet. With a larger water reservoir than that for the Arid-Venus planets, soil moisture collects throughout the tropics in the equilibrated climate, but still with a slight preference for the highland regions even though there is no standing water in the highlands in the initial condition. But unlike the Arid-Venus planets, soil moisture also collects at the poles, both for the 1 bar and the 250 mb atmosphere planets. Likewise, lakes in both simulations form in both the tropics and polar region, more so over the highlands for the 1 bar atmosphere but fairly uniformly distributed in longitude for the 250 mb atmosphere.

8 Conclusion

Whether Venus’ original water survived its initial magma ocean stage, or whether significant water was delivered afterwards, is unknown. It is therefore worth having a theoretical framework that considers the possibility of an early habitable Venus as a starting point for designing future observing strategies that might shed light on Venus’ past. In this spirit, we envision the following possible climatic evolution for Venus:

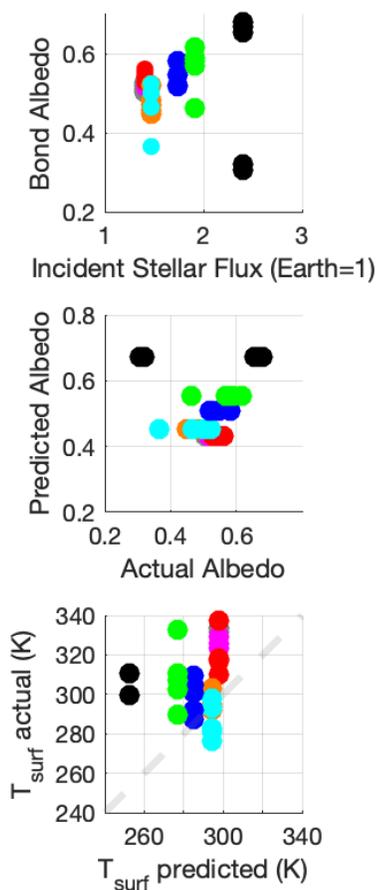


Figure 6. Upper panel: Bond albedo vs. incident solar flux for the planets in Figure 2. Middle panel: Predicted versus actual Bond albedos for the same set of planets after Del Genio et al. (2019). Lower panel: Predicted vs. observed surface temperatures after Del Genio et al. (2019).

974 1.) Toward the end of the accretion period (~ 4.2 Ga) Venus would have cooled rapidly
 975 as did Earth, as shown in the work of J. W. Valley et al. (2002). This would allow sur-
 976 face water to condense and early oceans to form. Early oceans in turn could create sig-
 977 nificant tidal dissipation that would spin down Venus' rotation rate on a relatively short
 978 time scale. Solid body dissipation may have also been effective, see Section 5. Slow ro-
 979 tation combined with an early ocean would provide the necessary ingredients for a day-
 980 side cloud deck to emerge, shielding the planet from high insolation and allowing at least
 981 some of the initial surface water to survive despite the planet being well inside the con-
 982 ventional inner edge of the habitable zone, as shown in the work of Yang et al. (2014);
 983 Way et al. (2018).

984 2.) The carbonate-silicate cycle in concert with interior volatile cycling would allow
 985 low CO_2 draw-down while N_2 was outgassed, eventually reaching a balance producing
 986 an N_2 -dominated atmosphere with trace amounts of CO_2 over gigayears with pressures
 987 ranging from several bars to hundreds of millibars.

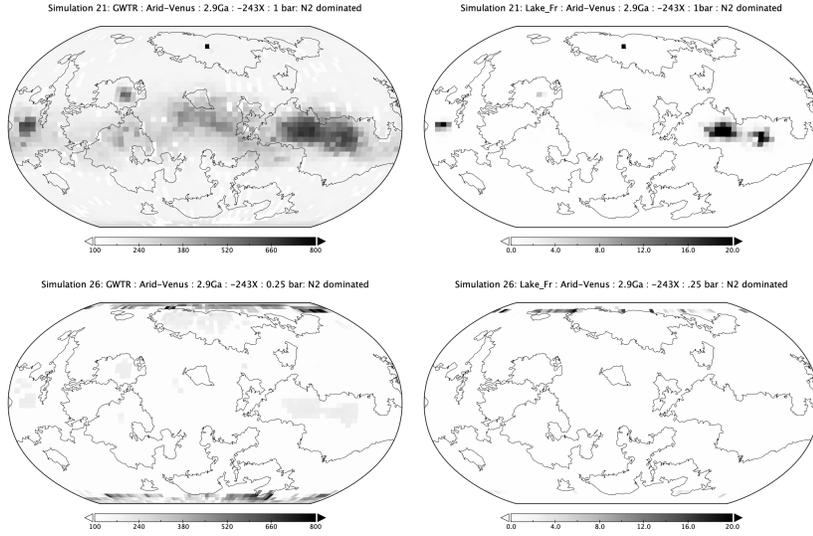


Figure 7. GWTR left panels, LakeFR right panels for Top to Bottom: Simulation 21 and 26 (Arid-Venus).

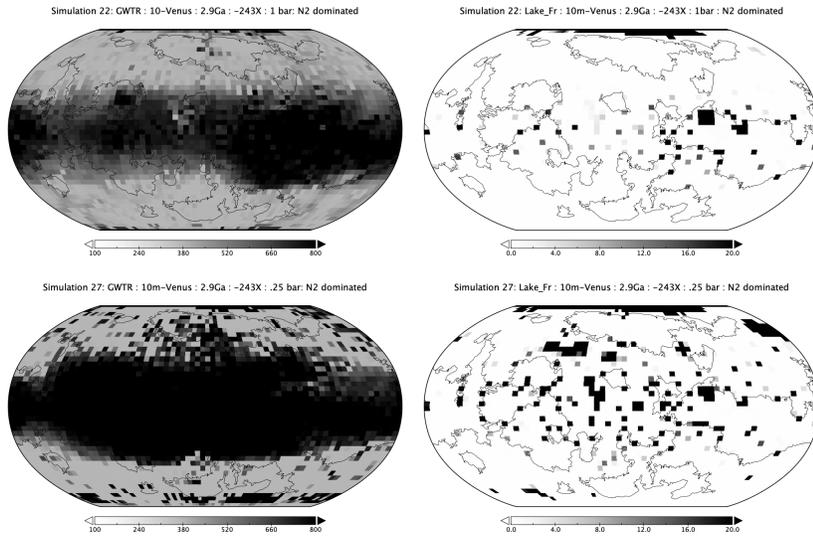


Figure 8. GWTR left panels, LakeFR right panels for Top to Bottom: Simulation 22 and 27 (10m-Venus).

988 3.) We propose that any stable Venusian climate period came to an end at around
 989 700Ma, as evidenced by the age of the surface from crater counts (McKinnon et al., 1997).
 990 We suggest that the ignition of multiple large igneous provinces (LIPs) became active
 991 at the same geological time. It is likely that the ignition of LIPs throughout Earth's his-
 992 tory (Ernst, 2014; Ernst et al., 2019) is a random stochastic process. This may imply
 993 that multiple large scale LIPs have not occurred simultaneously on Earth by purely ran-
 994 dom chance, which is fortuitous for life as we know it today. Venus may not have been
 995 as fortunate. Unfortunately, little is known about Venus' interior structure today, much
 996 less its initial state and subsequent evolution, so the question remains as to whether de-

997 terministic evolutionary processes in the interiors of Venus-like planets will inevitably
 998 lead to catastrophic changes and uninhabitable end states.

999 4.) Multiple large scale LIPs would have warmed the world markedly via the re-
 1000 lease of large amounts of CO₂ via one or more mechanisms and its greenhouse effect on
 1001 the climate (e.g., Ogden & Sleep, 2012).

1002 5.) This warming could have hastened a movement the onset of a moist (Kasting
 1003 et al., 1993) and possibly a runaway greenhouse state in which much of the surface wa-
 1004 ter was lost via photodissociation accompanied by hydrogen escape and an oxidized
 1005 surface. Recall that present day Venus D/H measurements (if they are accurate) imply
 1006 .6 to 16% of Earth’s present day surficial water stores (Donahue et al., 1997). Hence the
 1007 timescale of total loss would be at least an order of magnitude faster than that described
 1008 in Kasting et al. (1993), although there is still uncertainty about the stratospheric con-
 1009 ditions under which a moist greenhouse occurs (e.g., Chen et al., 2019) and when a run-
 1010 away greenhouse is achieved (e.g., Kasting & Ackerman, 1986; Goldblatt et al., 2013; Ramirez
 1011 et al., 2014).

1012 6.) The loss of water would in turn change the planet from an initial subductive
 1013 or mobile plate tectonic mode to more of a stagnant lid mode (as on present day Venus
 1014 & Mars) since it is currently believed that water plays a key role in plate tectonics on
 1015 Earth (e.g., Grove et al., 2012; Lécuyer, 2014). This scenario fits in very nicely with the
 1016 recent work of (Weller & Kiefer, 2019) who give a timescale of order 1Gyr for the tran-
 1017 sition from a mobile to a stagnant lid mode on Venus. Without a mechanism to efficiently
 1018 cycle volatiles in a stagnant lid mode (e.g., Tosi et al., 2017; Höning et al., 2019), out-
 1019 gassing would have continued without the major weathering and subduction surface sinks
 1020 that operate on Earth, hence CO₂ and N₂ would build up over time to reach the levels
 1021 we see on Venus today. Some studies have also shown that even in a stagnant lid mode
 1022 it is possible to cycle volatiles, possibly up to gigayears in time (e.g., Foley & Smye, 2018;
 1023 Godolt et al., 2019), but these mechanisms depend on the initial CO₂ budget and the
 1024 retention of at least some water after cooldown.

1025 7.) This stagnant lid mode may then allow very large mantle upwelling and/or down-
 1026 welling centers that would produce some of the features we see on Venus’ surface today
 1027 over hundreds of millions of years, as described most recently in the work of e.g Rolf et
 1028 al. (2018); Weller and Kiefer (2019).

1029 Clearly a great deal more modeling work and more observations are required to con-
 1030 firm or refute this hypothesis. Did Venus follow the ‘canonical’ path with Earth-like con-
 1031 ditions throughout most of its history, with its demise and present hothouse state the
 1032 consequence of a series of LIP events releasing CO₂ bound up in the crust as on Earth,
 1033 and/or released from the deep interior where CO₂ is more easily sequestered (e.g., Ku-
 1034 ramamoto & Matsui, 1996)? Or did it become bone dry in an extended magma ocean phase
 1035 in the first 100Myr, as described in Hamano et al. (2013) for Type II planets?

1036 We will not be able to distinguish between these two scenarios until we return to
 1037 Venus to make proper noble gas and other elemental and isotopic measurements at the
 1038 surface (Baines et al., 2013) and better constrain escape processes at the top of the at-
 1039 mosphere through time. The latter will also rely upon how such gases escape from present
 1040 day Earth given the possibility that Venus may have had a magnetic field in previous
 1041 epochs, even if it is not clear how important the magnetic field is to escape processes in
 1042 general (Gunell, Herbert et al., 2018). Likewise, whether the actual evolution of the one
 1043 Venus we can visit is the ultimate fate of all highly irradiated rocky planets or an ac-
 1044 cident of an evolutionary path that might have proceeded differently in other circum-
 1045 stances (Lenardic et al., 2016) is not known. The stakes are high for answering this ques-
 1046 tion, since many exoplanets have been discovered in the “Venus zone” just inside the tra-
 1047 ditional inner edge of the habitable zones of other stars (Kane et al., 2014). Efforts to

1048 simultaneously characterize the CO₂ concentrations and climates of a number of these
 1049 exoplanets, combined with a focused observational strategy for unveiling the history of
 1050 the “exoplanet next door” to Earth in our own solar system (Kane et al., 2019), will be
 1051 our best chance to understand whether the envelope for habitability and the emergence
 1052 of life is much broader than usually assumed.

Table 2. Results

ID ^a	Runtime years	Temp ^b C	Balance ^c W m ⁻²	Q _{top} ^e kg/kg	Q _{surf} ^d %	Albedo Planetary	Albedo Surface	Clouds High	Clouds Medium	Clouds Low
01	71	262	0.01	1.36e-05	0.08	54	29	54	5	0
02	111	151	27.21	2.67e-05	3.74	54	18	33	17	33
03	67	121	21.31	4.50e-04	7.87	52	11	45	23	37
04	52	120	51.28	1.74e-04	10.01	44	7	36	17	51
05	56	123	36.27	4.28e-04	9.16	49	10	46	24	44
06/01	100	96	0.06	1.73e-03	1.65	51	29	44	2	0
07/02	100	76	0.41	9.34e-03	7.70	53	16	71	32	21
08/03	53	61	5.92	6.49e-03	7.50	50	10	72	32	47
09/04	100	59	-0.48	6.23e-03	6.95	52	7	73	35	62
10/05	60	60	1.90	5.46e-03	7.19	52	9	69	37	53
11/01	200	73	0.05	9.03e-04	1.70	52	29	59	4	1
12/02	500	50	0.04	5.94e-04	2.07	55	16	57	10	10
13/03	200	45	0.87	1.56e-03	3.15	54	10	59	18	43
14/04	156	53	0.11	3.47e-03	4.81	51	7	66	25	60
15/05	200	56	0.44	4.58e-03	5.65	52	9	72	28	49
16/01	300	64	0.01	8.34e-04	0.80	53	28	49	4	1
17/02	500	44	0.06	5.68e-04	1.40	56	16	48	9	8
18/03	1000	37	0.13	8.18e-04	2.07	56	10	55	15	38
19/04	500	44	-0.50	1.59e-03	3.01	54	7	60	17	56
20/05	500	45	0.23	1.88e-03	3.15	54	9	63	19	44
21/01	500	22	0.02	2.52e-05	0.36	46	28	45	6	4
22/02	3000	21	-0.04	3.08e-05	0.61	47	17	50	10	12
23/03	1000	10	-0.05	2.53e-06	0.70	52	10	56	20	48
24/04	500	30	-0.12	3.51e-05	2.32	45	6	45	14	37
25/05	1000	19	0.05	5.87e-06	1.24	48	9	48	22	39
26/01	300	20	0.11	5.94e-03	0.53	37	29	29	5	9
27/02	1000	8	-0.14	9.98e-03	1.26	50	20	73	17	27
28/03	1000	3	-0.02	6.15e-03	1.94	52	13	73	8	13
29/04	1000	25	0.34	2.69e-02	6.39	47	6	66	1	0
30/05	500	10	0.08	8.58e-03	2.76	50	10	71	7	10
31/01	200	32	-0.03	1.12e-04	0.63	52	28	45	7	4
32/02	297	19	0.02	8.32e-05	0.87	2	1	48	20	26
33/03	1000	15	-0.03	1.18e-05	0.95	58	10	50	28	47
34/04	500	36	-0.03	6.41e-04	3.32	52	7	38	17	42
35/05	500	28	0.27	1.11e-04	2.04	55	9	42	26	40
36/01	50	60	-0.09	3.05e-03	0.58	46	29	34	2	1
37/02	2000	29	-0.01	4.33e-04	1.17	59	17	48	16	23
38/03	1000	16	0.12	3.52e-05	1.08	62	10	47	33	47
39/04	500	38	-0.07	1.37e-03	3.57	57	7	42	21	42
40/05	1000	34	0.35	1.10e-03	2.92	58	9	46	24	41
41/01	20	99	0.52	3.13e-03	0.43	32	29	8	0	0
42/02	8	288	90.21	1.00e-01	11.71	31	17	31	0	0
43/03	1000	26	0.01	3.35e-03	1.99	68	11	49	33	54
44/04	500	38	0.23	5.52e-03	3.75	67	8	51	44	54
45/05	35	83	32.64	1.05e-01	30.36	66	10	89	19	16
D	2000	56	0.19	5.31e-03	8.97	44	10	46	39	45

^aID: The colored numbers correspond to those in Figure 2.

^bTemp: Surface temperature in Celsius.

^cBalance: These numbers come from 50 year averages unless the Runtime (Column 2) is less than 150 years, then the average is 10 years.

^dSpecific humidity in top layer of the atmosphere.

^eSurface humidity as percentage of atmosphere.

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Appendix A Energy balance and temperature

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Different simulations reach radiative balance sooner than others, while some never reach it at all. Herein we plot the energy balance (in units of W m^{-2}) and surface temperature (in Celsius) as a function of simulation year. This should allow the reader to have a better grasp of which simulations are appropriate for a given interest.



Figure A1. Energy Balance (also called Net Radiative Balance) in the left columns and surface temperature in the right columns as a function of simulation year for all simulations in this study. Note that not all limits on the x or y axes are the same.

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Acknowledgments

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https://portal.nccs.nasa.gov/GISS_modelE/ROCKE-3D/publication-supplements/

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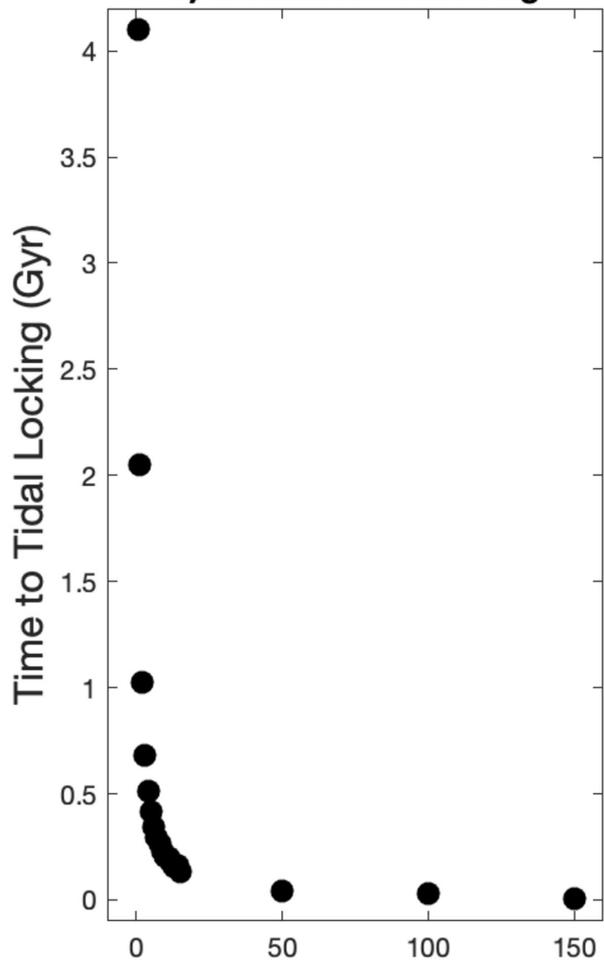
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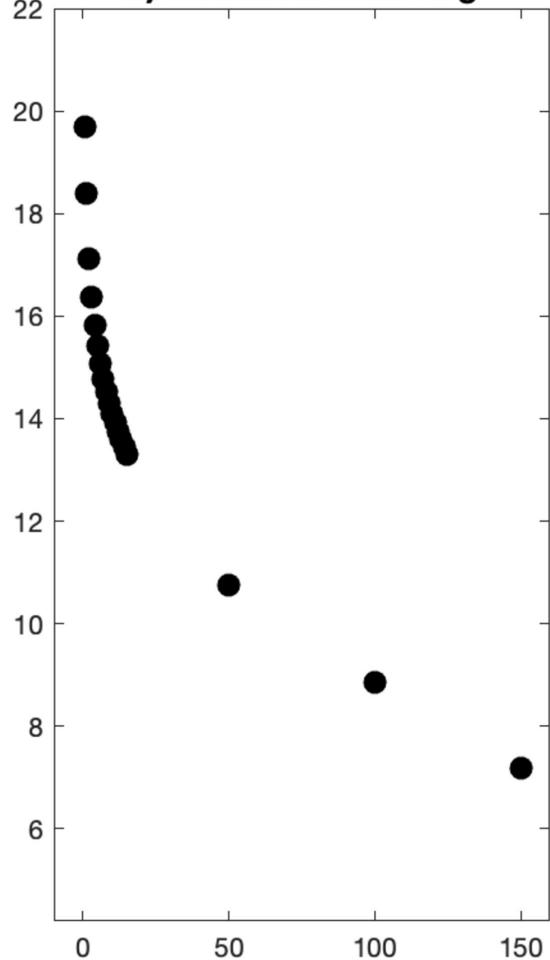
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Figure 1.

a) Constant Phase Lag



b) Constant Time Lag



Initial Prograde Rotation Rate in Earth Days

Figure 2.

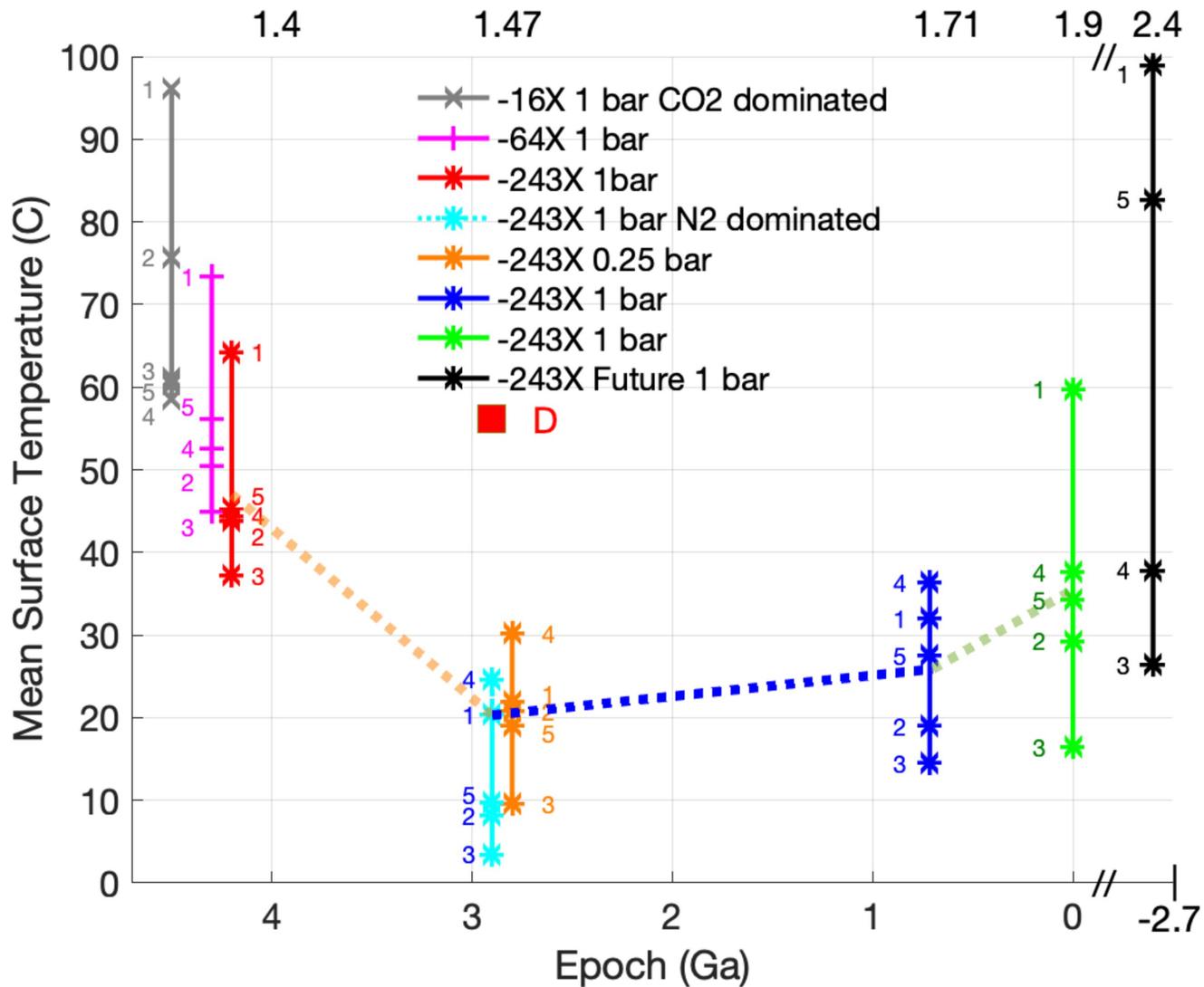


Figure 3.

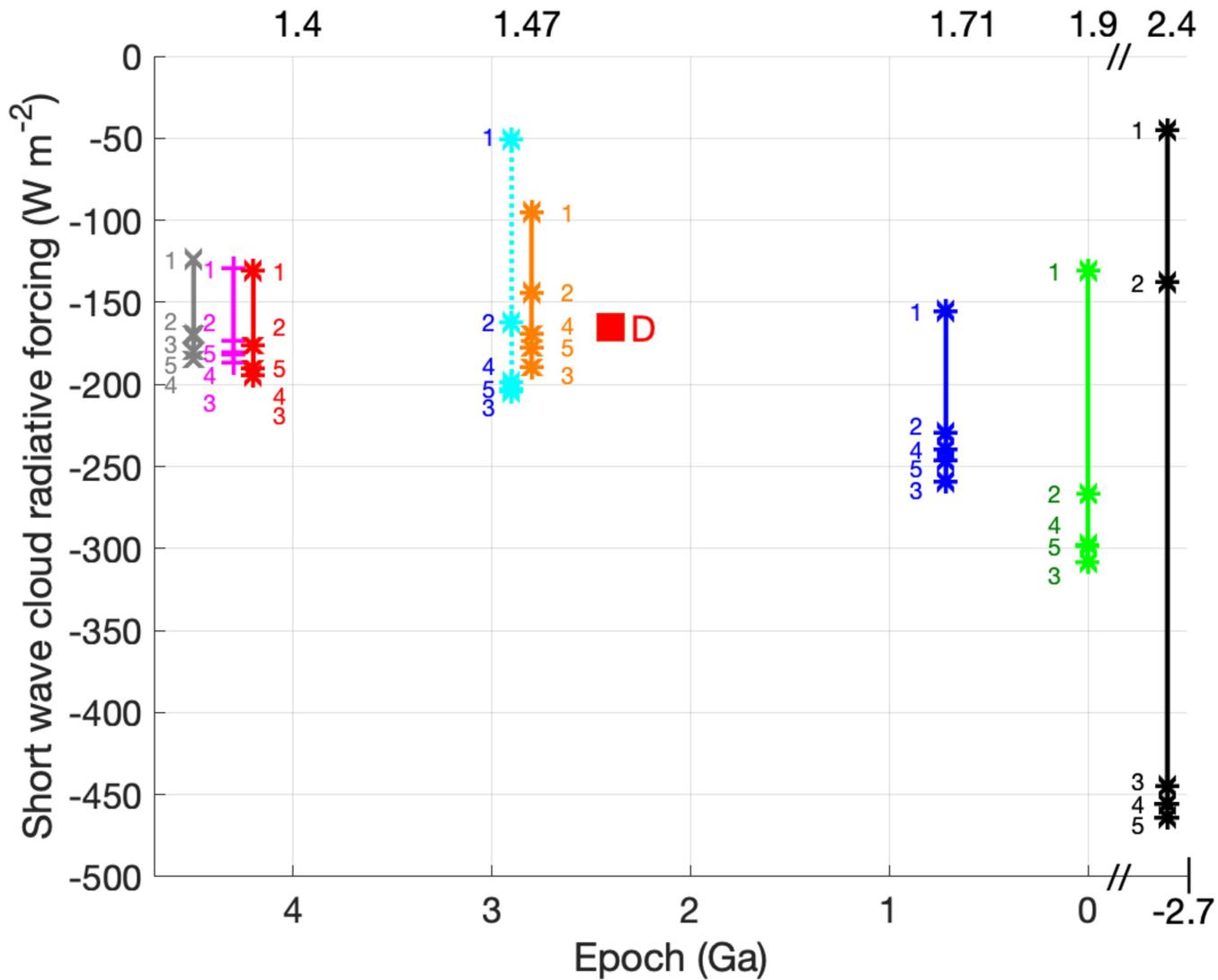


Figure 4.

Figure 5.

Simulation 28: 310m-Venus : 2.9Ga : -243X : 250mb : N2 dominated

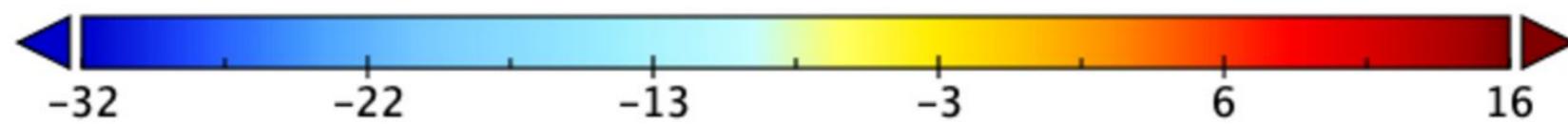
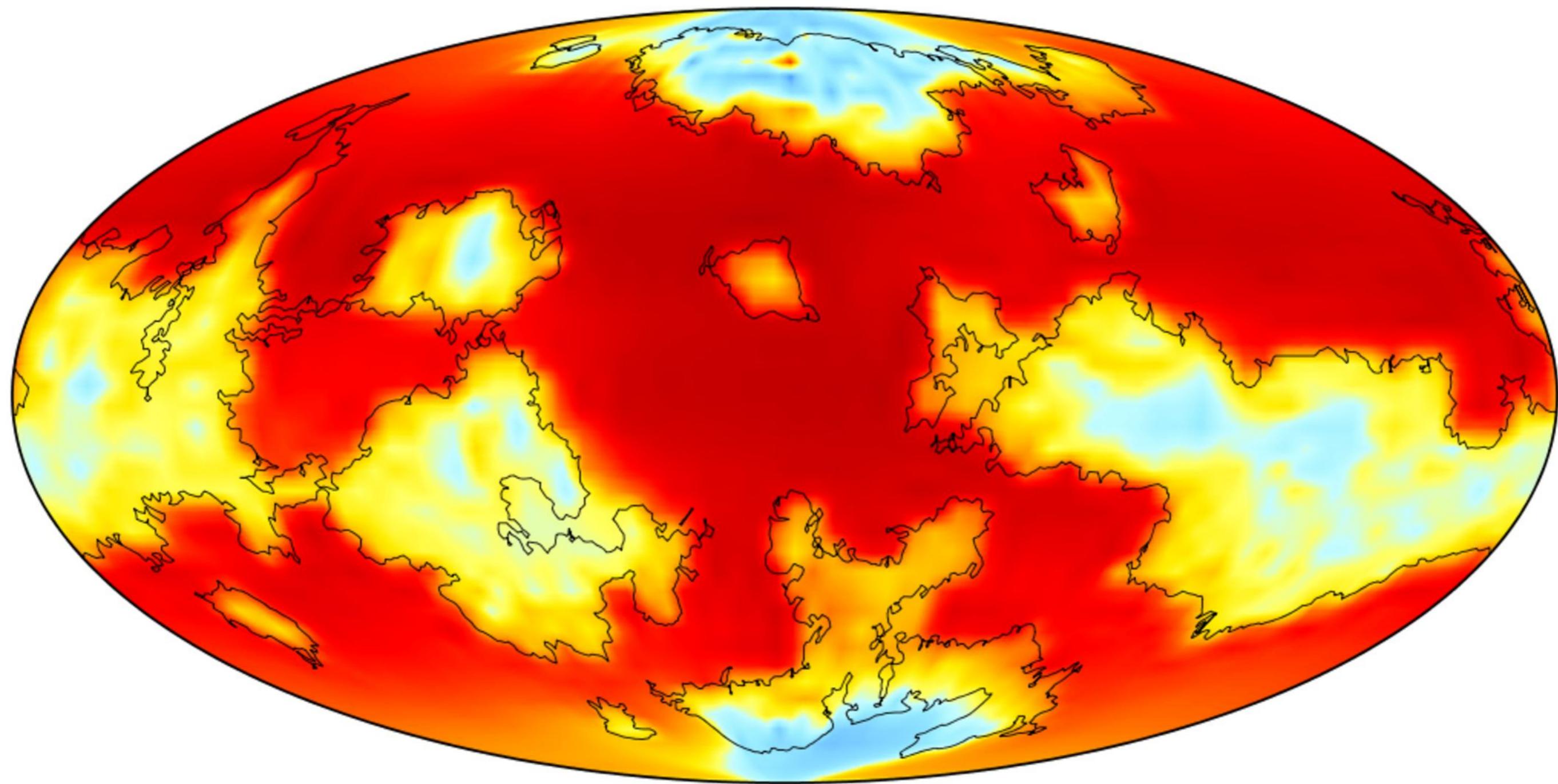


Figure 6.

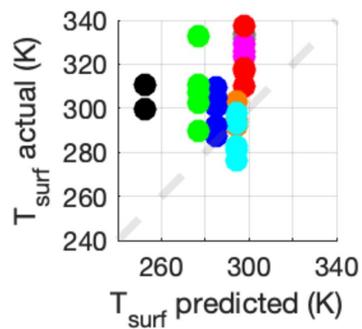
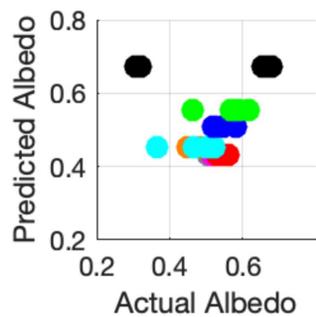
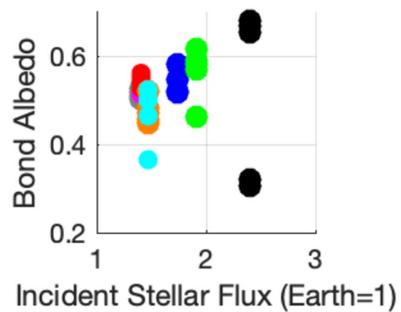


Figure 7a.

Simulation 21: GWTR : Arid-Venus : 2.9Ga : -243X : 1 bar: N2 dominated

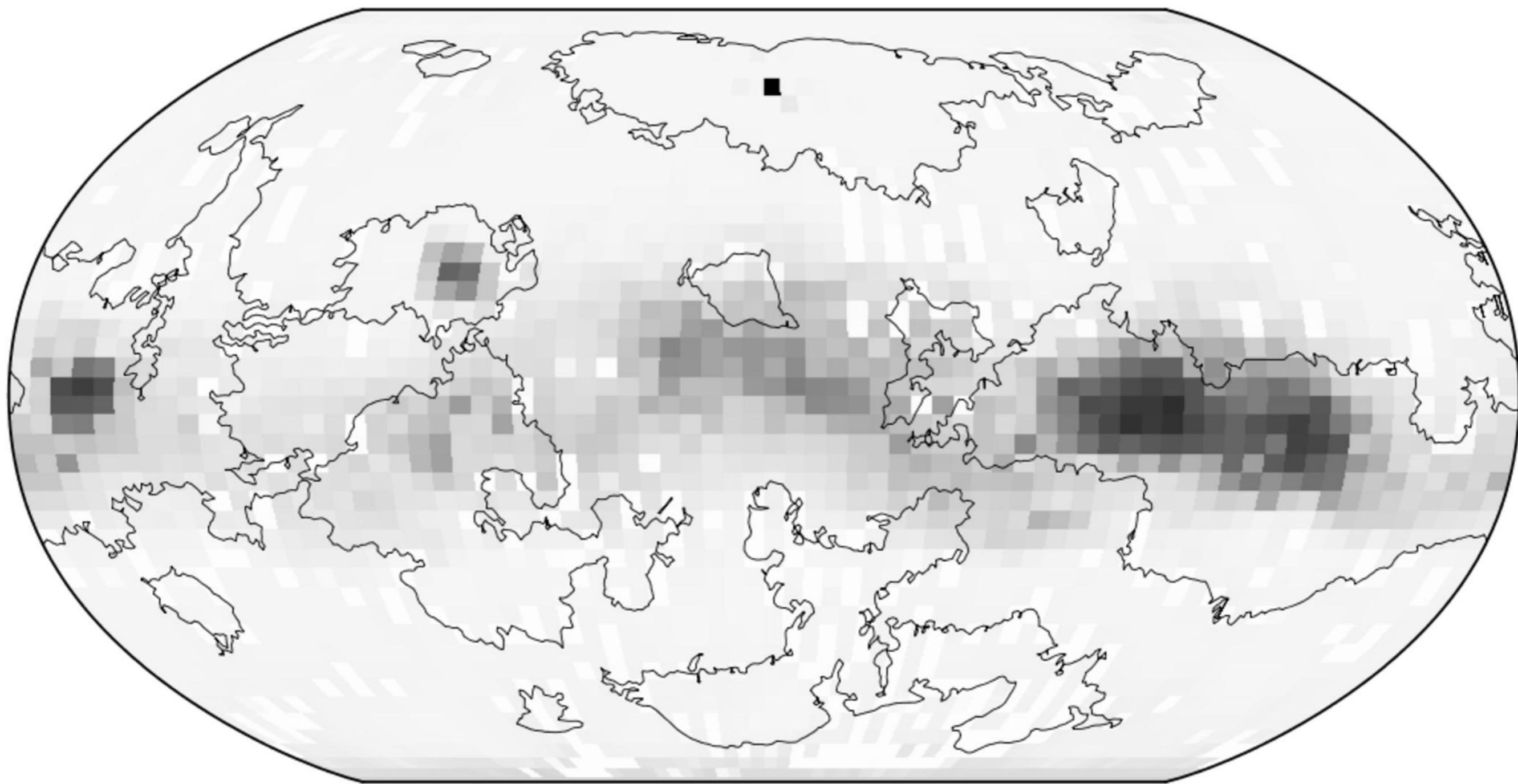


Figure 7b.

Simulation 21: Lake_Fr : Arid-Venus : 2.9Ga : -243X : 1bar : N2 dominated

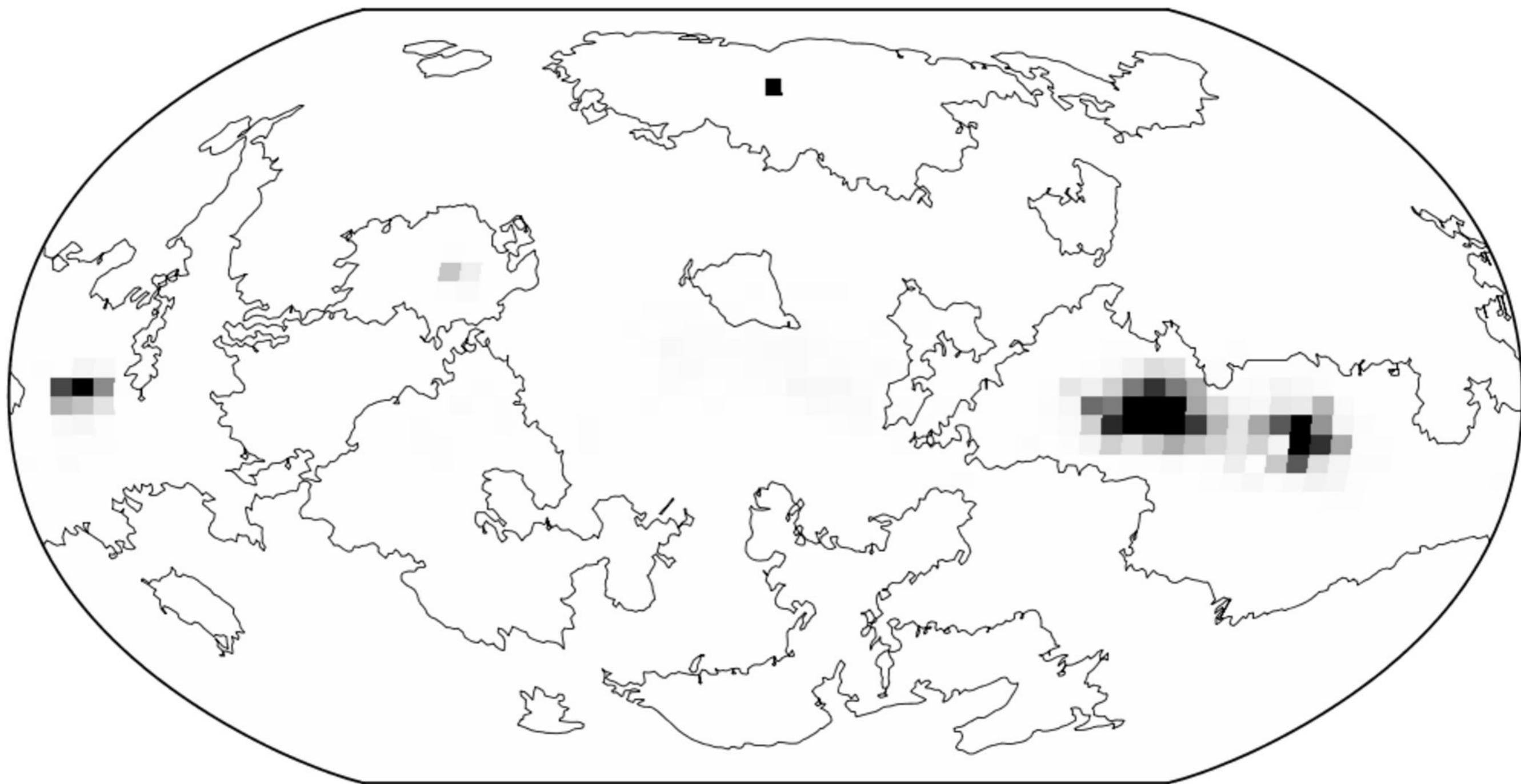


Figure 7c.

Simulation 26: GWTR : Arid-Venus : 2.9Ga : -243X : 0.25 bar: N2 dominated

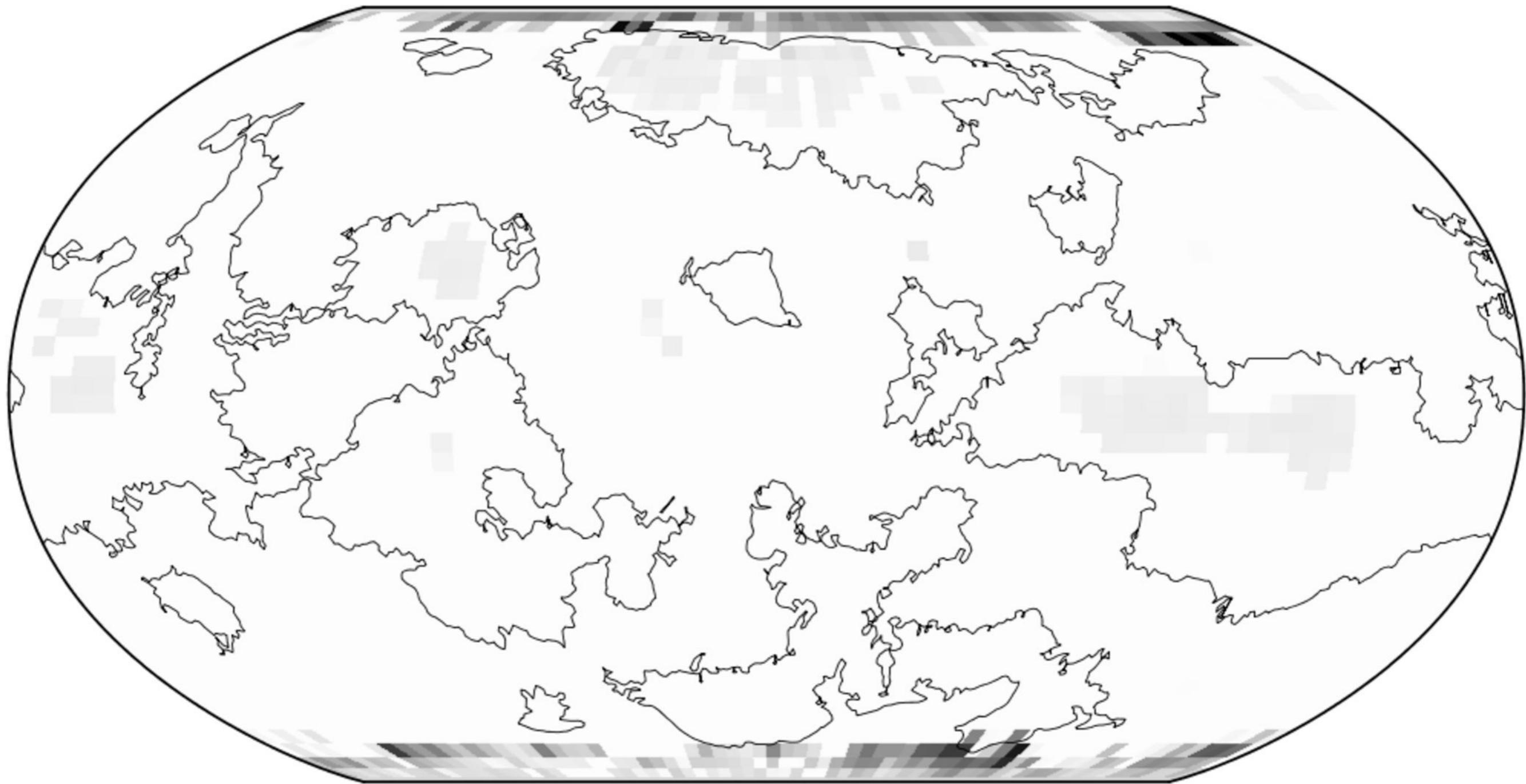


Figure 7d.

Simulation 26: Lake_Fr : Arid-Venus : 2.9Ga : -243X : .25 bar : N2 dominated

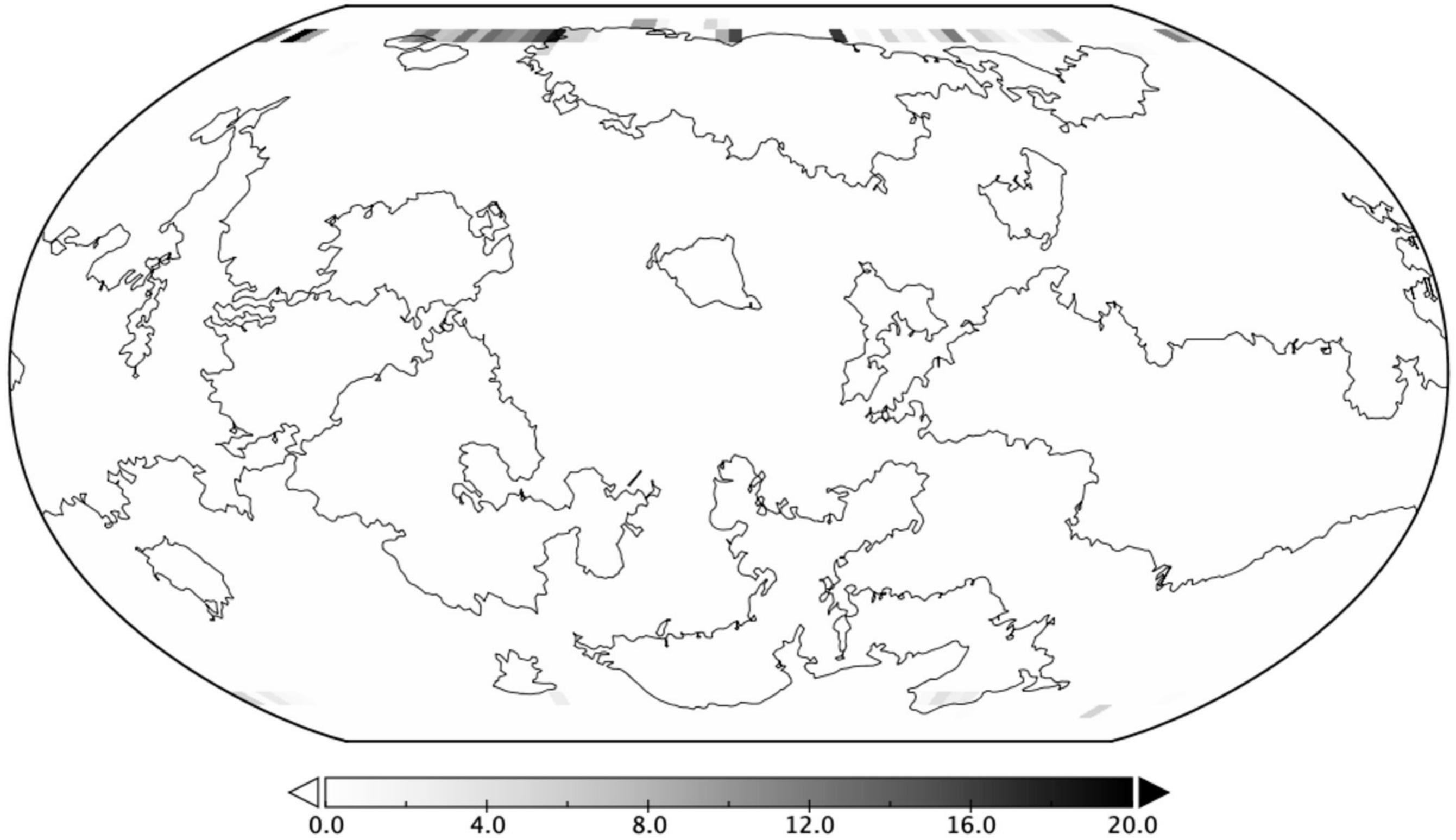


Figure 8a.

Simulation 22: GWTR : 10-Venus : 2.9Ga : -243X : 1 bar: N2 dominated

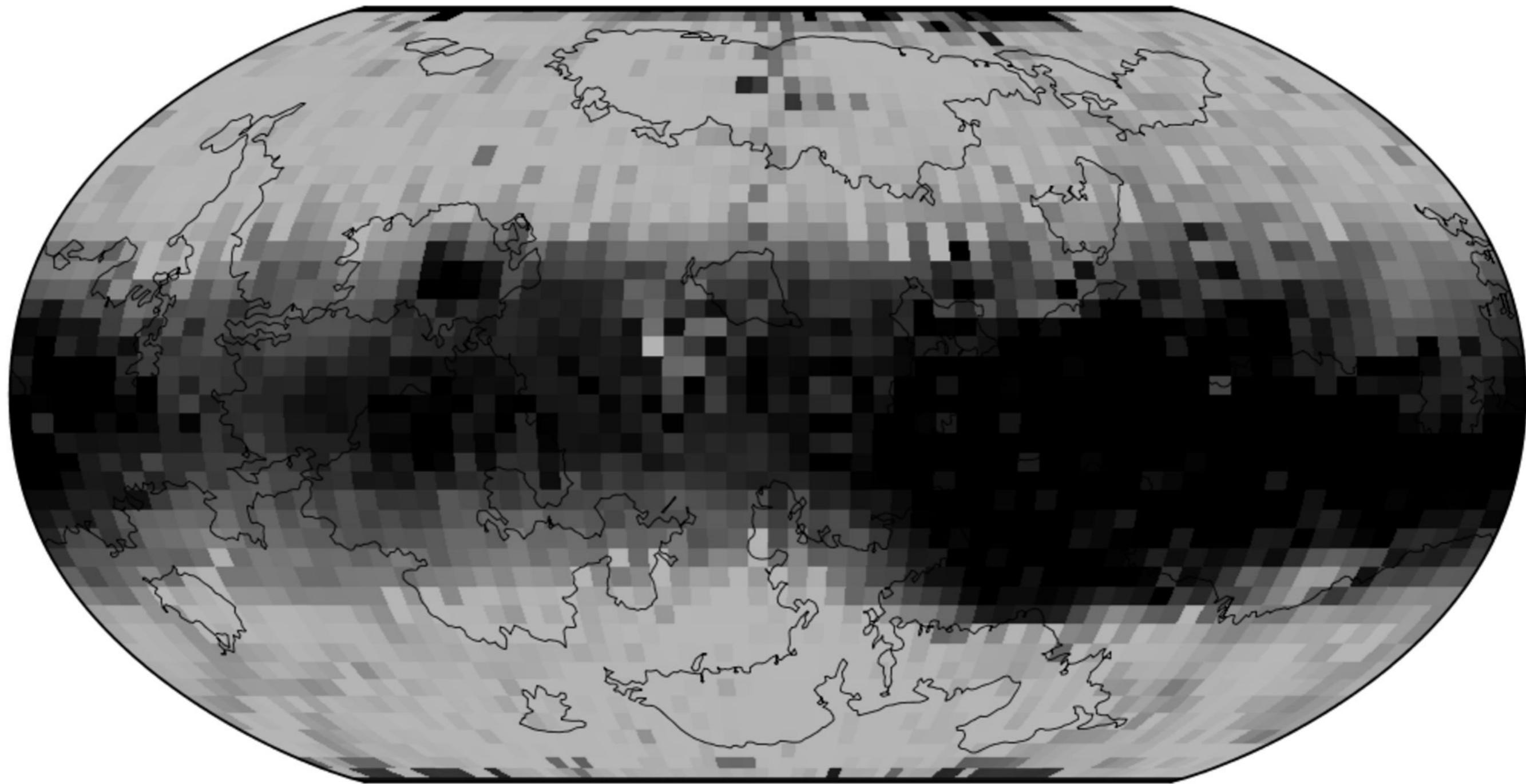


Figure 8b.

Simulation 22: Lake_Fr : 10m-Venus : 2.9Ga : -243X : 1bar : N2 dominated

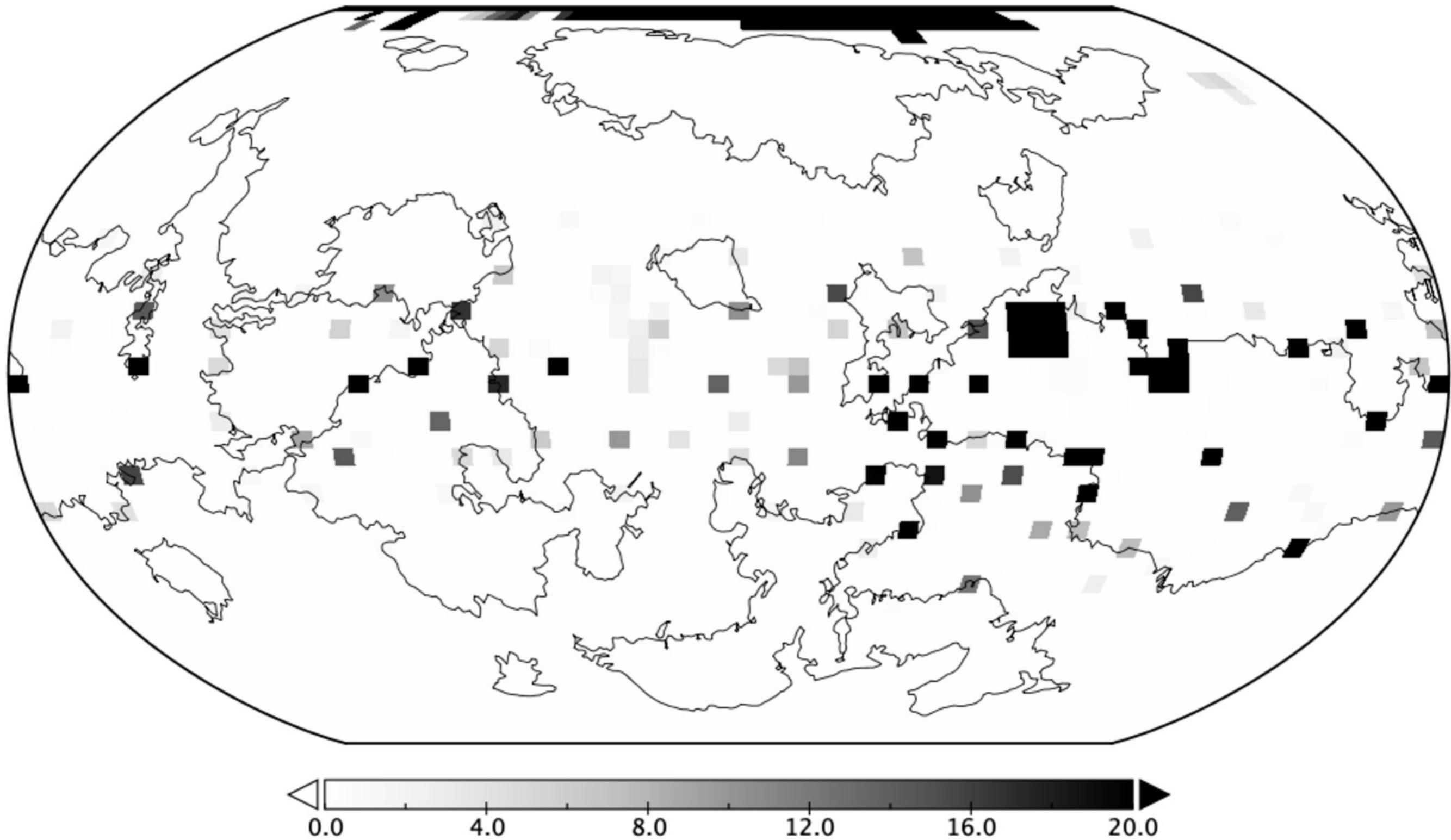


Figure 8c.

Simulation 27: GWTR : 10m-Venus : 2.9Ga : -243X : .25 bar: N2 dominated

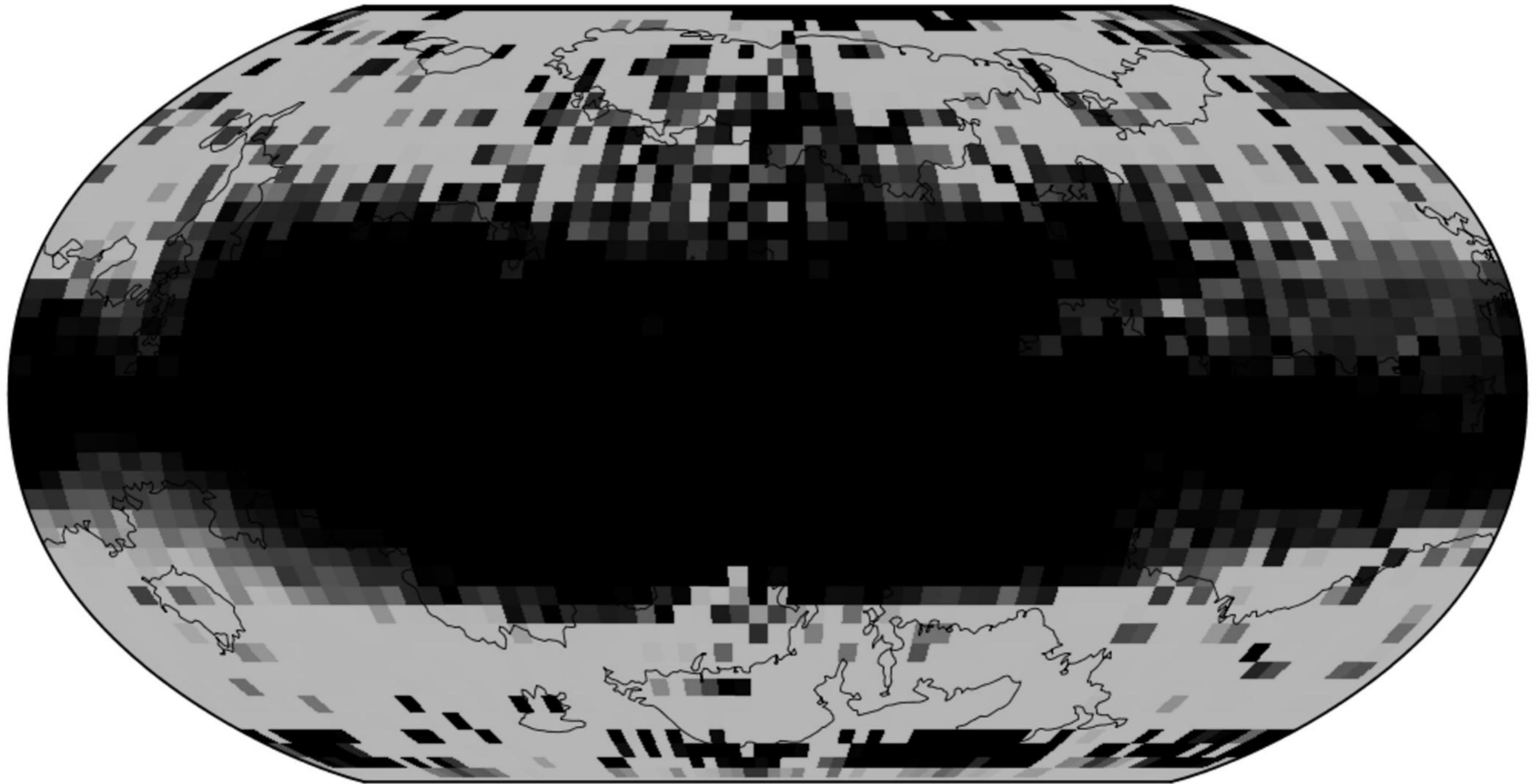


Figure 8d.

Simulation 27: Lake_Fr : 10m-Venus : 2.9Ga : -243X : .25 bar : N2 dominated

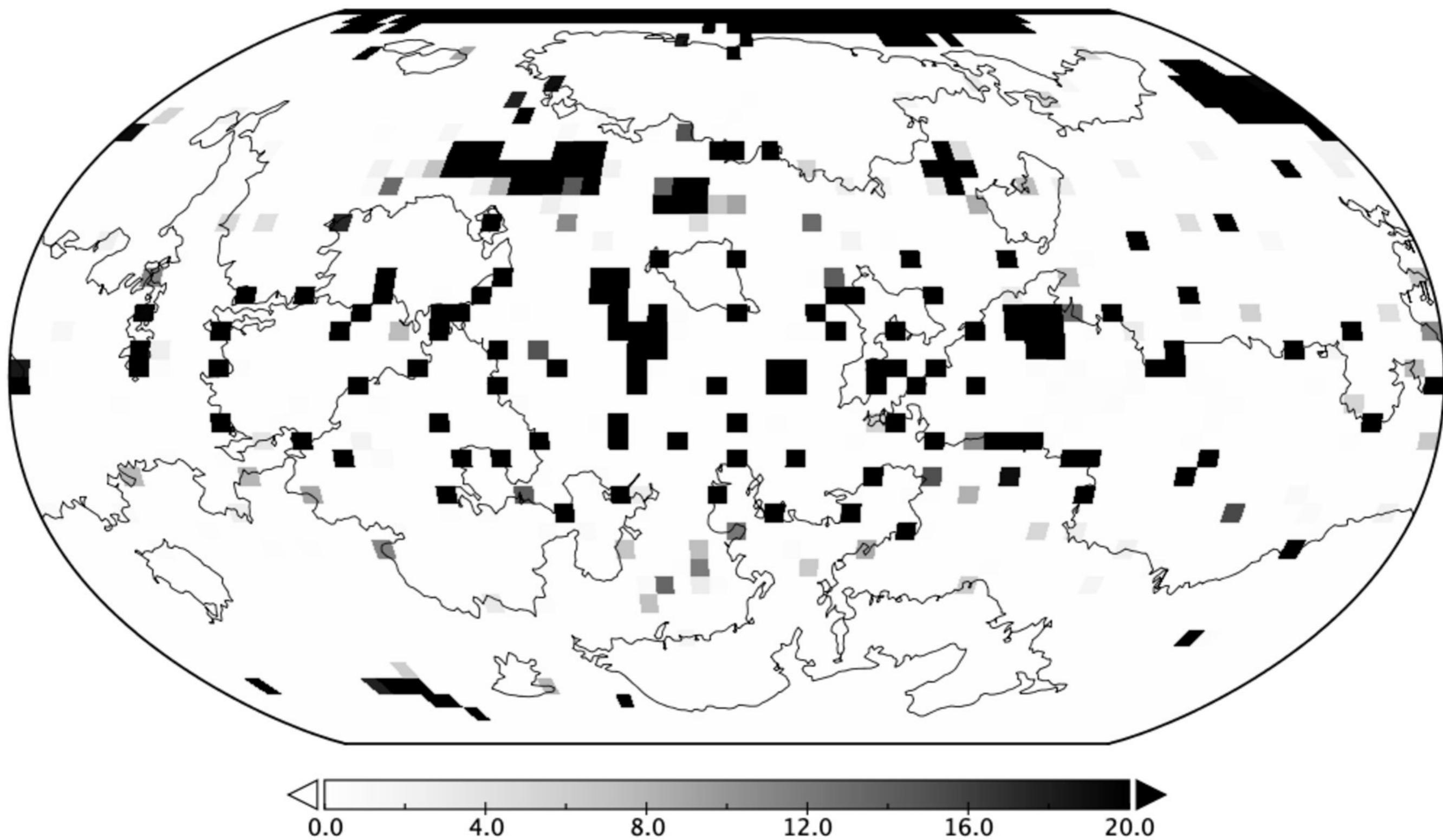


Figure 9a.

Sims 01-05 : -243x : 4.2Ga : 10 Bar : CO2 Dominated Atm

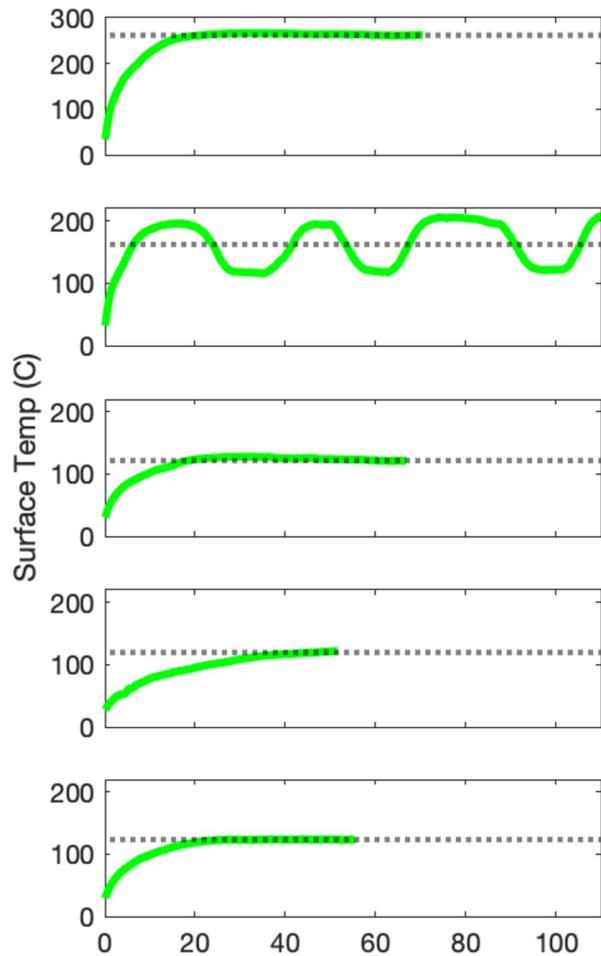
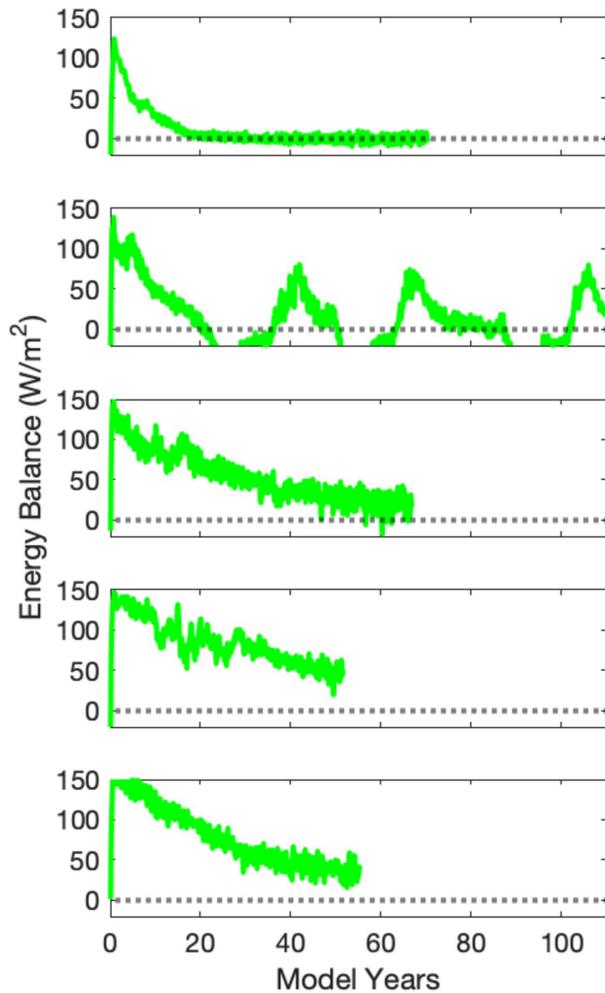


Figure 9b.

Sims 06-10 : -16x : 4.2Ga : 1 Bar : CO2 Dominated Atm

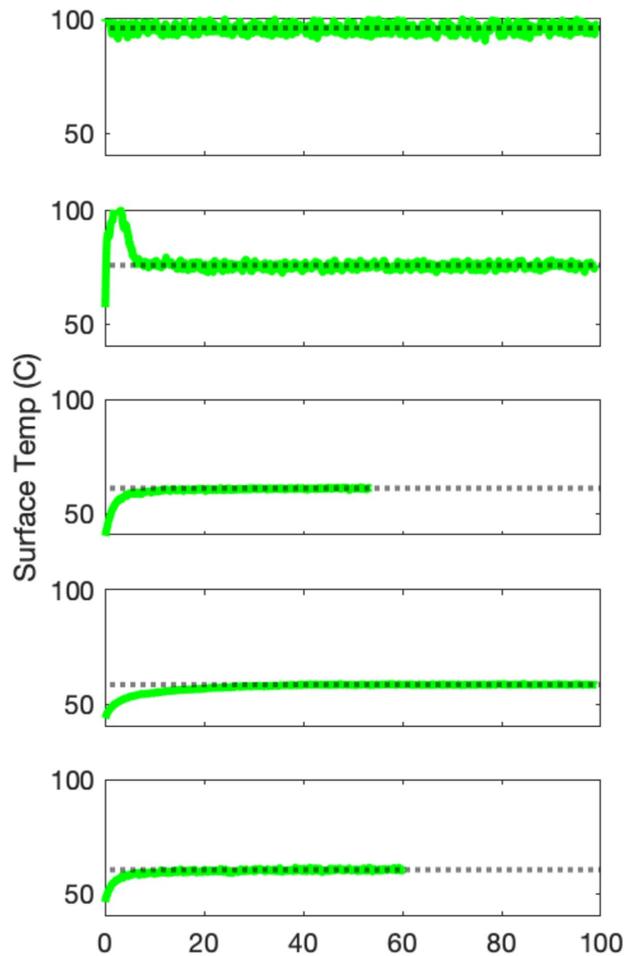
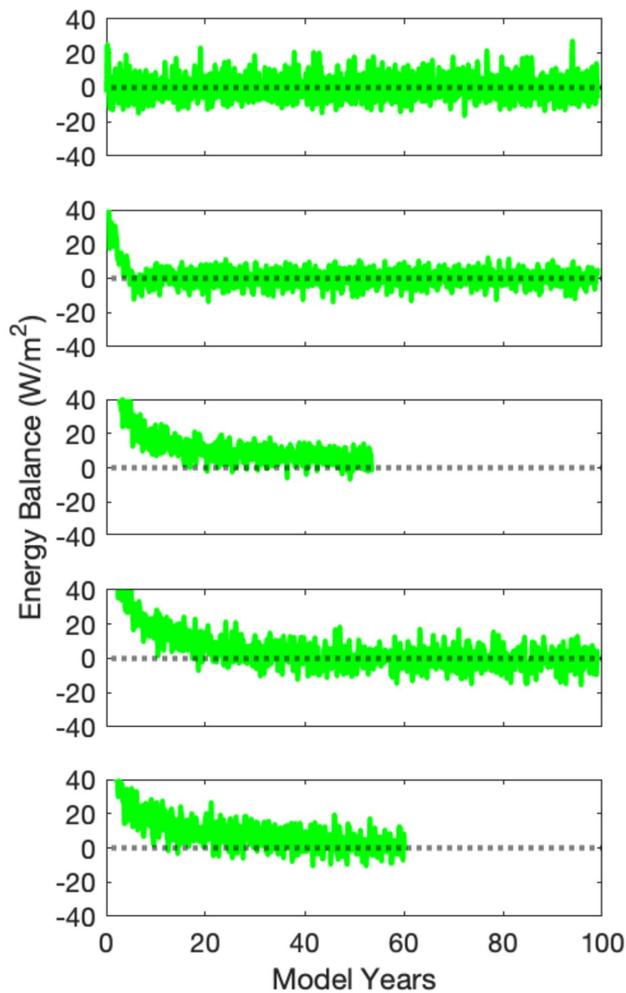


Figure 9c.

Sims 11-15 : -64x : 4.2Ga : 1 Bar : CO2 Dominated Atm

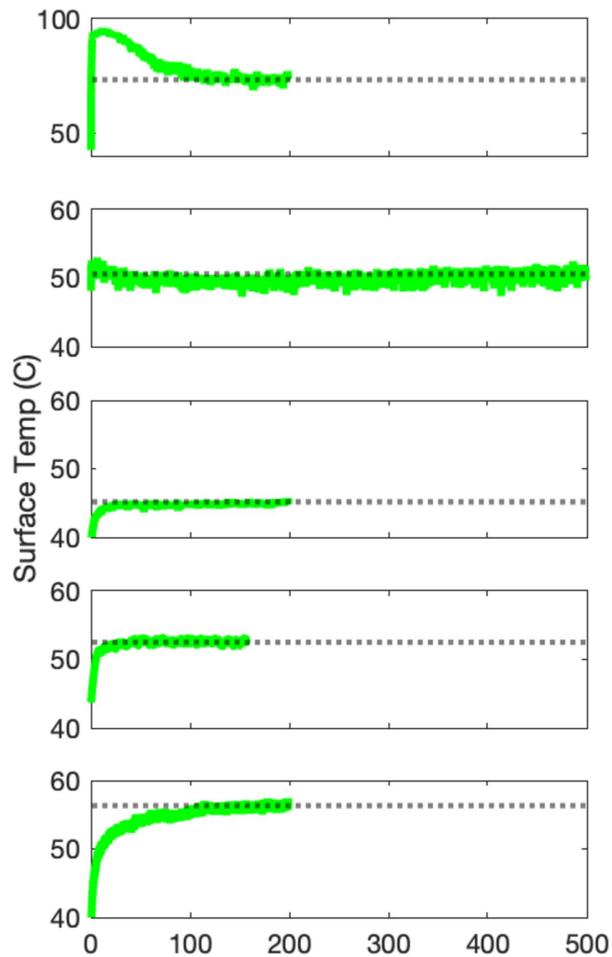
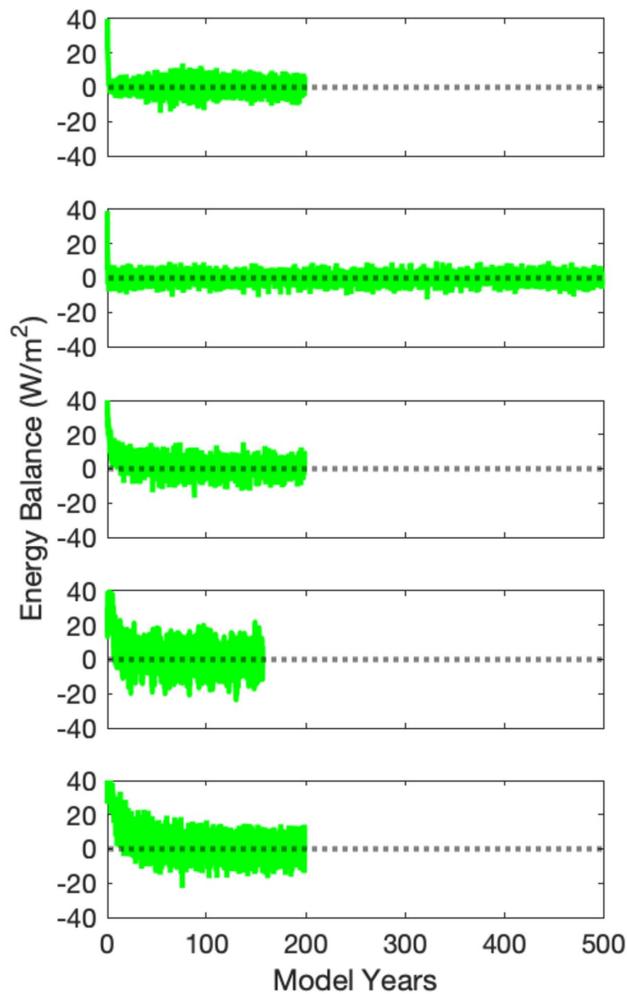


Figure 9d.

Sims 16-20 : -243 : 4.2Ga : 1 Bar : CO2 Dominated Atm

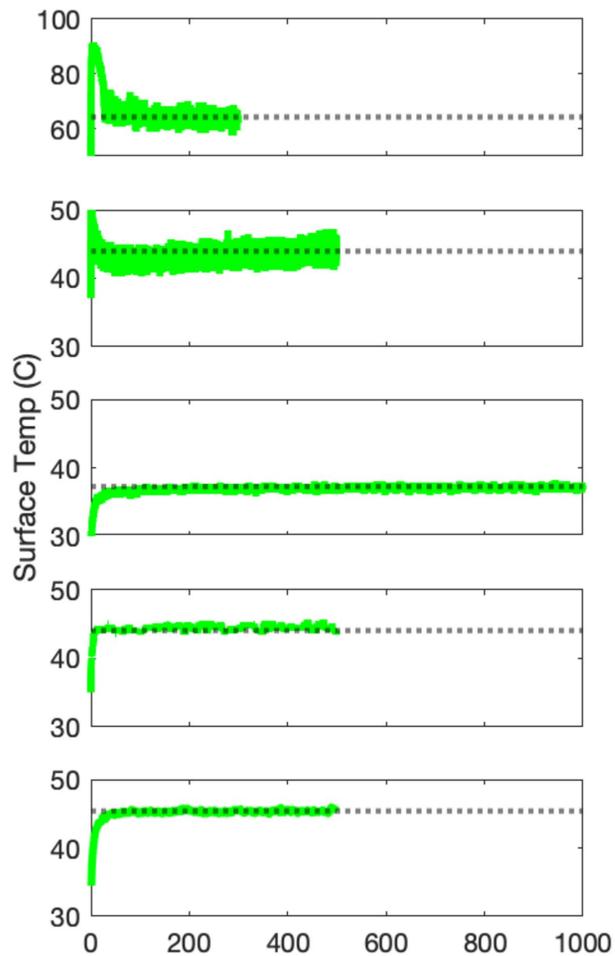
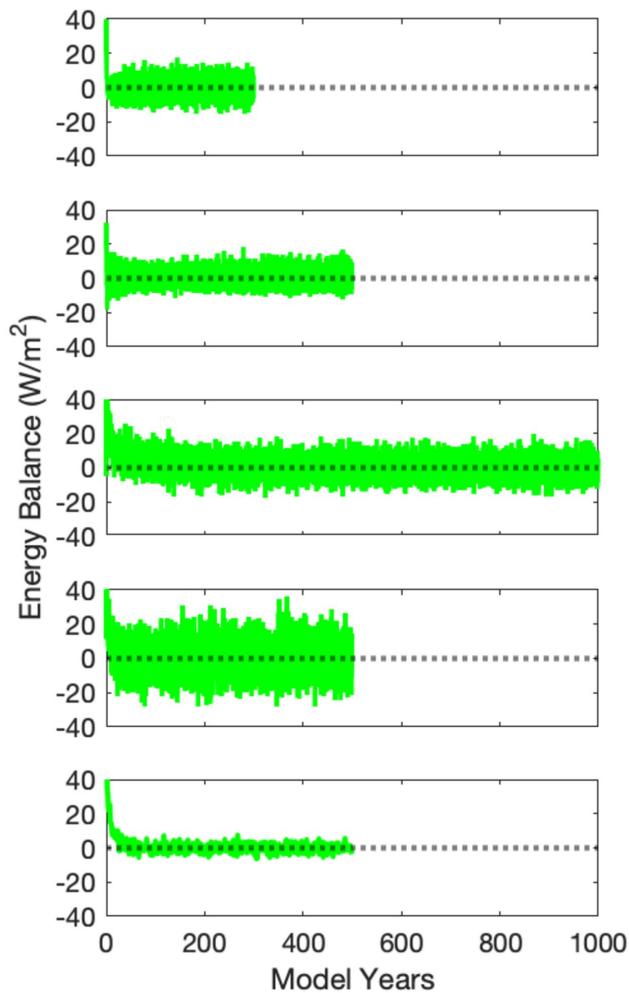


Figure 9e.

Sims 21-25 : -243x : 2.9Ga : 1 Bar : N2 Dominated Atm

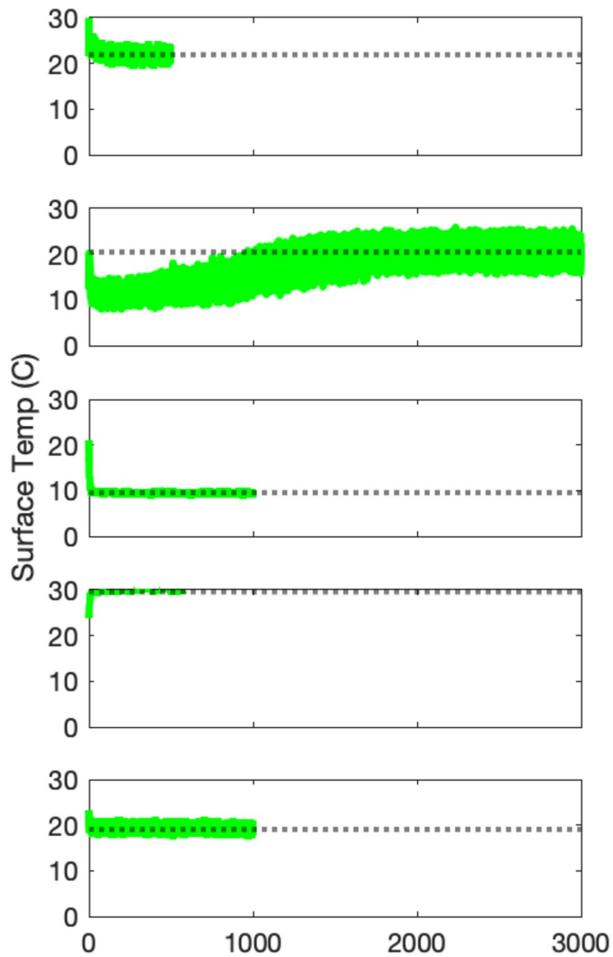
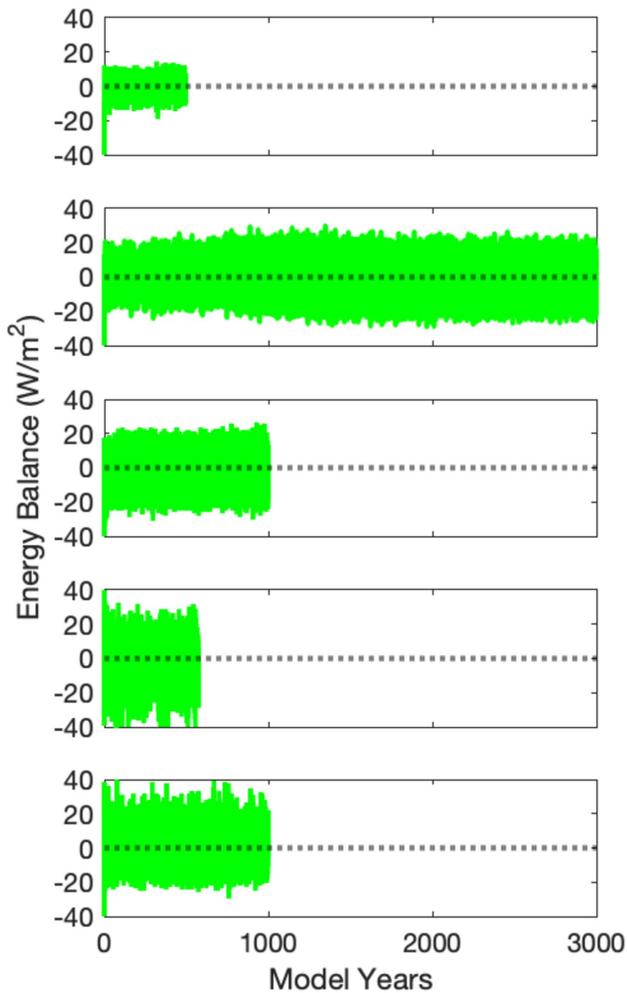


Figure 9f.

Sims 26-30 : -243x : 2.9Ga : 0.25 Bar : N2 Dominated Atm

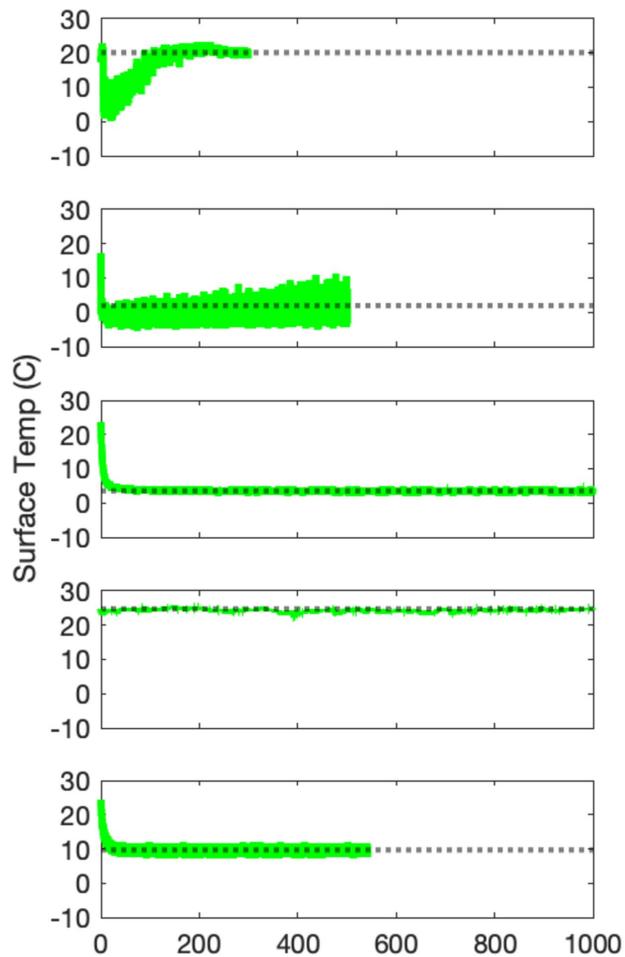
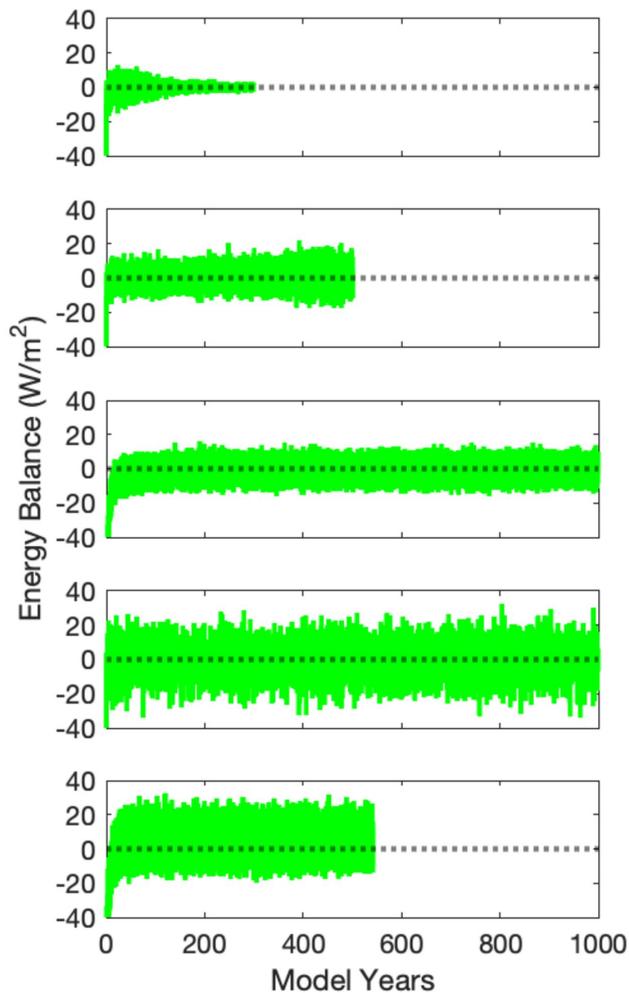


Figure 9g.

Sims 31-35 : -243x : 0.715Ga : 1 Bar : N2 Dominated Atm

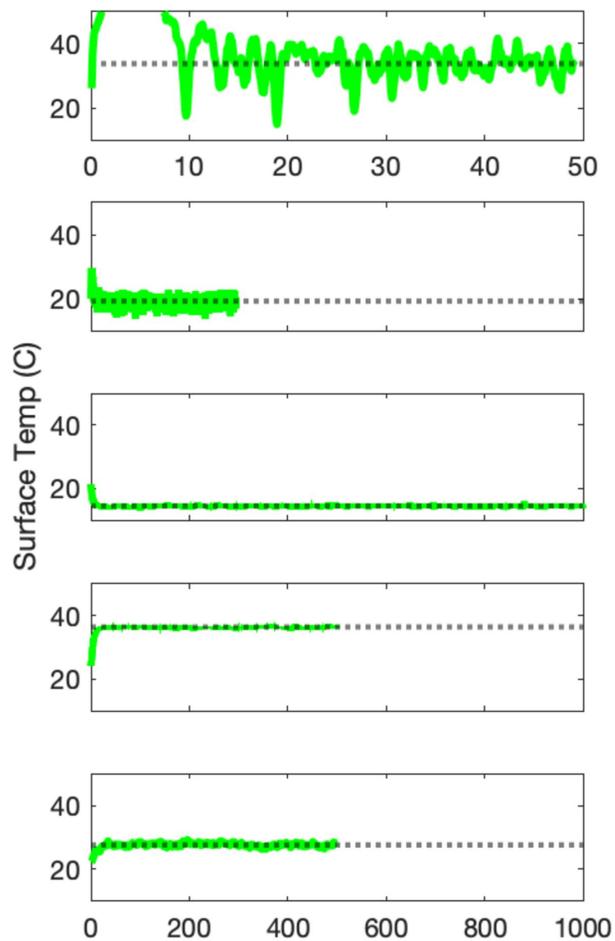
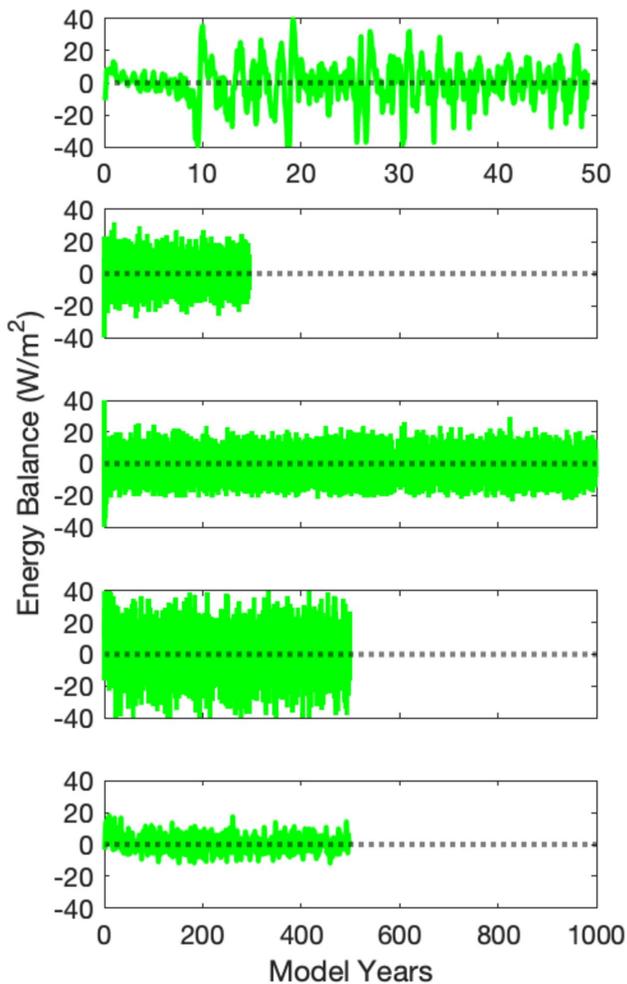


Figure 9h.

Sims 36-40 : -243x : 0.0Ga : 1 Bar : N2 Dominated Atm

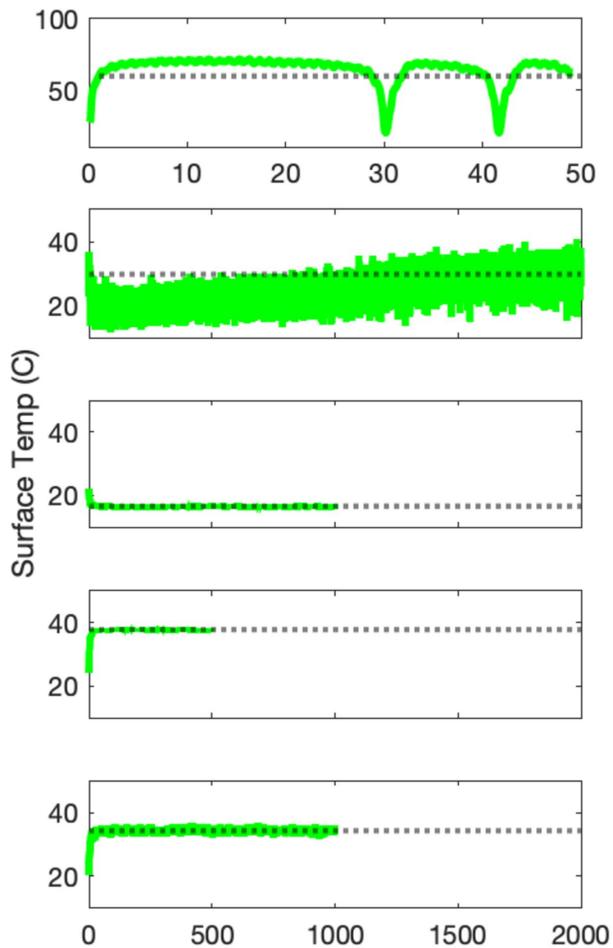
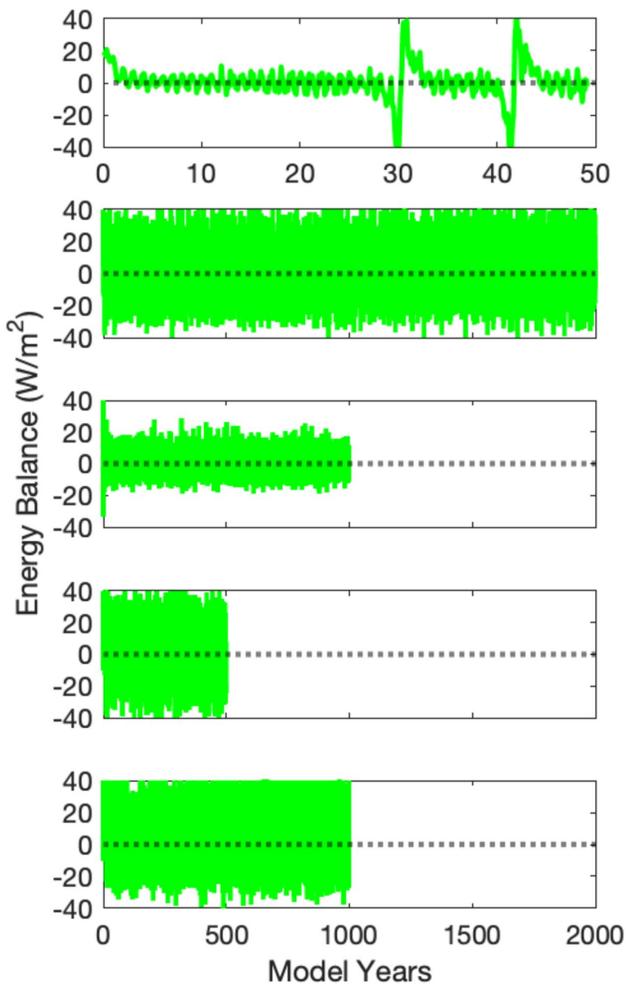


Figure 9i.

Sims 41-45 : -243x : Future : 1 Bar : N2 Dominated Atm

