

An ADI Scheme for Two-sided Fractional Reaction-Diffusion Equations and Applications to an Epidemic Model

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Reaction-diffusion equations are often used in epidemiological models. In this paper, we generalize the algorithm of Meerschaert and Tadjeran for fractional advection-dispersion flow equations to a coupled system of fractional reaction-diffusion like an equation that arises from vector borne disease modeling.

KEYWORDS

ADI scheme, fractional reaction-diffusion equations, fractional SIRUV model

1 | INTRODUCTION

The modeling and understanding of infectious diseases is for many decades an object of intensive study. Going back to the classical SIR model from Kermack and McKendrick [1] which describes the time evolution of the number of susceptible (S), infected (I) and recovered (R) individuals by a system of ordinary differential equations various refinements were developed and extended exhaustively in the last 90 years. Among those extensions are the introduction of new compartments e.g. to model vector-borne diseases such as Dengue or malaria, as well as more involved deterministic and stochastic models, see e.g. [2, 3, 4, 5, 6, 7, 8, 9, 10, 11, 12, 13]. Spatial disease spread can be modeled either in a discrete or continuous way. Popular space-discrete models are the metapopulation approach [14, 15, 16] and for Dengue recently [17, 18], Cellular automata [19, 20], epidemic spatial networks [21, 22, 23, 24], and lattice epidemic models [25, 26]. For space-continuous models integro-differential equation epidemic models [27, 28] and diffusion epidemic models [29, 30, 31], are studied. In the last 10 years numerous fractional epidemic models [32, 33, 34, 35] were established. A distinct feature of fractional derivatives is the capability to model long-range interactions. In a popular model, the second derivative in a classical diffusion model is substituted by α -order derivative.

Fractional differential equations are widely used to model non-local phenomena and it is nowadays an object of important studies. In interacting particle systems we expect different patterns of spreading of the particles in anomalous. Therefore, studying fractional diffusion approaches are more important when compared to the classical

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diffusion. This work is motivated by the work of Brockmann [32], where he has reported a developed technique which allows a rigid and quantitative study of human diffusion on a geographical region using the spread of dollar bills. The result was a super-diffusive behaviour caused by human mobility habits like flights and long car drives. The distribution of travelling distances hence decays according to a power law similar to in Levy flights. Therefore, in this work fractional *SIV* model is formulated taking the fact that the human dispersal using a fractional Laplacian and compare the dynamics to the classical diffusive case.

In this article, a fractional diffusion model is derived from the *SIRUV* compartmental model with migration. To simulate this equation numerically we adapt the Alternating Directions Implicit (ADI) scheme with a Crank-Nicholson discretization to the fractional case. For this purpose a shifted version of the typical Grünwald-Letnikov finite difference approximation is used. The ADI method joined together with a fractional Crank-Nicholson scheme for fractional diffusion examples was already implemented by Meerschaert et. al. [36, 37, 38, 39]. The novelty of this article is that it is generalized to a system of coupled fractional reaction-diffusion equations. For this we derive the ADI splits with the corresponding Grünwald-Letnikov operators. A numerical scenario and a comparison with the classical diffusion case for Dirichlet boundary conditions can be found at the end of the article.

2 | MODEL DEFINITION

In this article, the system of ordinary differential equations (ODEs) for *SIV* model is taken derive the fractional model. Instead of using the system of equations for *SIRUV* model as in [40], a reduced form is used by using the simplification $R(t) = N - S(t) - I(t)$ and $U(t) = M - V(t)$ is given by the system of equations (1). The corresponding system of ODEs is given as follows:

$$\begin{aligned}\frac{dS(t)}{dt} &= \mu \cdot (1 - S(t)) - \beta \cdot S(t) \cdot V(t) = g_S \\ \frac{dI(t)}{dt} &= \beta \cdot S(t) \cdot V(t) - (\mu + \gamma) \cdot I(t) = g_I \\ \frac{dV(t)}{dt} &= \vartheta \cdot (1 - V(t)) \cdot I(t) - \nu \cdot V(t) = g_V\end{aligned}\tag{1}$$

where β and ϑ are the infection rate from vectors to hosts and hosts to vectors respectively. The recovery rate from the compartment *I* is given by γ . The birth and death rates of the hosts are equal and denoted by μ in order to maintain constant population size. Similarly, a constant population of vectors is maintained by assuming that birth and death rates ν of the vectors to be equal. The initial conditions are given by $S(0)$, $I(0)$ and $V(0)$ for the corresponding compartments.

The susceptible and infected individuals are spatially distributed, where $S(x, y, t)$, $I(x, y, t)$ and $V(x, y, t)$ represent the three state variables for the compartments. The initial conditions are given by the notations $S(x, y, t_0)$, $I(x, y, t_0)$ and $V(x, y, t_0)$. The two-dimensional spatial variables are denoted by x and y . Model (1) is redefined and can be written in a form of a reaction-diffusion model as follows:

$$\begin{aligned}\frac{\partial S(x, y, t)}{\partial t} &= g_S + a^S \left\{ \frac{\partial^2 S(x, y, t)}{\partial x^2} \right\} + b^S \left\{ \frac{\partial^2 S(x, y, t)}{\partial y^2} \right\} \\ \frac{\partial I(x, y, t)}{\partial t} &= g_I + a^I \left\{ \frac{\partial^2 I(x, y, t)}{\partial x^2} \right\} + b^I \left\{ \frac{\partial^2 I(x, y, t)}{\partial y^2} \right\} \\ \frac{\partial V(x, y, t)}{\partial t} &= g_V + a^V \left\{ \frac{\partial^2 V(x, y, t)}{\partial x^2} \right\} + b^V \left\{ \frac{\partial^2 V(x, y, t)}{\partial y^2} \right\}\end{aligned}\tag{2}$$

on a finite rectangular domain $x_L < x < x_H$ and $y_L < y < y_H$. The fractional orders are given by $1 < \alpha_1 \leq 2$ and

$1 < \alpha_2 \leq 2$. Dirichlet boundary conditions are used on the boundary $x_L \leq x \leq x_H$ and $y_L \leq y \leq y_H$:

$$\begin{aligned} S(x = x_L, y, t) = S(x = x_R, y, t) = S(x, y = y_L, t) = S(x, y = y_R, t) &= 0 \\ I(x = x_L, y, t) = I(x = x_R, y, t) = I(x, y = y_L, t) = I(x, y = y_R, t) &= 0 \\ V(x = x_L, y, t) = V(x = x_R, y, t) = V(x, y = y_L, t) = V(x, y = y_R, t) &= 0. \end{aligned}$$

The fractional derivatives of the previous equations are replaced by two-sided fractional derivatives and hence, the two-sided fractional diffusion *SIV*-model yields,

$$\begin{aligned} \frac{\partial S(x, y, t)}{\partial t} &= g_S + a^S \left\{ (1 - r_1) \frac{\partial^{\alpha_1} S(x, y, t)}{\partial (-x)^{\alpha_1}} + r_1 \frac{\partial^{\alpha_1} S(x, y, t)}{\partial x^{\alpha_1}} \right\} \\ &\quad + b^S \left\{ (1 - r_2) \frac{\partial^{\alpha_2} S(x, y, t)}{\partial (-y)^{\alpha_2}} + r_2 \frac{\partial^{\alpha_2} S(x, y, t)}{\partial y^{\alpha_2}} \right\} \\ \frac{\partial I(x, y, t)}{\partial t} &= g_I + a^I \left\{ (1 - r_1) \frac{\partial^{\alpha_1} I(x, y, t)}{\partial (-x)^{\alpha_1}} + r_1 \frac{\partial^{\alpha_1} I(x, y, t)}{\partial x^{\alpha_1}} \right\} \\ &\quad + b^I \left\{ (1 - r_2) \frac{\partial^{\alpha_2} I(x, y, t)}{\partial (-y)^{\alpha_2}} + r_2 \frac{\partial^{\alpha_2} I(x, y, t)}{\partial y^{\alpha_2}} \right\} \\ \frac{\partial V(x, y, t)}{\partial t} &= g_V + a^V \left\{ (1 - r_1) \frac{\partial^{\alpha_1} V(x, y, t)}{\partial (-x)^{\alpha_1}} + r_1 \frac{\partial^{\alpha_1} V(x, y, t)}{\partial x^{\alpha_1}} \right\} \\ &\quad + b^V \left\{ (1 - r_2) \frac{\partial^{\alpha_2} V(x, y, t)}{\partial (-y)^{\alpha_2}} + r_2 \frac{\partial^{\alpha_2} V(x, y, t)}{\partial y^{\alpha_2}} \right\}, \end{aligned} \quad (3)$$

with weights $r_1, r_2 \in [0, 1]$, where $\frac{\partial^{\alpha_i} F(x, y, t)}{\partial (-x)^{\alpha_i}}$ and $\frac{\partial^{\alpha_i} F(x, y, t)}{\partial (-y)^{\alpha_i}}$ denote the negative (right) fractional derivatives.

3 | NUMERICAL SCHEME

A Crank-Nicholson type system of finite difference equations can be obtained by substituting the shifted Grünwald into the differential equation centered at time $t_{n+1/2} = \frac{1}{2}(t_{n+1} + t_n)$.

$$\begin{aligned} S_{ij}^{n+1} - S_{ij}^n &= \Delta t \left\{ \mu - \mu S_{ij}^{n+1/2} - \beta S_{ij}^{n+1/2} V_{ij}^{n+1/2} \right\} \\ &\quad + \frac{\Delta t}{2} \left\{ (1 - r_1) (\delta_{\alpha_1 x}^- S_{ij}^{n+1} + \delta_{\alpha_1 x}^- S_{ij}^n) + r_1 (\delta_{\alpha_1 x}^+ S_{ij}^{n+1} + \delta_{\alpha_1 x}^+ S_{ij}^n) \right\} \\ &\quad + \frac{\Delta t}{2} \left\{ (1 - r_2) (\delta_{\alpha_2 y}^- S_{ij}^{n+1} + \delta_{\alpha_2 y}^- S_{ij}^n) + r_2 (\delta_{\alpha_2 y}^+ S_{ij}^{n+1} + \delta_{\alpha_2 y}^+ S_{ij}^n) \right\} \\ I_{ij}^{n+1} - I_{ij}^n &= \Delta t \left\{ \beta S_{ij}^{n+1/2} V_{ij}^{n+1/2} - (\mu + \gamma) I_{ij}^{n+1/2} \right\} \\ &\quad + \frac{\Delta t}{2} \left\{ (1 - r_1) (\delta_{\alpha_1 x}^- I_{ij}^{n+1} + \delta_{\alpha_1 x}^- I_{ij}^n) + r_1 (\delta_{\alpha_1 x}^+ I_{ij}^{n+1} + \delta_{\alpha_1 x}^+ I_{ij}^n) \right\} \\ &\quad + \frac{\Delta t}{2} \left\{ (1 - r_2) (\delta_{\alpha_2 y}^- I_{ij}^{n+1} + \delta_{\alpha_2 y}^- I_{ij}^n) + r_2 (\delta_{\alpha_2 y}^+ I_{ij}^{n+1} + \delta_{\alpha_2 y}^+ I_{ij}^n) \right\} \\ V_{ij}^{n+1} - V_{ij}^n &= \Delta t \left\{ \vartheta I_{ij}^{n+1/2} - \vartheta V_{ij}^{n+1/2} I_{ij}^{n+1/2} - \nu V_{ij}^{n+1/2} \right\} \\ &\quad + \frac{\Delta t}{2} \left\{ (1 - r_1) (\delta_{\alpha_1 x}^- V_{ij}^{n+1} + \delta_{\alpha_1 x}^- V_{ij}^n) + r_1 (\delta_{\alpha_1 x}^+ V_{ij}^{n+1} + \delta_{\alpha_1 x}^+ V_{ij}^n) \right\} \\ &\quad + \frac{\Delta t}{2} \left\{ (1 - r_2) (\delta_{\alpha_2 y}^- V_{ij}^{n+1} + \delta_{\alpha_2 y}^- V_{ij}^n) + r_2 (\delta_{\alpha_2 y}^+ V_{ij}^{n+1} + \delta_{\alpha_2 y}^+ V_{ij}^n) \right\} \end{aligned} \quad (4)$$

After rearranging the terms, the previous equation can be written in the operator notations as (5):

$$\begin{aligned}
& \left(1 - \frac{\Delta t}{2} \{(1-r_1)\delta_{\alpha_1 x}^{S^-} + r_1\delta_{\alpha_1 x}^{S^+}\} - \frac{\Delta t}{2} \{(1-r_2)\delta_{\alpha_2 y}^{S^-} + r_2\delta_{\alpha_2 y}^{S^+}\}\right) S_{ij}^{n+1} \\
&= \left(1 + \frac{\Delta t}{2} \{(1-r_1)\delta_{\alpha_1 x}^{S^-} + r_1\delta_{\alpha_1 x}^{S^+}\} + \frac{\Delta t}{2} \{(1-r_2)\delta_{\alpha_2 y}^{S^-} + r_2\delta_{\alpha_2 y}^{S^+}\}\right) S_{ij}^n \\
&\quad + \Delta t \left\{ \mu - \mu S_{ij}^{n+1/2} - \beta S_{ij}^{n+1/2} V_{ij}^{n+1/2} \right\} \\
\\
& \left(1 - \frac{\Delta t}{2} \{(1-r_1)\delta_{\alpha_1 x}^{I^-} + r_1\delta_{\alpha_1 x}^{I^+}\} - \frac{\Delta t}{2} \{(1-r_2)\delta_{\alpha_2 y}^{I^-} + r_2\delta_{\alpha_2 y}^{I^+}\}\right) I_{ij}^{n+1} \\
&= \left(1 + \frac{\Delta t}{2} \{(1-r_1)\delta_{\alpha_1 x}^{I^-} + r_1\delta_{\alpha_1 x}^{I^+}\} + \frac{\Delta t}{2} \{(1-r_2)\delta_{\alpha_2 y}^{I^-} + r_2\delta_{\alpha_2 y}^{I^+}\}\right) I_{ij}^n \\
&\quad + \Delta t \left\{ \beta S_{ij}^{n+1/2} V_{ij}^{n+1/2} - (\mu + \gamma) I_{ij}^{n+1/2} \right\} \\
\\
& \left(1 - \frac{\Delta t}{2} \{(1-r_1)\delta_{\alpha_1 x}^{V^-} + r_1\delta_{\alpha_1 x}^{V^+}\} - \frac{\Delta t}{2} \{(1-r_2)\delta_{\alpha_2 y}^{V^-} + r_2\delta_{\alpha_2 y}^{V^+}\}\right) V_{ij}^{n+1} \\
&= \left(1 + \frac{\Delta t}{2} \{(1-r_1)\delta_{\alpha_1 x}^{V^-} + r_1\delta_{\alpha_1 x}^{V^+}\} + \frac{\Delta t}{2} \{(1-r_2)\delta_{\alpha_2 y}^{V^-} + r_2\delta_{\alpha_2 y}^{V^+}\}\right) V_{ij}^n \\
&\quad + \Delta t \left\{ \vartheta I_{ij}^{n+1/2} - \vartheta V_{ij}^{n+1/2} I_{ij}^{n+1/2} - \nu V_{ij}^{n+1/2} \right\} \quad (5)
\end{aligned}$$

Multi-dimensional diffusion equations are often solved with alternating directions implicit methods (ADI), where splitting is used to significantly reduce the computational work [41]. These techniques use a perturbation of Equation (5) in order to derive schemes that requires only the implicit numerical solution in one direction where the other spatial direction is computed iteratively. We obtain the equations

for S ,

$$\begin{aligned}
& \left(1 - \frac{\Delta t}{2} \{(1-r_1)\delta_{\alpha_1 x}^{S^-} + r_1\delta_{\alpha_1 x}^{S^+}\}\right) \left(1 - \frac{\Delta t}{2} \{(1-r_2)\delta_{\alpha_2 y}^{S^-} + r_2\delta_{\alpha_2 y}^{S^+}\}\right) S_{ij}^{n+1} \\
&= \left(1 + \frac{\Delta t}{2} \{(1-r_1)\delta_{\alpha_1 x}^{S^-} + r_1\delta_{\alpha_1 x}^{S^+}\}\right) \left(1 + \frac{\Delta t}{2} \{(1-r_2)\delta_{\alpha_2 y}^{S^-} + r_2\delta_{\alpha_2 y}^{S^+}\}\right) S_{ij}^n \\
&\quad + \Delta t \left\{ \mu - \mu S_{ij}^{n+1/2} - \beta S_{ij}^{n+1/2} V_{ij}^{n+1/2} \right\},
\end{aligned} \quad (6)$$

for I ,

$$\begin{aligned}
& \left(1 - \frac{\Delta t}{2} \{(1-r_1)\delta_{\alpha_1 x}^{I^-} + r_1\delta_{\alpha_1 x}^{I^+}\}\right) \left(1 - \frac{\Delta t}{2} \{(1-r_2)\delta_{\alpha_2 y}^{I^-} + r_2\delta_{\alpha_2 y}^{I^+}\}\right) I_{ij}^{n+1} \\
&= \left(1 + \frac{\Delta t}{2} \{(1-r_1)\delta_{\alpha_1 x}^{I^-} + r_1\delta_{\alpha_1 x}^{I^+}\}\right) \left(1 + \frac{\Delta t}{2} \{(1-r_2)\delta_{\alpha_2 y}^{I^-} + r_2\delta_{\alpha_2 y}^{I^+}\}\right) I_{ij}^n \\
&\quad + \Delta t \left\{ \beta S_{ij}^{n+1/2} V_{ij}^{n+1/2} - (\mu + \gamma) I_{ij}^{n+1/2} \right\},
\end{aligned} \quad (7)$$

for V ,

$$\begin{aligned}
& \left(1 - \frac{\Delta t}{2} \{(1-r_1)\delta_{\alpha_1 x}^{V^-} + r_1\delta_{\alpha_1 x}^{V^+}\}\right) \left(1 - \frac{\Delta t}{2} \{(1-r_2)\delta_{\alpha_2 y}^{V^-} + r_2\delta_{\alpha_2 y}^{V^+}\}\right) V_{ij}^{n+1} \\
&= \left(1 + \frac{\Delta t}{2} \{(1-r_1)\delta_{\alpha_1 x}^{V^-} + r_1\delta_{\alpha_1 x}^{V^+}\}\right) \left(1 + \frac{\Delta t}{2} \{(1-r_2)\delta_{\alpha_2 y}^{V^-} + r_2\delta_{\alpha_2 y}^{V^+}\}\right) V_{ij}^n \\
&\quad + \Delta t \left\{ \vartheta I_{ij}^{n+1/2} - \vartheta V_{ij}^{n+1/2} I_{ij}^{n+1/2} - \nu V_{ij}^{n+1/2} \right\}.
\end{aligned} \quad (8)$$

The equations (6), (7) and (8) form Peaceman-Rachford type matrix equations defining ADI method. This can be split as follows:

for S ,

$$\begin{aligned} \left(1 - \frac{\Delta t}{2} \{(1 - r_1)\delta_{\alpha_1 x}^- + r_1\delta_{\alpha_1 x}^+\}\right) S_{i,j}^* = & \left(1 + \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^- + r_2\delta_{\alpha_2 y}^+\}\right) S_{i,j}^n \\ & + \frac{\Delta t}{2} \left\{ \mu - \mu S_{i,j}^{n+1/2} - \beta S_{i,j}^{n+1/2} V_{i,j}^{n+1/2} \right\} \end{aligned} \quad (9)$$

$$\begin{aligned} \left(1 - \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^- + r_2\delta_{\alpha_2 y}^+\}\right) S_{i,j}^{n+1} = & \left(1 + \frac{\Delta t}{2} \{(1 - r_1)\delta_{\alpha_1 x}^- + r_1\delta_{\alpha_1 x}^+\}\right) S_{i,j}^* \\ & + \frac{\Delta t}{2} \left\{ \mu - \mu S_{i,j}^{n+1/2} - \beta S_{i,j}^{n+1/2} V_{i,j}^{n+1/2} \right\} \end{aligned} \quad (10)$$

for I ,

$$\begin{aligned} \left(1 - \frac{\Delta t}{2} \{(1 - r_1)\delta_{\alpha_1 x}^{I-} + r_1\delta_{\alpha_1 x}^{I+}\}\right) I_{i,j}^* = & \left(1 + \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{I-} + r_2\delta_{\alpha_2 y}^{I+}\}\right) I_{i,j}^n \\ & + \frac{\Delta t}{2} \left\{ \beta S_{i,j}^{n+1/2} V_{i,j}^{n+1/2} - (\mu + \gamma) I_{i,j}^{n+1/2} \right\} \end{aligned} \quad (11)$$

$$\begin{aligned} \left(1 - \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{I-} + r_2\delta_{\alpha_2 y}^{I+}\}\right) I_{i,j}^{n+1} = & \left(1 + \frac{\Delta t}{2} \{(1 - r_1)\delta_{\alpha_1 x}^{I-} + r_1\delta_{\alpha_1 x}^{I+}\}\right) I_{i,j}^* \\ & + \frac{\Delta t}{2} \left\{ \beta S_{i,j}^{n+1/2} V_{i,j}^{n+1/2} - (\mu + \gamma) I_{i,j}^{n+1/2} \right\}. \end{aligned} \quad (12)$$

for V ,

$$\begin{aligned} \left(1 - \frac{\Delta t}{2} \{(1 - r_1)\delta_{\alpha_1 x}^{V-} + r_1\delta_{\alpha_1 x}^{V+}\}\right) V_{i,j}^* = & \left(1 + \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{V-} + r_2\delta_{\alpha_2 y}^{V+}\}\right) V_{i,j}^n \\ & + \frac{\Delta t}{2} \left\{ \vartheta I_{i,j}^{n+1/2} - \vartheta V_{i,j}^{n+1/2} I_{i,j}^{n+1/2} - \nu V_{i,j}^{n+1/2} \right\} \end{aligned} \quad (13)$$

$$\begin{aligned} \left(1 - \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{V-} + r_2\delta_{\alpha_2 y}^{V+}\}\right) V_{i,j}^{n+1} = & \left(1 + \frac{\Delta t}{2} \{(1 - r_1)\delta_{\alpha_1 x}^{V-} + r_1\delta_{\alpha_1 x}^{V+}\}\right) V_{i,j}^* \\ & + \frac{\Delta t}{2} \left\{ \vartheta I_{i,j}^{n+1/2} - \vartheta V_{i,j}^{n+1/2} I_{i,j}^{n+1/2} - \nu V_{i,j}^{n+1/2} \right\}. \end{aligned} \quad (14)$$

To observe this multiply (9), (11) and (13) by $\left(1 + \frac{\Delta t}{2} \{(1 - r_1)\delta_{\alpha_1 x}^- + r_1\delta_{\alpha_1 x}^+\}\right)$ on both sides and then multiply (10), (12) and (14) by $\left(1 - \frac{\Delta t}{2} \{(1 - r_1)\delta_{\alpha_1 x}^- + r_1\delta_{\alpha_1 x}^+\}\right)$ on both sides to obtain the Equations (6), (7) and (8). Equations (9), (10), (11), (12), (13) and (14) calculates intermediate solutions $S_{i,j}^*$, $I_{i,j}^*$ and $V_{i,j}^*$ in order to develop the numerical solutions to S , I and V at time step n to the numerical solution $S_{i,j}^{n+1}$, $I_{i,j}^{n+1}$ and $V_{i,j}^{n+1}$ at time t_{n+1} .

The algorithm of ADI splitting method [36] to solve the SIV -fractional diffusion model is given by:

Algorithm 1 ADI Scheme

- 1: In order to acquire the intermediate solution slice $S_{i,j}^*$, $I_{i,j}^*$ and $V_{i,j}^*$, a set of $N_x - 1$ equations at the points x_i , where, $i = 1, 2, \dots, N_x - 1$ defined by equation (9) and (11) are solved initially on each fixed horizontal slice $y = y_k$ where $k = 1, 2, \dots, N_y - 1$.
 - 2: In addition, by alternating the spatial direction on every fixed verticle slice $x = x_k$ ($k = 1, 2, \dots, N_x - 1$) a set of $N_y - 1$ equations are solved at the points y_j where $j = 1, 2, \dots, N_y - 1$ defined by the equations (10) and (12) in order to obtain the solution for $S_{k,j}^{n+1}$, $I_{k,j}^{n+1}$ and $V_{k,j}^{n+1}$ at time $n + 1$.
-

The shifted Grünwald operators used in this model yield

$$\begin{aligned}
 \delta_{\alpha_1,x}^{S-} S_{ij}^n &= \frac{(a_{ij}^S)}{(\Delta x)^{\alpha_1}} \sum_{k=0}^{N_x-i+1} g_{\alpha_1,k} \cdot S_{i+k-1,j}^n, & \delta_{\alpha_1,x}^{S+} S_{ij}^n &= \frac{(a_{ij}^S)}{(\Delta x)^{\alpha_1}} \sum_{k=0}^{i+1} g_{\alpha_1,k} \cdot S_{i-k+1,j}^n \\
 \delta_{\alpha_2,y}^{S-} S_{ij}^n &= \frac{(b_{ij}^S)}{(\Delta y)^{\alpha_2}} \sum_{k=0}^{N_y-j+1} g_{\alpha_2,k} \cdot S_{i,j+k-1}^n, & \delta_{\alpha_2,y}^{S+} S_{ij}^n &= \frac{(b_{ij}^S)}{(\Delta y)^{\alpha_2}} \sum_{k=0}^{j+1} g_{\alpha_2,k} \cdot S_{i,j-k+1}^n \\
 \delta_{\alpha_1,x}^{I-} I_{ij}^n &= \frac{(a_{ij}^I)}{(\Delta x)^{\alpha_1}} \sum_{k=0}^{N_x-i+1} g_{\alpha_1,k} \cdot I_{i+k-1,j}^n, & \delta_{\alpha_1,x}^{I+} I_{ij}^n &= \frac{(a_{ij}^I)}{(\Delta x)^{\alpha_1}} \sum_{k=0}^{i+1} g_{\alpha_1,k} \cdot I_{i-k+1,j}^n \\
 \delta_{\alpha_2,y}^{I-} I_{ij}^n &= \frac{(b_{ij}^I)}{(\Delta y)^{\alpha_2}} \sum_{k=0}^{N_y-j+1} g_{\alpha_2,k} \cdot I_{i,j+k-1}^n, & \delta_{\alpha_2,y}^{I+} I_{ij}^n &= \frac{(b_{ij}^I)}{(\Delta y)^{\alpha_2}} \sum_{k=0}^{j+1} g_{\alpha_2,k} \cdot I_{i,j-k+1}^n \\
 \delta_{\alpha_1,x}^{V-} V_{ij}^n &= \frac{(a_{ij}^V)}{(\Delta x)^{\alpha_1}} \sum_{k=0}^{N_x-i+1} g_{\alpha_1,k} \cdot V_{i+k-1,j}^n, & \delta_{\alpha_1,x}^{V+} V_{ij}^n &= \frac{(a_{ij}^V)}{(\Delta x)^{\alpha_1}} \sum_{k=0}^{i+1} g_{\alpha_1,k} \cdot V_{i-k+1,j}^n \\
 \delta_{\alpha_2,y}^{V-} V_{ij}^n &= \frac{(b_{ij}^V)}{(\Delta y)^{\alpha_2}} \sum_{k=0}^{N_y-j+1} g_{\alpha_2,k} \cdot V_{i,j+k-1}^n, & \delta_{\alpha_2,y}^{V+} V_{ij}^n &= \frac{(b_{ij}^V)}{(\Delta y)^{\alpha_2}} \sum_{k=0}^{j+1} g_{\alpha_2,k} \cdot V_{i,j-k+1}^n
 \end{aligned} \tag{15}$$

Analogously finite difference schemes for the S , I and V compartments are obtained by substituting the shifted Grünwald operator into the equations from before. The corresponding ADI scheme reads:

(i). ADI split I:

$$\begin{aligned}
 S_{ij}^* - r_1 D_{ij}^S \sum_{k=0}^{i+1} g_{\alpha_1,k} \cdot S_{i-k+1,j}^* - (1-r_1) D_{ij}^S \sum_{k=0}^{N_x-i+1} g_{\alpha_1,k} \cdot S_{i+k-1,j}^* \\
 = S_{ij}^n + r_2 E_{ij}^S \sum_{k=0}^{j+1} g_{\alpha_2,k} \cdot S_{i,j-k+1}^n + (1-r_2) E_{ij}^S \sum_{k=0}^{N_y-j+1} g_{\alpha_2,k} \cdot S_{i,j+k-1}^n \\
 + \frac{\Delta t}{2} \left\{ \mu - \beta S_{ij}^{n+1/2} V_{ij}^{n+1/2} - \mu S_{ij}^{n+1/2} \right\},
 \end{aligned}$$

$$\begin{aligned}
 I_{ij}^* - r_1 D_{ij}^I \sum_{k=0}^{i+1} g_{\alpha_1,k} \cdot I_{i-k+1,j}^* - (1-r_1) D_{ij}^I \sum_{k=0}^{N_x-i+1} g_{\alpha_1,k} \cdot I_{i+k-1,j}^* \\
 = I_{ij}^n + r_2 E_{ij}^I \sum_{k=0}^{j+1} g_{\alpha_2,k} \cdot I_{i,j-k+1}^n + (1-r_2) E_{ij}^I \sum_{k=0}^{N_y-j+1} g_{\alpha_2,k} \cdot I_{i,j+k-1}^n \\
 + \frac{\Delta t}{2} \left\{ \beta S_{ij}^{n+1/2} V_{ij}^{n+1/2} - (\mu + \gamma) I_{ij}^{n+1/2} \right\},
 \end{aligned}$$

$$\begin{aligned}
 V_{ij}^* - r_1 D_{ij}^V \sum_{k=0}^{i+1} g_{\alpha_1,k} \cdot V_{i-k+1,j}^* - (1-r_1) D_{ij}^V \sum_{k=0}^{N_x-i+1} g_{\alpha_1,k} \cdot V_{i+k-1,j}^* \\
 = V_{ij}^n + r_2 E_{ij}^V \sum_{k=0}^{j+1} g_{\alpha_2,k} \cdot V_{i,j-k+1}^n + (1-r_2) E_{ij}^V \sum_{k=0}^{N_y-j+1} g_{\alpha_2,k} \cdot V_{i,j+k-1}^n
 \end{aligned}$$

$$+ \frac{\Delta t}{2} \left\{ \vartheta (1 - V_{ij}^{n+1/2}) I_{ij}^{n+1/2} - \nu V_{ij}^{n+1/2} \right\},$$

where $D_{ij}^X = \frac{\Delta t a_{ij}^X}{2(\Delta x)^{\alpha_1}}$, and $E_{ij}^X = \frac{\Delta t b_{ij}^X}{2(\Delta y)^{\alpha_2}}$. X represents the compartments S, I and V.

(ii). ADI split II:

$$\begin{aligned} S_{ij}^{n+1} - r_2 E_{ij}^S \sum_{k=0}^{j+1} g_{\alpha_2, k} \cdot S_{ij-k+1}^{n+1} - (1 - r_2) E_{ij}^S \sum_{k=0}^{N_y-j+1} g_{\alpha_2, k} \cdot S_{ij+k-1}^{n+1} \\ = S_{ij}^* + r_1 D_{ij}^S \sum_{k=0}^{i+1} g_{\alpha_1, k} \cdot S_{i-k+1, j}^* + (1 - r_1) D_{ij}^S \sum_{k=0}^{N_x-i+1} g_{\alpha_1, k} \cdot S_{i+k-1, j}^* \\ - \frac{\Delta t}{2} \left\{ \mu - \beta S_{ij}^{n+1/2} V_{ij}^{n+1/2} - \mu S_{ij}^{n+1/2} \right\} \end{aligned}$$

$$\begin{aligned} I_{ij}^{n+1} - r_2 E_{ij}^I \sum_{k=0}^{j+1} g_{\alpha_2, k} \cdot I_{ij-k+1}^{n+1} - (1 - r_2) E_{ij}^I \sum_{k=0}^{N_y-j+1} g_{\alpha_2, k} \cdot I_{ij+k-1}^{n+1} \\ = I_{ij}^* + r_1 D_{ij}^I \sum_{k=0}^{i+1} g_{\alpha_1, k} \cdot I_{i-k+1, j}^* + (1 - r_1) D_{ij}^I \sum_{k=0}^{N_x-i+1} g_{\alpha_1, k} \cdot I_{i+k-1, j}^* \\ - \frac{\Delta t}{2} \left\{ \beta S_{ij}^{n+1/2} V_{ij}^{n+1/2} - (\mu + \gamma) I_{ij}^{n+1/2} \right\} \end{aligned}$$

$$\begin{aligned} V_{ij}^{n+1} - r_2 E_{ij}^V \sum_{k=0}^{j+1} g_{\alpha_2, k} \cdot V_{ij-k+1}^{n+1} - (1 - r_2) E_{ij}^V \sum_{k=0}^{N_y-j+1} g_{\alpha_2, k} \cdot V_{ij+k-1}^{n+1} \\ = V_{ij}^* + r_1 D_{ij}^V \sum_{k=0}^{i+1} g_{\alpha_1, k} \cdot V_{i-k+1, j}^* + (1 - r_1) D_{ij}^V \sum_{k=0}^{N_x-i+1} g_{\alpha_1, k} \cdot V_{i+k-1, j}^* \\ - \frac{\Delta t}{2} \left\{ \vartheta (1 - V_{ij}^{n+1/2}) I_{ij}^{n+1/2} - \nu V_{ij}^{n+1/2} \right\} \end{aligned}$$

Before solving the system of equations defined by ADI split I and ADI split II, the intermediate solutions S_{ij}^* , I_{ij}^* and V_{ij}^* must be treated with care on the boundary in order to preserve the consistency of the set of equations defined by (9), (10), (11), (12), (13) and (14) with (6), (7) and (8). By subtracting (10) from (9), (12) from (11) and (14) from (13) we obtain,

$$\left(1 - \frac{\Delta t}{2} \{ (1 - r_2) \delta_{\alpha_2 y}^{S^-} + r_2 \delta_{\alpha_2 y}^{S^+} \} \right) S_{ij}^{n+1} + \left(1 + \frac{\Delta t}{2} \{ (1 - r_2) \delta_{\alpha_2 y}^{S^-} + r_2 \delta_{\alpha_2 y}^{S^+} \} \right) S_{ij}^n = 2S_{ij}^* \quad (16)$$

$$\left(1 - \frac{\Delta t}{2} \{ (1 - r_2) \delta_{\alpha_2 y}^{I^-} + r_2 \delta_{\alpha_2 y}^{I^+} \} \right) I_{ij}^{n+1} + \left(1 + \frac{\Delta t}{2} \{ (1 - r_2) \delta_{\alpha_2 y}^{I^-} + r_2 \delta_{\alpha_2 y}^{I^+} \} \right) I_{ij}^n = 2I_{ij}^* \quad (17)$$

$$(1 - \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{V^-} + r_2\delta_{\alpha_2 y}^{V^+}\})V_{ij}^{n+1} + (1 + \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{V^-} + r_2\delta_{\alpha_2 y}^{V^+}\})V_{ij}^n = 2V_{ij}^* \quad (18)$$

The boundary conditions for the intermediate solutions S_{ij}^* , I_{ij}^* and V_{ij}^* (i.e., $i = 0$ or $i = N_x$ for $j = 1, \dots, N_y - 1$) required to solve the set of equations (6), (7) and (8) are of the form

$$\begin{aligned} 2S_{0j}^* &= (1 - \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{S^-} + r_2\delta_{\alpha_2 y}^{S^+}\})S_{0j}^{n+1} + (1 + \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{S^-} + r_2\delta_{\alpha_2 y}^{S^+}\})S_{0j}^n \\ 2S_{N_x j}^* &= (1 - \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{S^-} + r_2\delta_{\alpha_2 y}^{S^+}\})S_{N_x j}^{n+1} + (1 + \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{S^-} + r_2\delta_{\alpha_2 y}^{S^+}\})S_{N_x j}^n \end{aligned} \quad (19)$$

$$\begin{aligned} 2I_{0j}^* &= (1 - \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{I^-} + r_2\delta_{\alpha_2 y}^{I^+}\})I_{0j}^{n+1} + (1 + \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{I^-} + r_2\delta_{\alpha_2 y}^{I^+}\})I_{0j}^n \\ 2I_{N_x j}^* &= (1 - \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{I^-} + r_2\delta_{\alpha_2 y}^{I^+}\})I_{N_x j}^{n+1} + (1 + \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{I^-} + r_2\delta_{\alpha_2 y}^{I^+}\})I_{N_x j}^n. \end{aligned} \quad (20)$$

$$\begin{aligned} 2V_{0j}^* &= (1 - \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{V^-} + r_2\delta_{\alpha_2 y}^{V^+}\})V_{0j}^{n+1} + (1 + \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{V^-} + r_2\delta_{\alpha_2 y}^{V^+}\})V_{0j}^n \\ 2V_{N_x j}^* &= (1 - \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{V^-} + r_2\delta_{\alpha_2 y}^{V^+}\})V_{N_x j}^{n+1} + (1 + \frac{\Delta t}{2} \{(1 - r_2)\delta_{\alpha_2 y}^{V^-} + r_2\delta_{\alpha_2 y}^{V^+}\})V_{N_x j}^n. \end{aligned} \quad (21)$$

Dirichlet boundary conditions are used and hence, S_{0j}^* , $S_{N_x j}^*$, I_{0j}^* , $I_{N_x j}^*$, V_{0j}^* and $V_{N_x j}^*$ becomes zero.

4 | EXAMPLE

The pattern of the spread of the infected hosts I and the infected vectors are quite similar in the figures. However, the intensities of the number of infected in the corresponding compartments are different. In the simulations we have considered the model parameters such that a disease outbreak will happen. The reproduction number of the model 1 is $9.9997 > 1$ ($R_0 = \frac{\beta\vartheta}{\nu(\gamma+\mu)}$, $\mu = 1/(72)$, $\beta = 365/(7)$, $\nu = 365/10$, $\vartheta = 5\nu$, $\gamma = 365/(14)$). During the considered time duration that is used in the the simulation therefore, there will be a increase in the disease infected and will eventually converge to the endemic equilibrium outside the considered time period and hence a reduction of the infected individuals can be seen. The model 3 is solved numerically by the ADI-CN scheme. To see the diffusion of the infection the following set of initial conditions are used,

$$I(0) = \begin{cases} 0.1 & \text{mid point of the finite grid} \\ 0 & \text{elsewhere} \end{cases} \quad (22)$$

$$S(0) = \begin{cases} 0.9 & \text{mid point of the finite grid} \\ 1 & \text{elsewhere} \end{cases} \quad (23)$$

$$V(0) = \begin{cases} 0 & \text{everywhere} \end{cases} \quad (24)$$

In this paper Dirichlet boundary conditions are of interest. Dirichlet conditions for both the S and I compartments on the rectangular region $x_L \leq x \leq x_H$ and $y_L \leq y \leq y_H$ are of the form:

$$S(0, y, t) = S(1, y, t) = S(x, 0, t) = S(x, 1, t) = 0$$

$$\begin{aligned} I(0, y, t) = I(1, y, t) = I(x, 0, t) = I(x, 1, t) = 0 \\ V(0, y, t) = V(1, y, t) = V(x, 0, t) = V(x, 1, t) = 0. \end{aligned} \quad (25)$$

The corresponding numerical solutions of the fractional diffusion *SIV*-model is compared with the classical diffusion *SIV*-model. Figure 1 and Figure 2 shows the results of the hosts and vectors corresponding to the fractional-order 1.2 compared with the corresponding classical case.

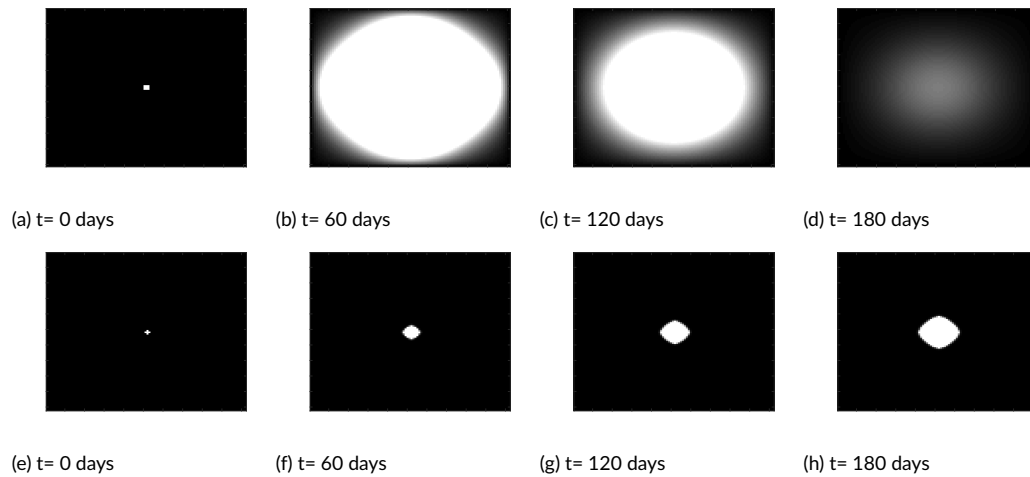


FIGURE 1 Spread of the infected hosts I by Fractional Diffusion *SIV*-model where $\alpha = 1.2$. and the classical model in the host compartment.

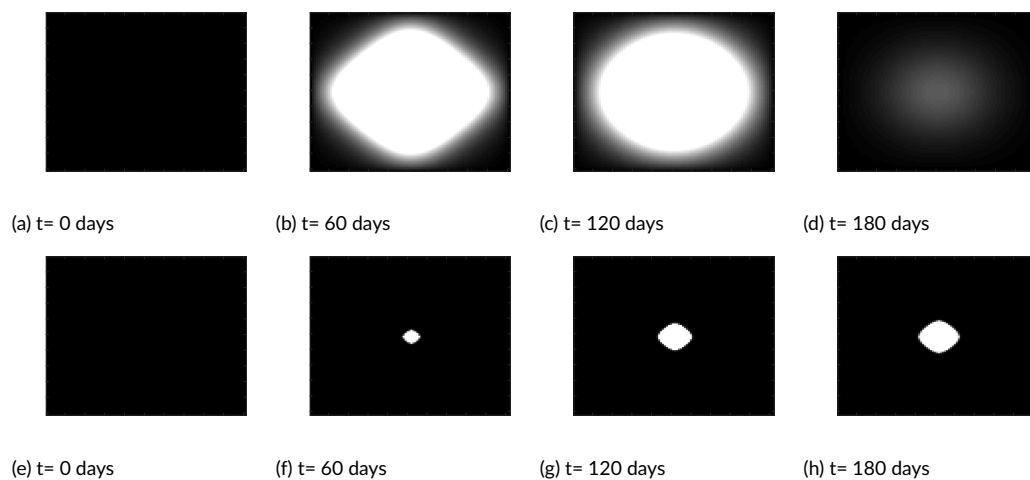


FIGURE 2 Spread of the infected vectors V by Fractional Diffusion *SIV*-model where $\alpha = 1.2$. and the classical model in the host compartment.

It can be seen that the spread of the classical diffusion is slower than that of the fractional diffusion. Hence, numerical results corresponding to the fractional model shows an anomalous diffusion which can be seen in the

infected hosts.

5 | CONCLUSION

In this article, the reaction-diffusion *SIV* partial differential equation model is derived by using the multi-patch system with the long term movements of the individuals. Further, here we introduce a model corresponding to the reaction-diffusion approach to the existing *SIV*-epidemic model. The second derivative of the consistent classical reaction-diffusion equations is substituted by using the α order fractional derivatives in the respective space derivatives. The model is simulated using the alternating directions implicit (ADI) scheme with a Crank-Nicholson discretization. The numerical results are compared with the results attained by the classical reaction-diffusion system. The results illustrate an anomalous diffusive behaviour compared to the classical diffusion approach.

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Conflict of interest declaration

The authors have no conflict of interest relevant to this article.

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