

# **Accurate determination of bubble size and expansion ratio for polymer foaming with non-isothermal PBB model based on additional energy conservation**

Yukai Ge, Zhiying Fang, Tao Liu\*

Shanghai Key Laboratory of Multiphase Materials Chemical Engineering, East China University of Science and Technology, Shanghai 200237, P. R. China.

\* Corresponding author: liutao@ecust.edu.cn (T. Liu)

**Abstract:** A non-isothermal pressure-balanced bubble-growth (PBB) model has been proposed based on mass, momentum and energy conservation, which additionally considered the decrease in the internal energy of gas due to the work done by gas expansion in bubble. The model could accurately predict the bubble size and expansion ratio for the melt foaming of four polymers for a wide range of cell densities from  $1.5 \times 10^{13}$  to  $1.9 \times 10^{15}$  cells/m<sup>3</sup>. Furthermore, the simulation results indicate that the bubble shell resisted bubble growth and consumed significant energy, preventing the growth of some small nucleations. During the melt foaming process, the energy cost of the linear polymer had a long-term effect, which reduced the bubble size, while that of the long-chain branched polymer had a short-term effect, thereby increasing the expansion ratio. Finally, we defined the gas efficiency of the foaming agent to evaluate the economic feasibility of the foaming agent in a foaming process.

**Keywords:** Non-isothermal PBB model; Bubble growth simulation; Energy conservation; Bubble size distribution; Gas efficiency

## 1. Introduction

Polymer melt foaming is an important process owing to its scientific and industrial applicability.<sup>1</sup> The density and bubble size distribution (BSD) of polymer foams are critical to foam quality. Several events, such as nucleation, bubble growth, possible coalescence and rupture<sup>2</sup> and final solidification, occur during the foaming process and determine the final density and BSD of the foam.<sup>3</sup> The first nucleation is the initial state of a bubble from which the bubble grows. Nucleation strongly depends on process parameters such as pressure and temperature, especially nucleation agents,<sup>4,5</sup> indicating that nucleation is adjustable in some aspects. Some effective nucleation agents have been used to control cell density, such as nano-CaCO<sub>3</sub>,<sup>6</sup> nano-clay<sup>7-9</sup>, nano-graphite<sup>10</sup> and multi-walled carbon nanotubes (MWCNTs)<sup>10-13</sup>. The nucleation and later bubble growth controlled the final density and BSD, which had considerable impact on the mechanical properties of the foam. For example, Bai et al.<sup>14</sup> found that the compressive strength and hardness of the prepared silicone rubber foams were improved as the cell diameter decreased from 4.97 to 1.12  $\mu\text{m}$ . Kabir et al.<sup>15</sup> used a Zwick/Rowell machine to test the tensile and fracture behaviours of polyvinyl chloride (PVC) foam and rigid polyurethane foam, and found that the tensile strength and modulus were strongly dependent on the foam density.

Visual observation directly facilitates the understanding of the growth of bubbles through visual foam growth experiments<sup>16</sup>. Guo et al.<sup>17</sup> observed the foaming process

of polystyrene in a visualized autoclave and found that bubble growth was extremely sensitive to pressure drop. It only took 0.05 s to complete the bubble growth at an extremely high rate of pressure drop. Wong et al.<sup>18</sup> observed that the bubble growth behaviour changed with the macroscopic matrix deformation. As observed by Tammaro et al.<sup>19</sup>, the main bubble growth process of a bulk sample could last for 20–30 s. Although these observations can help us observe the specific bubble growth process, the experimental phenomena could only provide a limited improvement in the understanding of the foaming mechanism.

Bubble growth determines the structure and properties of foams<sup>1,20</sup>, thus being a crucial step for preparing polymer the foam with a uniform microcell structure. To deeply understand the mechanism of bubble growth, a few numerical models have been proposed.<sup>1,21–27</sup> The bubble growth model can be distinguished according to the region of a single bubble, that is, finite shell and infinite shell. Some scholars<sup>24,26,28</sup> studied diffusion-induced bubble growth in infinite mediums. Because the infinite shell provided an infinite amount of foaming agent, endless bubble growth could be observed. In the case of a finite bubble shell, the bubble stops growing after growing to a certain extent, which depends on the predetermined value of the shell volume. Li et al.<sup>29</sup> simulated the bubble growth of polypropylene (PP) with a finite shell, and their simulation showed that the bubble growth could stop and reach equilibrium. However, in their simulation, the final radius of a bubble was approximately 1250 times larger

than the initial value, which appeared extremely large compared to the actual situation. Feng et al.<sup>1</sup> attempted to predict the bubble size using the finite shell model and described the shell melt using the Oldroyd-B model. The fitting results were in good agreement with the experimental results, but a rapid pressure drop was not investigated, and the coverage was small. Ge et al.<sup>20</sup> proposed a pressure-balanced bubble-growth (PBB) model following the concept of finite volume. The model was unable to determine the final bubble radius effectively if the actual bubble density was used as the input.

Almost all the bubble growth models available assume that the bubble growth process is isothermal, and no in-situ measurement technology is available to challenge this assumption. The terminated bubble size considerably depends on its inner pressure<sup>5,20,30–32</sup>, which is related to the gas temperature and inner energy. The bubble size calculation deviates from the real values if the gas temperature in the bubble is assumed to be unchanged. Thus, it is meaningful to consider the temperature change in the inner bubble<sup>22,33,34</sup>. Moreover, the bubble growth also determines the expansion ratio, and a correct bubble size calculation can result in an accurate theoretical expansion ratio. Therefore, an accurate bubble size calculation is of great significance, wherein the inner bubble temperature is one of the key factors.

In the present study, we proposed a non-isothermal PBB model from the viewpoint of mass, momentum, and energy conservation, which considered the

decrease in the internal energy of the gas owing to the work done performed by gas expansion. The model was used to predict the bubble size and expansion ratio for the melt foaming of four different polymers in a wide range of cell densities. The gas efficiency of the foaming agent was defined as Gas efficiency was defined as the changing rate of expansion ratio to the foaming pressure, providing a method to evaluate the economic feasibility of the foaming agent in the foaming process.

## **2. Experiments**

### **2.1 Materials**

The foaming behaviours of four types of materials were studied in this work. Pellets of low-density polyethylene (LDPE) (DFDA-1253NT) were supplied by Dow Chemical Co. Ltd., China. linear (L) structure polypropylene (PP) (T1903) was provided by Sinopec Maoming Petrochemical Co. Ltd., China. The long-chain branched (LCB-) PP (WB140) with a high melt strength was supplied by Borealis Co. Ltd. The LCB-PET was produced from a fibre-grade PET supplied by Xinshan Petrochemical Co. Ltd., China<sup>20</sup>. CaCO<sub>3</sub> nanospheres 40 nm in radius, tube-shaped MWCNTs 20 nm in length, lamellar nano graphite 40 nm in thickness, and nano clay 25 nm in thickness were supplied by XFNANO Chemical Co, Nanjing, China. CO<sub>2</sub> (purity: 99.9%, w/w %) and N<sub>2</sub> (purity: 99.9%, w/w %) were obtained from Air Production Co., Shanghai, China. All the chemical agents were used as received.

## 2.2 Molar mass characterization

The molar mass distribution was determined by size exclusion chromatography (Waters 1515). A linear polystyrene PS with a molar mass of  $75 \text{ kg}\cdot\text{mol}^{-1}$  was used for calibration. The samples were pulverized and dissolved in hexafluoroisopropanol (HFLP) at a concentration of 4 g/L. All measurements were performed at 35 °C and a constant flow rate of 1.0 mL/min. **Table 1** lists the molecular information of the four polymers, and the details are given in the **Supporting Information**.

**Table 1.** Molecular structure parameters, CO<sub>2</sub> solubility, and interface parameters of the four polymers

Sample	M <sub>w</sub>	M <sub>n</sub>	M <sub>w</sub> /M <sub>n</sub>	LCB	γ	H <sub>a</sub> (g CO <sub>2</sub>
	(kg/mol)	(kg/mol)			(Pa/m)	/g melt/Pa)
LDPE <sup>a</sup>	107	25	11	very high	0.015	$3.9 \times 10^{-9}$
LCB-PP <sup>a</sup>	609	80	8	high	0.013	$5.6 \times 10^{-9}$
L-PP <sup>a</sup>	419	73	6	none	0.013	$5.6 \times 10^{-9}$
LCB-PET <sup>b</sup>	65	28	2	very high	0.020	$1.8 \times 10^{-9}$

<sup>a</sup>Molecular weight information was measured in this work.

<sup>b</sup>Molecular weight information was got from the reference 20.

## 2.3 Rheological characterization in shear and elongation

The linear viscoelasticity of the four polymers was measured using a Haake Mars III rheometer with a 35 mm plate fixture. Oscillation tests at a frequency range of 100–0.1 rad/s and at different temperatures were conducted in a nitrogen atmosphere.

The testing temperatures were 110, 130, and 160 °C for LDPE; 150, 180, and 200 °C for PP; and 265 °C for PET. The extensional viscosity of the four polymers was measured using a Sentmanat extensional rheometer (SER) fixture. Samples used in the measurements were cut into a shape of 20 mm length and 8 mm width.

The rheological behaviours of each polymer were also carefully modelled. In **Table 2**, three important parameters of each polymer are listed, including the zero shear viscosity and the two free factors that were used to describe strain hardening. In addition, the specific viscoelasticity analyses are presented in the **Supporting Information**. The molecular stress function (MSF) was used to describe the tensile behaviour of the melt. The details of the MSF model are available in the **Supporting Information**.

**Table 2.** Zero shear rate viscosity and non-linear parameters in MSF of the four polymers

Sample	$\eta_0^i$ (Pa·s)	$B$	$f_{max}^2$
LDPE	141000	1.8	80
LCB-PP	128000	1.6	100
Linear-PP	31959	1.0	4
LCB-PET	64970	2.4	120

## **2.4 Physical mixing of LDPE and different nucleation agents**

The physical mixing of LDPE and 0.5 wt.% nucleation agents was conducted in a twin-screw extruder (Nanjing Giant SHJ-20). The screw length was 600 mm and the diameter was 20 mm. The speed of rotation was 350 rpm to obtain a strong shear for good blending. The temperature was set at 200 °C from the feeding to the die, and the material was fed at a rate of 10 kg/h. The nucleation agents added were CaCO<sub>3</sub> nanospheres, tube-shaped MWCNTs, lamellar nanographite, and nanoclay, respectively.

## **2.5 Batch foaming process at melt state**

Batch foaming was conducted using a high-pressure vessel equipped with a pressure transducer. The high-pressure CO<sub>2</sub> used for foaming was provided by a TELIDE S-486-JN-60 pump. The temperature of the foaming system was controlled with an oil bath ( $\pm 0.5$  °C). The samples used for foaming were first heated to 150 °C for LDPE, 180 °C for LCB-PP, and 265 °C for LCB-PET. Then, the samples were pressurised with CO<sub>2</sub> at a pressure range of 10–20 MPa at the same temperature held for 30 min. Subsequently, the system pressure was kept constant, and the processing temperature was reduced to the foaming temperature at which it was held for 60 min. The foaming temperatures were 110 °C for LDPE, 150 °C for LCB-PP, and 265 °C for LCB-PET. Thereafter, the pressure was rapidly released with a depressurization rate of approximately 100 MPa/s, leading to the formation of foamed cells. All

foaming experiments were conducted thrice under the same conditions.

## **2.6 Morphological analysis of polymer foams**

Foaming samples were cooled in liquid nitrogen for 10 min and then fractured to achieve a cross-section. The cell morphology of the foam was observed using a scanning electron microscope (SEM) NOVE NanoSEM450 (FEI, USA). The information on cell size distribution was obtained through SEM photographs with the software Image-Pro Plus.<sup>35</sup>

## **3 Numerical simulation**

### **3.1 Non-isothermal PBB model**

The non-isothermal PBB model originated from the cell model<sup>25</sup>, where the bubble shell was assumed to be a viscoelastic fluid having a finite volume. The bubble shell could hold a certain amount of foaming gas. The dissolved gas was transferred into the bubble and drove the bubble growth. In this process, three conservations including mass, momentum, and energy were involved, based on which, the non-isothermal PBB model was established. The isothermal assumption in the previous cell model required the total system including the bubble shell and inner gas to be isothermal. However, to achieve an accurate cell size calculation, a non-isothermal bubble growth model was necessary in which the change in gas temperature inside the bubble had to be considered. Meanwhile, the bubble shell

(solid phase) could be assumed to be isothermal for two reasons. First, in a polymer foaming bubble system, the mass of gas was more than one order of magnitude smaller than that of the polymer. For example, the solubility of CO<sub>2</sub> in LDPE was only 0.068 g gas/g LDPE<sup>32</sup>. Thus, the total inner energy of the gas was much smaller than that of the polymer, indicating that the temperature of the gas phase is more sensitive. Second, the polymer matrix deformation and friction also generated heat, which could offset the cooling caused by the gas. Moreover, the internal energy loss of the gas caused immediate temperature dropping of gas. When a gas expands and performs external work, the internal energy and temperature reduce instantly. Thus, it was reasonable to consider only the temperature change in the gas phase while the bubble shell remained isothermal. All the model assumptions are listed here:

- (a) The bubble shell was isothermal, but the inner gas was non-isothermal.
- (b) The bubble was spherical.<sup>3</sup>
- (c) The polymer melt was incompressible.<sup>3,24,25,27,29</sup>
- (d) The effect of plasticisation was ignored.<sup>29</sup>
- (e) The effects of inertial forces and gravity could be neglected.<sup>24</sup>
- (f) The bubble had no mass transfer with the environment, indicating that the bubble was a closed system.<sup>29</sup>
- (g) The bubble growth rate followed the form of the Hencky strain, and the main

growth process was at a constant rate.<sup>20</sup>

- (h) The inner pressure of the bubble depended on the amount and temperature of the gas inside the bubble, and the equilibrium vapor pressure of the polymer-gas solution was the driving force of mass transfer.<sup>20</sup>

### 3.2 Mass conservation

The mass of CO<sub>2</sub> in a closed bubble system follows the principle of mass conservation during bubble growth.

$$n_{CO_2(\text{system})} = n_{CO_2(\text{inner bubble})} + n_{CO_2(\text{dissolved})} \quad (1)$$

$$n_{CO_2(\text{system})} = H_a \cdot P_0 \cdot V_0 \quad (2)$$

$$n_{CO_2(\text{inner bubble})} = (P_i \cdot \left(\frac{4}{3} \pi R_i^3\right)) / (R T_i) \quad (3)$$

$$n_{CO_2(\text{dissolved})} = c \cdot V_0 = H_a \cdot P_{eo} \cdot V_0 \quad (4)$$

$n_{CO_2(\text{system})}$  represents the entire CO<sub>2</sub> in the bubble system,  $n_{CO_2(\text{dissolved})}$  the CO<sub>2</sub> dissolved in

the bubble shell, and  $n_{CO_2(\text{inner bubble})}$  the CO<sub>2</sub> molar inner the bubble.  $R_i$  is the inner bubble

radius,  $V_0$  is the volume of the finite bubble shell,  $P_0$  is the foaming pressure, and  $P_{in}$  is the bubble inner pressure.  $T_{in}$  is the bubble inner temperature.  $P_{eo}$  represents the

equilibrium vapor pressure of the CO<sub>2</sub>/polymer solution, which can be described by Henry's law:

$$c = H_a \cdot P_{eo} \quad (5)$$

where  $H_a$  is the Henry's coefficient of the CO<sub>2</sub>/polymer solution, and  $c$  is the instantaneous concentration of CO<sub>2</sub>. The final mass conservation equation is as follows:<sup>5</sup>

$$H_a \cdot P_0 \cdot V_0 = \frac{P_i}{RT} \cdot \frac{4}{3} \pi R_i^3 + H_a \cdot P_{eo} \cdot V_0 \quad (6)$$

### 3.3 Momentum conservation

The principle of momentum conservation was satisfied during the bubble growth process. There were two differences compared to the traditional cell model<sup>23–25,29</sup>. First, the bubble growth rate was determined as a parameter of 5 s<sup>-1</sup>.<sup>20</sup> In the rapid pressure dropping situation, the bubble growth time was short and its rate could be assumed to be a constant so that the corresponding bubble growth time was about 0.2–0.3 s.<sup>20</sup> This time scale was close to the actual situation<sup>17</sup> in that, the bubble growth stage was less than 0.4 s under a rapid pressure relief rate. Second, the MSF model was used to calculate the elongational stress  $\sigma$  of the bubble shell. Despite the existence of small differences, both the form and meaning of the momentum equation were almost unchanged compared to the previous work.<sup>23–25,29</sup> The momentum equation of bubble growth is given by:

$$P_i - P_e - \frac{2\gamma}{R_i} = 2 \int_{R_i}^{R_{out}} \frac{\sigma_r + \eta^i \dot{\gamma}_r}{r} dr \quad (7)$$

where  $P_e$  is the environmental pressure. The interface tension  $\gamma$  of the two phases lead to increased pressure inside the bubble.  $R_{out}$  is the outer radius of the bubble, where the shell thickness is included.  $\eta^i$  represents the melt viscosity, and  $\dot{\gamma}$  represents the shear rate.

The LHS of the equation represents the total driven force, and the RHS is the resistance caused by the bubble shell. The growth resistance contains two different terms, that is, elongational stress and shear friction. Shear friction can be calculated by multiplying the viscosity and shear rate. The elongational stress  $\sigma_r$  of the bubble shell was calculated using the MSF<sup>36–39</sup>. Although other viscoelastic models were also available in stress calculations and were used in previous work, such as the upper-convected Maxwell<sup>28</sup> and Oldroyd-B models<sup>1</sup>, the MSF possessed a few advantages—first, strain hardening was well described, which is a crucial property in foaming<sup>5,40,41</sup>. Second, a good connection between the molecular structure and the elongational properties was available. Härth et al.<sup>42</sup> used the MSF to model the elongational behaviour of LCB-PET and obtained a quantitative relationship between the branches and strain hardening. Ahirwal et al.<sup>43</sup> modeled the elongational behaviour of both PP and PE, and the MSF in both polymer systems showed accurate results. Therefore, the MSF model was selected for modelling the shell resistance of the matrix elongation.

### 3.4 Energy conservation

Bubble growth is a process of energy consumption, in which the gas expansion follows the first law of thermodynamics and can be regarded as a Carnot cycle<sup>44</sup>. When a high temperature gas is expanded outward, the gas temperature is decreased. The gas performs external work during bubble growth, and this energy was consumed in matrix friction and stretching. Simultaneously, the inner energy of the bubble was supplemented because of the continuous evaporation of the foaming agent gas from the polymer matrix to the inner bubble. The foaming agent that was dissolved in the bubble shell was the energy source for bubble growth, and the stored energy decreased continuously because of the diffusion and evaporation. The bubble stopped growing when the system energy approached zero. The energy consumed in the bubble growth process can be regarded as the energy cost to change the polymer shape, that is, from the solid state to the foam form. The expanding energy in the bubble blowing process can be directly provided by the gas inside the bubble:

$$dW_{\text{blow}} = P_i dV \quad (8)$$

$$P_i dV = dU = nc_v dT_i \quad (9)$$

$$c_v = \frac{i}{2} R_g \quad (10)$$

where  $d$  indicates a differential.  $W_{\text{blow}}$  is the work provided by the inner gas because of its expansion, and  $V$  is the bubble volume. The loss in gas inner energy  $U$  leads to a

decrease in the gas temperature  $T_{in}$ , and  $R_g$  is the gas constant.  $c_v$  is the specific heat capacity of the ideal gas, and for CO<sub>2</sub> with three atoms in one molecule, the parameter  $i$  is 6 according to the ideal gas theory.

A polymer/gas solution can provide high-pressure vapor, and, evidently, a certain amount of energy. The internal energy of the gas in the bubble was supplemented because of gas evaporation and diffusion from the bubble shell. During this period, the concentration of the gas/polymer solution as well as its energy, decreased. We quantified this energy and referred to it as the system energy  $E_{system}$ .

$$E_{system} = c H_a V_0 \quad (11)$$

$H_a$  is the Henry's coefficient,  $c$  is the concentration, and  $V_0$  is the shell volume. The system energy decreased with a reduction in the gas concentration:

$$dE_{system} = H_a V_0 dc \quad (12)$$

The mass transfer stops when the system energy becomes zero. Energy can be represented by pressure. The system energy  $E_{system}$  related to the vapor pressure  $P_{eo}$  is given by:

$$dE_{system} \cdot V_0 = P_{eo} \quad (13)$$

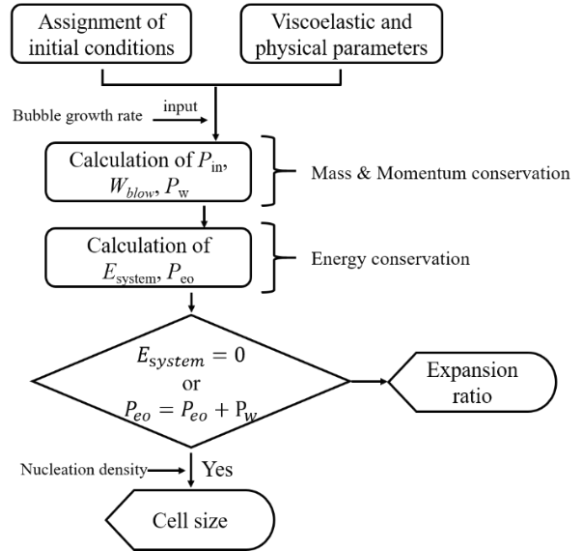
The blowing work ( $W_{blow}$ ) leads to a temperature drop and pressure drop represented as  $P_w$ .

$$W_{\text{blow}} = \Delta n c_v T_c = \frac{c_v}{R_g} V \Delta P = \frac{c_v}{R_g} V P_w \quad (14)$$

$P_w$  was ignored in traditional bubble growth models with an isothermal assumption. These models believed that the bubble stopped growing when the inner pressure equalled the vapor pressure i.e.  $P_{\text{eo}} = P_{\text{in}}$ . However, in the non-isothermal situation, the criterion for stopping the growth is  $P_{\text{eo}} = P_{\text{in}} + P_w$ .

### 3.5 Calculation procedure of bubble growth using non-isothermal PBB model

A flow chart of the non-isothermal PBB model is plotted for the calculation process specification. The initial conditions include the foaming temperature, pressure, and basic physical parameters, including Henry's constant and surface tension. The relaxation properties and strain hardening parameters in MSF are also required to describe the bubble shell deformation. Bubble growth simulation was conducted to obtain the values of  $P_{\text{in}}$ ,  $W_{\text{blow}}$ ,  $P_w$ ,  $P_{\text{eo}}$ , and  $E_{\text{system}}$ . The expansion ratio and cell size were determined at end of the growth process. Veracious nucleation information was needed to determine the accurate cell size. Notably, the expansion ratio can be determined by using an assumed nucleation condition as well.



**Figure 1.** Computational flow chart of non-isothermal PBB model.

## 4. Results and discussion

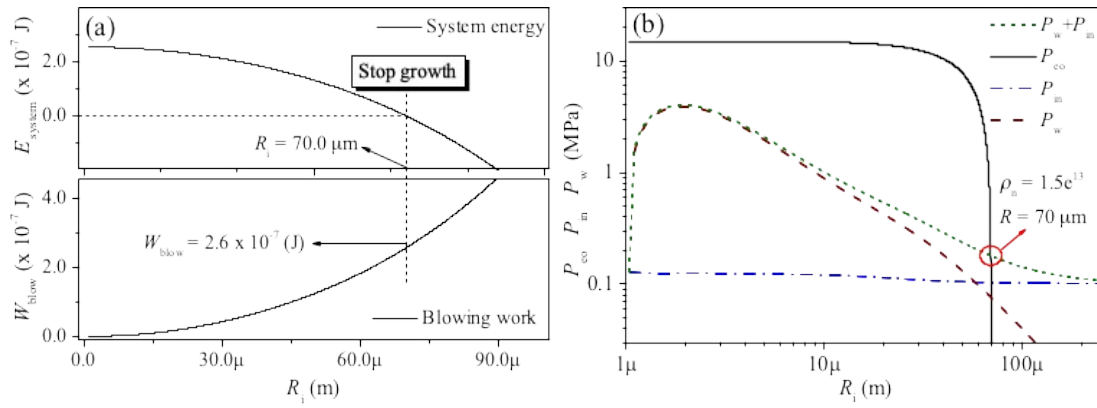
### 4.1 Bubble growth simulation of LDPE foaming using non-isothermal PBB model

#### 4.1.1 Bubble growth simulation of neat LDPE foaming

The batch foaming simulation and experiments using LDPE were conducted at 110 °C and 15 MPa CO<sub>2</sub>. All the inputs of the simulation were determined by actual experiments. The cell density was determined to be  $1.5 \times 10^{13}$  cells/m<sup>3</sup> according to the SEM image shown in **Figure S4(a)** in the **Supporting Information**.

**Figure 2 (a)** shows the variations in system energy and blowing work during bubble growth. The simulation results were suitable for cell size calculation because the bubble growth stopped when the system energy reached zero. This occurred when the bubble grows to a radius of 70.0 μm. Accordingly, the theoretical bubble radius

was 70  $\mu\text{m}$ , and the total blowing energy of one bubble was  $2.6 \times 10^{-7}$  (J). In **Figure 3**, the experimental results showed that the actual bubble diameter was 148 ( $\mu\text{m}$ ), i.e.,  $R_i = 74 \mu\text{m}$ , which agreed well with the theoretical calculation. If the bubble system comprised more energy, the growth could continue and obtain a higher expansion ratio. For example, Zhang et al.<sup>32</sup> used a mixture of  $\text{CO}_2$  and  $\text{i-C}_4\text{H}_{10}$  as foaming agents. The cell diameter and expansion ratio increased with the amount of  $\text{i-C}_4\text{H}_{10}$ . It was because  $\text{i-C}_4\text{H}_{10}$  could store more gas and energy compared to  $\text{CO}_2$  leading to a further bubble growth.



**Figure 2.** Simulation results of neat LDPE using non-isothermal PBB model. (a) Energy variations. (b) Pressure variations.

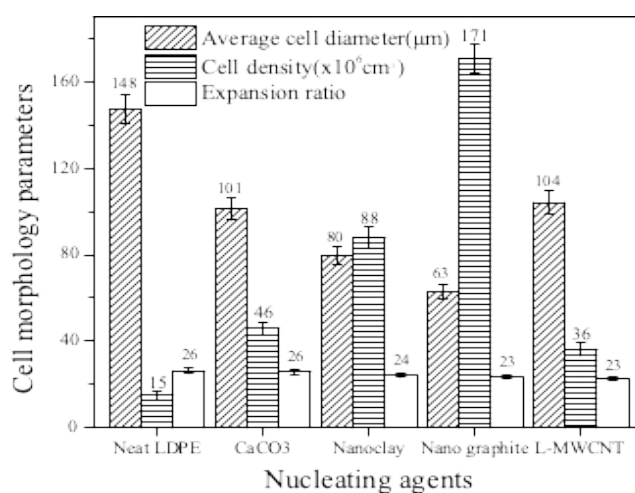
In **Figure 2(b)**, the simulation of bubble pressure can provide a detailed analysis of the bubble growth process. Three transfer behaviours, that is, energy, mass, and momentum, work together and determine the bubble growth behaviour. Three types of pressures are included—bubble inner pressure ( $P_{\text{in}}$ ), gas phase equilibrium pressure ( $P_{\text{co}}$ ), and decreased pressure because of the gas expansion work ( $P_w$ ). The  $P_{\text{co}}$  curve is

mainly related to mass conservation, and  $P_{eo}$  declines after evaporation and the diffusion of  $\text{CO}_2$  from the bubble shell. The  $P_{in}$  is controlled by momentum conservation, and energy conservation determines the  $P_w$  curve. A theoretical bubble radius can be found at the crossover point of the two curves, that is, the  $P_w + P_{in}$  and  $P_{eo}$  curves. The radius calculated in the pressure form is the same as that calculated in the energy form.

As shown in **Figure 2(b)**,  $P_w$  increases rapidly at the initial bubble growth stage and then decreases slowly and continuously. The existence of the peak in the  $P_w$  curve at approximately 2  $\mu\text{m}$  in radius has two reasons. First, the initial bubble was of micro dimensions, which meant that the bubble volume was small and the bubble shell was very thick, indicating that the resistance of bubble growth would be extremely large. Second, bubble growth consumed the internal energy of the gas, and the energy cost was related to resistance. The small amount of the internal energy of the gas was required to provide a large amount of energy at the initial bubble-growth stage, which led to a fast pressure drop corresponding to the peak in the  $P_w$  curve. Both the attenuation of the bubble shell and the increase in the gas amount weakened the influences of the expansion work. Thus, the  $P_w$  curve shows that the effects of gas inflation work and internal energy change were significant at the initial bubble growth stage, but would be mitigated later. Additionally, the gas temperature inside the bubble reveals the effects, as shown in **Figure S8** in the **Supporting Information**.

#### 4.1.2 Bubble growth simulation of LDPE foaming under different nucleation conditions

Four different nucleation agents were used to change the density of the LDPE foams in the experiments. Foaming experiments were conducted under the same conditions as the neat LDPE, and the BSD of each foam was characterized according to the SEM images shown in **Figure S4** in the **Supporting Information**. **Figure 3** shows the morphology statistics of the different LDPE foams. The cell density varies more than one order of magnitude from  $1.5 \times 10^{13}$  cells/m<sup>3</sup> to  $1.7 \times 10^{14}$  cells/m<sup>3</sup> because of the use of different nucleation agents. Although the cell size changes significantly with cell density, the expansion ratio appears to be unaffected. The expansion ratios of the different LDPE foams vary around 24.

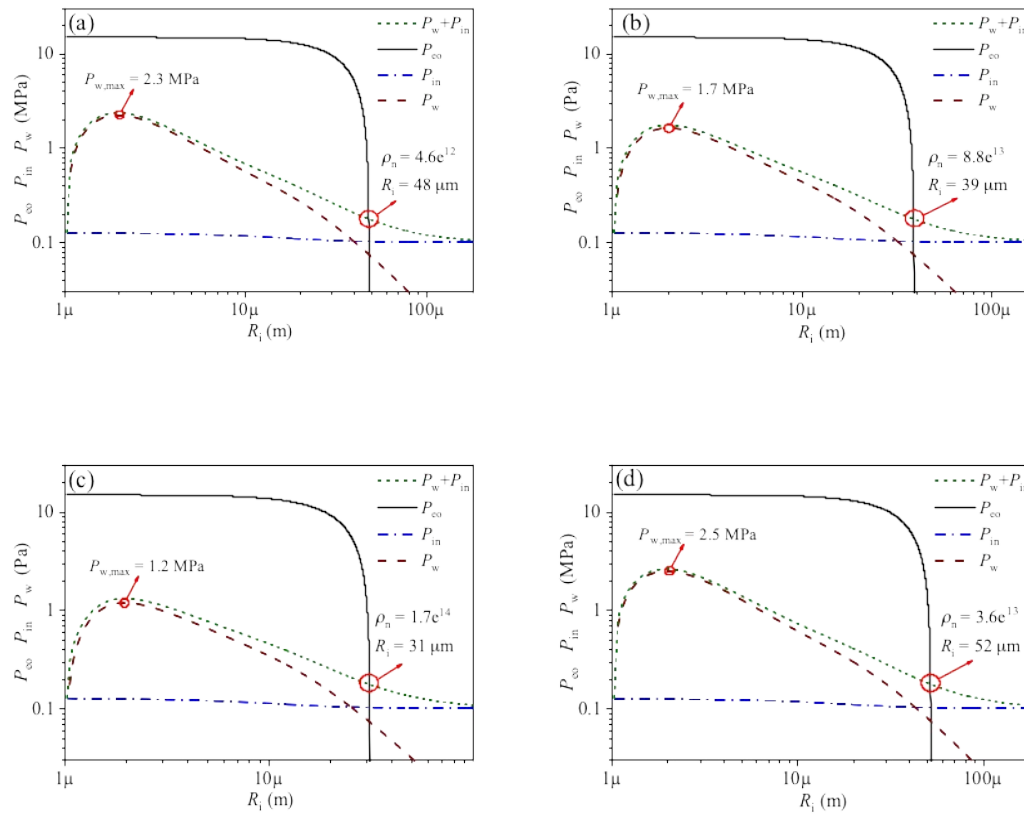


**Figure 3.** Morphology statistics of LDPE foams with different nucleation agents.

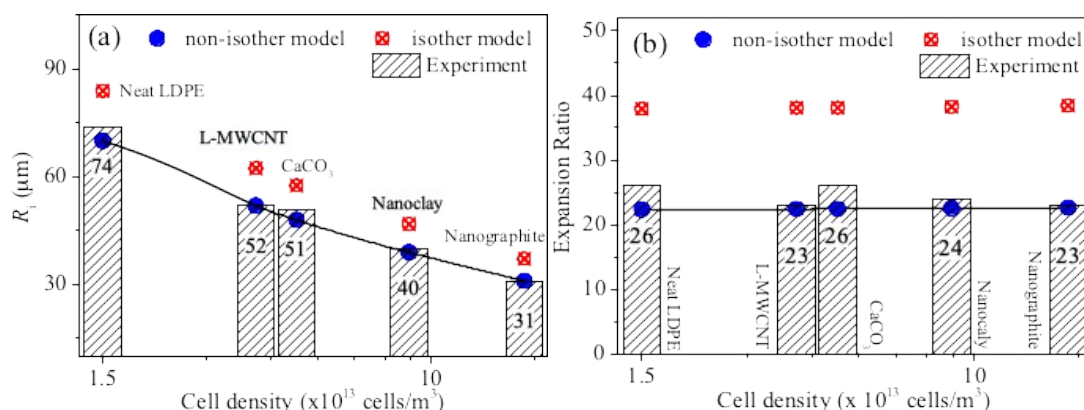
The cell density was updated in the simulation according to the SEM images.

The theoretical bubble growth processes were calculated, and the results are plotted in **Figure 4**. In addition to predicting the bubble size, two more meaningful results can be obtained by bubble growth simulation. First, according to the calculated bubble radius, the expansion ratio calculation is also available. The calculation results are consistent with the foaming experiments for both the average bubble size and expansion ratio, as shown in **Figure 5**. If the energy conservation is not considered, that is, in the isothermal situation, the calculated radius and expansion ratio will be much larger than the actual values. The calculated results, without considering energy conservation, contained significant errors in the expansion ratio up to approximately 58 %. However, the accuracy in both the bubble size and expansion ratio predictions was greatly improved by considering energy conservation. It was proven that the bubble growth is a non-isothermal process, and the energy cost of growth should be considered. Second, a high peak in the  $P_w$  curve appears at low nucleation density, indicating a considerable resistance in bubble growth. The nucleation density not only affects the potential cell number but also corresponds to the bubble growing resistance, which was not found in previous bubble growth simulations because the energy conservation was not considered. Correspondingly, a large nucleation size favours bubble growth and increases the final cell density. Accordingly, the ideal nucleation agents should satisfy two requirements, that is, providing a large number of nucleation points and a large nucleation size. Some early experimental phenomena could be explained by the simulation results. Chen et al.<sup>45</sup> found that the nucleation

agent, MWCNTs, with a small size had low nucleation efficiency at the low foaming pressure, but the density was improved after increasing the foaming pressure. Meanwhile, the large size of MWCNTs had a high nucleation efficiency at all foaming pressures. According to the simulation, a significant peak in the  $P_w$  curve appears if the nucleation size is small, indicating that it is challenging for the nucleation point to grow. In addition, the effects of  $P_w$  decrease after increasing the foaming pressure (related to  $P_{e0}$ ) as the entire driven force is  $P_{e0} - P_w$ . Therefore, in the experiments conducted by Chen et al.<sup>45</sup>, the high growth resistance, instead of fewer nucleation points, may be attributable for the low cell density when small MWCNTs were used.



**Figure 4.** Bubble growth simulation for LDPE samples with nucleation agents. (a)  $\text{CaCO}_3$ /LDPE, (b) Nano-Clay/LDPE, (c) Nano-graphite/LDPE and (d) MWCNTs/LDPE.



**Figure 5.** Comparison of foaming results and simulations with isothermal and non-isothermal models. (a) bubble radius, (b) expansion ratio. The lines only indicated the trend.

#### 4.2 Bubble growth simulations of L-PP, LCB-PP, and LCB-PET foaming using non-isothermal PBB model

To explore the validity of the non-isothermal PBB model, bubble growth simulations for foaming of three different polymers— L-PP, LCB-PP, and LCB-PET — were also performed. Furthermore, foaming experiments were conducted, and the micromorphology of each foam was carefully analysed. The cell sizes were measured according to the SEM images shown in **Figure S5** of the **Supporting Information**. It was found that the BSD met a nearly normal distribution, as shown in **Figure S6** in the **Supporting Information**. Bubble growth simulations were conducted, in which all the inputs were equal to the exact conditions for the foaming experiments. The results, including the predicted bubble radius  $R_p$ , the exact cell size  $R_i$ , expansion ratio, and cell density  $\rho_n$  are listed in **Table 3**.

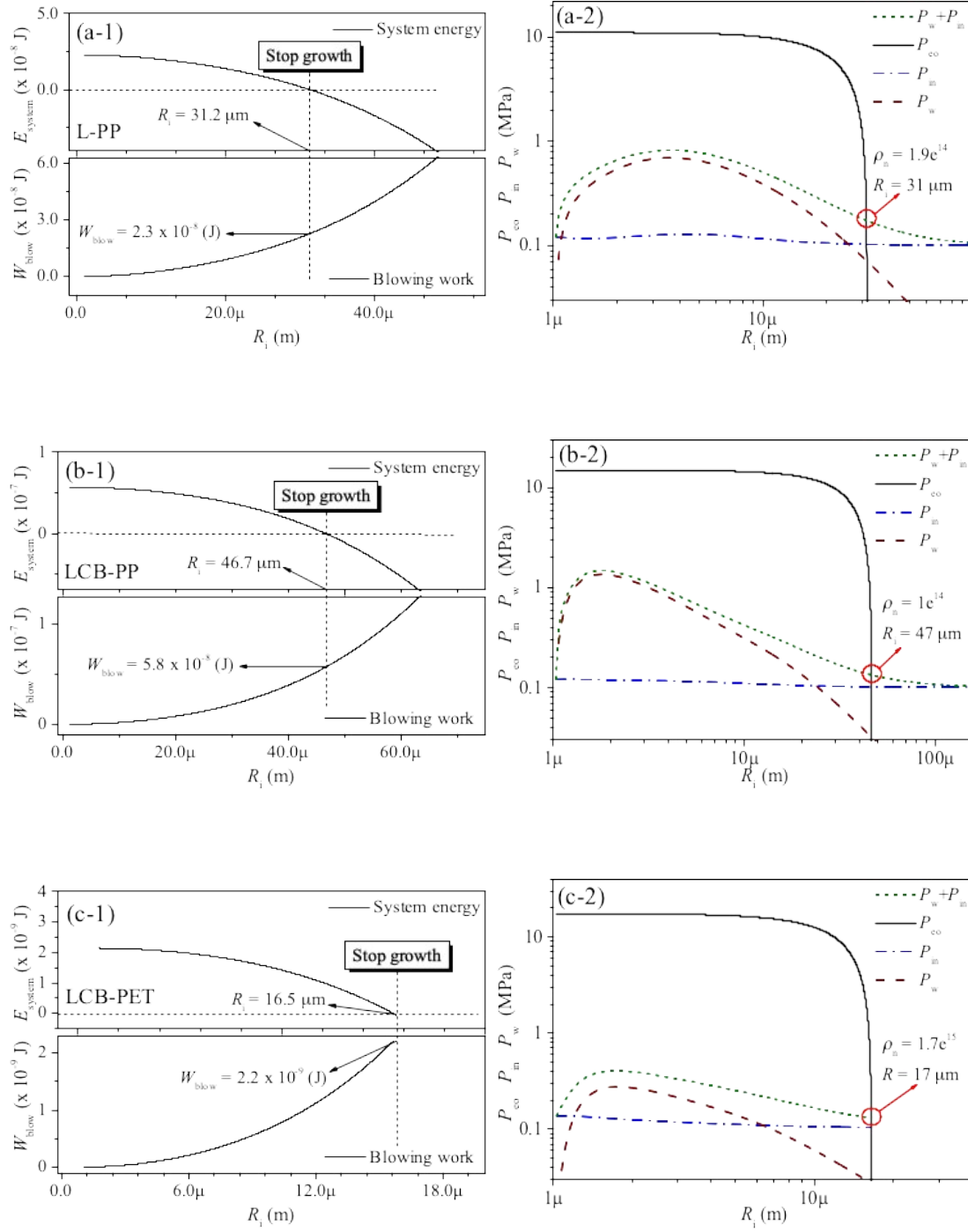
**Table 3.** Foaming conditions, morphology statistics of foams, and the model calculated results

Sample	T	P	$\rho_n^a$	$R_i^a$	$R_p^b$	Expansion	Gas efficiency <sup>b</sup>
	(°C)	(MPa)	(cells/m <sup>3</sup> )	( $\mu\text{m}$ )	( $\mu\text{m}$ )	Ratio <sup>a&amp;b</sup>	(MPa <sup>-1</sup> )
LDPE	110	15	$1.5 \times 10^{13}$	71	70	26 & 22	1.8
LCB-PP	150	15	$1.0 \times 10^{14}$	47	47	40 & 45	3.5
L-PP	150	12	$1.9 \times 10^{14}$	30	31	26 & 26	2.0 <sup>1st</sup> , 11 <sup>2st</sup>
LCB-PET	265	18	$1.7 \times 10^{15}$	17	17	32 & 33	2.1

<sup>a</sup>Results obtained from experiments.

<sup>b</sup>Results obtained from simulations.

The termination of bubble growth can be easily determined when the system energy becomes zero, as shown in **Figures 6(a-1), (b-1), and (c-1)**. The blowing energy  $W_{\text{blow}}$  of each bubble is obtained. As shown in **Figure 6(c-1)**, LCB-PET has the lowest system energy and its blowing work is only  $2.2 \times 10^{-9}$  J/cell, which is one magnitude smaller than that of L-PP and LCB-PP, as shown in **Figures 6(a-1) and (b-1)**. The energy of LCB-PET is limited by the solubility of CO<sub>2</sub> in PET. As shown in **Table 1**, it has the lowest Henry's coefficient. Therefore, more system energy was required to obtain PET foam with a larger expansion ratio. The second reason for the low  $W_{\text{blow}}$  was that the cell density of PET was extremely high, and the bubble size was small.



**Figure 6.** Simulations of polymer foaming using non-isothermal PBB model in energy form: (a-1) L-PP, (b-1) LCB-PP, (c-1) LCB-PET; and in pressure form: (a-2) L-PP, (b-2) LCB-PP, (c-2) LCB-PET.

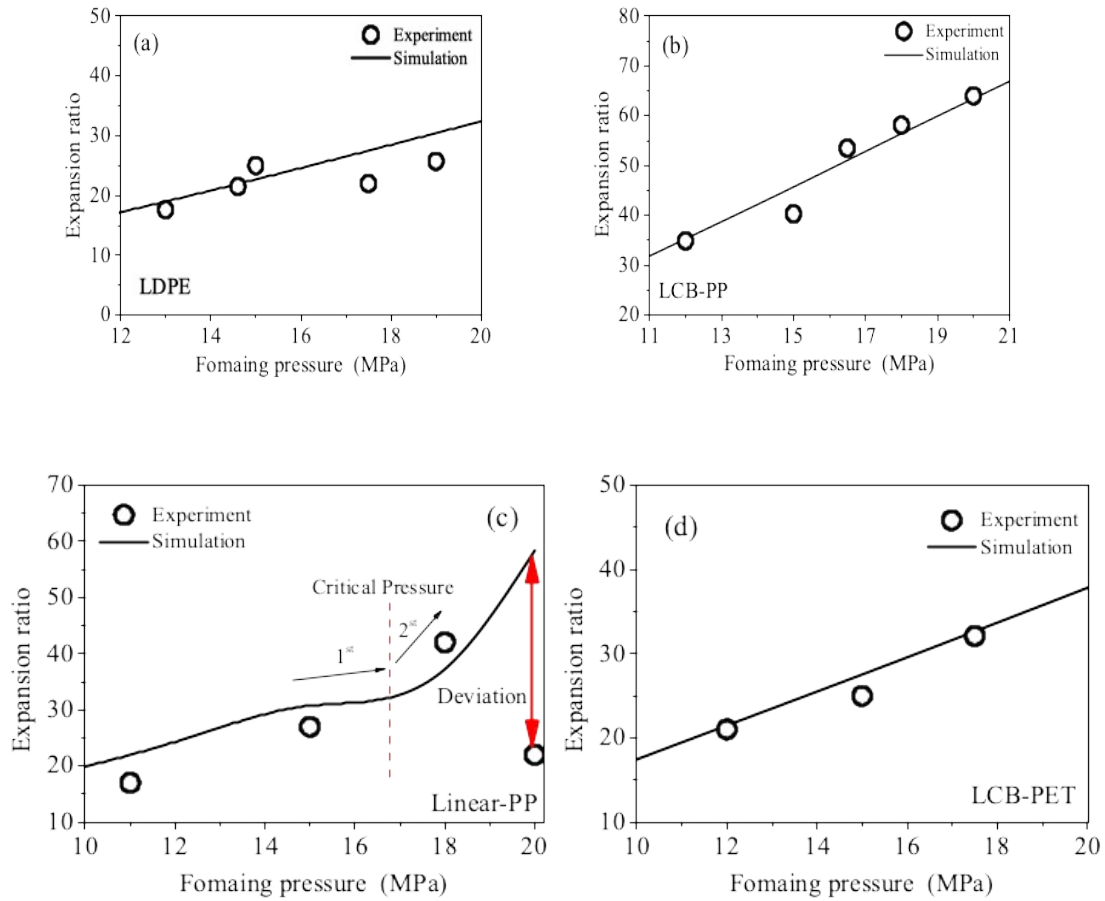
The  $P_w$  curve has different shapes for the three polymers, which affects the

bubble growth behaviours.  $P_w$  was related to the energy loss owing to the growth resistance. For LCB-PP and LCB-PET, the corresponding  $P_w$  curves decrease rapidly, indicating that bubble growth faced a weak resistance. L-PP shows a near-elliptic  $P_w$  curve, which means that the inner energy cost affects bubble growth for a long period. The difference in the  $P_w$  curve is related to the polymer chain structure because linear and LCB structures have different elongational behaviours, as shown in **Figure S3** in the **Supporting Information**. The L-PP is more viscous and has a relatively flat elongational viscosity, while the LCB-PP has a low viscosity at the initial bubble growth stage and shows a significant increase in viscosity after stretching. Low initial viscosity indicates a weak resistance at the beginning of bubble growth. The subsequent increase in the elongational viscosity is beneficial to the foaming stability when the bubble wall becomes thin and weak after bubble growth. Therefore, a relatively low initial viscosity and a significant viscosity increase are the ideal rheological properties for the foaming process. Therefore, the LCB structure is more suitable for foaming applications than linear structure in both process stability and obtaining a high expansion ratio.

#### 4.3 Gas efficiency in melt foaming of four polymers

In the foam industry, it is usually desirable to produce foams with a high expansion ratio under the lowest foaming pressure, so as to achieve profit and safety. However, the relationship between the expansion ratio and foaming pressure is often

unknown. To achieve an overall understanding of this relationship, gas efficiency (CO<sub>2</sub> in this work) was proposed in this work. Gas efficiency was defined as the changing rate of expansion ratio to the foaming pressure.



**Figure 7.** Determination of expansion ratio at different foaming pressures in both experiments and simulations. (a) LDPE, (b) LCB-PP, (c) L-PP, (d) LCB-PET.

Both the simulations and experiments were conducted in a pressure range of 11 to 20 MPa. In the simulations, only the expansion ratio was focused on, while the inputs of nucleation density, an ordinary value of  $1 \times 10^{13}$  cells/m<sup>3</sup>, was assumed for all the samples. The simulated expansion ratios as a function of foaming pressure are

plotted in **Figure 7**. The results of the experiment and simulation show good agreement, which proves that the non-isothermal PBB model is an effective method for analysing the expansion ratio. It should be noted that the prediction at extreme high foaming pressures, such as above 20 MPa, contains errors because of plasticization<sup>4,46</sup>.

The curve of the expansion ratio shows an almost linear shape for all the LCB-polymers, that is, LDPE, LCB-PP, and LCB-PET, as shown in **Figure 7(a), (b), and (c)**, indicating that their gas efficiencies are constant. However, the gas efficiencies of the three polymers are different: 1.8, 3.5, and 2.1 for LDPE, LCB-PP, and LCB-PET, respectively. The expansion ratio of LDPE is relatively insensitive to the foaming pressure, indicating that it is not an economical to obtain a high expansion PE foam using high-pressure CO<sub>2</sub>. The gas efficiency of LCB-PET is similar to that of LDPE, which means that CO<sub>2</sub> alone is insufficient to achieve a satisfactory expansion ratio. However, the reason for the low gas efficiencies is different for the two polymers. For LDPE, the main reason is the high melt viscosity, which consumes a large amount of energy. For LCB-PET, it has a low solubility of CO<sub>2</sub>, leading to a low expansion ratio. The gas efficiency of LCB-PP is favourable, which means that the use of CO<sub>2</sub> in PP foaming is economical, and 20 MPa CO<sub>2</sub> is sufficient to produce PP foam with a high expansion ratio of 65.

The L-PP shows a considerably different outcome, and its expansion ratio curve

exhibits nonlinearity. By increasing the foaming pressure to approximately 17 MPa, an accelerated increase in the expansion ratio is observed. The gas efficiency at low foaming pressure is small and approximately  $2.0 \text{ (MPa}^{-1}\text{)}$ , but increases to  $11 \text{ (MPa}^{-1}\text{)}$  at high foaming pressure. The appearance of the second stage, in which the gas efficiency increases, is caused by the linear structure of L-PP. The initial viscosity of L-PP is high, which is the reason for the low gas efficiency at the first stage. Compared with the LCB-polymers, the L-PP has a low elongational viscosity after stretching because of the lack of strain hardening, indicating low growth resistance at the late bubble growth stage. A deviation between the experiments and simulations appears at a foaming pressure of 20 MPa for the L-PP system, as shown in **Figure 7(c)**. The SEM images shown in **Figure S7** in the **Supporting Information** indicate that the L-PP foam produced at 20 MPa has cell wall rupture behaviour, which led to an open-cell structure. This morphology did not satisfy the model assumption that the bubble system was a sealed sphere. The gas and corresponding energy leaked at the rupture point, leading to a low expansion and an open-cell structure. This phenomenon also indicated that the bubble shell might be broken under an extremely high foaming pressure.

## **Conclusion**

A non-isothermal PBB model was proposed from the viewpoint of mass, momentum, and energy conservation, which additionally considered the decrease in

gas internal energy owing to the work performed by gas expansion in the bubble. Accurate predictions of bubble size and expansion ratio for the polymer foaming process were made with the model calculation. The model correction was verified in four kinds of polymers and in a wide range of cell densities from  $1.5 \times 10^{13}$  cells/m<sup>3</sup> to  $1.9 \times 10^{15}$  cells/m<sup>3</sup>. The simulation results indicated that some nucleation points were difficult to grow because of shell resistance, leading to a significant energy cost. Furthermore, the effects of gas expansion work on bubble growth were different for linear and LCB polymers; a relatively prolonged decrease in the inner pressure appeared for the linear situation. Finally, the gas efficiency of foaming was theoretically calculated. The gas efficiency of CO<sub>2</sub> for PP foaming was satisfactory, while for PE and PET, single CO<sub>2</sub> demonstrated a low gas efficiency.

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