

Flow structures with high Lagrangian coherence promote diatom blooms in oligotrophic waters

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Key Points:

- We provide an objective measure of the range of dynamical coherence scales in geophysical flows
- We propose the Lagrangian coherence scale as a new metric to classify biologically relevant flow structures
- Observations evidence that turbulent flow structures with a high Lagrangian coherence are conducive to diatom blooms in oligotrophic waters.

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Abstract

Diatoms are among the most efficient marine organisms for primary production and carbon sequestration, absorbing at least 10 billion tonnes of carbon dioxide every year. Yet, the spatial distributions of these planktonic organisms remain puzzling and the underlying physical processes poorly known. Here we investigate what dynamical conditions are conducive to episodic diatom blooms in oligotrophic waters based on Lagrangian diagnosis and satellite-derived phytoplankton functional types and ocean currents. The Lagrangian coherence of the flow is diagnosed in space and time simultaneously to identify which structures favor diatom growth. Observations evidence that flow structures with a high degree of coherence (40 days or longer) in high turbulent kinetic energy and vorticity sustain high concentrations of diatoms in the sunlit layers. Our findings show that the integration of Eulerian kinematic variables into a Lagrangian frame allows revealing new dynamical aspects of geophysical turbulence and unveil transport properties having large biological impacts.

Plain Language Summary

Marine diatoms is a photosynthetic plankton group that plays an important role in the generation of oxygen and the removal of carbon dioxide from the atmosphere. However, their abundance is predicted to decline as a result of the increasing ocean stratification induced by global warming. Understanding the underlying physical processes that are conducive to diatom blooms is crucial to accurately predict the evolution of primary production and carbon export. Here using an objective measure of the range of dynamical coherence scales and satellite observations, we provide evidences that flow structures with a high degree of coherence in high turbulent kinetic energy and vorticity track the favorable dynamical conditions sustaining diatoms growth in stratified and oligotrophic waters.

1 Introduction

Earth System model simulations as well as mesocosm experiments indicate that in response to increased ocean stratification and reduced vertical mixing caused by global warming, diatoms are likely to decline or be replaced by smaller phytoplankton (Barton et al., 2016; Frémont et al., 2022). However, these model predictions do not account for the role of the physical environment at scales smaller than 100 km in structuring the distribution of phytoplankton functional types. Understanding the role played by the ocean circulation at these scales in the development of diatoms in stratified and oligotrophic waters is essential to accurately predict the future evolution of primary production and carbon export (Falkowski et al., 1998; Field et al., 1998; Tréguer et al., 2018).

In contrast to other groups of smaller plankton, such as dinoflagellates, which are favored under stratified (low-turbulent), low-nutrient conditions (Glibert, 2016), diatoms are known to thrive in high nutrient-high turbulent systems (e.g. polar, upwelling and coastal areas). They are, however, not exclusive to these systems, and recent observations revealed that short-living spatially-restricted diatoms blooms may occur in oligotrophic, nutrient limited waters (Malviya et al., 2016), fueled by transient meso- and submesoscale dynamical structures such as eddies, fronts or filaments (Tréguer et al., 2018; Hernández-Carrasco et al., 2020). However it is not yet possible to link specific dynamical structures to functional phytoplankton groups. On one hand, model-based studies are strongly dependent on the modeling formulation and how the interactions between biological and physical processes are parameterized (Barton et al., 2010; D. J. McGillicuddy, 2016). Observational studies, on the other hand, remain scarce due to the lack of simultaneous measurements of velocity fields and phytoplankton taxonomy at synoptic scale (Hernández-Carrasco et al., 2020). Furthermore, most monitoring studies occur at a fixed location, linking the characteristics of the local phytoplankton communities to some prop-

70 erties of the Eulerian (“frozen” in space) flow (Cotti-Rausch et al., 2016). Contrasting
 71 with the Eulerian view, an important issue when studying bio-physical interactions in
 72 fluid environments is that any physically-driven phytoplankton community shift or change
 73 in activity (Lévy et al., 2014; Wilkins et al., 2013) would reflect the cumulative effects
 74 of coherent transport dynamics affecting nutrient supply and other growth limiting fac-
 75 tors.

76 In this study we propose a methodology to classify the flow structures according
 77 to their Lagrangian coherence scales. The method is based on the Lagrangian assessment
 78 F of Eulerian kinematic variables f , which provides information on the cumulative ef-
 79 fect of f along trajectories, as it integrates the observed variable over the time-evolution
 80 of moving water parcels arriving to a specific point. The finite-time dynamical systems
 81 theory has shown that relevant information regarding the dynamical properties of the
 82 flow can be inferred from F , especially when f is chosen to be a function related to the
 83 velocity field (Haller & Poje, 1998; Mezić et al., 2010). We assess coherent transport pro-
 84 cesses through the finite-time Lagrangian integration of vorticity and eddy kinetic en-
 85 ergy (EKE), denoted by Ω_T and K_T , respectively. The proposed methodology allows for
 86 investigating the impact of Lagrangian persistent lateral turbulent properties of the un-
 87 derlying flow structures on the phytoplankton community structure. In particular, we
 88 can accurately assess the spatial and temporal coherence scales of turbulent structures
 89 that do or do not favor a sustained growth of diatoms, providing a robust diagnostic to
 90 identify, among all the eddies, those associated with a biological response. Previously
 91 developed Lagrangian flow diagnostics, such as Lyapunov exponents based Lagrangian
 92 coherent structures (Hernández-Carrasco et al., 2011) have shown to be very useful to
 93 identify dynamical boundaries separating different functional phytoplankton groups (d’Ovidio
 94 et al., 2010). However, the proposed methodology allows further identifying not only the
 95 boundaries but also the dynamical characteristic scales of the underlying flow associated
 96 with a particular functional phytoplankton group. We apply this framework to data from
 97 two independent sources of satellite observations. Specifically, we use data of dominant
 98 phytoplankton functional type (PFT) from a regional adaptation of the PHYSAT algo-
 99 rithm (Navarro et al., 2017) and daily velocity fields at $1/8^\circ$ of spatial resolution derived
 100 from a Ssalto/Duacs multimission altimeter regional product.

101 We focus on the Mediterranean Sea, which is characterized by intense mesoscale
 102 activity (d’Ovidio et al., 2009; Morales-Márquez et al., 2022) and by high productivity
 103 only during spring time, between February and April (Basterretxea et al., 2018), with
 104 dominance of large phytoplanktonic cells such as diatoms (Navarro et al., 2017). Dur-
 105 ing the rest of the year is on average oligotrophic and highly stratified, disrupted by lo-
 106 cal - in time and space - events of elevated phytoplankton growth. The Mediterranean
 107 Sea is therefore a good proxy for ocean deserts, which cover most of the global ocean,
 108 offering an ideal laboratory to explore the biological response of an oligotrophic and highly
 109 stratified environment to intermittent, dynamical changes in the physical structure of
 110 the upper ocean.

111 2 Finite-time Lagrangian diagnostics

112 Given a Eulerian descriptor f , its Lagrangian counterpart F can be obtained by
 113 computing the path-integral of f . They can be expressed in a general way by the ma-
 114 terial integration of f along the fluid parcel trajectory, $\mathbf{R}_T(\mathbf{r}_0, t_0)$, initially located at
 115 $\mathbf{r}_0(t_0)$ over a finite time interval of integration T as, $F_T(\mathbf{r}_0, t_0) = \frac{1}{T} \int_{t_0}^{t_0+T} f(\mathbf{R}_t(\mathbf{r}, t_0), t) dt$.
 116 The ergodic theory of dynamical systems asserts that when $T \rightarrow \infty$ such Lagrangian func-
 117 tions should be constant along invariant manifolds, independently of the Eulerian func-
 118 tion chosen. This property has been used to extract the flow geometry in autonomous
 119 or periodic dynamical systems (Haller & Poje, 1998). Nevertheless the ocean flow is dy-
 120 namically an aperiodic and finite system; consequently such functions can only be path-
 121 integrated (or averaged in a Lagrangian sense) over a finite time period ($T \ll \infty$). Even

122 so, relevant information regarding the dynamical properties of the flow can still be in-
 123 ferred from F , especially when f is chosen to be a function related to the velocity field,
 124 \mathbf{v} (Haller & Poje, 1998; Mezić et al., 2010). Here we quantify transport coherent pro-
 125 cesses extending the choice of these functions to relative vorticity (ω) and eddy kinetic
 126 energy (EKE), since EKE informs on the turbulent component of the flow associated with
 127 eddy activity, while ω takes into account the shear and the rotation of the flow.

128 The Lagrangian description of the vorticity was originally reported in Casey, 1991
 129 (Casey & Naghdi, 1991) and recently used to develop an objective identification of co-
 130 herent vortex boundaries based on well-defined contours of Lagrangian-averaged vortic-
 131 ity deviation (LAVD) (Haller et al., 2016). Considering the motion of a fluid parcel on
 132 the time interval $[t_0, t_0+T]$ in an incompressible two dimensional flow, we define the finite-
 133 time Lagrangian vorticity (FTLV), denoted by Ω_T , at the position $\mathbf{r}_0 = (x_0, y_0)$ at time
 134 t_0 as,

$$\Omega_T(\mathbf{r}_0, t_0) = \frac{1}{T} \int_{t_0}^{t_0+T} \omega(\mathbf{R}_t(\mathbf{r}, t_0), t) dt, \quad (1)$$

135 where ω is the Eulerian relative vorticity: $\omega(\mathbf{r}, t) = \frac{\partial v(\mathbf{r}, t)}{\partial x} - \frac{\partial u(\mathbf{r}, t)}{\partial y}$.

136 Furthermore we introduce a new metric to analyze the evolution of the EKE along
 137 fluid parcel trajectories. Specifically, we compute the finite-time Lagrangian EKE (FTLK),
 138 denoted by K_T , as follows,

$$K_T(\mathbf{r}_0, t_0) = \frac{1}{T} \int_{t_0}^{t_0+T} \kappa(\mathbf{R}_t(\mathbf{r}, t_0), t) dt, \quad (2)$$

139 where κ is the Eulerian EKE given by, $\kappa(\mathbf{r}, t) = (u(\mathbf{r}, t) - \langle u(\mathbf{r}) \rangle_{\Delta T})^2 + (v(\mathbf{r}, t) - \langle v(\mathbf{r}) \rangle_{\Delta T})^2$,
 140 obtained through the instantaneous anomaly of the zonal and meridional components
 141 of the velocity field with respect to the steady mean of the total kinetic energy, which
 142 is achieved averaging over a time period of $\Delta T=10$ years. This diagnostic allows clas-
 143 sifying the flow, in the domain under consideration, in subregions according to the Lagrangian-
 144 persistent level of the turbulent mesoscale activity. High values of K_T identify therefore
 145 regions where the flow is concentrating turbulent energy. The selection of the time in-
 146 tegration, T , depends on the typical time scales of the process to be studied.

147 3 Satellite data

148 The identification from satellite observations of the dominant phytoplankton func-
 149 tional type (PFT), is performed using the PHYSAT algorithm (Alvain et al., 2005). Here
 150 we use a regional adaptation of this algorithm to the specific bio-optical characteristics
 151 of the Mediterranean Sea, PHYSAT-Med (Navarro et al., 2014, 2017). The PHYSAT-
 152 Med identifies nanoeukaryotes, coccolithophorids, *Prochlorococcus*, *Synechococcus*-like
 153 cyanobacteria, diatoms and Phaeocystis-like phytoplankton through the analysis of the
 154 specific signatures in the normalized water leaving radiance (nLw) spectra measured by
 155 ocean color sensors. For this study, we use the PHYSAT-Med version published in 2017
 156 (Navarro et al., 2017), in which the authors used the OC-CCI v3.0 dataset generated from
 157 merged normalized remote-sensing reflectance derived from four satellite sensors (Storm
 158 et al., 2013): SeaWiFS, MODIS, MERIS, and VIIRS. PHYSAT-Med retrieves the domi-
 159 nant group for a given satellite image pixel (4 km for the Mediterranean Sea) by iden-
 160 tifying which phytoplankton group is the major contributor to the radiance anomaly. In
 161 other words, it associates a given phytoplankton type with its main bio-optical signa-
 162 ture, while taking into account biases induced by other optically-active components, such
 163 as dust deposition events and rivers discharge. From this high-resolution database (about
 164 6,600 daily images), 10-day and monthly maps of dominant phytoplankton groups were
 165 obtained by calculating the phytoplankton group that was dominant during the integra-
 166 tion period (10-day or monthly, respectively) at each geographical pixel, not including
 167 “unidentified” pixels. The regional PHYSAT-Med product was validated by comparing

168 the outputs of the algorithm with more than 5,000 in situ measurements analyzed by high-
 169 performance liquid chromatography (HPLC) (Navarro et al., 2014, 2017).

170 The current flow field is derived from daily absolute geostrophic surface velocities
 171 at $1/8^\circ$ of spatial resolution derived from a Ssalto/Duacs multimission altimeter regional
 172 product released in 2016 specific for the Mediterranean Sea. This product is currently
 173 distributed by the Copernicus Marine Environment Monitoring Service (CMEMS).

174 4 Results

175 Dynamical co-coherence of oceanic flow features

176 We first discuss how the flow characterization inferred from a given kinematic quan-
 177 tity differs when evaluated through either Eulerian or Lagrangian functions. The latter
 178 are not directly related to instantaneous measurements of such functions, unless some
 179 kind of dynamic equilibrium or ergodicity-type property is established. In this case, the
 180 time-integrated effect can be effectively related to the instantaneous or averaged spatial
 181 patterns (for instance, if the spatial arrangement of eddies at a given time provides in-
 182 formation about the typical time evolution of a water parcel). Using daily gridded geostrophic
 183 velocities derived from altimetry (see section 3) we compare daily snapshots of the Eu-
 184 lerian vorticity and EKE, denoted by Ω_0 and K_0 , respectively, (Fig. 1A and 1C), and
 185 their Lagrangian counterpart evaluated for $T=40$ days, Ω_{40} and K_{40} (Fig. 1B and 1D)
 186 (see Section 2 for a detailed description of these metrics). Spatial patterns of Ω_T and
 187 K_T differ significantly from their Eulerian equivalence whenever the integration time is
 188 greater than 15 days (see Fig. S1 in the Supporting Information). Filamentary struc-
 189 tures of intense accumulated vorticity and EKE around and within the eddies that could
 190 originate from eddy-flow or eddy-eddy interactions are unveiled by the Lagrangian di-
 191 agnostic but not by the Eulerian one. Furthermore, some structures with high values of
 192 vorticity and EKE identified by Ω_0 and K_0 vanish in Ω_{40} and K_{40} , suggesting that these
 193 structures are not dynamically persistent, while only the highly coherent flow structures
 194 remain as T increases (Fig. S2 in the Supporting Information). The above analysis sug-
 195 gests that the Lagrangian characterization of the marine turbulent flow unveils super-
 196 coherent structures, i.e singular structures of high dynamical coherence, that can not be
 197 identified in the classical Eulerian framework.

198 To assess how the Ω_T and K_T scale with the duration of integration, we analyze
 199 their statistical behavior comparing the probability density functions (PDF) of Ω_T and
 200 K_T for different T values (Fig. 1E and 1F). As T increases the peak of the PDFs in-
 201 creases while their width narrows, and the PDF becomes asymmetrical, with a more pro-
 202 nounced tail for negative vorticity values and high EKE values. This PDF behaviour -
 203 a time-dependent histogram whose peak becomes sharper and higher as the time inte-
 204 gration lengthens - is characteristic of a multifractal scaling (Falconer, 1990; Hernández-
 205 Carrasco et al., 2011), often used to interpret spatial scales but extended here to tem-
 206 poral scales.

207 Given that in the ocean nearly 80% of the total kinetic energy is captured in vor-
 208 tical coherent structures (Klein et al., 2019), a relationship between Eulerian EKE and
 209 vorticity may be expected. However when comparing the spatial patterns of both Eu-
 210 lerian quantities (Fig. 1A, C), we do not find a clear correspondence. Conversely, the
 211 spatial distribution is very similar when comparing their Lagrangian counterparts (Fig.
 212 1B, D). To further study the relationship between Ω_T and K_T we examine the histograms
 213 of Ω_T conditioned by the value of K_T for different values of T (Fig. 2). It allows explor-
 214 ing which aspects of the flow dynamical coherence is well-captured by the Lagrangian
 215 assessment but missed by the Eulerian approach. We observe a dispersed cloud of points
 216 for $T=0$ days (Fig. 2A), while a clear modal line (e.g. line of maximum conditioned prob-
 217 abilities) formed by two different, nearly straight, segments is evident for $T=40$ days (Fig.

218 2B). These segments are associated with a slowly increasing and faster decreasing de-
 219 pendency of Ω_T with K_T , respectively. In order to verify if a functional dependence be-
 220 tween vorticity and EKE could hold for our dataset, we compute various conditioned av-
 221 erages $\hat{\Omega}_T(K_T)$ for different integration times T (Fig. 2C). While the plot shows con-
 222 ditioned averages, the conditioned standard deviation is not negligible as reflected by the
 223 correlation coefficients R^2 and associated slopes reported in the caption of Fig. 2. While
 224 there is a very weak relationship among Eulerian quantities (black curve), we observe
 225 more robust relationships as T increases (green curves). The resulting conditional mean
 226 is a constant value for K_T smaller than $0.13 \text{ m}^2/\text{s}^2$ suggesting that Ω_T and K_T are sta-
 227 tistically independent, and there is not a functional relationship between the two vari-
 228 ables at low values of EKE. A functional relationship, on the other hand, exists for val-
 229 ues of K_T greater than $0.13 \text{ m}^2/\text{s}^2$. Overall, Fig. 2C suggests that we can approximate
 230 Ω_T in terms of $(K_T)^{1/2}$ allowing to infer a spatial scale, L , from the slopes A after fit-
 231 ting: $\Omega_T = A \cdot (K_T)^{1/2}$ ($L = 1/A$) for high values of $(K_T)^{1/2}$. L has units of the in-
 232 verse of length (1/m) and provides an estimation of the spatial scale associated with the
 233 coherence of these Lagrangian structures and their evolution as a function of the time
 234 integration. It could be interpreted as the typical length-scales of the vortical structures
 235 that contribute most, and for longer times, to the turbulent energy. Values of L decreases
 236 as T increases, with values ranging between [60 - 15 km] for T = [0 - 80 days], indicat-
 237 ing that when T increases smaller structures are captured. The minimum L value in the
 238 flow considered converges to 15 km, which represents the smallest length-scale of the La-
 239 grangian dynamical structures that can be inferred from this velocity field.

240 The high correlation between Ω_T and K_T for T > 40 days, indicates that spatial
 241 variations of Lagrangian coherent vorticity are closely related to spatial variations of La-
 242 grangian coherent mesoscale activity. This suggests that the persistent anticyclonic ed-
 243 dies in the portion of Mediterranean Sea analyzed emanate fluid parcels associated with
 244 accumulated mesoscale turbulent energy. Consequently, Ω_T for T > 40 days is a more
 245 suitable Lagrangian diagnostic to link turbulence properties to phytoplankton dynam-
 246 ics than its Eulerian counterpart (T = 0 days), whenever the analyzed flow is derived
 247 from altimetry.

248 **Influence of oceanic coherence on the phytoplankton community com- 249 position**

250 Here we dive on the response of diatoms to the ocean physical conditions focus-
 251 ing on the Lagrangian coherence scales of the marine flow structures unveiled by these
 252 novel metrics. Figure 3A shows a patchy spatial distribution of chlorophyll-a concentra-
 253 tion, a proxy of total phytoplankton, with local maxima dominated by diatoms (Fig. 3B),
 254 located at $38.5^\circ\text{N}/5.5^\circ\text{E}$. In order to compare with PHYSAT-Med images, which corre-
 255 spond to a 10 days averaged distribution of phytoplankton types, we show the spatial
 256 patterns of time averaged snapshots of Eulerian vorticity, Ω_0 , and EKE, K_0 , over the same
 257 10 days in Fig. 3C and Fig. 3E, revealing different structures of high vorticity and EKE
 258 distributed throughout the entire area. This map is significantly different from the time
 259 average over the same 10 days of daily snapshots of Lagrangian vorticity (Ω_{40} , Fig. 3D)
 260 and Lagrangian EKE (K_{40} , Fig. 3F) for T=40 days. Only one intense eddy-like struc-
 261 ture stands out among all other structures inferred from Ω_0 and K_0 . It indicates again
 262 that while Eulerian flow structures assess coherence in the spatial domain, Lagrangian
 263 structures can assess coherence in space and time simultaneously. As such, most-eddy
 264 like structures identified by Ω_0 and K_0 and vanishing in the Ω_T and K_T fields for T >
 265 40 days exhibit spatial coherence only; conversely, the central eddy is unique exhibiting
 266 high coherence over both space and time. By combining both Lagrangian metrics we can
 267 identify singular supercoherent vortical objects associated with high mesoscale turbu-
 268 lent energy which should impact plankton communities.

269 Comparing the maps of Ω_0 and K_0 (Fig. 3C and E, respectively) with the spatial
 270 distribution of the dominant phytoplankton groups (Fig. 3B) we observe that the con-
 271 fined and circular patch dominated by diatoms is associated with local maxima of EKE
 272 and extrema (both positive and negative) of vorticity. However, the reciprocal statement
 273 is not true: not all local maxima of both Eulerian quantities are characterized by diatoms.
 274 Looking at the Ω_T and K_T diagnostics (Fig. 3D and F), we found the singular struc-
 275 ture associated with high Lagrangian turbulent kinetic energy and vorticity values uniquely
 276 identifies the diatoms bloom. In the Eulerian framework, many structures of high EKE
 277 and vorticity are present in the domain without a clear correspondence with any dom-
 278 inant phytoplankton group. Conversely, our novel Lagrangian diagnostics distinguish sin-
 279 gular coherent structures that support specific biological response from the other tran-
 280 sient features that have no clear biological signal. Further examples of the correspon-
 281 dence between these coherent dynamical objects and diatom blooms are shown in Sup-
 282 porting Information (Fig. S3 to S11).

283 The co-occurrence between diatoms blooms and these singular dynamical struc-
 284 tures suggests a relationship between the abundance of diatoms and the degree of co-
 285 herence of the oceanic features. To further explore this relationship we compute the prob-
 286 ability density functions of the diatom abundances, henceforth referred to as $P(\text{diatoms})$,
 287 conditioned by different discretized values of Ω_T and K_T (50 bins are taken) and for dif-
 288 ferent values of T averaged over 15 diatom blooms x 30 snapshots/event distributed in
 289 the stratification period (May-September) among the analyzed 10 years of PHYSAT-Med
 290 data. Area-normalized PDFs (Fig. 4A and B) have a characteristic signature that be-
 291 comes flatter at lower values of Ω_T and K_T as the time integration lengthens. When T
 292 = 0 higher values of $P(\text{diatoms})$ are found in regions characterized by low values of vor-
 293 ticity and turbulence, and high $P(\text{diatoms})$ values shift toward higher values of vortic-
 294 ity and turbulent kinetic energy as T increases. The drifting fluid parcels that largely
 295 contribute to promoting diatom growth are those that preserve high negative vorticity
 296 and turbulent kinetic energy while they move with the flow for a sufficiently long amount
 297 of time. The shape of the histograms persists nearly unaltered from T = 40 to T = 80
 298 days. While the existence of this 'collapsed' shape of the histogram confirms the need
 299 to use a Lagrangian approach to identify the flow properties that matters for biological
 300 variables, it further suggests that 40-80 days is the characteristic time-scale of bio-physical
 301 interactions most relevant to predict/explain localized diatom blooms.

302 In order to estimate the optimal time scale of the flow coherence that maximizes
 303 bio-physical interactions, we compute the area-normalized number of pixels where di-
 304 atoms dominate (i.e. the number of pixels identified as diatoms normalized by the to-
 305 tal area of high Ω_T and high K_T , that is a proxy of diatoms relative abundance) over
 306 regions characterized by high values of Ω_T and K_T fields (defined as $\Omega_T/f > 0.1$ or <
 307 -0.1 for negative vorticity and $K_T > 0.1 \text{ m}^2/\text{s}^2$) as a function of T. The abundance of
 308 diatoms increases in regions of high negative Lagrangian vorticity (Fig. 4C) while de-
 309 creases in regions of positive Lagrangian vorticity as T increases (Fig. 4E). This asym-
 310 metric behavior is opposite to that recovered through Eulerian analyses (Ω_0) in which
 311 diatoms appear associated predominantly with coherent (in space only) positive vortic-
 312 ity. This counter-intuitive behavior is explained by the fact that fluid parcels with pos-
 313 itive/negative vorticity at time t_0 may have had opposite sign vorticity when the whole
 314 history of the parcel is considered (i.e. integrating the relative vorticity backward in time
 315 from t_0 to $t_0 - T > 0$ days). For example fluid parcels coming from an anticyclonic/cyclonic
 316 eddy can be absorbed by an cyclonic/anticyclonic eddy (Hernández-Carrasco et al., 2020).
 317 For both high negative and positive Ω_T we observe a plateau in $P(\text{diatoms})$ for $T > 70$
 318 days. A similar behavior is identified when looking at $P(\text{diatoms})$ in regions of high La-
 319 grangian turbulent kinetic energy ($K_T > 0.1 \text{ m}^2/\text{s}^2$) for different T. The probability to
 320 find high abundance of diatoms increases with T and reaches its maximum value for T
 321 = 30 days (Fig. 4D). By simultaneously conditioning the abundance of diatoms in re-
 322 gions of high Ω_T (either positive or negative) together with K_T (Fig. 4F), we find that
 323 $P(\text{diatoms})$ associated with these turbulent flow conditions initially increases over in-

324 tegration time and reaches a plateau for $T > 40$ days. From this T threshold, one can
 325 infer the optimal coherence time of vortical structures for diatoms blooms under strat-
 326 ified, oligotrophic conditions.

327 5 Discussion and conclusions

328 To understand the effects of global warming on global ocean physical and biogeo-
 329 chemical properties, it is essential to identify the scales that are relevant for dynamical
 330 and biological interactions, and then predict their evolution. The coherence analysis of
 331 the surface ocean dynamics through the Lagrangian descriptors, K_T and Ω_T provides
 332 for the first time a theoretically-grounded and objective measure of the range of dynam-
 333 ical coherence time and spatial scales associated with mesoscale ocean features. We show
 334 that the relationship between EKE (K_T) and vorticity (Ω_T) emerges for topological struc-
 335 tures characterized with high dynamical coherence. As T increases the Lagrangian vort-
 336 icity and the EKE field emerge organized in filaments of high values. In the case of the
 337 altimetry velocity field analyzed here, we find a dynamical resolution barrier with La-
 338 grangian flow structures size converging and collapsing at spatial scale $L=[15-20$ km] for
 339 time scales of coherence $T > 40$ days. At these scales the filamentary dynamics associ-
 340 ated with the fine-mesoscale turbulence becomes relevant. This is likely due to the vor-
 341 tex filaments induced by the (mesoscale) eddy-eddy interactions and their contribution
 342 to the energy cascade from large to small scales (in nature down to the submesoscales,
 343 which are not well captured by current altimetry, but will be partially resolved by the
 344 future SWOT mission (Morrow et al., 2019)). This Lagrangian similarity between dif-
 345 ferent kinematic flow properties can be used to assess the finest effective scales of the dy-
 346 namics captured in a given velocity field, as well as to evaluate the capability of mod-
 347 els or observations to resolve different processes.

348 The Lagrangian description of vorticity and turbulent energy through K_T and Ω_T
 349 does offer some advantages as shown by comparing it with the conventional (Eulerian)
 350 description. The Lagrangian objects uniquely identified as high-coherent turbulent ed-
 351 dies track the favorable dynamical conditions, namely a turbulent environment that is
 352 sufficiently coherent in space and time, favouring diatoms growth. This allows to differ-
 353 entiate the mesoscale structures that effectively promote diatoms blooms from those that
 354 only stir existing communities (d’Ovidio et al., 2010). Analyzing the accumulated effect
 355 of the flow dynamical properties along fluid parcel trajectories hosting phytoplankton
 356 cells allows for estimating the typical time and length scales at which the interactions
 357 between the dynamical flow and the fast-growing opportunist planktonic groups (diatom-
 358 like in our case) are most significant. We observe that $T = 40$ days is a critical coher-
 359 ence time scale at which the mesoscale turbulent conditions (i.e. vorticity and EKE) in-
 360 fluence phytoplankton community composition by promoting diatoms, through persist-
 361 ent cell re-suspension and nutrient vertical fluxes favoring their encounter in the euphotic
 362 zone. This value is of the same order of magnitude but slightly smaller than the 60 days
 363 of mean life associated with non-linear eddies that display the maximum correlation with
 364 high concentrations of surface Chl a reported in previous studies (Chelton et al., 2011).

365 The limited resolution of the altimetry currents implies that we cannot quantify
 366 the biological response of rapidly fluctuating small scale structures but only the slow dy-
 367 namics, i.e. temporal scales longer than 1 week. It should be noted that using a coarse
 368 velocity field (coarser than in the real ocean) may introduce spurious Lagrangian dynam-
 369 ics at small scales (Haza et al., 2016), but does not impact the dynamics at scales larger
 370 than resolution, that remains robust (Beron-Vera et al., 2019; Hernández-Carrasco et al.,
 371 2011). This ensures that the resolution in the altimetry dynamical fields is not affect-
 372 ing the evaluation of the temporal scales at or above 40 days at which diatoms ‘feel’ the
 373 flow coherence. This is evident in Fig. 4C, where the concentration of diatoms in struc-
 374 tures of high vorticity increases as a function of Lagrangian integration up to $T=70$
 375 days, which is longer than the critical scale of 40 days identified by the Lagrangian ef-

376 fective resolution analysis. Our findings, based on satellite observations, provide a ro-
 377 bust target against which models can be now tested to verify their ability to properly
 378 represent physical-biogeochemical interactions. Mesoscale resolving regional simulations
 379 should indeed capture the diatom preferences, at the relevant spatial and temporal co-
 380 herence scales, that we identified. In interpreting our results, we excluded the possibil-
 381 ity that the fluid parcels with high diatom concentrations are composed of nutrient rich
 382 coastal waters by analyzing the origin maps shown in the Supporting Information (Fig.
 383 S12-S14). This suggests that diatom growth occurs whenever diatoms in resting stages,
 384 likely present at low concentrations in most of the surface ocean (Ryneerson et al., 2013)
 385 (even oligotrophic), are entrained in coherent physical structures and can exploit ver-
 386 tical mixing events (e.g. ephemeral hot spots of intense vertical fluxes (D. McGillicuddy
 387 et al., 1998; Falkowski et al., 1998; Klein & Lapeyre, 2009; Lévy et al., 2018)) induced
 388 by specific mesoscale eddies that persist for enough time (more than one month). Indeed,
 389 our analysis allows for identifying these short-living and moving diatoms habitats de-
 390 spite high-stratification conditions typical of "oceanic deserts", as the major oligotrophic
 391 gyres of the world ocean.

392 6 Open Research

393 Absolute geostrophic velocity data are available on the CMEMS web platform [https://](https://data.marine.copernicus.eu/products)
 394 data.marine.copernicus.eu/products and at DOI:10.48670/moi-00141. PHYSAT-
 395 Med data are referenced at Navarro et al., 2017 and are available at [https://nimbus](https://nimbus.imedea.uib-csic.es/index.php/s/6w6YzWBpTmWD8BW)
 396 [.imedea.uib-csic.es/index.php/s/6w6YzWBpTmWD8BW](https://nimbus.imedea.uib-csic.es/index.php/s/6w6YzWBpTmWD8BW). The code to compute the par-
 397 ticle trajectories and the Finite-Time Lagrangian diagnostics are available at [https://](https://doi.org/10.5281/zenodo.7705122)
 398 doi.org/10.5281/zenodo.7705122 with Creative Commons Attribution 4.0 Interna-
 399 tional licence. Figures are plotted using Matlab v2021 software.

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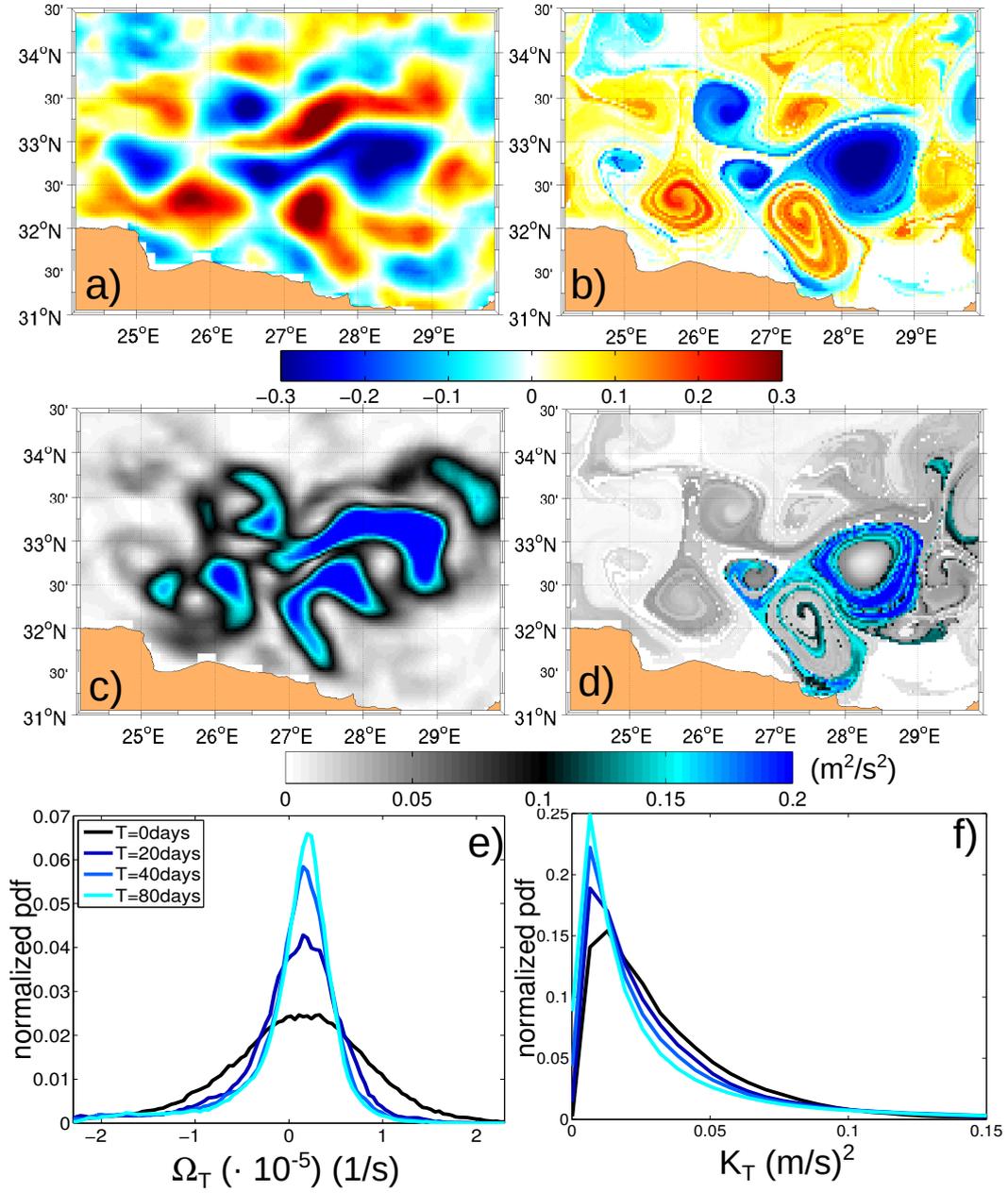


Figure 1. (A) and (B) show snapshots corresponding to March 22, 2016 of Ω_0 and Ω_{40} divided by Coriolis parameter. (C) and (D) show K_0 and K_{40} for the same date. Panels (E) and (F) show the PDF for Ω_T and K_T , respectively, using values from weekly snapshots over 2006-2016 for different integration times, T.

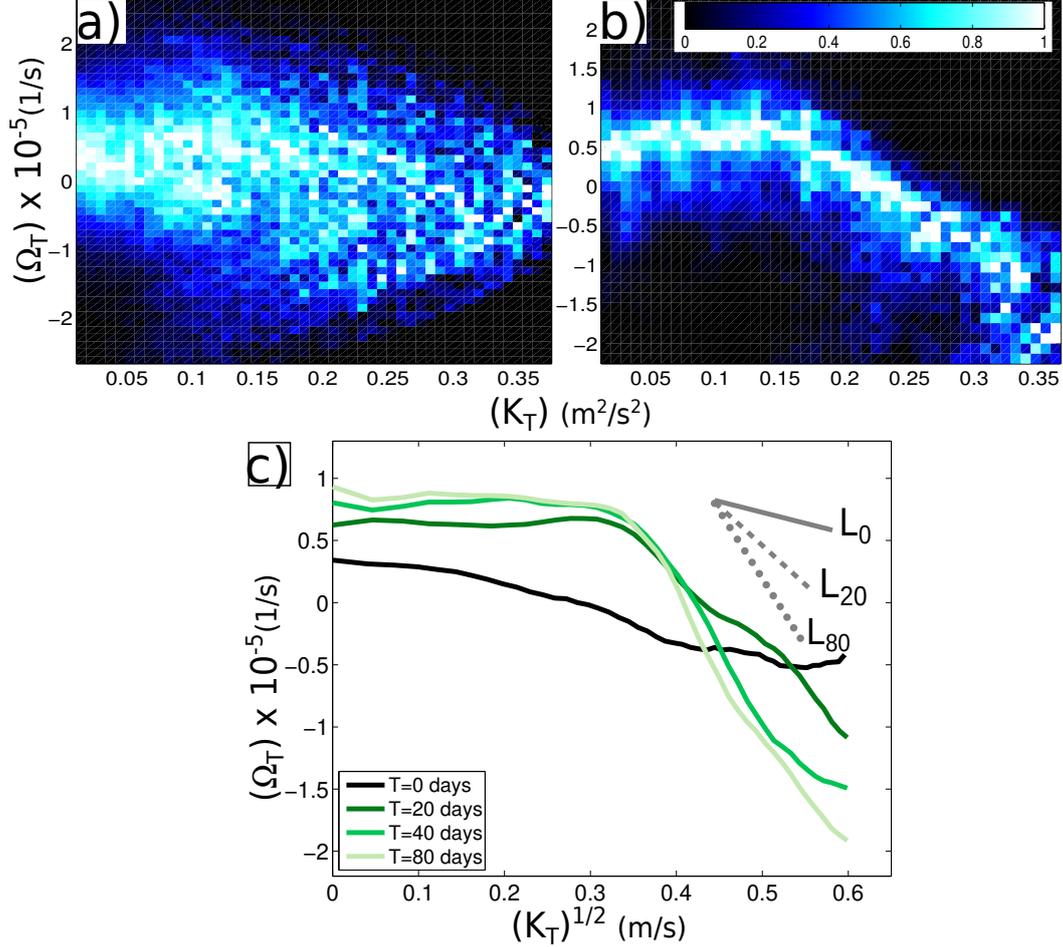


Figure 2. Probability distributions of Ω_T conditioned by K_T for two different time periods of integration, $T = 0$ days (panel A) and $T = 40$ days (panel B). The brightest color (light blue) corresponds to the maximum probability at each column; the darkest color (pure black) corresponds to zero. (C) shows functional dependence between Ω_T vs $K_T^{1/2}$ for different integration times T . Correlation coefficient are computed from co-variance dependence, obtaining $R^2=0.19$ for $T=0$ days; $R^2=0.52$ for $T=20$; $R^2=0.69$ for $T=40$ and $R^2=0.76$ for $T=80$ days. Slopes resulting from the linear regression fit analysis of the functional dependencies between Ω_T vs $K_T^{1/2}$ for $T=0, 20$ and 80 days are shown as grey lines in C. Their corresponding length scales are $L_0=56$ km $L_{20}=29$ km and $L_{80}=15$ km.

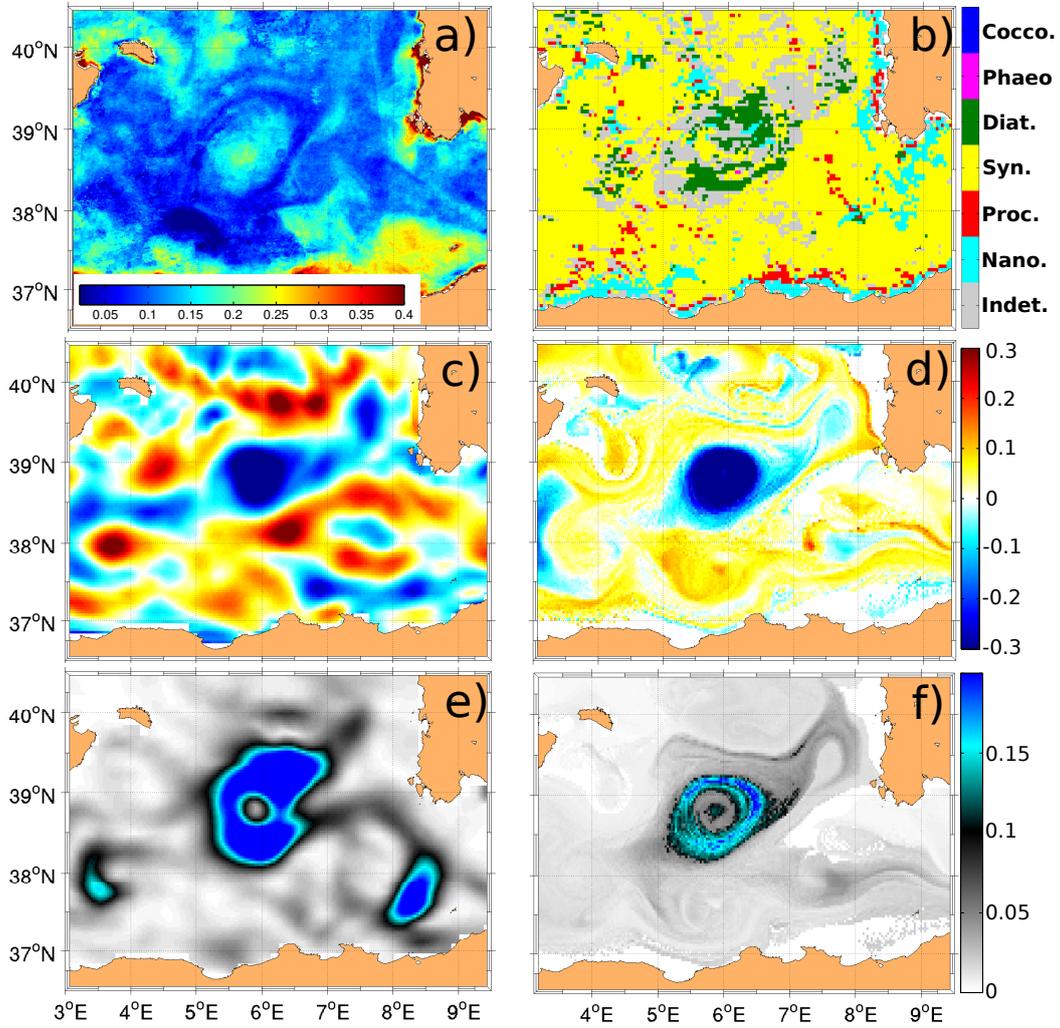


Figure 3. (A) Chlorophyll (a) (in mg/m^3) derived from Ocean color satellite data; (B) map of dominant phytoplankton functional types retrieved from the PHYSAT-Med product corresponding to the period: May 28, 2009 - June 07, 2009; (C) and (E) are the Eulerian vorticity, Ω_0 , divided by Coriolis, and Eulerian EKE, K_0 (in m^2/s^2), both averaged over 10 daily snapshots (May 28, 2009 - June 07, 2009); (D) and (F)) are the Lagrangian vorticity, Ω_{40} , and Lagrangian EKE, K_{40} (in m^2/s^2), for $T=40$ days, and averaged over the same period.

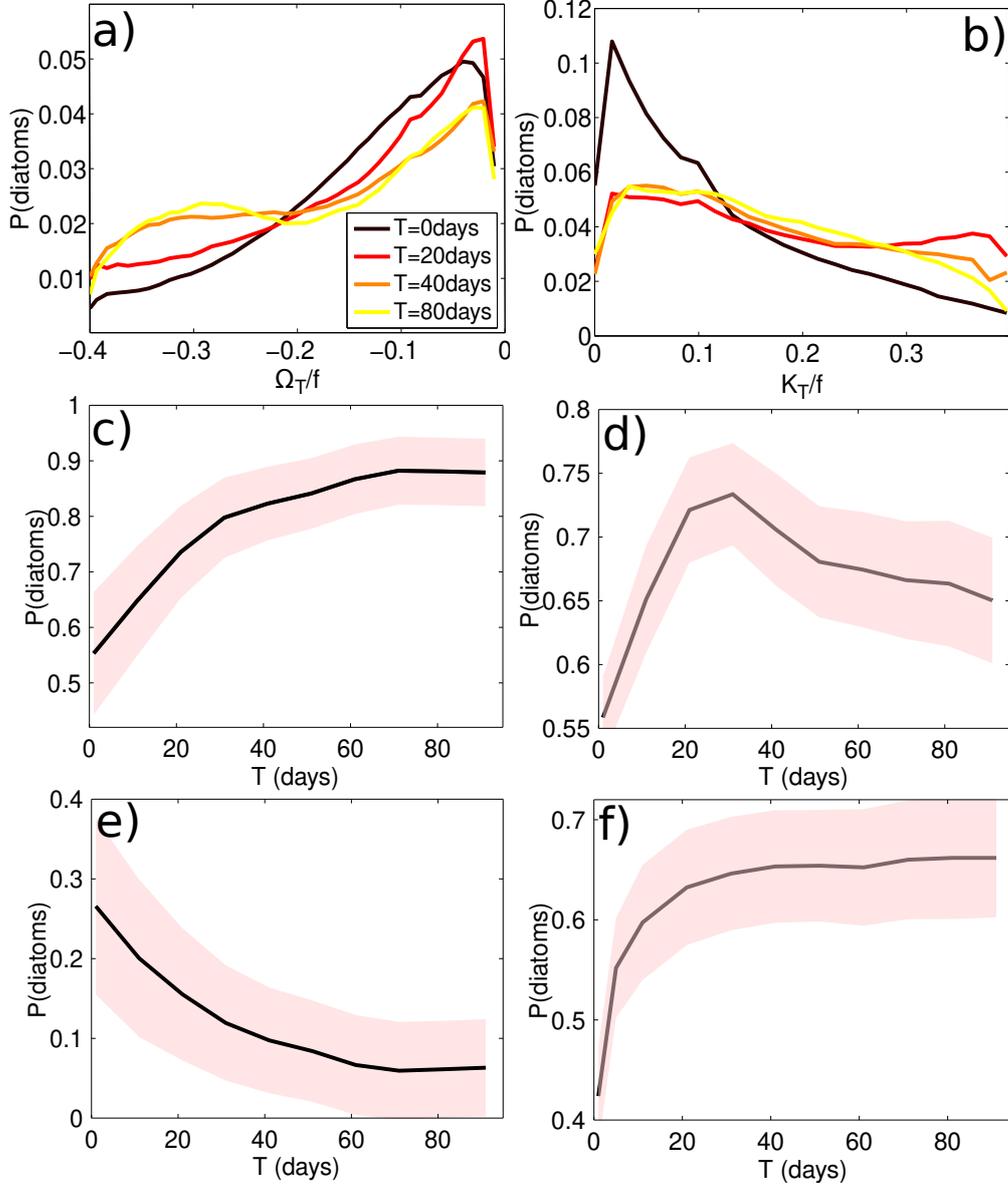


Figure 4. Probability density function of the normalized number of diatoms, or diatoms abundance, (to unity area) versus values of Ω_T (A) and K_T (B). Normalized diatom abundance as a function of the time integration T found in regions of high negative (C), positive Ω_T (E) and high K_T (D); and in regions high negative Ω_T and high K_T simultaneously (F).