

1 **Adaptive covariance hybridization for the assimilation**
2 **of SST observations within a coupled Earth system**
3 **reanalysis**

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8 **Key Points:**

- 9 • Hybrid covariance handles sampling error and improves the update of deep wa-
10 ter masses when assimilating surface observation with an EnKF.
11 • The method is well suited to provide a long coupled reanalysis of the past century
12 • Hybrid covariance with adaptive hybrid coefficients explicitly estimated in space
13 and time achieved the best performance

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Abstract

Ensemble data assimilation methods, such as the Ensemble Kalman Filter (EnKF), are well suited for climate reanalysis because they feature flow-dependent covariance. However, because Earth System Models are heavy computationally, the method uses a few tens of members. Sampling error in the covariance matrix can introduce biases in the deep ocean, which may cause a drift in the reanalysis and in the predictions. Here, we assess the potential of the hybrid covariance approach (EnKF-OI) to counteract sampling error. The EnKF-OI combines the flow-dependent covariance computed from a dynamical ensemble with another covariance matrix that is static but less prone to sampling error. We test the method within the Norwegian Climate Prediction Model (NorCPM), which combines the Norwegian Earth System Model (NorESM) and the EnKF. We test the performance of the reanalyses in an idealised twin experiment, where we assimilate synthetic sea surface temperature observations monthly over 1980-2010. The dynamical and static ensembles consist respectively of 30 members and 315 seasonal members sampled from a pre-industrial run. We compare the performance of the EnKF to an EnKF-OI with a global hybrid coefficient, referred to as standard hybrid, and an EnKF-OI with adaptive hybrid coefficients estimated in space and time. Both hybrid covariance methods cure the bias introduced by the EnKF at intermediate and deep water. The adaptive EnKF-OI performs best overall by addressing sampling noise and rank deficiencies issues and can sustain low analysis errors by doing smaller updates than the standard hybrid version.

Plain Language Summary

Data assimilation is a statistical method that reduces uncertainty in a model, based on observations. Because of their ease of implementation, the ensemble data assimilation methods, that rely on the statistics of a finite ensemble of realisations of the model, are popular for climate reanalysis and prediction. However, observations are sparse – mostly near the surface – and the sampling error from data assimilation method introduces a degradation in the deep ocean. We use a method that complements this ensemble with a pre-existing database of model states to reduce sampling error. We show that the approach substantially reduces error at the intermediate and deep ocean. The method typically requires the tuning of a parameter, but we show that it can be estimated online, achieving the best performance.

1 Introduction

Data assimilation estimates the state of a model (called the analysis) that approaches the “unknown true state” of the system based on observations, a prior model estimate, and statistical information on their uncertainties. Data assimilation is applied sequentially/recursively to provide a reanalysis, which can also be used to understand the mechanisms of variability and initialise predictions. Data assimilation has been one of the key ingredients in the progress of numerical weather prediction (Bauer et al., 2015) and is now used in a wide range of geosciences applications (Carrassi et al., 2018), including climate prediction. The ensemble Kalman Filter (EnKF, Evensen, 2003), is an advanced data assimilation method that provides flow-dependent covariance – *i.e.*, that can evolve in time and space with a transient state or a regime shift – and the ensemble provides a quantification of the uncertainty of the system. These properties are well suited for climate reanalysis and predictions, and the method is becoming increasingly popular in that field (*e.g.*, Zhang et al., 2009; Counillon et al., 2014; Brune et al., 2015; Karspeck et al., 2018; O’Kane et al., 2019).

Ensemble-based data assimilation methods (*e.g.*, the EnKF) approximate the forecast error covariance matrix using a finite-size ensemble from a Monte Carlo integration step. Rank deficiencies and sampling errors emerge and deteriorate the performance of

64 the system, causing an artificial reduction of the ensemble spread and may even lead to
 65 filter divergence. Several ad-hoc methods have been introduced to counteract sampling
 66 errors. Localization (Houtekamer & Mitchell, 2001; Evensen, 2003; Ott et al., 2004; An-
 67 derson, 2007) limits the spatial extent of the corrections, based on the approximation
 68 that the covariance function decays as a function of the distance – and can be seen as
 69 a way to effectively reduce the degree of freedom of the system. Inflation (Anderson, 2001;
 70 Whitaker & Hamill, 2012; Raanes et al., 2019) counteracts the spread-collapse by arti-
 71 ficially inflating the ensemble spread at every assimilation step. The last method is the
 72 covariance hybridization method, which is the topic of our paper.

73 Covariance hybridization (Hamill & Snyder, 2000) linearly combines the flow-dependent
 74 covariance computed from a finite Monte-Carlo ensemble with another covariance ma-
 75 trix that is less prone to sampling error. The static matrix can be parameterized (Hamill
 76 & Snyder, 2000; Weaver & Courtier, 2001), computed from a long model simulation (Counillon
 77 et al., 2009), computed as the average of the background error covariance matrices from
 78 a previous data assimilation run (Carrió et al., 2021) or computed from a dynamical en-
 79 semble at a lower resolution (Rainwater & Hunt, 2013). The hybrid covariance method
 80 achieves better performance than the standalone EnKF, particularly for small ensem-
 81 bles, and performance converges to that of the EnKF for large ensembles (X. Wang et
 82 al., 2007; Counillon et al., 2009; Raboudi et al., 2019). The computational cost of the
 83 hybrid covariance methods is customizable to the desired cost.

84 The linear coefficients combining the static and the dynamic covariance are called
 85 the “hybridization coefficients”, which optimally balance the superior but noisy sample
 86 covariance with that of less noisy but static covariance. To achieve optimal performance,
 87 it is crucial to tune these coefficients (X. Wang et al., 2007; Counillon et al., 2009; Raboudi
 88 et al., 2019; Gharamti et al., 2014). The optimal values of these coefficients depend on
 89 the non-stationarity of the dynamical system as well as the data assimilation settings,
 90 such as the dynamical ensemble size, localisation and inflation settings. As such, it is ex-
 91 pected that the optimal value of the hybridization coefficients should vary in space and
 92 time. A first attempt to estimate spatial and time-varying hybridization coefficient has
 93 been developed (Gharamti, 2020) with a Bayesian framework (using fixed localisation
 94 settings). They found that a spatially heterogeneous hybridization coefficient formula-
 95 tion outperforms a homogeneous formulation. Ménétrier and Auligné (2015) and Ménétrier
 96 et al. (2015) formulated the problem of hybridization as a linear filtering problem of the
 97 background error covariance matrix to optimize both the localization and the hybridiza-
 98 tion coefficients simultaneously.

99 In this work, we aim to investigate the benefit of background error covariance hy-
 100 bridization for climate reanalysis and climate prediction systems, as for example with
 101 the CMIP6 Decadal Climate Prediction Project (DCPP, Boer et al., 2016). We use the
 102 Norwegian Climate Prediction Model (NorCPM, Counillon et al., 2014, 2016) that pro-
 103 vides coupled reanalysis and contributed to CMIP6 DCPP (Bethke et al., 2021). Here,
 104 we focus on long coupled reanalysis as NorCPM will produce such a reanalysis from 1850
 105 to the present. Sea surface temperature (SST) dominates the ocean observation network
 106 before the emergence of altimetry in the 90s and Argo data in the 2000s. Using only SST,
 107 NorCPM can control the upper ocean heat content, and major indices of climate vari-
 108 ability in the North Atlantic well (Counillon et al., 2016). Two features of NorCPM, flow-
 109 dependent assimilation and assimilation in isopycnal coordinates, were found to be par-
 110 ticularly important in that success. However, it also yields an unrealistic update of the
 111 intermediate water masses in the North Atlantic Subpolar Gyre (SPG) (typically below
 112 1000 m) (Counillon et al., 2016), which subsequently causes a drift in the multi-year pre-
 113 dictions in the North Atlantic SPG region (Bethke et al., 2018). We aim to address this
 114 limitation in our current assimilation system.

115 We test the performance of NorCPM for coupled reanalysis in the framework of ide-
 116 alised twin experiments with the assimilation of SST. We assess whether 1) hybrid co-

117 variance can solve this issue and compare the performance of a 31-year coupled reanal-
 118 ysis produced with the EnKF (currently used in NorCPM) and hybrid covariance meth-
 119 ods, 2) compare robustness and optimality of two flavours of hybrid covariance meth-
 120 ods: one using an empirically tuned globally constant hybridization coefficient and one
 121 where the hybridization coefficients are estimated adaptively in space and time (Ménétrier
 122 et al., 2015; Ménétrier & Auligné, 2015; Ménétrier, 2024).

123 This paper is organised as follows. Section 2 presents the Earth System Model (ESM)
 124 used in this work, the Norwegian Earth System Model (NorESM). Section 3 presents the
 125 deterministic EnKF, and its practical implementation within the NorCPM. Section 4 de-
 126 scribes hybridization with a static covariance matrix. Section 5 describes the adaptive
 127 covariance hybridization method and its practical implementation within NorCPM for
 128 SST assimilation. In section 6, we introduce the experimental design and the evaluation
 129 metrics. The numerical results are presented in section 7. Section 8, provides a conclu-
 130 sion, discussion and future perspective to this work.

131 **2 Model system: the Norwegian Earth System Model**

132 NorESM1-ME (Bentsen et al., 2013) is based on version 1.0.4 of the Community Earth
 133 System Model (Hurrell et al., 2013). Its atmosphere component is the CAM4-OSLO, the
 134 ocean component is the Bergen Layered Ocean Model, BLOM, (Bentsen et al., 2013),
 135 the land component is the Community Land Model, CLM4, (Lawrence et al., 2011), the
 136 sea ice component is the Los Alamos Sea Ice Model, CICE4, (Bitz et al., 2012), and the
 137 coupler is CPL7 (Craig et al., 2012).

138 The atmosphere and the land components have $1.9^\circ \times 2.5^\circ$ latitude-longitude res-
 139 olution. The atmosphere component has 26 hybrid sigma-pressure levels ranging from
 140 the surface up to 3 hPa. The ocean and the sea ice components have a 1° horizontal res-
 141 olution in both latitude and longitude with a bipolar grid. BLOM comprises in the ver-
 142 tical a stack of 51 isopycnic layers and two layers for representing the bulk mixed layer.
 143 Before 2005, the forcings are the CMIP5 historical forcings (Taylor et al., 2012), while
 144 after 2005, they are the representative Concentration Pathway 8.5 forcings (van Vuuren
 145 et al., 2011).

146 **3 The deterministic ensemble Kalman filter**

147 Let $\mathbf{X} \in \mathbb{R}^{n \times N}$ an ensemble of N model states ($\mathbf{x}_1, \mathbf{x}_2, \dots, \mathbf{x}_N$), $\mathbf{x} \in \mathbb{R}^n$ the en-
 148 semble mean and $\mathbf{A} \in \mathbb{R}^{n \times N}$ the ensemble anomalies. We note n the model state di-
 149 mension. \mathbf{x} and \mathbf{A} are given by Eq. (1) and (2):

$$150 \quad \mathbf{x} = \frac{1}{N} \mathbf{X} \mathbf{1}, \quad (1)$$

$$151 \quad \mathbf{A} = \mathbf{X} \left(\mathbf{I} - \frac{1}{N} \mathbf{1} \mathbf{1}^T \right), \quad (2)$$

152 where $\mathbf{I} \in \mathbb{R}^{N \times N}$ is the identity matrix and $\mathbf{1} \in \mathbb{R}^N$ is a vector with all elements equal
 153 to 1. In the following equations, the superscripts f and a stand respectively for the fore-
 154 casted and analysed states of the mean and the anomalies.

155 We note \mathbf{x}^t the true state of the system. The observations \mathbf{y} are defined by Eq. (3):

$$156 \quad \mathbf{y} = \mathbf{H} \mathbf{x}^t + \varepsilon, \quad \varepsilon \sim \mathcal{N}(\mathbf{0}, \mathbf{R}), \quad (3)$$

157 where $\mathbf{H} \in \mathbb{R}^{p \times n}$ is the observation operator, $\mathbf{R} \in \mathbb{R}^{p \times p}$ is the observation error co-
 158 variance matrix, and p is the number of observations.

159 In this study, the deterministic EnKF (DEnKF) introduced by Sakov and Oke (2008)
 160 is used. The DEnKF is a square-root (deterministic) flavour of the EnKF that solves the
 161 analysis without the need for perturbation of the observations. It inflates the error by
 162 construction and is intended to perform well in applications where corrections are small
 163 (Sakov & Oke, 2008). The scheme has been robustly tested and validated (Sakov et al.,
 164 2012; Counillon et al., 2016; Bethke et al., 2021). The DEnKF decomposes into two steps:
 165 a forecast step and an analysis step. In the forecast step, each analyzed member \mathbf{x}_i^a is
 166 integrated forward in time, which becomes the prior \mathbf{x}_i^f at the following assimilation cycle:
 167

$$168 \quad \mathbf{x}_i^f = \mathcal{M}(\mathbf{x}_i^a), \quad i = 1, \dots, N, \quad (4)$$

169 where \mathcal{M} is an operator that stands for the model integration.

170 The analysis step of the DEnKF proceeds in two steps, the update of the ensemble
 171 mean, Eq. (5), and the update of the ensemble anomalies, Eq. (6):

$$172 \quad \mathbf{x}^a = \mathbf{x}^f + \mathbf{K}(\mathbf{y} - \mathbf{H}\mathbf{x}^f), \quad (5)$$

$$173 \quad \mathbf{A}^a = \mathbf{A}^f - \frac{1}{2}\mathbf{K}\mathbf{H}\mathbf{A}^f, \quad (6)$$

174 where:

$$175 \quad \mathbf{K} = \mathbf{P}^f \mathbf{H}^T (\mathbf{H} \mathbf{P}^f \mathbf{H}^T + \mathbf{R})^{-1}, \quad (7)$$

$$176 \quad \mathbf{P}^f = \frac{\mathbf{A}^f (\mathbf{A}^f)^T}{N - 1}, \quad (8)$$

177 are respectively the Kalman gain matrix and the background error covariance matrix es-
 178 timated from the ensemble anomalies.

179 In the following, the DEnKF will be referred to as EnKF since general conclusions
 180 of this work are independent of the flavour of the EnKF analysis scheme used.

181 Applying an EnKF with a large dimensional system requires few ad-hoc implemen-
 182 tations. In order to avoid a too abrupt start of assimilation, the variance of the obser-
 183 vation error is multiplied by a factor 8 at the first assimilation cycle and is then reduced
 184 by 1 every two months until it reaches 1 over the course of 14 months. We use the rfac-
 185 tor inflation scheme (Sakov et al., 2012), for which the observation error is inflated by
 186 2 when updating the ensemble anomaly in Eq. 6. We also use pre-screening of the ob-
 187 servation; *i.e.*, the observation error variance is inflated so that the analysis remains within
 188 2 standard deviations of the forecast error from the ensemble mean of the forecasts (Sakov
 189 et al., 2012). We also used the upscaling method (Y. Wang et al., 2016) that handles the
 190 update of the water layers thickness (truncated Gaussian) and better preserve mass, heat
 191 and salt. For producing long reanalysis (from 1850) with SST assimilation, we use the
 192 Hadley Centre Sea Ice and Sea Surface Temperature (HadISST2.1.0.0) available from
 193 1850–present. This type of product is practical because it handles the biases between
 194 different data set and provide a grided ensemble SST that can be used to quantify the
 195 uncertainty. Still, it is constructed by statistical interpolation/extrapolation from the
 196 raw data and the neighbouring observation errors are highly correlated. Our assimila-
 197 tion code assumes the observation error to be independent (*i.e.*, \mathbf{R} is diagonal) and it
 198 was therefore decided to only retain the nearest SST observation to update the water
 199 column (we speak of “strong localization”). For more details about the implementation
 200 of the EnKF within NorCPM, see Counillon et al. (2014) and Counillon et al. (2016).

4 Background error covariance matrix hybridization

The dynamical covariance matrix \mathbf{P}_d^f is estimated from the dynamic ensemble \mathbf{X}_d . The size of \mathbf{X}_d is limited to 30 members in the current version of NorCPM. Such a small ensemble size results in spurious covariances (Anderson, 2007; Bishop & Hodyss, 2007) and rank deficiencies (Oke et al., 2007). Background error covariance hybridization was initially introduced by Hamill and Snyder (2000) to combine an EnKF with a 3DVar and bring some flow-dependency in variational data assimilation. Covariance hybridization has been used in sequential ensemble data assimilation by X. Wang et al. (2007), Counillon et al. (2009) and Gharamti et al. (2014) (hereafter referred to as EnKF-OI) as a way to limit the impact of under-sampling and rank deficiency. The background error covariance matrix combines linearly a dynamical covariance matrix \mathbf{P}_d^f with another covariance matrix \mathbf{P}_s^f (where the subscript s stands for static) computed from a climatological ensemble of size N_s , \mathbf{X}_s (where $N_s \gg N_d$). That static ensemble is constructed by gathering model outputs before running the assimilation experiment. As such, the EnKF-OI does not increase the computational cost of the integration step and has only a limited impact on the computational cost of the analysis step (Counillon et al., 2009). In the special case of the EnKF-OI in NorCPM the CPU-time of the EnKF-OI is 7% larger than that of the standard EnKF; we discuss in Section 8 possible ways to reduce the computational cost of the EnKF-OI in NorCPM. We denote \mathbf{P}_h^f the hybrid covariance matrix:

$$\mathbf{P}_h^f = \alpha_d \mathbf{P}_d^f + \alpha_s \mathbf{P}_s^f, \quad \alpha_d, \alpha_s \geq 0 \quad (9)$$

Unless explicitly mentioned, the sum of the coefficients α_d and α_s can be different from 1.

The update of the mean and the anomalies with the EnKF-OI writes:

$$\mathbf{x}_d^a = \mathbf{x}_d^f + \mathbf{K}_h (\mathbf{y} - \mathbf{H} \mathbf{x}_d^f), \quad (10)$$

$$\mathbf{A}_d^a = \mathbf{A}_d^f - \frac{1}{2} \mathbf{K}_h \mathbf{H} \mathbf{A}_d^f, \quad (11)$$

where \mathbf{K}_h is the hybrid Kalman gain,

$$\mathbf{K}_h = \mathbf{P}_h^f \mathbf{H}^T \left(\mathbf{H} \mathbf{P}_h^f \mathbf{H}^T + \mathbf{R} \right)^{-1}. \quad (12)$$

In practice, we do not compute explicitly the hybrid covariance matrix \mathbf{P}_h :

$$\mathbf{P}_h = \frac{\mathbf{A}_h (\mathbf{A}_h)^T}{N_h - 1}, \quad (13)$$

where \mathbf{A}_h stands for the hybrid anomalies (Counillon et al., 2009):

$$\mathbf{A}_h = \sqrt{N_h - 1} \left[\sqrt{\frac{\alpha_d}{N_d - 1}} \mathbf{A}_d, \sqrt{\frac{\alpha_s}{N_s - 1}} \mathbf{A}_s \right], \quad (14)$$

and $[\cdot, \cdot]$ stand for the concatenation of two sets of vectors. Therefore, the EnKF-OI is implemented within NorCPM following Evensen (2003) and Sakov et al. (2010).

5 Adaptive covariance hybridization

5.1 Explicit optimality of the hybridization coefficients

In this section, we build on the work of Ménétrier et al. (2015); Ménétrier and Auligné (2015) and adapt a method proposed in Ménétrier (2024) to determine the optimal hybridization coefficients.

We take advantage of the strong localization setting (see Section 3) to consider only the covariance between a single point at the surface of the ocean and the multivariate state vector of the water column below so that the covariance matrix in each point reduces to a vector. In the following, we note \mathbf{P} the asymptotic covariance vector at a given point (which would be obtained with an infinite ensemble) from which we only have a noisy estimation \mathbf{P}_d . We use the subscript i to refer to the i^{th} element of the covariance vector (*e.g.*, \mathbf{P}_i for the i^{th} element of the asymptotic covariance vector). We define the scalar product of two random vectors \mathbf{X} and \mathbf{Y} as $\mathbb{E}[\mathbf{X} \cdot \mathbf{Y}] = \sum_{i=1}^p \mathbb{E}[\mathbf{X}_i \mathbf{Y}_i]$. We assume that \mathbf{P}_d and \mathbf{P} are realizations of two independent random processes and that the sampling error of \mathbf{P}_d , *i.e.*, $\mathbf{P}_d - \mathbf{P}$, is unbiased, and orthogonal to the asymptotic covariance, (see Ménétrier et al. (2015), section 5.a and Ménétrier (2024), Eqs. (108)-(114)). Therefore:

$$\mathbb{E}[\mathbf{P}_d - \mathbf{P}] = 0, \quad (15a)$$

$$\mathbb{E}[(\mathbf{P}_d - \mathbf{P}) \cdot \mathbf{P}] = 0. \quad (15b)$$

The optimal hybridization coefficients (α_d, α_s) are defined as those minimizing, in a statistical sense, the square of the error between \mathbf{P}_h and \mathbf{P} , *i.e.*, (α_d, α_s) minimize the function e :

$$e(\alpha_d, \alpha_s) = \mathbb{E}[\|\mathbf{P}_h - \mathbf{P}\|^2] = \mathbb{E}[\|\alpha_d \mathbf{P}_d + \alpha_s \mathbf{P}_s - \mathbf{P}\|^2], \quad (16)$$

where $\|\cdot\|$ stands for the L^2 -norm of a vector. It can be shown (see Appendix A) that minimizing the function e is equivalent to solving a system of two equations, with unknown α_d and α_s , whose solution is given by:

$$(\alpha_d, \alpha_s) = \left(\frac{n_d}{\Delta}, \frac{n_s}{\Delta} \right), \quad (17)$$

where:

$$\Delta = \|\mathbf{P}_s\|^2 \mathbb{E}[\|\mathbf{P}_d\|^2] - \mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s]^2, \quad (18)$$

$$n_d = \|\mathbf{P}_s\|^2 \mathbb{E}[\|\mathbf{P}\|^2] - \mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s]^2, \quad (19)$$

$$n_s = \left(\mathbb{E}[\|\mathbf{P}_d\|^2] - \mathbb{E}[\|\mathbf{P}\|^2] \right) \mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s]. \quad (20)$$

5.2 Properties

The properties highlighted by Ménétrier et al. (2015) in the case of the dual optimization of localization and hybridization hold here:

1. **Behavior of the hybridization coefficients:** if the static covariance is multiplied by a factor λ , then α_s is divided by λ , while α_d remains unchanged. As such,

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it is not necessary to tune the static covariance with a scalar a priori, as done in Evensen (2003), Oke et al. (2008) and Counillon et al. (2009).

2. **Asymptotic behavior:** with an infinite ensemble, $\mathbb{E} [\|\mathbf{P}_d\|^2] = \mathbb{E} [\|\mathbf{P}\|^2]$. We can replace $\mathbb{E} [\|\mathbf{P}_d\|^2]$ by $\mathbb{E} [\|\mathbf{P}\|^2]$ in Eqs. (19)-(20), and obtain $(\alpha_d, \alpha_s) = (1, 0)$ as expected – there is no need for hybridization.

3. **Benefits of hybridization:** whatever the choice of the static covariance (see Appendix B),

$$e(1, 0) \geq e(\alpha_d, \alpha_s), \quad (21)$$

showing the superiority of the hybrid scheme over the standalone EnKF.

4. **Optimality condition:** at optimality the following equivalence is verified, see Ménétrier et al. (2015); Ménétrier and Auligné (2015):

$$\begin{cases} \frac{\partial e}{\partial \alpha_d} = 0 \\ \frac{\partial e}{\partial \alpha_s} = 0 \end{cases} \Leftrightarrow \begin{cases} \mathbb{E} [(\mathbf{P}_h - \mathbf{P}) \cdot \mathbf{P}_d] = 0 \\ \mathbb{E} [(\mathbf{P}_h - \mathbf{P}) \cdot \mathbf{P}_s] = 0 \end{cases}, \quad (22)$$

which means that there is optimality if and only if the error of the hybrid to the asymptotic covariance is independent from both \mathbf{P}_d and \mathbf{P}_s .

Here are some remarks:

1. if $\alpha_d \neq 1$, \mathbf{P}_h can be interpreted as the linear interpolation between \mathbf{P}_d and $\lambda \mathbf{P}_s$, where:

$$\lambda = \frac{\alpha_s}{1 - \alpha_d}, \quad (23)$$

thus, $\lambda \geq 1$ (resp ≤ 1) is equivalent to $\alpha_d + \alpha_s \geq 1$ (resp. ≤ 1). λ acts as an inflation or deflation term for the matrix \mathbf{P}_s and \mathbf{P}_h is the linear interpolation between \mathbf{P}_d and the inflated/deflated \mathbf{P}_s .

2. The numerator of α_d , n_d , can be interpreted as a measure of the collinearity of the static covariance \mathbf{P}_s , and the expectation of the asymptotic covariance \mathbf{P} . Hence, n_d is equal to 0 if and only if $\text{Var}(\mathbf{P}) = 0$ and \mathbf{P} and \mathbf{P}_s are collinear. Similarly, the denominator of α_d and α_s is a measure of the collinearity of \mathbf{P}_s and $\mathbb{E}[\mathbf{P}_d]$, see Appendix C. Conversely, $\mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s] = 0$ ($\mathbb{E}[\mathbf{P} \cdot \mathbf{P}_s] = 0$ respectively) if and only if $\mathbb{E}[\mathbf{P}_d]$ and \mathbf{P}_s ($\mathbb{E}[\mathbf{P}]$ and \mathbf{P}_s) are orthogonal.

3. It follows from Eq. (15b) that (see Appendix B, Eqs. (B1)-(B2)):

$$\mathbb{E} [\|\mathbf{P}_d - \mathbf{P}\|^2] = \mathbb{E} [\|\mathbf{P}_d\|^2] - \mathbb{E} [\|\mathbf{P}\|^2] \geq 0. \quad (24)$$

Therefore, the difference $\mathbb{E} [\|\mathbf{P}_d\|^2] - \mathbb{E} [\|\mathbf{P}\|^2]$ can be interpreted as a measure of the optimality (or the non-optimality) of the covariance function computed from the dynamic ensemble \mathbf{P}_d : the smaller the difference, the smaller the distance $\|\mathbf{P}_d - \mathbf{P}\|$ in a statistical sense. Conversely, the larger the difference, the larger the distance $\|\mathbf{P}_d - \mathbf{P}\|$.

4. It follows from remarks 3 and 4 that the hybridization coefficients α_d and α_s are the combination of the optimality of the dynamic covariance function \mathbf{P}_d and the collinearity/orthogonality of the static covariance \mathbf{P}_s and the expectation of the asymptotic covariance \mathbf{P} /the dynamic covariance \mathbf{P}_d .

5. As a consequence of Eq.(24) and Appendix C, $0 \leq \alpha_d \leq 1$. We can not give similar upper and lower boundaries for the values of α_s , as the term $\mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s]$ can be negative and we do not know its lower bound. Numerical simulations showed that this term is almost always positive (not shown). We can just say that α_s is maximal when Δ is minimal and therefore $\mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s]$ is maximal. In that case, α_d is minimal.

- 315 6. The hybrid scheme exhibited in this work does not consider model bias and takes
 316 only into account the sampling noise in the computation of the hybridization co-
 317 efficient. However, following (Ménétrier, 2024), section 10.8, we can replace \mathbf{P} in
 318 Eq. (16) by an hybrid target $\bar{\alpha}_d \mathbf{P} + \bar{\alpha}_s \mathbf{P}_s$. This reduces to defining a new hybrid
 319 matrix $\mathbf{P}_h = \bar{\alpha}_d \mathbf{P}_h + \bar{\alpha}_s \mathbf{P}_s$, see Appendix D, where \mathbf{P}_h is the hybrid matrix com-
 320 puted with the approach followed in this paper. The coefficients $\bar{\alpha}_d$ and $\bar{\alpha}_s$ can
 321 be computed by a method that takes into account the model bias, see for exam-
 322 ple (Gharamti, 2020).
- 323 7. In this study, we have used one coefficient per water column in order to preserve
 324 the dynamical consistency of the model *e.g.* the relationship between the ocean
 325 bottom pressure and the layer thicknesses. The method could be adapted to use
 326 different hybridization coefficients in the vertical over one water column in order
 327 to further minimize the distance between the hybrid covariance \mathbf{P}_h and the asymp-
 328 totic covariance \mathbf{P} . In that case, a strong post-processing is needed to deal with
 329 the dynamical consistency in the vertical.

330 5.3 Practical implementation

331 Quantities in Eq. (19) and Eq. (20) can not be computed directly as they are a function
 332 of $\mathbb{E} [\|\mathbf{P}\|^2]$, $\mathbb{E} [\mathbf{P}_d \cdot \mathbf{P}_s]$, and $\mathbb{E} [\|\mathbf{P}_d\|^2]$ that are unknown.

333 Nonetheless, the sampling theory developed in Ménétrier (2024) allows us to express
 334 $\mathbb{E} [\mathbf{P}_i^2]$, $i = 1, \dots, p$ as a function of the covariance and variance of the dynamic ensem-
 335 ble. Using Eq. (123a) of Ménétrier (2024) one can write:

$$336 \mathbb{E} [\mathbf{P}_i^2] = \frac{(N_d - 1)^2}{N_d (N_d - 3)} \mathbb{E} [\mathbf{P}_{di}^2] + \frac{N_d - 1}{N_d (N_d - 2) (N_d - 3)} \mathbb{E} [\mathbf{v}_{di} \mathbf{v}_{d1}] - \frac{N_d}{(N_d - 2) (N_d - 3)} \mathbb{E} [\Xi_{d1i}], \quad (25)$$

337 where where \mathbf{v}_{di} is the variance of the dynamic ensemble for the i -th element of the model
 338 state and:

$$339 \Xi_{d1i} = \frac{1}{N_d} \sum_{k=1}^{N_d} \mathbf{A}_{d1k}^2 \mathbf{A}_{dik}^2, \quad (26)$$

340 where \mathbf{A}_{dik} is the i -th element of the anomaly of the k -th member of the dynamic en-
 341 semble.

342 In the case where the ensemble is Gaussian, one can write using Eq. (127a) of Ménétrier
 343 (2024):

$$344 \mathbb{E} [\mathbf{P}_i^2] = \frac{(N_d - 1)^2}{(N_d - 2) (N_d + 1)} \mathbb{E} [\mathbf{P}_{di}^2] - \frac{N_d - 1}{(N_d - 2) (N_d + 1)} \mathbb{E} [\mathbf{v}_{di} \mathbf{v}_{d1}]. \quad (27)$$

345 The quantities n_d , n_s , and Δ are then fully expressed as a function of the static covari-
 346 ance and the expectation of the sample variance, the sample covariance and the sam-
 347 ple fourth-order centered moment, Ξ_d , of the dynamic ensemble.

348 The expectation terms in Eq. (18), (19), (20), and (27) are estimated under a sim-
 349 plifying assumption of “local homogeneity”: it is assumed that in an area surrounding
 350 the water column, the vertical covariance functions are representative of the covariance
 351 function of the water column. The expectation terms are then estimated as the average
 352 of the surrounding vertical covariance functions. Numerical tests have shown that at least
 353 500 covariance functions are necessary to compute reliable statistics. In this work, we

354 consider covariance functions in a radius up to 1000 km around the water column, which
 355 usually provides between 500 and 1000 covariance functions to compute the expectation
 356 terms. Numerical experiments (not shown) indicated that the hybridization coefficients
 357 converge to the values exhibited in this study for a radius larger than 750 km.

358 In order to limit the computational burden of estimating the hybridization coeffi-
 359 cients, they are computed on a subgrid of the domain (every 5 grid cells). The hybridiza-
 360 tion coefficients are then interpolated to the remaining wet points using linear interpo-
 361 lation of the neighbouring wet points.

362 The hybridization coefficients are estimated based on both temperature and salin-
 363 ity as they are key oceanic variables. Doing so yields, on one hand, a lower root mean
 364 square error than when computing the hybridization coefficients solely based on the tem-
 365 perature (not shown) and, on the other hand, a lower computational cost than when adding
 366 variables like *e.g.* the velocity. In other words, this is the best compromise between ac-
 367 curacy and computational efficiency.

368 6 Experimental design and evaluation metrics

369 6.1 Experimental design

370 The free ensemble run (hereafter referred to as FREE), consists of 30 members run
 371 with transient forcing from 1850 to 2014. The true run (hereafter referred to as TRUE)
 372 is created by spawning one member (adding noise to surface temperature) on member
 373 1 of FREE in 1960 and running it up to 2010. It was verified in Y. Wang et al. (2022)
 374 that TRUE and member 1 of FREE were fully de-synchronised at the start of the ex-
 375 periment in 1980. The synthetic observations of SST are generated by adding white noise
 376 to the monthly SST of TRUE. The amplitude of the noise is set equal to the observa-
 377 tion uncertainty (in space and time) of HadISST2. As in the real framework for assim-
 378 ilation of SST, we do not use SST data under sea ice.

379 We produce reanalyses with monthly assimilation of SST observations from January
 380 1980 to December 2010. All experiments start with the same initial dynamic ensemble
 381 (taken from FREE in January 1980). The static ensemble is made from the monthly restarts
 382 of a 315 years stable pre-industrial run. It is different for each calendar month and we
 383 use the same static ensemble for the same month of two different years. This implies that
 384 we have in total 12 different static ensembles, one for each calendar month. The exper-
 385 iments are separated into three categories:

- 386 • **EnKF:** the standard EnKF used in NorCPM (Counillon et al., 2016).
- 387 • **Standard hybrid:** a constant and global hybridization coefficients (see Section
 388 4). The sum of α_d and α_s is 1. We performed 7 reanalyses with $\alpha_d = 0, 0.2, 0.4, 0.6, 0.8, 0.9, 1$.
 389 The case where $\alpha_d = 0$ is equivalent to an ensemble of EnOI, and the case where
 390 $\alpha_d = 1$ is equivalent to the standard EnKF.
- 391 • **Adaptive hybrid:** the hybridization coefficients are estimated at each assim-
 392 ilation cycle, they vary spatially and their sum is not imposed equal to 1 (see Sec-
 393 tion 5).

394 6.2 Evaluation metrics

395 The accuracy of the reanalyses is estimated based on the root mean square error (RMSE).
 396 The RMSE is computed as:

$$397 \text{RMSE} = \sqrt{\sum_{i=1}^N \omega_i (\mathbf{x}_i - \mathbf{x}_i^t)^2}. \quad (28)$$

398 In the following, the RMSE is computed either over a time series at a given point
 399 (in which case $\omega_i = \frac{1}{N}$), or over the whole domain at a given time (in which case ω_i
 400 is the relative size of the grid cell).

401 In order to easily compare RMSE between the nine different schemes (see Section 6.1),
 402 we introduce the Mean Skill Score of one configuration i , MSS_i . It is the relative reduc-
 403 tion of RMSE compared to the mean of the RMSE of the nine configurations, Eq. (29):

$$404 \quad MSS_i = 1 - \frac{RMSE_i}{\frac{1}{9} \sum_{j=1}^9 RMSE_j}, \quad (29)$$

405 where $RMSE_i$ is the RMSE of one of the schemes. The MSS is 1 if the scheme is per-
 406 fect (RMSE is equal to 0), between 0 and 1 if the scheme performs better than the mean
 407 of the other schemes and negative otherwise.

408 Another important metric to evaluate the relative efficiency of different data assim-
 409 ilation schemes is to consider the “degrees of freedom for signal” (DFS, Cardinali et al.,
 410 2004; Wahba et al., 1995). It can be interpreted as the number of modes of variability
 411 reduced from the ensemble by the assimilation (*i.e.* the assimilation change). The DFS
 412 is defined as follows:

$$413 \quad DFS = \frac{\partial \mathbf{H} \mathbf{x}_d^a}{\partial \mathbf{y}} = \text{Tr}(\mathbf{KH}) \quad (30)$$

414 The DFS is between 0 (*i.e.*, the observations have no impact on the ensemble), and
 415 the total number of degrees of freedom (*i.e.*, observations has collapsed the number of
 416 modes of variability into a single one, Xie et al., 2018). The total number of degrees of
 417 freedom is the minimum between the ensemble size and the number of observations used
 418 for the local assimilation. In NorCPM, in the context of strong localisation (where we
 419 retain only the nearest observation, see Section 3), it implies that the DFS is between
 420 0 and 1 (independently of the ensemble size). This allows for an inter-comparison of the
 421 DFS even though the schemes have different ensemble sizes.

422 7 Results

423 7.1 Stability of the adaptive covariance hybridization

424 The adaptive covariance hybridization method (see Section 5) estimates adaptive hy-
 425 bridization coefficients both in space and time.

426 Figure 1 shows the time series of globally averaged α_d , α_s , and $\alpha_d + \alpha_s$ (sea ice-covered
 427 points where there are no SST data are masked). After a spin-up period of approximately
 428 three years, the averaged values of the hybridization coefficients converge to a global av-
 429 erage of 0.7 for α_d , and 0.175 for α_s . This shows that the mean values of α_d and α_s
 430 are stable in time and display a limited temporal variability despite an important spatial vari-
 431 ability, and so does the sum $\alpha_d + \alpha_s$. Hence, the mean values of α_d and α_s computed
 432 in specific basins show similar behaviour and converge within 3 years (not shown). The
 433 global averaged value of $\alpha_d + \alpha_s$ is roughly 0.875. Following remark 1 in Section 5.2,
 434 it implies that the static ensemble has a larger error variance than the error growth within
 435 one assimilation cycle and needs to be reduced - in agreement with Oke et al. (2008),
 436 Counillon and Bertino (2009), and Evensen (2003).

437 Figures 2 shows the pointwise averaged map of α_d and α_s computed over the years
 438 1983 to 2010 for January and July. The values of α_d display an important spatial vari-
 439 ability with values ranging approximately from 0.4 up to 1. Regions where α_d is small
 440 coincide with places where α_s is large. The spatial variations of the values of α_d and α_s

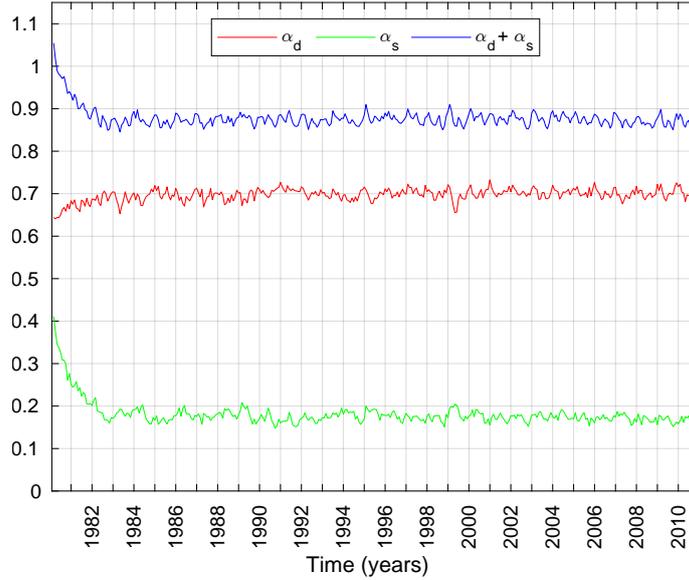


Figure 1. Time series of the global average values of α_d , α_s , and $\alpha_d + \alpha_s$. Ice-covered regions are masked.

441 (Section 5.2) can be explained from the perspective of the optimality of the dynamic covar-
 442 iance \mathbf{P}_d (depending on the sampling error in the dynamical ensemble), and the collinear-
 443 ity between the static covariance \mathbf{P}_s and the asymptotic covariance \mathbf{P} (meaning that static
 444 covariance is sufficient). Larger values of α_d are found in locations where the dynamic
 445 is non-stationary and internal variability is large; *e.g.*, in the Northern part of the At-
 446 lantic Ocean (Gulf Stream pathway, Subpolar Gyre, near the ice edge), the North Pa-
 447 cific, El Niño–Southern Oscillation and in the Southern Ocean. Conversely, there are rel-
 448 atively low values of α_d in the Indian Ocean where variability is primarily externally forced
 449 (Guemas et al., 2013), the decadal fluctuations are less pronounced than in the Atlantic
 450 or the Pacific Ocean and where the Pacific Ocean teleconnections dominate the region-
 451 ally driven variability (Frankcombe et al., 2015). In the tropical Atlantic, the model is
 452 performing very poorly and has no skill (Counillon et al., 2021); it is thus not surpris-
 453 ing that α_d is also low. In the Southern Ocean, in summer, high values of α_d and α_s co-
 454 exist. This implies that the sampling noise of the dynamic covariance is small while the
 455 numerator of α_s is small too. This means that $\mathbb{E}[\mathbf{P}_d]$ and \mathbf{P}_s are close to collinearity and
 456 that $\text{Var}(\mathbf{P}_{di})$ for all i is small according to remark 3, therefore the dynamic covariance
 457 is stationary.

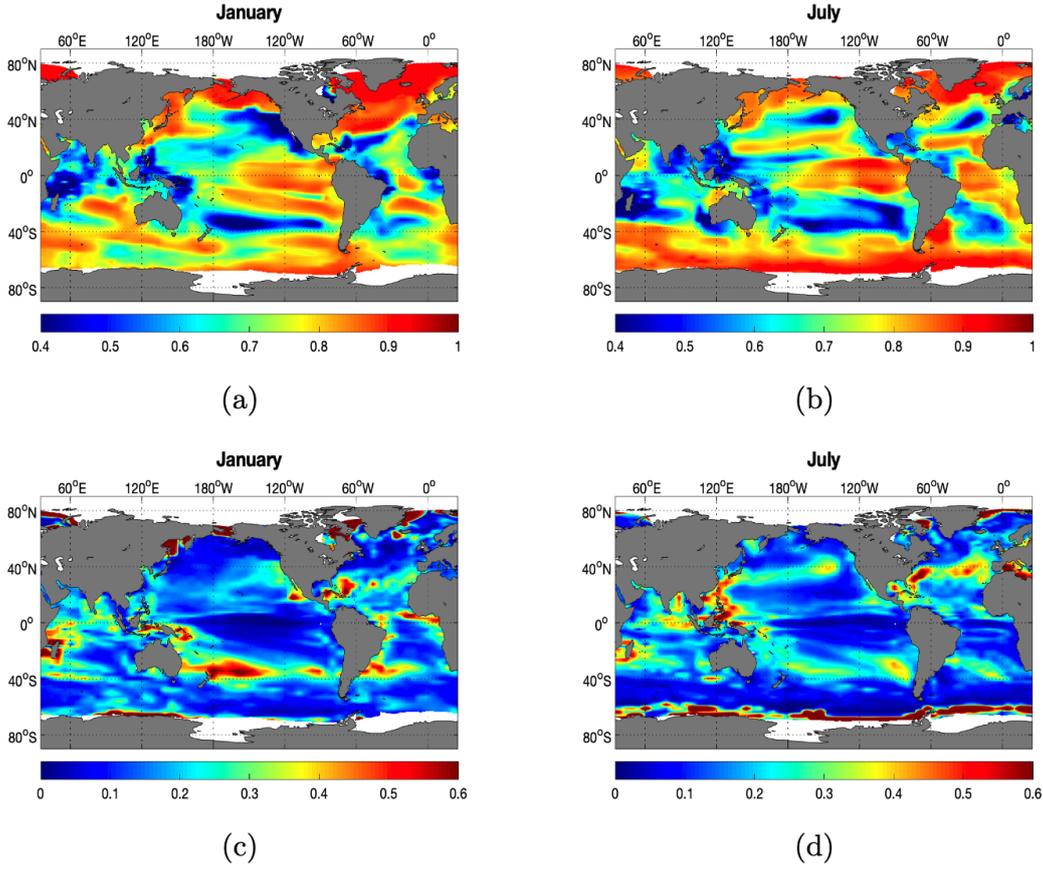


Figure 2. Pointwise averaged estimate of α_d (top row), and of α_s (bottom row) computed over 1983–2010 for the months of January and July.

458 In Fig. 3, we analyse the interannual de-seasoned standard deviation of the hybrid
 459 coefficient beyond year 3 (once it has converged). We can see that the variability is very
 460 small, except in a few places, *e.g.* in the Arctic, in Indian Ocean and in the tropical At-
 461 lantic and Pacific Gyre. In those places, the performance between the standard hybrid
 462 coefficient method is relatively small (not shown).

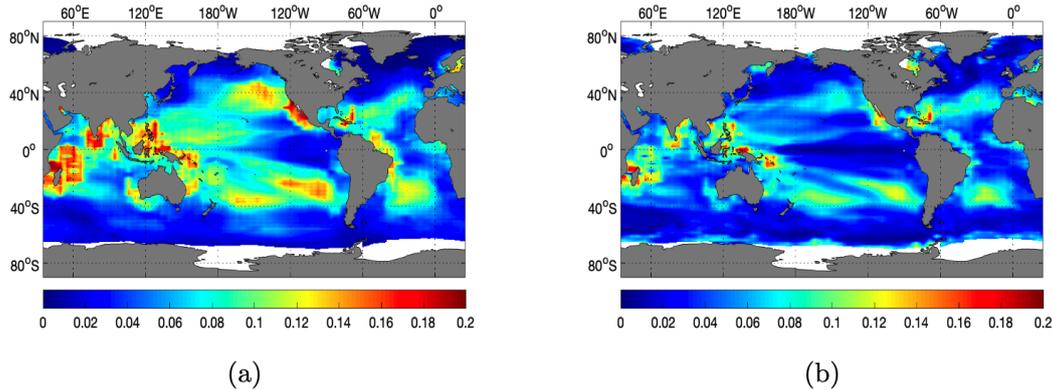


Figure 3. Standard-deviation of the de-seasoned values of α_d (a), and of α_s (b) computed over 1983–2010.

7.2 Intercomparison of the performance of the EnKF and the hybrid covariance schemes

Figure 4 shows the MSS (see Section 6.2) of ocean heat and salt content for the different schemes at different depth ranges (0-200 m, 200-500 m, ..., 2000-4000 m). We include the EnKF ($\alpha_d = 1$ and $\alpha_s = 0$) and the ensemble of EnOI ($\alpha_d = 0$ and $\alpha_s = 1$) as particular cases of the standard hybrid covariance method. A red cell (resp. blue cell) indicates that the scheme provides a reduction (resp. an increase) of RMSE compared to the average performance of all the schemes for a given depth range. For example, the adaptive hybrid and the standard hybrid scheme with $\alpha_d = 0.9$ reduces the RMSE of the temperature at depth 500-1000 m by 10% compared to the average performance, while the standard hybrid with $\alpha_d = 0$ increases the RMSE at the depth 200-500 m by 15%. The results for heat and salt content are very similar. As expected, the EnKF is outperforming the ensemble of EnOI (*i.e.*, $\alpha_d = 0$), showing the superiority of flow-dependent covariance over static covariance. It also shows the importance of tuning the hybrid coefficient as for a large span of standard hybrid coefficient values, the hybrid covariance methods perform poorer than the EnKF. When α_d is larger than 0.8 ($\alpha_d = 0.9$ being optimal), the standard hybrid covariance outperforms the EnKF; notably between 2000 and 4000 m. In the latter, the error is gradually increasing due to spurious covariance at depth (Y. Wang et al., 2022; Bethke et al., 2018). In the former, the improvements are mostly the consequence of the filtering of the sampling noise with the static covariance, especially in the North Pacific and the North Atlantic. In the Weddell Sea and the Southern Ocean, there is virtually no hybridization (with small values of α_s associated with a very limited variability), and the improvement is associated to deflation/regularization of the dynamic covariance with the dynamic coefficient α_d . The adaptive hybrid covariance method performs best at nearly all depth levels for heat and salt content. In the following, we will therefore present the adaptive hybrid and assess the spatial distribution of the improvements over the EnKF, but results with the best standard hybrid are nearly comparable (not shown).

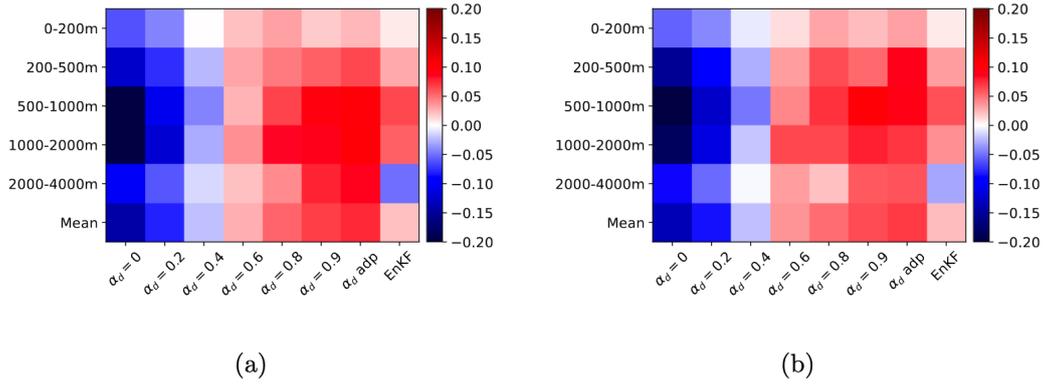


Figure 4. MSS of all the schemes for temperature (a) and salinity (b) at different depth intervals. $\alpha_d = 0, \alpha_d = 0.2, \dots, \alpha_d = 0.9$ refer to the hybridization coefficient of the standard hybrid. $\alpha_d = 0$ is the ensemble of EnOI, full static case, and $\alpha_d = 1$ is the EnKF – the default scheme used in NorCPM. "α adp" stands for the adaptive hybridization scheme. The warm colour indicates that the scheme performs better than the average skill of all systems.

491 The adaptive hybrid and the EnKF achieve similar performance in the top 1000 m
 492 (Figure 4), and we focus on performance below this depth range. We compare the re-
 493 duction of RMSE of the EnKF and the adaptive hybrid compared to that of FREE for
 494 two depths range.

495 Between 1000 and 2000 m (Fig. 5), the EnKF reduces the error overall (warmer value
 496 is predominant) compared to FREE. Still, there are few places where it increases the RMSE
 497 of temperature, *e.g.*: in the North Pacific, the subtropical Atlantic, and near the Wed-
 498 dell Sea. Results are relatively comparable for salt content (see Fig. 5-b). The overall
 499 pattern is similar with the adaptive hybrid. Still, it yields further improvement, as in
 500 the North Atlantic subpolar gyre and it mitigates the degradation in the aforementioned
 501 regions. The degradation in the Weddell Sea is nearly completely removed.

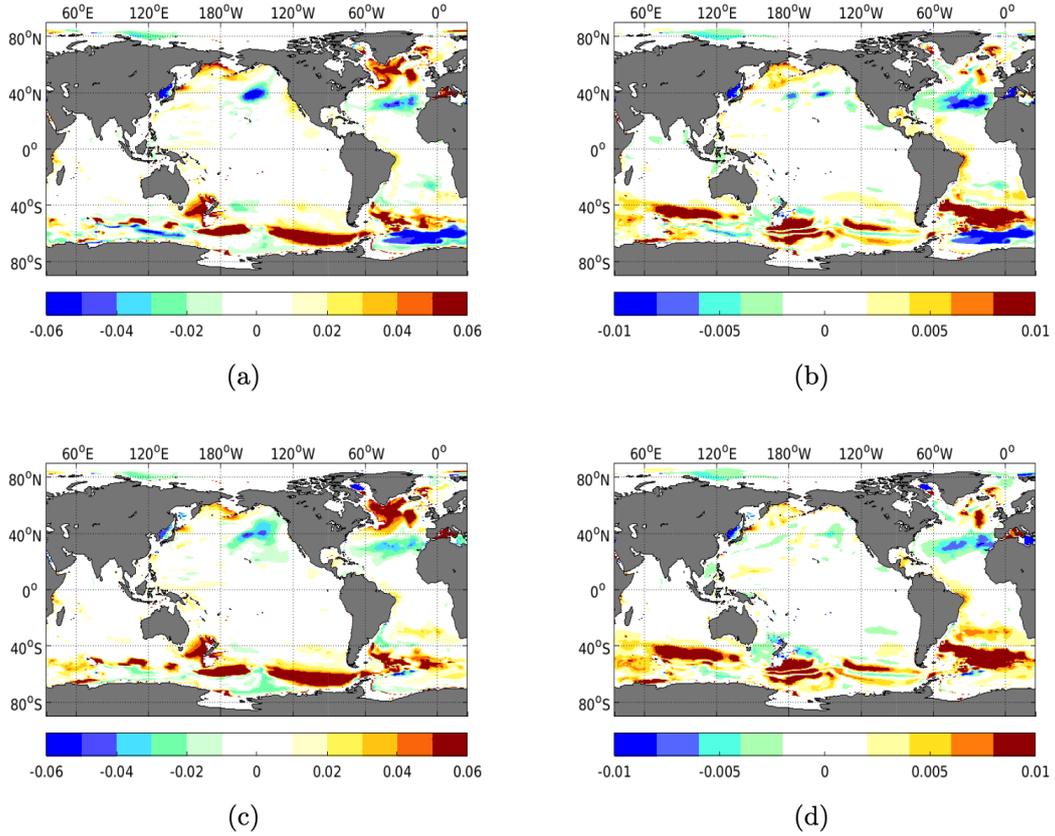


Figure 5. RMSE difference between FREE and the EnKF (a,b), and FREE and the adaptive hybrid (c,d) between 1000 and 2000 m depth for the temperature (left column) and the salinity (right column). Warm colour indicates that assimilation reduces error compared to FREE.

502 Between 2000 and 4000 m (Fig. 6), the EnKF degrades overall performance compared
 503 to FREE. The degradation is larger in the North Pacific, the North Atlantic, and the
 504 Southern Ocean for both the temperature and the salinity. The improvements are also
 505 limited to the South Atlantic Ocean. The adaptive hybrid corrects or mitigates these bi-
 506 ases. Some degradation remains (in the North Atlantic subpolar gyre, the Sea of Japan
 507 in particular for salinity), but the assimilation yields an overall improvement over FREE.

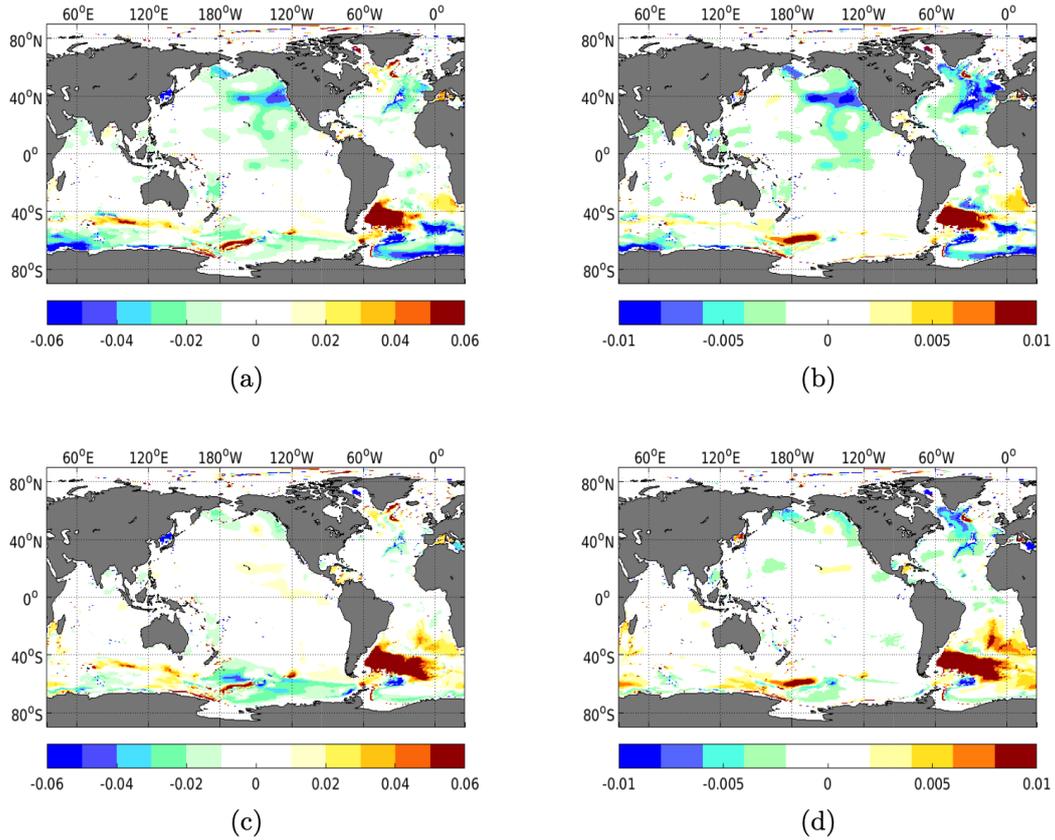


Figure 6. Same as fig. 5 but for 2000–4000 m depth interval.

508 An ideal assimilation system achieves minimal error while making the smallest change
 509 possible during the assimilation. Figure 7 shows the difference of DFS (that quantifies
 510 the assimilation change) between the EnKF and the best standard hybrid ($\alpha_d = 0.9$)
 511 (panel a) and between the EnKF and the adaptive hybrid (panel b). The standard hy-
 512 brid has a larger DFS value than the EnKF (negative values), implying that the data
 513 assimilation induces more change. This is most notable in the Southern Ocean and the
 514 tropical Pacific. In the Southern Ocean, the standard hybrid covariance method performs
 515 better than the EnKF, so it can be argued that the larger corrections are beneficial. How-
 516 ever, in the tropical Pacific, the Δ RMSE of the two remains quite close, meaning the anal-
 517 ysis induces more changes without improving performance. On the contrary, the adap-
 518 tive hybrid, Fig. 7-(b), has a DFS close to that of the EnKF. There are some slight dif-
 519 ferences (in the North Pacific, the North Atlantic, and the Southern Ocean), with a max-
 520 imum in the Irminger Sea, where it strongly outperforms the EnKF (*e.g.*, 1000–2000 m).
 521 It implies that the adaptive hybrid induces only change where this yields improved per-
 522 formance.

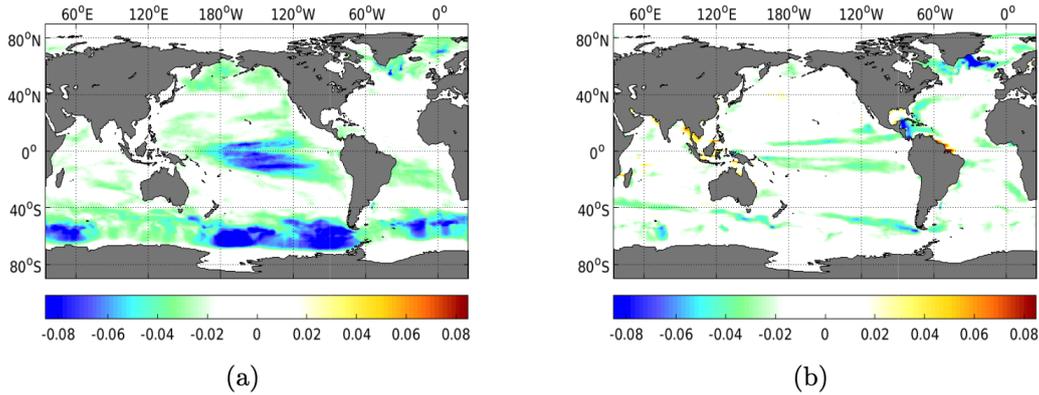


Figure 7. Difference of DFS between the EnKF and the standard hybrid covariance with $\alpha_d = 0.9$ (a), and between EnKF and adaptive hybrid covariance (b). The cold colour indicates that the hybrid covariance yields a larger reduction of DFS than with the EnKF.

8 Discussion and conclusion

In this work, we compare two different versions of hybrid covariance data assimilation with the standard EnKF for producing climate reanalysis. We use the Norwegian Climate Prediction Model (NorCPM) and work in an idealised twin experiment framework. The reanalyses are performed with sole assimilation of SST for the period 1980-2010. In the first hybrid coefficient method, the hybridization coefficients are tuned empirically to optimize the performance, while in the second, the hybridization coefficients are estimated adaptively, both in space and time. The two hybrid coefficient methods outperform the standard EnKF and mitigate the degradation it introduces in the intermediate and deep ocean compared to unassimilated simulations. The adaptive performs best and is doing so by making smaller corrections than the standard hybrid. The hybridization coefficients with the adaptive hybrid are converging quickly (less than 3 years) to stable values and only show small seasonal variations.

Other alternatives have been developed in parallel to address the sampling error with the EnKF in NorCPM – namely the isopycnal vertical localization (Y. Wang et al., 2022). The latter limits the assimilation update of temperature and salinity to a fixed isopycnal level and was shown to mitigate the degradation seen in the standard EnKF. Combining the two approaches is straightforward and will be tested in the future. However, the isopycnal vertical localization detailed in Y. Wang et al. (2022) was tuned for an ensemble size of 30 members, while now the ensemble size is much larger (350 members). The vertical tapering will thus need to be revised.

The adaptive hybrid coefficients method is slightly more expensive than the standard hybrid as it requires additional computation related to the estimation of the hybrid coefficient at each assimilation step and corresponds to an increase of 7.5% of the total CPU-time of the simulation. Here, the hybridization coefficients are estimated at every 5 grid cells, but we could have estimated them at every 10 grid cells (reducing the cost by 4) with a comparable solution (not shown). In that case, the increase of the total CPU-time of the simulation would be only of 0.5% to 1%. Furthermore, as these coefficients converge rapidly to stable estimates (within 3 model years, showing only a weak seasonal variability). They could be stored and directly used instead of being recalculated every time. As such, we do not consider that the additional computational cost would be much larger than the standard hybrid, which also has an additional cost (empirical estimation of the global coefficient).

556 In this study, the estimation of the hybridization coefficients in the adaptive method
 557 is constant in the vertical. Nonetheless, adapting the method to estimate different hy-
 558 bridization coefficients for different vertical levels or variables would be relatively triv-
 559 ial. Furthermore, we tested the method for the particular case of assimilation of SST ob-
 560 jective analysis, where we update a single water column with a single observation, (*i.e.*,
 561 "strong localization"). For the assimilation of temperature/salinity profile data, the ob-
 562 servation error is uncorrelated, and a larger localisation radius is used in NorCPM, the
 563 method can be adjusted following Ménérier (2024).

564 A consequence of the perfect twin experiments setting is that models are unbiased
 565 and that for an infinite ensemble size, the ensemble covariance matrix converges with the
 566 true covariance matrix. These assumptions fall apart with Earth System Models with
 567 considerable biases (Palmer & Stevens, 2019). Handling these biases in data assimila-
 568 tion remains challenging, particularly with sparse and inhomogeneous observation net-
 569 works and where model bias can change, e.g. as a response to external forcings. There-
 570 fore, the method presented in this work would benefit from further developments before
 571 it can be tested in a realistic system. Currently, several methods are being developed and
 572 tested within NorCPM to handle climate biases directly, namely: anomaly coupling (Counillon
 573 et al., 2021), multivariate parameter estimation (Singh et al., 2022), super-resolution data
 574 assimilation (Barthélémy et al., 2022), and supermodelling (Counillon et al., 2023; Scheven-
 575 hoven et al., 2023).

576 Appendix A Minimization of the function e

577 The function e is defined as:

$$578 e(\alpha_d, \alpha_s) = \mathbb{E} \left[\|\alpha_d \mathbf{P}_d + \alpha_s \mathbf{P}_s - \mathbf{P}\|^2 \right]. \quad (\text{A1})$$

579 By linearity of the expectation operator and by definition of the L^2 -norm $\|\cdot\|$, we can
 580 write:

$$581 e(\alpha_d, \alpha_s) = \alpha_d^2 \mathbb{E} \left[\|\mathbf{P}_d\|^2 \right] + \alpha_s^2 \|\mathbf{P}_s\|^2 + \mathbb{E} \left[\|\mathbf{P}\|^2 \right] \\ 582 + 2\alpha_d \alpha_s \mathbb{E} [\mathbf{P}_d \cdot \mathbf{P}_s] - 2\alpha_d \mathbb{E} [\mathbf{P}_d \cdot \mathbf{P}] - 2\alpha_s \mathbb{E} [\mathbf{P} \cdot \mathbf{P}_s]. \quad (\text{A2})$$

583 It follows from Eq. (15b) that:

$$584 \mathbb{E} [\mathbf{P}_d \cdot \mathbf{P}] = \mathbb{E} \left[\|\mathbf{P}\|^2 \right] \quad (\text{A3})$$

585 Replacing Eq. (A3) in Eq. (A2), we obtain the following expression of e :

$$586 e(\alpha_d, \alpha_s) = \alpha_d^2 \mathbb{E} \left[\|\mathbf{P}_d\|^2 \right] + \alpha_s^2 \|\mathbf{P}_s\|^2 + (1 - 2\alpha_d) \mathbb{E} \left[\|\mathbf{P}\|^2 \right] + 2\alpha_s (\alpha_d - 1) \mathbb{E} [\mathbf{P}_d \cdot \mathbf{P}_s]. \quad (\text{A4})$$

587 e being a quadratic function of two variables, α_d and α_s , with positive coefficients
 588 associated to α_d^2 and α_s^2 , its has a unique minimum where both the partial derivatives
 589 with respect to α_d and α_s are null. Hence, minimizing the function e is equivalent to solve
 590 the following system of two equations:

$$591 \begin{cases} \frac{\partial e(\alpha_d, \alpha_s)}{\partial \alpha_d} = 0 \\ \frac{\partial e(\alpha_d, \alpha_s)}{\partial \alpha_s} = 0 \end{cases} \quad (\text{A5})$$

592 The partial derivatives of e with respect to α_d and α_s are given by Eq. (A6) and (A7):

$$593 \quad \frac{\partial e(\alpha_d, \alpha_s)}{\partial \alpha_d} = 2\alpha_d \mathbb{E} \left[\|\mathbf{P}_d\|^2 \right] + 2\alpha_s \mathbb{E} [\mathbf{P}_s \cdot \mathbf{P}_d] - 2\mathbb{E} \left[\|\mathbf{P}\|^2 \right], \quad (\text{A6})$$

$$594 \quad \frac{\partial e(\alpha_d, \alpha_s)}{\partial \alpha_s} = 2\alpha_d \mathbb{E} [\mathbf{P}_s \cdot \mathbf{P}_d] + 2\alpha_s \|\mathbf{P}_s\|^2 - 2\mathbb{E} [\mathbf{P}_s \cdot \mathbf{P}_d]. \quad (\text{A7})$$

595 From which it follows that minimizing the function e is equivalent to solving the sys-
596 tem of two equations:

$$597 \quad \begin{cases} \alpha_d \mathbb{E} \left[\|\mathbf{P}_d\|^2 \right] + \alpha_s \mathbb{E} [\mathbf{P}_d \cdot \mathbf{P}_s] = \mathbb{E} \left[\|\mathbf{P}\|^2 \right] \\ \alpha_d \mathbb{E} [\mathbf{P}_d \cdot \mathbf{P}_s] + \alpha_s \|\mathbf{P}_s\|^2 = \mathbb{E} [\mathbf{P}_d \cdot \mathbf{P}_s], \end{cases} \quad (\text{A8})$$

598 Appendix B Benefits of hybridization

599 By definition of the function e :

$$600 \quad e(1, 0) = \mathbb{E} \left[\|\mathbf{P}_d - \mathbf{P}\|^2 \right]. \quad (\text{B1})$$

601 By linearity of the expectation operator and by definition of the L^2 -norm, we can
602 show that

$$603 \quad e(1, 0) = \mathbb{E} \left[\|\mathbf{P}_d\|^2 \right] + \mathbb{E} \left[\|\mathbf{P}\|^2 \right] - 2\mathbb{E} [\mathbf{P}_d \cdot \mathbf{P}]. \quad (\text{B2})$$

604 Replacing Eq. (A3) in Eq. (B2) we get:

$$605 \quad e(1, 0) = \mathbb{E} \left[\|\mathbf{P}_d\|^2 \right] - \mathbb{E} \left[\|\mathbf{P}\|^2 \right] \quad (\text{B3})$$

606 From which we can express the difference of the errors between the EnKF and the
607 hybrid scheme:

$$608 \quad e(1, 0) - e(\alpha_d, \alpha_s) = (1 - \alpha_d^2) \mathbb{E} \left[\|\mathbf{P}_d\|^2 \right] - 2(1 - \alpha_d) \mathbb{E} \left[\|\mathbf{P}\|^2 \right] \\ 609 \quad - \alpha_s^2 \|\mathbf{P}_s\|^2 - 2\alpha_s(\alpha_d - 1) \mathbb{E} [\mathbf{P}_s \cdot \mathbf{P}_d]. \quad (\text{B4})$$

610 For the sake of simplicity, we note:

$$611 \quad \alpha = \mathbb{E} \left[\|\mathbf{P}\|^2 \right], \quad (\text{B5})$$

$$612 \quad \beta = \|\mathbf{P}_s\|^2, \quad (\text{B6})$$

$$613 \quad \gamma = \mathbb{E} [\mathbf{P}_d \cdot \mathbf{P}_s], \quad (\text{B7})$$

$$614 \quad \delta = \mathbb{E} \left[\|\mathbf{P}_d\|^2 \right]. \quad (\text{B8})$$

615 Eq. (B4) rewrites:

$$e(1, 0) - e(\alpha_d, \alpha_s) = (1 - \alpha_d^2)\delta - 2(1 - \alpha_d)\alpha - \alpha_s^2\beta - 2\alpha_s(\alpha_d - 1)\gamma. \quad (\text{B9})$$

Given that $\alpha_d = \frac{n_d}{\Delta}$ and $\alpha_s = \frac{n_s}{\Delta}$:

$$e(1, 0) - e(\alpha_d, \alpha_s) = \frac{(\Delta^2 - n_d^2)\delta - 2(\Delta^2 - \Delta n_d)\alpha - n_s^2\beta - 2n_s(n_d - \Delta)\gamma}{\Delta^2}. \quad (\text{B10})$$

Δ^2 being positive, showing that $e(1, 0) - e(\alpha_d, \alpha_s) \geq 0$ is equivalent to show that:

$$(\Delta^2 - n_d^2)\delta - 2(\Delta^2 - \Delta n_d)\alpha - n_s^2\beta - 2n_s(n_d - \Delta)\gamma \geq 0. \quad (\text{B11})$$

On the other hand, n_d , n_s , and Δ write:

$$n_d = \alpha\beta - \gamma^2, \quad (\text{B12})$$

$$n_s = \gamma\delta - \alpha\gamma, \quad (\text{B13})$$

$$\Delta = \beta\delta - \gamma^2. \quad (\text{B14})$$

Replacing n_d , n_s , and Δ by their expression given by Eqs. (B12), (B13), and (B14) in the left hand side of Eq. (B11), and developing all the terms we can show that Eq. (B11) is verified if and only if:

$$\beta(\beta\delta - \gamma^2)(\delta - \alpha)^2 \geq 0. \quad (\text{B15})$$

$\beta \geq 0$ as a sum of squares and $(\delta - \alpha)^2 \geq 0$ as a square. Showing that $e(1, 0) - e(\alpha_d, \alpha_s) \geq 0$ is then equivalent to show that $\beta\delta \geq \gamma^2$ i-e $\|\mathbf{P}_s\|^2 \mathbb{E}[\|\mathbf{P}_d\|^2] - \mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s]^2 \geq 0$. This inequality holds, see Appendix C, Eq. (C9), and replacing \mathbf{P} by \mathbf{P}_d . As a consequence, $e(1, 0) \geq e(\alpha_d, \alpha_s)$.

Appendix C Collinearity/orthogonality of $\mathbb{E}[\mathbf{P}]$ and \mathbf{P}_s

Because \mathbf{P}_d is an unbiased estimation of \mathbf{P} , see Eq. (15a), we have: $\mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s] = \mathbb{E}[\mathbf{P} \cdot \mathbf{P}_s]$. Therefore:

$$\|\mathbf{P}_s\|^2 \mathbb{E}[\|\mathbf{P}\|^2] - \mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s]^2 = \|\mathbf{P}_s\|^2 \mathbb{E}[\|\mathbf{P}\|^2] - \mathbb{E}[\mathbf{P} \cdot \mathbf{P}_s]^2. \quad (\text{C1})$$

By definition of $\|\mathbf{P}_s\|^2$ and $\mathbb{E}[\|\mathbf{P}\|^2]$:

$$\|\mathbf{P}_s\|^2 \mathbb{E} [\|\mathbf{P}\|^2] = \sum_{i=1}^p \mathbf{P}_{si}^2 \mathbb{E} [\mathbf{P}_i^2] + \sum_{1 \leq i \neq j \leq p} \mathbf{P}_{si}^2 \mathbb{E} [\mathbf{P}_j^2], \quad (\text{C2})$$

$$\begin{aligned} \|\mathbf{P}_s\|^2 \mathbb{E} [\|\mathbf{P}\|^2] &= \sum_{i=1}^p \mathbf{P}_{si}^2 \left(\text{Var}(\mathbf{P}_i) + \mathbb{E} [\mathbf{P}_i^2] \right) \\ &+ \sum_{1 \leq i \neq j \leq p} \mathbf{P}_{si}^2 \left(\text{Var}(\mathbf{P}_j) + \mathbb{E} [\mathbf{P}_j^2] \right), \end{aligned} \quad (\text{C3})$$

$$\begin{aligned} \|\mathbf{P}_s\|^2 \mathbb{E} [\|\mathbf{P}\|^2] &= \sum_{i=1}^p \mathbf{P}_{si}^2 \mathbb{E} [\mathbf{P}_i^2] + \sum_{1 \leq i \neq j \leq p} \mathbf{P}_{si}^2 \mathbb{E} [\mathbf{P}_j^2] + \sum_{i=1}^p \mathbf{P}_{si}^2 \text{Var}(\mathbf{P}_i) \\ &+ \sum_{1 \leq i \neq j \leq p} \mathbf{P}_{si}^2 \text{Var}(\mathbf{P}_j). \end{aligned} \quad (\text{C4})$$

On the other hand:

$$\mathbb{E} [\mathbf{P} \cdot \mathbf{P}_s]^2 = \sum_{i=1}^p \mathbf{P}_{si}^2 \mathbb{E} [\mathbf{P}_i]^2 + 2 \sum_{1 \leq i < j \leq p} \mathbf{P}_{si} \mathbb{E} [\mathbf{P}_i] \mathbf{P}_{sj} \mathbb{E} [\mathbf{P}_j]. \quad (\text{C5})$$

Therefore:

$$\begin{aligned} \|\mathbf{P}_s\|^2 \mathbb{E} [\|\mathbf{P}\|^2] - \mathbb{E} [\mathbf{P} \cdot \mathbf{P}_s]^2 &= \sum_{i=1}^p \mathbf{P}_{si}^2 \text{Var}(\mathbf{P}_i) + \sum_{1 \leq i \neq j \leq p} \mathbf{P}_{si}^2 \text{Var}(\mathbf{P}_j) \\ &+ \sum_{1 \leq i \neq j \leq p} \mathbf{P}_{si}^2 \mathbb{E} [\mathbf{P}_j]^2 \\ &- 2 \sum_{1 \leq i < j \leq p} \mathbf{P}_{si} \mathbb{E} [\mathbf{P}_i] \mathbf{P}_{sj} \mathbb{E} [\mathbf{P}_j], \end{aligned} \quad (\text{C6})$$

$$\begin{aligned} \|\mathbf{P}_s\|^2 \mathbb{E} [\|\mathbf{P}\|^2] - \mathbb{E} [\mathbf{P} \cdot \mathbf{P}_s]^2 &= \sum_{i=1}^p \mathbf{P}_{si}^2 \text{Var}(\mathbf{P}_i) + \sum_{1 \leq i \neq j \leq p} \mathbf{P}_{si}^2 \text{Var}(\mathbf{P}_j) \\ &+ \sum_{1 \leq i < j \leq p} \mathbf{P}_{si}^2 \mathbb{E} [\mathbf{P}_j]^2 + \mathbf{P}_{sj}^2 \mathbb{E} [\mathbf{P}_i]^2 \\ &- 2 \sum_{1 \leq i < j \leq p} \mathbf{P}_{si} \mathbb{E} [\mathbf{P}_i] \mathbf{P}_{sj} \mathbb{E} [\mathbf{P}_j], \end{aligned} \quad (\text{C7})$$

$$\begin{aligned} \|\mathbf{P}_s\|^2 \mathbb{E} [\|\mathbf{P}\|^2] - \mathbb{E} [\mathbf{P} \cdot \mathbf{P}_s]^2 &= \sum_{i=1}^p \mathbf{P}_{si}^2 \text{Var}(\mathbf{P}_i) + \sum_{1 \leq i \neq j \leq p} \mathbf{P}_{si}^2 \text{Var}(\mathbf{P}_j) \\ &+ \sum_{1 \leq i < j \leq p} (\mathbf{P}_{si} \mathbb{E} [\mathbf{P}_j] - \mathbf{P}_{sj} \mathbb{E} [\mathbf{P}_i])^2 \end{aligned} \quad (\text{C8})$$

As a consequence,

$$\|\mathbf{P}_s\|^2 \mathbb{E} [\|\mathbf{P}\|^2] - \mathbb{E} [\mathbf{P} \cdot \mathbf{P}_s]^2 \geq 0, \quad (\text{C9})$$

as the sum of positive terms, and $\|\mathbf{P}_s\|^2 \mathbb{E} [\|\mathbf{P}\|^2] - \mathbb{E} [\mathbf{P} \cdot \mathbf{P}_s]^2$ is equal to 0 if and only if $\text{Var}(\mathbf{P}_i) = 0$ for all $i = 1, \dots, p$ and $(\mathbf{P}_{si} \mathbb{E} [\mathbf{P}_j] - \mathbf{P}_{sj} \mathbb{E} [\mathbf{P}_i])^2 = 0$ for all $1 \leq i < j \leq p$. In particular, this condition is equivalent to:

$$\mathbf{P}_{s1}\mathbb{E}[\mathbf{P}_j] = \mathbf{P}_{sj}\mathbb{E}[\mathbf{P}_1], \quad j = 2, \dots, p, \quad (\text{C10})$$

which means that \mathbf{P}_s and $\mathbb{E}[\mathbf{P}]$ are collinear.

Appendix D Hybrid target

Following (Ménétrier, 2024), we can replace the asymptotic covariance \mathbf{P} in Eq. (16) by an hybrid target $\bar{\alpha}_d\mathbf{P} + \bar{\alpha}_s\mathbf{P}_s$, where $\bar{\alpha}_d$ and $\bar{\alpha}_s$ are given by the user. Following Appendix A, we can show that the new optimal hybridization coefficients (α'_d, α'_s) are solutions to the following system of equations (D1):

$$\begin{cases} \alpha_d\mathbb{E}[\|\mathbf{P}_d\|^2] + \alpha_s\mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s] = \bar{\alpha}_d\mathbb{E}[\|\mathbf{P}\|^2] + \bar{\alpha}_s\mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s] \\ \alpha_d\mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s] + \alpha_s\|\mathbf{P}_s\|^2 = \bar{\alpha}_d\mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s] + \bar{\alpha}_s\|\mathbf{P}_s\|^2. \end{cases} \quad (\text{D1})$$

The solution of the system (D1) is given by:

$$(\alpha'_d, \alpha'_s) = \left(\frac{n'_d}{\Delta'}, \frac{n'_s}{\Delta'} \right), \quad (\text{D2})$$

where:

$$\Delta' = \|\mathbf{P}_s\|^2\mathbb{E}[\|\mathbf{P}_d\|^2] - \mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s]^2, \quad (\text{D3})$$

$$n'_d = \bar{\alpha}_d \left(\|\mathbf{P}_s\|^2\mathbb{E}[\|\mathbf{P}\|^2] - \mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s]^2 \right), \quad (\text{D4})$$

$$\begin{aligned} n'_s &= \bar{\alpha}_d \left(\mathbb{E}[\|\mathbf{P}_d\|^2] - \mathbb{E}[\|\mathbf{P}\|^2] \right) \mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s] \\ &\quad + \bar{\alpha}_s \left(\|\mathbf{P}_s\|^2\mathbb{E}[\|\mathbf{P}_d\|^2] - \mathbb{E}[\mathbf{P}_d \cdot \mathbf{P}_s]^2 \right). \end{aligned} \quad (\text{D5})$$

Therefore:

$$\begin{cases} \alpha'_d = \bar{\alpha}_d \times \alpha_d \\ \alpha'_s = \bar{\alpha}_d \times \alpha_s + \bar{\alpha}_s \end{cases} \quad (\text{D6})$$

The new hybrid covariance $\bar{\mathbf{P}}_h$ is given by:

$$\bar{\mathbf{P}}_h = \alpha'_d\mathbf{P}_d + \alpha'_s\mathbf{P}_s \quad (\text{D7})$$

$$\bar{\mathbf{P}}_h = \bar{\alpha}_d \times \alpha_d\mathbf{P}_d + \bar{\alpha}_d \times \alpha_s\mathbf{P}_s + \bar{\alpha}_s\mathbf{P}_s \quad (\text{D8})$$

$$\bar{\mathbf{P}}_h = \bar{\alpha}_d\mathbf{P}_h + \bar{\alpha}_s\mathbf{P}_s \quad (\text{D9})$$

Open Research

All the data on which this study is based were synthetically generated by (Barthélémy, 2023) and are available from: <https://zenodo.org/record/8037714> (DOI: 10.5281/zenodo.8037714)

The figures were made using:

- 685 • Matplotlib version 3.5.2 [Software], Hunter (2007), available from: [https://doi](https://doi.org/10.5281/zenodo.6982547)
686 [.org/10.5281/zenodo.6982547](https://doi.org/10.5281/zenodo.6982547)

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