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Asymmetric Responses of the Western Tropical Pacific Sea Level to El Niño and La Niña

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1 **Key Points.**

- 2 1) The responses of the western tropical Pacific sea level to ENSO are
3 obviously asymmetric, and the response to El Niño is evidently stronger.
- 4 2) The different surface wind anomaly structure between El Niño and La Niña
5 is critical for the asymmetric response of sea level.
- 6 3) The asymmetric ocean responses to ENSO may contribute to the ENSO
7 asymmetry.

8 **Abstract**

9 The western tropical Pacific (WTP) exhibits large interannual sea level
10 anomalies (SLAs), and the sea level falling in El Niño is evidently stronger than the
11 rising in La Niña. The asymmetry is most prominent near 160°E with the response to
12 El Niño larger by three times and becomes less obvious near the western boundary.
13 Sensitivity experiments of a simplified ocean model suggest that the asymmetry in
14 surface wind forcing structure between El Niño and La Niña is critical. The El Niño's
15 westerly wind anomaly patch locates more east than the La Niña's easterly wind patch
16 during the mature stage, and its upwelling effects are accumulated over a wider
17 longitude range and cause stronger negative SLAs in the WTP. Near the western
18 boundary, however, upwelling effects are attenuated by easterly wind anomalies
19 during El Niño conditions. The asymmetric ocean responses to ENSO winds may
20 participate in the asymmetry of ENSO cycle.

21 **Plain Language Summary**

22 El Niño–Southern Oscillation (ENSO) is the most influential climate variability
23 mode of the Pacific and cause strong interannual sea level anomalies (SLAs) in the
24 western tropical Pacific (WTP). We notice that the WTP's sea level falling in El Niño
25 condition is stronger than its sea level rising in La Niña. This difference is most
26 prominent near 160°E, where the falling in El Niño is stronger by three times. This
27 phenomenon becomes much less evident near the western boundary of the Pacific
28 basin. We further show that the difference in surface wind anomaly structures between
29 El Niño and La Niña is the primary cause. El Niño's westerly wind anomaly center

30 locates more east than La Niña's easterly wind anomaly center in their mature phase,
31 and there are easterly wind anomalies emerging near the western boundary during El
32 Niño condition. As a result, effect of El Niño's westerly wind anomaly is accumulated
33 over a wider longitude range and causes stronger sea level falling in the WTP. But this
34 effect becomes weaker near the western boundary, due to easterly wind anomalies
35 there. By contrast, effect of La Niña's easterly wind anomaly strengthens
36 monotonically from east to west, but the produced sea level rising signatures are
37 mainly confined to the WTP.

38 **Keywords**

39 Western Tropical Pacific, Sea Level, ENSO, ENSO asymmetry, Interannual

40 Variability

41 **1. Introduction**

42 Regional sea level change, as an essential aspect of climate change, is attracting
43 increasing attention of scientific communities and the general public, because of its
44 threats on eco-systems and people of coastal residence (e.g., Nicholls & Cazenave,
45 2010; Cazenave & Cozannet, 2014). There is a paramount demand for improved
46 understanding of regional sea level changes on various timescales (e.g., Milne et al.,
47 2009; Church et al., 2013; Stammer et al., 2013). Sea level changes over the tropical
48 Pacific Ocean are particularly pronounced on interannual timescale, as largely
49 modulated by El Niño-Southern Oscillation (ENSO) (e.g., Cazenave et al., 2008;
50 Antonov et al., 2005; Cheng et al., 2008). It is evident in Figure 1a that the western
51 tropical Pacific (WTP) shows stronger interannual sea level anomalies (SLAs) than
52 other regions, as quantified the standard deviation of low-passed SLAs. The WTP
53 shows strong sea level rising (falling) during the La Niña (El Niño) condition, as a
54 result of prevailing easterly (westerly) wind anomalies over the tropical Pacific basin
55 (Wyrtki, 1975; Zebiak, 1984; Alory & Delcroix, 2002; Gu & Li, 2009; Nerem et al.,
56 2010; Merrifield, 2011; Zhang & Church, 2012; Chang et al., 2013; Becker et al.,
57 2016; Hamlington et al., 2016). The sea level rising during La Niña conditions is
58 expected to aggravate coastal erosion, extreme marine flooding, or saltwater intrusion
59 in coastal aquifers (Nicholls & Tol, 2006; Nicholls et al., 2007; Nicholls & Cazenave,
60 2010), thus becomes a major threat for the densely populated coastal regions.
61 However, the sea level falling during El Niño causes the coral reef exposure and could
62 damage the region's ecosystem. There are many heavily-populated coasts and islands

63 in the WTP, so the sea level changes during ENSO are well worth studying.

64 El Niño and La Niña are, however, not mirrors of each other. They exhibit
65 obvious asymmetry in magnitude, duration time, and occurrence frequency (e.g.,
66 Burgers & Stephenson, 1999; Ohba & Ueda, 2007; Gergis & Fowler, 2009; An &
67 Choi, 2009; Okumura & Deser, 2010). Their signatures on sea level are also
68 asymmetric, owing to not only ENSO's asymmetry but also the nonlinearity in ocean
69 response (Niedzielski & Kosek, 2010; Swierczynska et al., 2013; Im et al., 2015; An
70 & Kim, 2017). Previous studies have demonstrated that ENSO asymmetry is
71 contributed by nonlinear dynamical heating of ocean current advection (e.g., Kang &
72 Kug, 2002; Jin et al., 2003; An & Jin, 2004; Su et al., 2010). As dynamically
73 associated with the upper-ocean currents through pressure gradient, the asymmetric
74 SLAs induced by ENSO may affect ocean current advection of heat and thereby
75 participate in the ENSO asymmetry.

76 In this study, we aim to investigate the asymmetric response of the WTP sea
77 level to ENSO and the underlying dynamical processes. This effort is of paramount
78 need for the prediction and projection of regional sea level change and adaptations of
79 the low-lying coasts and islands, as well as for understanding the dynamics of ENSO
80 asymmetry. The rest of the paper is structured as follows. Section 2 describes the data
81 and models. Section 3 describes the asymmetric SLAs of the WTP in response to
82 ENSO. Section 4 explores the underlying dynamics through a simplified ocean model.
83 Section 5 presents concluding remarks.

84 **2. Data and Models**

85 **2.1. Data**

86 In this study we use $0.25^\circ \times 0.25^\circ$, monthly satellite SLA data of 1993-2016
87 from Archiving Validation, and Interpretation of Satellite Oceanography (AVISO; Le
88 Traon et al., 1998) and sea level records of three tidal gauge stations in the WTP:
89 Malakal at 7.33°N , 134.45°E for 1979-2016, Kapingamrangi at 1.1°N , 154.78°E for
90 1929-2016, and Lombrum at 2.04°S , 147.47°E for 1995-2016 (Figure S1 in the
91 Supporting Information). The Hadley Centre Sea Ice and Sea Surface Temperature
92 (HadISST) data set (Kennedy et al., 2011) of the Met Office during 1979-2016 with
93 horizontal resolution of $1^\circ \times 1^\circ$ is used to compute Niño-3.4 index and identify El
94 Niño and La Niña conditions. Surface winds and other surface atmospheric fields of
95 1979-2016 are taken from the 0.75° monthly dataset of the European Centre for
96 Medium-Range Weather Forecasts (ECMWF) ERA-interim (Dee et al., 2011).

97 **2.2. HYCOM**

98 The HYbrid Coordinate Ocean Model (HYCOM) version 2.2.18 (Bleck, 2002)
99 are used to simulate interannual SLAs in the tropical Pacific Ocean. HYCOM is
100 configured to the Pacific Ocean basin between 48°S - 48°N , 110°E - 70°W , with a
101 horizontal resolution of $1/3^\circ \times 1/3^\circ$ and 26 hybrid vertical layers (Li et al., 2015; Ren et
102 al., 2020). Surface atmospheric forcing fields (winds, heat fluxes, precipitation, etc.)
103 are taken from ERA-interim. Three sponge layers are applied to the western, southern
104 and northern open-ocean boundaries, where model temperature and salinity are
105 related to the World Ocean Atlas 2009 (WOA09) climatology (Antonov et al., 2010).

106 More details of model configuration are described in Ren et al. (2020). Subsequent to
107 the spin-up run of 30 years under monthly climatologic forcing, two parallel
108 experiments are performed under daily ERA-Interim fields for the period of 1979-
109 2016. The control run (HYCOM-CTL) is forced with the original daily atmospheric
110 fields and assumed to contain the complete processes. This experiment can well
111 reproduce the amplitude and spatial distribution of interannual SLA and upper-ocean
112 circulation variation in observation (Figure 1b; Ren et al., 2020). Another experiment,
113 HYCOM-TAU, uses daily wind stress forcing as HYCOM-CTL, but all the other
114 forcing fields (heat and freshwater fluxes) are fixed to monthly climatology.
115 HYCOM-TAU is used to evaluate the effects of wind forcing on sea level.

116 **2.3. Reduced-Gravity Ocean Model**

117 To achieve more in-depth understanding, a series of experiments are performed
118 with a 1.5-layer nonlinear reduced-gravity ocean (RGO) model. This model mainly
119 represents the 1st-mode baroclinic response of the ocean to surface wind forcing,
120 which is the dominant source of large-scale interannual variability in sea level and
121 upper-ocean circulation of the WTP (e.g., Qiu & Chen, 2010, 2012). The model is
122 configured to the Pacific Ocean basin between 40°S-40°N, 100°E-70°W with
123 horizontal resolutions of 0.25°×0.25° and forced by monthly ERA-Interim surface
124 winds. Readers are referred to Duan et al. (2019) for more details of the model
125 configuration. After a spin-up of 10 years under climatological wind forcing, the
126 control run of RGO model (RGO-CTL) is forced by realistic monthly winds. RGO
127 model experiments (Table S1) are forced by idealized wind forcing to examine the

128 role of wind forcing asymmetry and are described in Section 4.

129 **3. Asymmetric Responses to El Niño and La Niña**

130 To highlight the interannual variations associated with ENSO, we analyze the 13-
131 month low-pass filtered anomaly fields with the monthly climatology removed.

132 Skewness S is a measure of the distribution asymmetry with $S = 0$ indicating a normal
133 distribution (White, 1980). Figure 1c shows the skewness of the observed SLA over
134 the tropical Pacific for 1993-2016. The eastern tropical Pacific is positively skewed
135 with the maximum S of ~ 2.0 , while the central-western tropical Pacific is negatively
136 skewed with the minimum S of -2.0 and a horseshoe structure extending from the
137 equator to extratropical regions in both hemispheres. This distribution of SLA
138 skewness resembles that of SST anomaly (An & Jin, 2004; Niedzielski & Kosek,
139 2010) and is likely associated with the positive skewness of ENSO (Nerem et al.,
140 2010). The El Niño condition is characterized by positive SLAs in the eastern Pacific
141 and negative SLAs in the WTP, and these anomalies are stronger in amplitude than
142 the opposite SLAs occurring in La Niña condition (Niedzielski & Kosek, 2010;
143 Figure S1).

144 In addition to the asymmetry residing in ENSO (e.g., as quantified by the
145 skewness of Niño-3.4 index), the asymmetric responses of sea level to El Niño and La
146 Niña also contribute to the SLA asymmetry shown in Figure 1c. We regress SLAs
147 onto the normalized Niño-3.4 index separately for the El Niño condition (Niño-3.4 >
148 0) and the La Niña condition (Niño-3.4 < 0). The corresponding regression
149 coefficients, namely $k_{\text{Niño}}$ and $k_{\text{Niña}}$, are used to quantify the responses of SLA to El

150 Niño and La Niña, respectively (Figures 1d and 1e). The response time of sea level to
151 ENSO shows spatial variation, as indicated by the lead-lag correlation (Figure S2a).
152 For each grid point, the lead-lag time of the maximal correlation is used to compute
153 the regression coefficient. The results are not dramatically different from those of
154 simultaneous regression (Figure S2b). Figures 1d and 1e show distributions of k_{Nino}
155 and k_{Nina} . In the WTP, the maximal k_{Nino} of ~ -0.2 m is located 20° - 30° away from the
156 western boundary, and k_{Nino} decreases in magnitude as approaching the western
157 boundary. By contrast, the peak k_{Nina} values of ~ -0.15 m are close to the western
158 boundary. We further use the ratio of k_{Nino} to k_{Nina} to quantify the response asymmetry,
159
$$R_k = \frac{k_{Nino}}{k_{Nina}}. \quad (1)$$

160 $R_k = 1$ denotes symmetric response of SLAs to El Niño and La Niña, while $R_k > 1$ and
161 $R_k < 1$ indicates stronger and weaker response to El Niño than to La Niña,
162 respectively. As shown in Figure 1f, R_k reaches the largest value near 160° E with
163 values exceeding 3.0, indicating that the response of SLA to El Niño is stronger by at
164 least 3 times than the response to La Niña. R_k is weakened to ~ 1.0 near the western
165 boundary, implying SLA responses there are nearly symmetric. This interesting
166 distribution of R_k in the WTP and underlying dynamics are worthy of systematic
167 investigation. One may notice that R_k is < 1 in the central Pacific and is > 1 in the
168 eastern Pacific, indicative of prevailing asymmetry over the tropical Pacific basin. In
169 the following, we focus on explaining the R_k distribution in the WTP.

170 Simulations of HYCOM and RGO model have faithfully reproduced the
171 observed interannual variations of sea level at three tidal gauge stations (Figure S3).

172 The upper-layer thickness (ULT) anomaly of the 1.5-layer RGO model is a good
173 proxy of SLA in the tropical Pacific (e.g., Qiu & Chen, 2010, 2012; Chang et al.,
174 2013; Duan et al., 2019). The correlations among AVISO, HYCOM-CTL, HYCOM-
175 TAU and RGO-CTL at the three tidal gauge stations are all above 0.85. The
176 asymmetric responses of SLAs to El Niño and La Niña in the WTP during 1993-2016
177 can be realistically reproduced by HYCOM-CTL, HYCOM-TAU, and RGO-CTL
178 (Figures 2a-2c), although RGO-CTL fails to capture the features in the eastern
179 Pacific. The good performance of HYCOM-TAU and RGO-CTL indicates that the
180 interannual SLAs in the WTP and their asymmetric features are primarily the results
181 of ENSO wind forcing, and the underlying dynamics can be explored by sensitive
182 model experiments with prescribed wind forcing fields.

183 **4. Dynamics**

184 To include more ENSO events, we use the period of 1979-2016 to perform
185 model experiments (Table S1), although R_k of 1979-2016 shows detailed differences
186 from that of 1993-2016 in the northwest Pacific (Figures 2d-2f). According to existing
187 studies of ENSO asymmetry, the different spatial and temporal characteristics of wind
188 anomalies are essential to cause the SST asymmetry between El Niño and La Niña
189 (Kang & Kug, 2002; An & Kim, 2017). To examine which aspect of ENSO's wind
190 forcing is critical in regulating the asymmetric responses of SLAs, we adopt a
191 statistical model based on the singular value decomposition (SVD) of wind stress and
192 SST (Kang & Kug, 2000), which is expressed as

$$193 \quad \boldsymbol{\tau}'(x, y, t) = \sum_n^N c(n) [\sum_{x,y} V_{\text{SST}}(x, y, n) T(x, y, t)] V_{\boldsymbol{\tau}}(x, y, n), \quad (2)$$

194 where x , y , and t represent longitude, latitude, and time, respectively, V_{SST} and V_{τ} are
 195 the SVD singular vectors for SST and wind stress, T is the SST anomaly field, and n
 196 $= 1, 2, \dots, N$ indicates the n th mode of SVD, $\sum_{x,y}$ indicates spatial integration over
 197 the region of 100° - 290° E, 40° S- 40° N. $c(n)$ represents the correlation between SST
 198 and wind stress anomalies,

$$199 \quad c(n) = \frac{\sum_t t_{\text{SST}}(t,n)t_{\tau}(t,n)}{\sum_t t_{\text{SST}}(t,n)^2}, \quad (3)$$

200 where t_{SST} and t_{τ} are the corresponding time series, the numerator is the covariance
 201 of t_{SST} and t_{τ} and the denominator is the variance of t_{SST} , and \sum_t indicates temporal
 202 integration from January 1979 to December 2016. Zonal and meridional components
 203 of τ' are separately computed. All the RGO sensitive experiments are forced by
 204 monthly wind stress anomalies constructed by Eq. 2 (representing ENSO wind
 205 forcing) plus monthly climatological winds.

206 We first perform two experiments, namely EXP1 and EXP2. EXP1 uses different
 207 wind stress anomaly fields for El Niño and La Niña conditions. Specifically, $\tau'_{\text{Niño}}$ and
 208 $\tau'_{\text{Niña}}$ are reconstructed separated for Niño-3.4 ≥ 0 and Niño-3.4 < 0 conditions
 209 (Figures S4a and S4b) using Eq. 2, respectively, so that the synthesized τ' still retains
 210 the difference in spatial structure between El Niño and La Niña conditions. By
 211 contrast, EXP2 does not distinguish El Niño and La Niña conditions and uses
 212 reconstructed τ' for the entire model period (Figure S4c). As such, the difference
 213 between EXP1 and EXP2 represents the effect of different wind anomaly structures
 214 between El Niño and La Niña on SLAs. In EXP1 and EXP2, we use only the leading
 215 mode ($n = 1$) of SVD to reconstruct τ' (Figure S4), which explains $> 85\%$ of the total

216 covariance and mainly represents ENSO's mature phase (Figure S5). The higher SVD
217 modes largely represent the transition stages between El Niño and La Niña polarities
218 (Figure S5) and have limited impacts on the asymmetry of SLAs (Figure S6).

219 Figures 3a and 3b show R_k distributions of ULT produced by EXP1 and EXP2,
220 respectively. EXP1 is able to reproduce large R_k values in the WTP as RGO-CTL,
221 whereas R_k is generally close to 1.0 in EXP2. These results suggest that the
222 asymmetry in surface wind structures between El Niño and La Niña is largely
223 responsible for asymmetric responses of the WTP sea level. Note that EXP2 still
224 retains some asymmetric characteristics of ENSO winds, such as the asymmetries in
225 intensity, frequency, and temporal evolution. Figure 3b however indicates that these
226 factors have little contributions to SLA asymmetry. This is confirmed by an additional
227 experiment EXP3, which adopts an idealized sine time series for τ' (Table S1) and
228 achieves similar results to EXP2 (Figure S7).

229 In the equatorial zone, zonal component of wind stress τ^x is much more
230 influential for the ocean than the meridional component. We repeat EXP1 and EXP2
231 using only τ^x anomaly, and τ^y is fixed to monthly climatology (EXP1-TAUX and
232 EXP2-TAUX). The results of two experiments achieve are broadly consistent with
233 EXP1 and EXP2 (Figures 3c and 3d). Therefore, the spatial structure of τ^x is critical
234 for the asymmetric responses. The typical structures of ULT and τ^x in EXP1-TAUX
235 for warm and cold phases are shown in Figures 3e and 3f, respectively, as represented
236 by the leading SVD mode of τ^x and ULT. Notice that westerly wind anomaly patch of
237 El Niño locates more east than the easterly wind anomaly patch of La Niña in their

238 mature phase, and correspondingly the zero value of ULT anomaly of El Niño also
 239 locates more east (Jin, 1997; Kang & Kug, 2002). During El Niño there are easterly
 240 wind anomalies near the western boundary. The strong negative ULT anomalies in El
 241 Niño are seen over a wide longitude range and weaken as approaching the western
 242 boundary, while the positive ULT anomalies in La Niña are confined to the far WTP
 243 and generally strengthen westward.

244 To better elucidate how the wind forcing structure cause asymmetric SLAs, we
 245 show in Figure 4 the zonal distributions of τ^x and ULT in the equatorial band (5°S-
 246 5°N). It is clearly discernible in Figure 4a that the El Niño's westerly wind anomaly
 247 patch locates more east than the La Niña's easterly wind anomaly patch. The forcing
 248 effect of equatorial zonal wind stress on sea level and ULT slopes can be roughly
 249 expressed in a linear relationship (Sverdrup, 1947; McCreary, 1977; Alory &
 250 Delcroix, 2002; Palanisamy et al., 2014),

$$251 \quad \frac{dh}{dx} = \tau^x, \quad (4)$$

252 where h is ULT anomaly or SLA. Therefore, $h(x)$ at a given longitude x can be
 253 determined by the integration of Eq. (4) from the eastern boundary $x_E = 70^\circ\text{W}$ along
 254 equator,

$$255 \quad h(x) = h(x_E) + \int_{x_E}^x \tau^x dx. \quad (5)$$

256 Since $h(x_E)$ diffuses quickly away from the eastern boundary as free Rossby waves
 257 and has little impact on the interior Pacific (Qiu et al., 2013), $h(x)$ is primarily
 258 determined by the zonal integral of τ^x from x_E to x (second term). As such, the
 259 negative ULT anomaly in El Niño in the WTP is much stronger than the positive ULT

260 anomaly in La Niña owing to the much wider westerly wind anomaly patch to its east.
 261 West of 160°E, easterly wind anomaly near the western boundary causes downwelling
 262 of the ocean, resulting in the attenuation of negative ULT anomalies there. By
 263 contrast, the La Niña's positive ULT anomalies are continuously strengthened by
 264 easterly wind anomalies. As a result, the ULT anomalies of El Niño and La Niña are
 265 comparable in amplitude in the far WTP, and the asymmetry of response is no longer
 266 evident there (right panel of Figure 4a). In EXP2-TAUX (Figure 4b), without the
 267 difference in wind forcing structure (left panel), the asymmetric responses cannot be
 268 reproduced (middle and right panels).

269 Figure 4c shows the zonal distributions of ULT anomaly and regression
 270 coefficients predicted by the linear theory, which compare favorably with Figure 4a. It
 271 indicates that the linear theory can to a large extent capture the processes causing
 272 asymmetric responses to ENSO winds and confirms the critical role played by the
 273 structure of zonal wind anomaly. Under this theoretical framework, the response
 274 asymmetry R_k can be theoretically expressed as,

$$275 \quad R_k(x) \approx \frac{\int_{x_E}^x \tau^x_{Nino} dx}{\int_{x_E}^x \tau^x_{Nina} dx}, \quad (6)$$

276 where τ^x_{Nino} and τ^x_{Nina} are the zonal wind stress anomaly in El Niño and La Niña
 277 conditions, respectively. Eq. (6) clearly suggests the sensitivity of $R_k(x)$ to the
 278 distribution of τ^x from x to the eastern boundary.

279 **5. Concluding Remarks**

280 The WTP exhibits large interannual variations of sea level, and the sea level

281 falling in El Niño is stronger than the rising in La Niña. Here we show that this
282 asymmetry is most prominent near 160°E with the response to El Niño larger by ~3
283 times and becomes much less obvious near the western boundary. RGO model
284 experiments suggest that the asymmetric surface wind anomaly structure between El
285 Niño and La Niña conditions is critical. El Niño's westerly wind anomaly patch
286 locates more east than La Niña's easterly wind anomaly patch in their mature stages.
287 As such, the upwelling effects of westerly wind anomalies are accumulated over a
288 wider longitude range and cause stronger negative SLAs in the WTP. As approaching
289 further toward the western boundary, positive SLAs in La Niña continue to amplify,
290 while negative SLAs of El Niño are attenuated by easterly wind anomalies in the far
291 WTP.

292 Here we reveal the sensitivity of the asymmetric SLAs in the WTP to the ENSO
293 wind structures. It is interesting to investigate whether the asymmetric SLAs in turn
294 contributes to the ENSO asymmetry in amplitude and temporal evolutions. This can
295 be investigated through careful heat budget analysis that evaluate the effects of
296 asymmetric current advection on SST variability. In addition, the zonal surface wind
297 patch dominates the SLA asymmetry along the equator, and its off-equatorial structure
298 may affect the asymmetry beyond the equator by modifying the wind stress curl (An
299 & Bong, 2016). In addition to wind forcing, the effects of local nonlinear processes,
300 such as mesoscale eddies, on sea level in the WTP are not resolved by the RGO
301 model experiments (Chen et al., 2015; Qiu et al., 2015), which can be also examined
302 in the future study.

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307 <http://marine.copernicus.eu/services-portfolio/access-to-products/>. Tidal gauges data
308 are available at <https://www.psmsl.org/>. ERA-Interim wind data are available at
309 <https://apps.ecmwf.int/datasets/>. HadISST data are downloaded from Met-Office
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474

475 **Figure captions**

476 **Figure 1.** Standard deviation (STD) of sea level anomalies (SLAs, m) of 1993-2016
477 from (a) AVISO sea level product and (b) HYCOM-CTL. (c) Skewness of SLAs. (d)
478 Regression coefficient $k_{\text{Niño}}$ (m) of SLAs onto the normalized Niño-3.4 index for El
479 Niño condition (Niño-3.4 > 0). (e) Same as (b) but for La Niña condition (Niño-3.4 <
480 0). (f) Ratio of $k_{\text{Niño}}$ to $k_{\text{Niña}}$, i.e., $R_k = k_{\text{Niño}}/k_{\text{Niña}}$. SLA data in (c)-(f) are derived from
481 AVISO sea level product of 1993-2016.

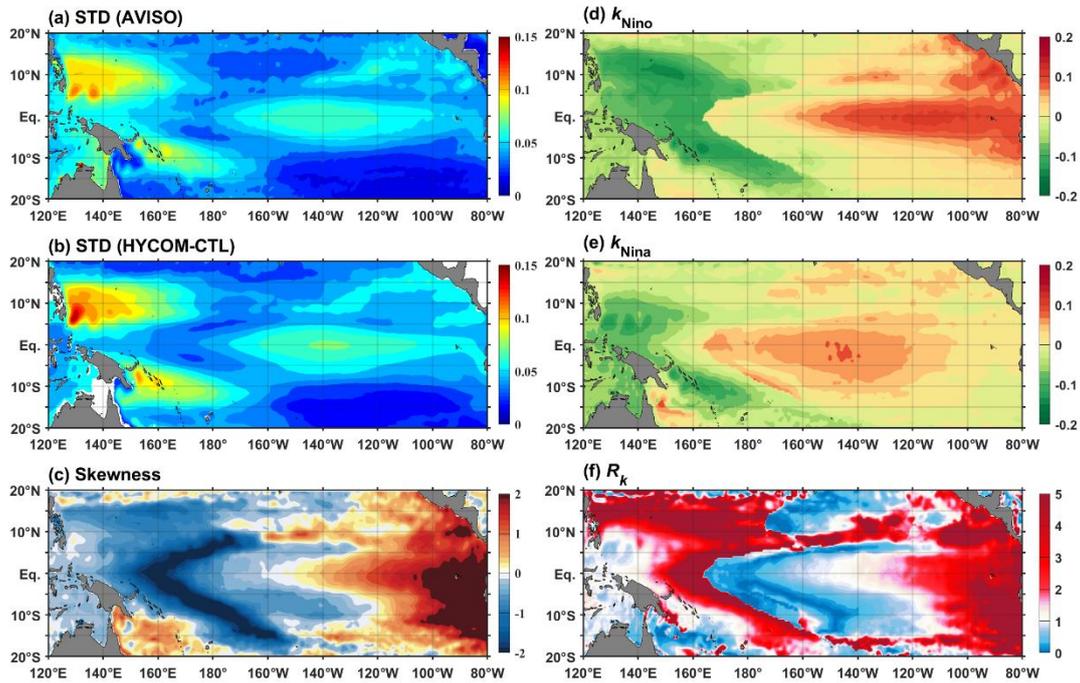
482 **Figure 2.** R_k distributions during 1993-2016 simulated by (a) HYCOM-CTL, (b)
483 HYCOM-TAU, and (c) RGO-CTL. (d)-(f) are the same as (a)-(c), but for the 1979-
484 2016 period.

485 **Figure 3.** R_k distributions from (a) EXP1, (b) EXP2, (c) EXP1-TAUX and (d) EXP2-
486 TAUX. (e) and (f) show the 1st singular value decomposition (SVD) modes for zonal
487 wind stress τ^x (N m^{-2} ; solid and dashed contours for positive and negative values)
488 and upper-layer thickness ULT (m; color shading) of EXP1-TAUX1 for El Niño and
489 La Niña conditions, respectively.

490 **Figure 4.** (a) Zonal structure of the 5°S-5°N average τ^x (left), and ULT (middle) of
491 the 1st SVD mode, and regression coefficients ($k_{\text{Niño}}$ and $k_{\text{Niña}}$; right) in EXP1-TAUX,
492 computed separated for the for El Niño and La Niña conditions. (b) Same as (a) but
493 for EXP2-TAUX. (c) is the same as (a), but for τ^x of the 1st SVD mode in RGO-
494 CTL, and theoretically-predicted ULT and regression coefficients (see the text for
495 details), and they are normalized by the standard deviation.

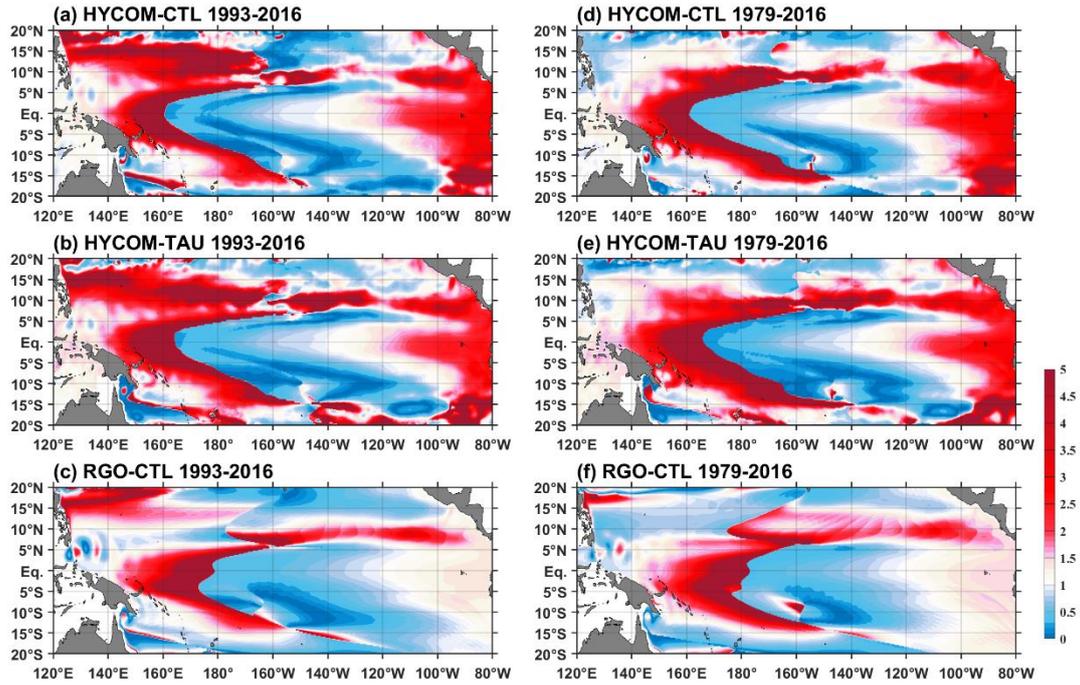
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497 **Figures**



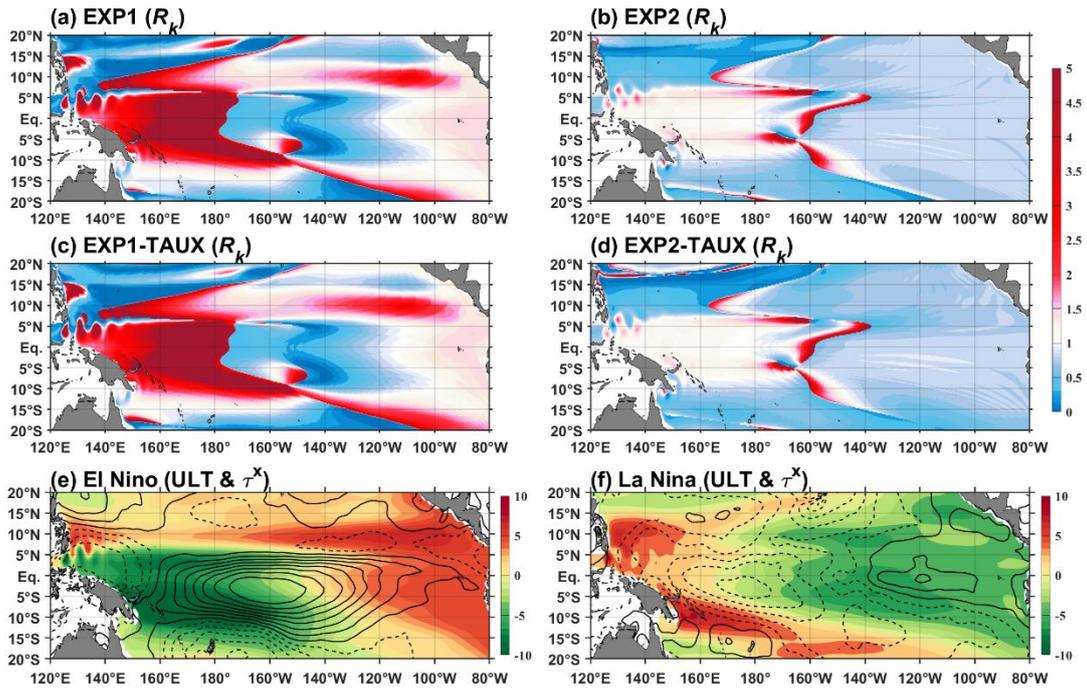
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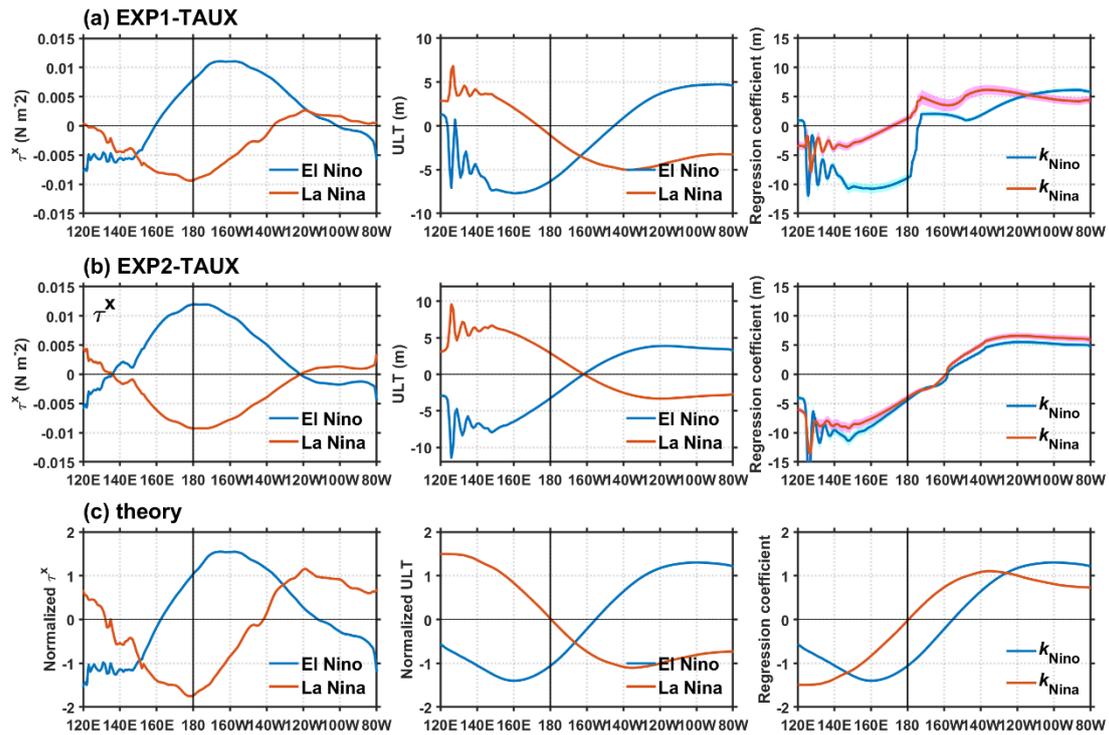
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