Modeling the Distribution of Iron-oxides in Basalt by combining FIB-SEM and MicroCT Measurements

Frenk Out¹, Rosa A de Boer¹, John Walmsley², and Lennart Vincent de Groot¹

¹Utrecht University ²University of Cambridge

September 11, 2023

Abstract

Micromagnetic tomography (MMT) aims to go beyond paleomagnetic measurements on bulk samples by obtaining magnetic moments for individual iron-oxide grains present in a sample. To obtain accurate MMT results all magnetic sources and all their magnetic signals should be known. Small particles ($<<1 \mu$ m) are often not detected by MicroCT analyses, but do have a magnetic signal, and therefore hamper obtaining reliable MMT results. Currently it is unknown how many of these small 'ghost grains' are present in basaltic samples. Here we aim to obtain a realistic grain-size distribution for iron-oxides in a typical Hawaiian basalt. We characterize the entire grain-size range of interest to paleomagnetism, from the superparamagnetic threshold of 40 nm to multidomain grains with sizes up to 10 μ m. This requires a combination of FIB-SEM slice-and-view and MicroCT techniques: FIB-SEM characterizes the grains between 20 nm and 1 μ m and MicroCT detects iron-oxides >750 nm. The FIB-SEM and MicroCT data are combined through normalizing the grain-size distribution is acquired for the entire grain-size range. Our dataset enables future studies to populate (MMT) models with a realistic distribution of even the smallest iron-oxide grains, which ultimately may shed light on the confounding influence of such ghost grains on MMT results.

Modeling the Distribution of Iron-oxides in Basalt by combining FIB-SEM and MicroCT Measurements 2

Frenk Out¹, Rosa A. de Boer¹, John Walmsley², Lennart V. de Groot¹

4	$^1\mathrm{Paleomagnetic}$ laboratory Fort Hoofddijk, Department of Earth Sciences, Utrecht University,
5	Budapestlaan 17, 3584 CD Utrecht, The Netherlands.
6	$^{2} \mathrm{Department} \text{ of Materials Science \& Metallurgy, University of Cambridge, 27 Charles Babbage Road, CB3}$
7	0FS Cambridge, United Kingdom.

Key Points: 8

1

3

9	• We produced a grain-size distribution of iron-oxides in a typical Hawaiian basalt
10	from the superparamagnetic threshold (~40 nm) to 10 micron
11	• We combined FIB-SEM and MicroCT data from sister specimens by normalizing
12	them to the mineral surface area of the non-magnetic minerals
13	• Our grain-size distribution can be used to populate realistic models of iron-oxides
14	in a Hawaiian basalt

Corresponding author: Frenk Out, f.out@uu.nl

15 Abstract

Micromagnetic tomography (MMT) aims to go beyond paleomagnetic measurements on 16 bulk samples by obtaining magnetic moments for individual iron-oxide grains present in 17 a sample. To obtain accurate MMT results all magnetic sources and all their magnetic 18 signals should be known. Small particles ($<<1 \mu m$) are often not detected by MicroCT 19 analyses, but do have a magnetic signal, and therefore hamper obtaining reliable MMT 20 results. Currently it is unknown how many of these small 'ghost grains' are present in 21 basaltic samples. Here we aim to obtain a realistic grain-size distribution for iron-oxides 22 in a typical Hawaiian basalt. We characterize the entire grain-size range of interest to 23 paleomagnetism, from the superparamagnetic threshold of ~ 40 nm to multidomain grains 24 with sizes up to 10 μ m. This requires a combination of FIB-SEM slice-and-view and Mi-25 croCT techniques: FIB-SEM characterizes the grains between 20 nm and 1 μ m and Mi-26 croCT detects iron-oxides >750 nm. The FIB-SEM and MicroCT data are combined through 27 normalizing the grain-size distribution using the surface area of non-magnetic minerals 28 that are characterised in both datasets. Then, a lognormal-like grain-size distribution 29 is acquired for the entire grain-size range. Our dataset enables future studies to popu-30 late (MMT) models with a realistic distribution of even the smallest iron-oxide grains, 31 which ultimately may shed light on the confounding influence of such ghost grains on 32 MMT results. 33

34

Plain Language Summary

Micromagnetic tomography (MMT) is a promising new technique that extracts high-35 quality magnetic information from lavas. Magnetic particles in lavas obtain a magnetic 36 signal while cooling in presence of Earth's magnetic field. However, not all particles store 37 the signal well, meaning that both good and bad recorders are present. Classical pale-38 omagnetic techniques measure the magnetic signal of all recorders together, i.e. the bulk 39 signal. MMT, however, acquires the signal from individual recorders in the lava, enabling 40 selecting only the good recorders and rejecting the signal of bad recorders. MMT needs 41 two pieces of information: (1) the surface magnetic field produced by recorders and (2) 42 the location of all recorders in the lava. Unfortunately, the position of recorders smaller 43 than $\sim 1 \ \mu m$ are often unknown due to measurement limitations. If these small parti-44 cles are not detected, they may scramble the calculated signal of larger recorders. To in-45 vestigate how many disruptive undetected grains are present, we scanned a small vol-46

-2-

- 47 ume of lava on high resolution to extract all magnetic recorders and obtain a grain-size
- 48 distribution. With this distribution we may estimate if these small recorders distort the
- ⁴⁹ signal of larger detected grains.

50 1 Introduction

Volcanic rocks are often perceived to be excellent recorders of paleomagnetic fields, although only a small portion of all minerals in a volcanic rock has magnetic properties. These minerals, iron-oxide grains, present in a lava acquire a magnetization by cooling in the Earth's magnetic field. Their magnetization is thus a thermo-remanent magnetization (TRM), and these magnetic signals are regarded to be stable, often over millions or billions of years (e.g. Dunlop & Özdemir, 1997; Tauxe & Yamazaki, 2015).

Small particles, or single domain (SD) grains, have two configurations for storing 57 their magnetic signal: either parallel or anti-parallel to the grain's easy axis. One of these 58 two options is selected at random but with a slight preference to align with the prevail-59 ing magnetic field (e.g. Tauxe, 2010; Berndt et al., 2016). In slightly larger particles, which 60 are often referred to as pseudo-single domain (PSD) grains, the magnetic signal is stored 61 in a more complex vortex structure. Fortunately, this signal usually represents the orig-62 inal imparting magnetic field accurately (Nagy et al., 2017). The last group of particles 63 are iron-oxides with diameters $>1 \ \mu m$ and belong to the multi-domain (MD) class. Al-64 though these MD grains may store magnetic signals for millions of years, their expres-65 sion of the past field is often disturbed, overprinted, unstable, or lost (e.g. De Groot et 66 al., 2014). This instability is caused by magnetic reassembling of the original magnetic 67 structure of a grain over time to achieve energy minimization (Néel, 1955). MD grains 68 are therefore mostly perceived as bad paleomagnetic recorders, while SD and PSD grains 69 are considered to be more reliable recorders of paleomagnetic fields. To make it even more 70 difficult, most natural rock samples consist of a mixture of SD, PSD and MD grains. The 71 combination of these good and bad recorders in rocks is the main reason that generally 72 only 10 to 20% of all paleointensity experiments pass sufficient selection criteria to yield 73 usable results (e.g. Tauxe & Yamazaki, 2015; Nagy et al., 2017). 74

To circumvent the problem of acquiring signals simultaneously from good and bad recorders, de Groot et al. (2018) proposed a new paleomagnetic technique to obtain magnetic moments of individual iron-oxide grains in a sample: micromagnetic tomography (MMT). MMT infers magnetic moments by first acquiring a magnetic surface image of a (cut-out of a) thin-section of a sample using a surface magnetometry technique (e.g. a Quantum Diamond Microscope: Glenn et al., 2017). Then, the position of each ironoxide grain in that sample is obtained with X-Ray micro computed tomography (MicroCT).

-4-

These two pieces of information are combined in a mathematical inversion to retrieve individual magnetic moments by minimizing residuals in the magnetic surface map. Both de Groot et al. (2021) and Kosters et al. (2023) have shown the capabilities of MMT on Hawaiian rock samples. The mathematical uncertainties of the technique have been modeled in Out et al. (2022).

Theoretically, MMT is able to retrieve the magnetic moment of even the smallest 87 SD iron-oxide particles. These particles are unfortunately often not detected by present-88 day MicroCT equipment, because these machines have resolutions down to 350-500 nm. 89 As a grain can only be reliably resolved if it consists of a couple of voxels, it is only pos-90 sible to detect particles with a diameter of >1 μ m. Consequently, SD and PSD parti-91 cles may produce a detectable magnetic signal in the magnetic surface scan, while they 92 do not show up in the MicroCT data. This problem challenges MMT, because the method 93 can only deliver unique and reliable results if all measured magnetic signals originate from 94 detected iron-oxides in the sample (Fabian & De Groot, 2019). Currently it is unknown 95 how many small iron-oxides that go undetected by the MicroCT analyses (i.e. ghost par-96 ticles) are present in (Hawaiian) basaltic lavas. Here we aim to image and characterize 97 the distribution of these ghost particles using Focused Ion Beam Scanning Electron Mi-98 croscopy (FIB-SEM), following the slice-and-view approach of Nikolaisen et al. (2020). 99 With FIB-SEM it is possible to image a sample of micrometer size with a resolution of 100 ~ 20 nm by slicing (or milling) through the sample and imaging the cleared surface. In 101 this way, a 3D reconstruction of the sample volume with these previous undetected iron-102 oxide particles can be reconstructed. Since iron-oxide grains reflect electrons well, the 103 particles are simply detected through backscattered electron images obtained with scan-104 ning electron microscopy. Nikolaisen et al. (2020) used the grain shapes to model mi-105 cromagnetic properties of the individual grains that were characterized by the FIB-SEM 106 analyses. Here we will use the three dimensional characterisation of the small iron-oxides 107 present in our sample to reconstruct a grain-size distribution. We combine the grain-size 108 distribution obtained using the FIB-SEM with the MicroCT analyses on the same sam-109 ple material to obtain a grain-size distribution spanning both the nanometer and microm-110 eter scale. Linking these two datasets, however, is not straightforward because iron-oxides 111 are not uniformly present in a rock sample but cluster around large mineral interfaces. 112 We therefore use the surface area of relatively large, non-magnetic, grains in both the 113 FIB-SEM and MicroCT data to normalize and combine the grain-size distributions ob-114

-5-



Figure 1. Overview of the study area within a Hawaiian basalt examined by slice-and-view through FIB-SEM. After the ion beam removed a 20.18 nm slice from the bottom of the study area, a backscattered and secondary reflected electron image is recorded. Instrumental drift is constrained by calibrating the FIB-SEM every 100 slices using the cross-correlation mark in the top left corner.

tained from both analyses. Lastly, we will demonstrate how these particles can be placed
in numerical models to simulate a realistic basalt sample. With these models, it might
ultimately be possible to investigate the effect of resolution limitations of MicroCT leading to the presence of ghost particles in MMT analyses.

¹¹⁹ 2 Grain-size distribution

Pivotal in obtaining a realistic grain-size distribution is acquiring dimensional properties of all iron-oxide grains present in basaltic rocks. The sample we have used in this

study was taken from the 1907 lava flow from Hawaii (HW-03; de Groot et al., 2013; ter 122 Maat et al., 2018)). The MicroCT data of this sample had already been studied by de 123 Groot et al. (2021) and Kosters et al. (2023), facilitating the detection of grains larger 124 than $\sim 1 \ \mu m$, but leaving most of the SD and PSD particles undetected. Fortunately, with 125 the help of FIB-SEM (e.g. Einsle et al., 2016; Nikolaisen et al., 2020) these remaining 126 SD and PSD grains could finally be made visible. FIB-SEM allows a pixel size of 10 nm 127 and a field of view close to 20 by 20 μ m, which enables the detection of both the small-128 est SD grains above the superparamagnetic threshold of 40 nm (Dunlop & Özdemir, 1997) 129 and large MD grains of $>1 \ \mu m$, which are normally at the lower detection limit of Mi-130 croCT. 131

132

2.1 Image acquisition with FIB-SEM

We applied FIB-SEM on a sample of HW-03, following the slice-and-view proce-133 dure described in Einsle et al. (2016) and Nikolaisen et al. (2020). A 60 μ m thick slice 134 with a diameter of 2 mm was extracted from the specimen and coated with a nanome-135 ter thick layer of gold. Around the edges of the sample, silver was applied with a tooth-136 pick so that electrons could easily pass through the sample. The sample was placed in 137 a Zeiss Crossbeam 540 after which we searched the top of the sample for a 20 x 20 μ m 138 area containing many small iron-oxides. This area was prepared for the slice-and-view 139 procedure by first applying a 1 μ m thick layer of platinum while maintaining an ion beam 140 current of 700 pA and an accelerating voltage of 30 kV. This accelerating voltage did 141 not change unless explicitly stated otherwise. 142

Then, five 0.5 µm deep fiducials were engraved with a current of 50 pA (Fig. 1). 143 The three central fiducials were created parallel to each other to enable image alignment 144 during data processing. The outer two fiducials were created under an angle of 25 de-145 grees with respect to the three central fiducials. These two fiducials ensured that the real 146 thickness of the individual image slices could later be checked. To make these fiducials 147 visible, a 0.6 μ m thick carbon layer was deposited on top of the platinum layer with an 148 ion beam current of 700 pA and a dwell-time of 400 seconds. Then, three trenches with 149 a depth of 25 μ m and an area of 27 by 46 μ m² were eroded at three sides of the selected 150 area. For this erosion procedure, we used an ion beam current kept at 15 nA with a back-151 and-forth milling pattern (bidirectional). Using the same bidirectional routine, the cur-152 rent was reduced from 15 nA to 7 nA to create smooth trenches next to the studied area. 153

Instrumental drift during measurements was reduced to a minimum by applying a cross correlation mark outside the studied area (Fig. 1).

The data acquisition process was initiated after a 20.18 nm slice of the sample area 156 was removed with FIB (beam current at 1.5 nA with an accelerating voltage between 157 1 to 6 kV). Backscattered and secondary reflected electrons images where obtained us-158 ing SEM, after which a 20.18 nm layer of sample was removed using FIB. After this pro-159 cedure was repeated a hundred times instrumental drift was corrected by checking align-160 ment with the cross-correlation mark. The entire process was finished in one day, in which 161 601 grayscale images of both backscattered and secondary reflected electrons were saved 162 in TIFF format. Each file encompassed a pixel area of 3072 x 2104 pixels with a pixel 163 size of 10.09 nm and a 8-bit grayscale. The total volume of sample analysed was there-164 fore 31.0 x 21.2 x 12.1 μm^3 . 165

166

2.2 Dataprocessing

167

2.2.1 Extracting particles from FIB-SEM data

After image acquisition, a grain-size distribution could be extracted from the data. Every data processing step was executed with Python 3 and the OpenCV library (Bradski, 2000). We initiated our data processing by loading the 601 backscattered electron images and aligning them using the three parallel fiducials and the matchTemplate function of OpenCV. After alignment, we removed all duplicate images. To remove the effects of both curtaining and the platina and carbon deposited layers, every image was cropped to 2600×1000 pixels.

Through denoising and thresholding, SD and PSD iron-oxides could be extracted 175 from FIB-SEM images as shown by Nikolaisen et al. (2020). However, we designed the 176 extraction process in such a way that both the surface and the centre of each iron-oxide 177 grain were properly retrieved. We initiated denoising by 2×2 binning of the 549 remain-178 ing images, which resulted into 1300×500 pixels per image. A non-local means denois-179 ing filter from OpenCV (fastNlMeansDenoisingMulti) was used with a search window-180 size of 20 pixels, a filter strength of 20, and template windowsize of 3 pixels applied in 181 the x, y, and z-direction. Each image was then thresholded at a pixelvalue of 220. Af-182 ter thresholding, all images are stacked together into one 3D-volume. In this 3D-volume, 183

-8-

we grouped connected pixels into grains. With this high threshold, we have only obtained
the brightest central iron-oxide pixels.

However, the iron-oxide grains are not properly imaged yet, because the sides of 186 the grains frequently have a lower pixel value than their center and thus have dropped 187 out in the previous thresholding step. To correctly image the boundary of the iron-oxides 188 we applied a lower threshold value of 160 instead of 220 to the previous set of images. 189 Then we applied the canny edge algorithm of OpenCV (Canny). This algorithm detects 190 spatial changes in pixel value and draws, based on that change, boundaries between pixel 191 values. After drawing boundaries, small gaps in the boundaries were closed using the mor-192 phology close operation in order to create enclosed areas. Again, all images were com-193 bined into a 3D-volume and connected pixels were grouped into grains. However, the low 194 threshold outputted both well defined grain shapes of iron-oxides and poorly defined shapes 195 of other unwanted minerals. To get rid of these unwanted minerals, we mapped the grains 196 obtained with high thresholding (poorly defined boundaries, most certainly iron-oxides) 197 on top of the grains obtained with low thresholding (well defined boundaries, possible 198 iron-oxides) to only remain with actual iron-oxides. Lastly, we removed all iron-oxides 199 that extend beyond the edges of the FIB-SEM domain. We calculated the volume of 1,558 200 remaining iron-oxides, and transformed these volumes into an equivalent diameter as-201 suming spherical grains. 202

203

2.2.2 Scaling FIB-SEM and MicroCT data

To obtain a grain-size distribution spanning the entire range of naturally occur-204 ring iron-oxides, we combined our FIB-SEM data with MicroCT data that was obtained 205 on a sister specimen of HW-03 by Kosters et al. (2023). Since iron-oxides are not homo-206 geneously distributed throughout the rock, it is not straightforward to link MicroCT data 207 one-to-one to FIB-SEM data based on volume. Most iron-oxides are positioned on the 208 interface between larger grains (mainly plagioclase and pyroxene); especially the small-209 est iron-oxides attach themselves to this 'honeycomb' structure of other minerals (Fig. 210 2). Because this honeycomb structure of relatively large minerals can be imaged prop-211 erly using MicroCT, it allows for extrapolating the number of undetected small iron-oxides 212 from the FIB-SEM analysis to the MicroCT data and hence produce a continuous grain-213 size distribution over both analyses. 214



Figure 2. Overview of the $26.2 \times 10.1 \times 10.5 \ \mu m^3$ volume exposed to the slice-and-view procedure with FIB-SEM. The large grains are indicated by a grayish color. The iron-oxides residing in this volume are indicated in yellow-brown. Note that most iron-oxides reside on the large mineral interfaces.

Characterizing the mineral surface area of the larger minerals, mainly plagioclase 215 and pyroxene, required reprocessing the FIB-SEM data, with this specific goal in mind. 216 First, the FIB-SEM dataset was denoised using a non-local means filter; then we applied 217 a median filter with a spherical footprint of 13 pixels in diameter to smooth out irreg-218 ularities at the boundaries. After this pre-processing we applied a K-means algorithm 219 to create five groups of pixels. For each individual image, we removed the first group be-220 cause it corresponded to drilling artifacts. Additionally, we removed all areas smaller than 221 3000 pixels, because these areas do not represent grain interfaces of the larger minerals 222 present in the sample that are typically occupied by iron-oxides. Finally, the mineral sur-223 face area was calculated for the remaining pixel values using a 'Lindblad' algorithm (Lindblad, 224 2005). The mineral surface area consisted of about 3.8×10^6 squared pixels that cor-225 respond to an area of $1.53 \times 10^{-3} \text{ mm}^2$. Fig. 2 shows the position of the mineral sur-226 face area together with the iron-oxides. 227

To calculate the mineral surface area in the MicroCT dataset, we started with raw MicroCT data and processed it using Dragonfly software. After loading the DICOM files, we cropped about 100 μ m on both sides of the sample to remove scanning residuals. To align the sample with the coordinate axes, we rotated the sample 0.85° clockwise around

-10-

the z-axis, 0.50° anticlockwise around the x-axis, and resampled the image with 0.75 μ m 232 pixel size. Then the image was cropped further, which resulted into a region of 1979 \times 233 44×1901 voxels. We denoised the image by applying a non-local means filter with a 234 3D spherical kernel size of 11 pixels. To create a larger contrast we applied a median fil-235 ter with a 3D spherical kernel size of 13 pixels. The different minerals were separated 236 by applying a K-means filter of 3 groups, of which the group with the highest pixel value 237 corresponds to iron-oxides. The other two groups chiefly consisted of plagioclase and py-238 roxene. Since we observed that SD and PSD iron-oxides tend to cluster on the interfaces 239 of these minerals, we applied a 'sobel edge detection' filter to these groups. The result 240 of this filter enabled us to create a honeycomb structure of the sample, on which we could 241 populate SD and PSD iron-oxides to create realistic models of basalts. This mineral sur-242 face area of 4.00 mm^2 is approximately 2,620 times larger than the surface area found 243 in the FIB-SEM dataset. This means that we need to multiply the number of FIB-SEM 244 grains by 2,620 to scale correctly to the MicroCT dataset. This resulted into a combined 245 dataset of 4.7×10^6 FIB-SEM and 1.6×10^3 MicroCT iron-oxide grains. The FIB-SEM 246 and MicroCT datasets can now be combined to produce a continuous grain-size distri-247 bution of iron-oxides for the entire range of interest (Fig. 3). The median grain-size has 248 a diameter of 70 nm. From 30 to 70 nm the occurrence rapidly increases, after which the 249 occurrence of grain diameters between 70 nm and 10 μ m follows are more gradually de-250 caying trend. 251

252

2.3 Constructing the grain-size distribution

The trend in grain-size distribution can be described using a lognormal-like distri-253 bution between 20 nm and 10 μ m. We have chosen the lognormal distribution, because 254 this distribution is frequently used to describe grain-sizes for e.g. magnetite powders (Smirnov, 255 2006; Yu et al., 2002). Nevertheless, the lognormal trend line was originally fitted to the 256 10th-logarithm of the data. This logarithm ensures a better fit to grain-sizes larger than 257 70 nm. Through an iterative procedure reducing the least-squares error of our fit to the 258 data, we found the best fitting distribution after which we transformed it back to lin-259 ear space: 260

$$p = 3.51 * 10^{-6 + \frac{5.61}{d\sqrt{2\pi}}e^{-\frac{\ln^2(0.60d)}{6.29}}} \text{ if } d : [0, 10]$$
(1)



Figure 3. Overview of the non-normalized grain-size distribution. The histograms show the number of iron-oxides obtained from FIB-SEM or MicroCT data. The datasets are scaled to each other using mineral surface area (see main text). The logarithmic trend line shows the fit to these scaled distributions.

with d the diameter of the grain in μ m and p the occurrence or probability of finding that grain-size. Note that this probability density function, or pdf, is no longer a lognormal distribution, yet still accurately describes the data (Fig. 3). The function is only valid for grain-sizes between 0 and 10 μ m, although it could be extended to larger grain-sizes through rescaling.

To obtain a function from which grain-sizes can be sampled, the probability den-266 sity function is integrated into a cumulative probability density function (cdf). This func-267 tion takes a grain-size as input and returns the probability of finding that grain-size or 268 smaller as output. Then to actually create a sampling function, this cdf has to be inverted. 269 By inverting the function, a grain-size could be obtained as function of inputted cumu-270 lative probability. However, this inverted function has no analytical expression, so we 271 created a look-up table to be able to sample the function. To create this table we inputted 272 100,000 diameters between 0 and 10 μ m into the cdf, which returned 100,000 cumula-273 tive probabilities. Each grain-size-probability pair was then put into the table. To sam-274 ple grain-sizes from this look-up table we would generate a pool of random numbers, ob-275

-12-

tained from a uniform distribution between 0 and 1 (U(0,1)), equal to the number of samples requested. For each random number we would then extract the corresponding grain-size from the look-up table. Grain-sizes or probabilities not present in the table would be obtained through cubic interpolation of the nearest values. All Python codes to create and sample the grain-size distribution, and the look-up table are available in our repository (Out et al., 2023).

²⁸² 3 Discussion

283

3.1 Dataset limitations and validation

For this study only one FIB-SEM sample was used from a specific, but typical Hawai-284 ian lava flow. How well this sample represents the sample material of site HW-03, let alone 285 material from other sites, flows, and/or locations is a big unknown. Obtaining data from 286 rock samples using FIB-SEM, however, is not a straightforward task; sample prepara-287 tion, machine handling, and availability of resources complicate the scanning process. 288 To get the most out of our data, we decided during initial phase of the FIB-SEM anal-289 yses to target an area with the highest amount of iron-oxides in view. This implies that 290 relatively many small grains will be present in our FIB-SEM study, leading to an over-291 representation of small grains in our final grain-size distribution. 292

The overestimation, however, might be somewhat damped through how we con-203 structed the scaling factor. Our scaling factor is dependent on the amount of mineral 294 surface area which probably scales with the number of grains. This might mean that our 295 scaling factor is less sensitive to grain density variations in the studied FIB-SEM area, 296 than, for example, a scaling factor based on volume. In case of volume scaling, the FIB-297 SEM area would be scaled by a factor of 24,000, independent of the number of grains 298 or amount of mineral surface area present in the studied area. This would most defini-299 tively result in a severe overestimation of particles imaged by FIB-SEM compared to par-300 ticles imaged by MicroCT, because this volume scaling factor is one order of magnitude 301 larger than our scaling factor. We, therefore, consider scaling by grains per mineral in-302 terface area a stable protocol to ensure comparability between FIB-SEM and MicroCT 303 results as it dampens local variations of the studied area. 304

To obtain an independent verification for our method of combining FIB-SEM and MicroCT analyses, we validated our dataset through a scanning electron microscopy (SEM)

-13-

image on a sister specimen from site HW03 (de Groot et al., 2013; ter Maat et al., 2018). Two areas of $48.2 \times 71.3 \ \mu\text{m}^2$ and $33.2 \times 49.0 \ \mu\text{m}^2$ within the sample were imaged with a resolution of 80 and 55 nm per 8-bit pixel respectively. After noise was removed, the images were thresholded at a pixel value of 153 to only retain iron-oxides. Then the major axis *m* and minor axis *n* of each grain were determined and grain diameters, *d*, were calculated using (Yu et al., 2002):

$$d = 2\sqrt{m\frac{n}{\pi}} \tag{2}$$

The diameters of the grains were sorted to produce a normalized histogram with bin-313 sizes of 0.1 μ m ranging between 0.2 and 3.0 μ m - straddling the transition from the FIB-314 SEM realm to the MicroCT range. We discarded grains smaller than 0.2 μ m from the 315 SEM data because they would be constructed from two pixels or less. On top of the SEM 316 grain-size data, we overlaid the FIB-SEM and MicroCT normalized data (Fig. 4). The 317 SEM data misses some of the smaller grains, as expected. The trends in the FIB-SEM+MicroCT 318 and SEM datasets for grain-sizes $>0.3 \mu m$, however, correspond well. Taking into ac-319 count that SEM and FIB-SEM are very different scanning techniques compared to Mi-320 croCT, we are confident that our constructed grain-size distribution is a proper descrip-321 tion of the actual grain-size distribution of iron-oxides in our samples, although it may 322 emphasize small, sub-micron, grains. 323

324

3.2 Building models of the physical distribution of iron-oxides

MMT struggles with the smallest particles as they are not detected by MicroCT acquisitions. Although these grains are invisible for MicroCTs, they might still produce a signal in a magnetic surface scan, possibly leading to incorrect magnetic moments of other detected particles (Fabian & De Groot, 2019; Out et al., 2022). However, with the grain-size distribution presented here, it is now possible to study the effect of these ghost grains through generating realistic grain models of basalt and simulating their effect on MMT results in future studies.

To generate a realistic model of the physical distribution of iron-oxide grains in a basaltic sample, we aim to populate a honeycomb structure of the larger minerals obtained from MicroCT analyses with a simulated distribution of small iron-oxides that the MicroCT would miss (Fig. 5). The first step is to extract a honeycomb structure and

-14-



Figure 4. Verification of MicroCT and FIB-SEM data using separate SEM imaged iron-oxides using normalized histograms. Note that the distribution is similar for the two data-series except for the smallest grain-sizes.

the large iron-oxides from a scanned basaltic sample and determine the resolution of the 336 MicroCT. The MicroCT resolution determines how many ghost grains are missing per 337 mm^2 honeycomb structure. By multiplying this ghost grain density (Figure 6) with the 338 honeycomb structure surface area extracted from MicroCT images, we can calculate the 339 total number of missing grains. Then we use the grain-size distribution obtained in this 340 study to make a list of the missing ghost grains and their diameters by feeding the num-341 ber of missing ghost grains together with the MicroCT resolution into the grain-size dis-342 tribution sampling function (available in the repository of Out et al., 2023). Finally, we 343 simulate grains from the list of ghost grains and place them on the honeycomb struc-344 ture obtained from the MicroCT analysis besides the larger iron-oxide particles that were 345 already detected. Hence, by extracting the honeycomb structure and MicroCT resolu-346 tion, we are able to simulate a realistic distribution of all iron-oxides in a basaltic sam-347 ple, with diameters down to the superparamagnetic range. 348

As an example we consider a hypothetical MicroCT analysis with a resolution of $0.1 \ \mu m$ that produced a honeycomb structure with a surface area of 4 mm². Then we would need to populate this honeycomb structure with 9×10^5 particles per mm² (dashed line in Fig. 6). This implies that we miss 3.6×10^6 ghost grains in the entire sample with

-15-



Figure 5. Overview of the protocol to build realistic models of the physical distribution of iron-oxides in a basalt. By segmenting the MicroCT image, we only retain the honeycomb structure and the visible iron-oxides. The MicroCT resolution is required as input for the relation to sample the grain-size distribution of iron-oxides (Fig. 6). Multiplying the ghost grain density with the honeycomb surface area extracted from the MicroCT image, we are able to determine how many grains are left undetected by MicroCT. Given the MicroCT resolution and the number of ghost grains, we sample the grain-size distribution function (see Out et al., 2023) to obtain a list of grain diameters. These grain diameters are used to construct particles that are randomly placed on the extracted honeycomb structure. A graphical overview can be found in the supplementary information (Figure S1).



Figure 6. Maximum probability and particles per mm² mineral interface area as function of MicroCT resolution. Based on this figure we can determine how many grains have to be sampled from the provided grain-size distribution. By providing the maximum cumulative probability to the grain-size distribution, we can select the maximum grain-size that should be sampled. For example, for 4 mm² mineral interface area and a MicroCT resolution of 0.1 μ m, we would sample 900.000 particles/mm² × 4 mm² = 3.6 million particles with a maximum cumulative probability of 0.57. For a resolution larger than 1.0 μ m, we would sample 1.20 million grains per mm² mineral interface area.

a surface area of 4 mm². To obtain the list of grains that we need to add to the honeycomb structure we sample the grain-size distribution 3.6 million times using a maximum cumulative probability of 0.57 (dashed line in Fig. 6) for grains with a diameter between 0 and 0.1 μ m. These grains are then randomly placed on the honeycomb structure together next to the larger iron-oxide particles that were already detected with MicroCT.

358

3.3 Implications for Micromagnetic Tomography

The results of previous MMT based on synthetic samples in which the grain size 359 distribution was controlled (e.g. (most) results in de Groot et al., 2018, 2021)) are most 360 likely not disturbed by ghost grains, because the iron-oxide grains have been sieved for 361 a size larger than the MicroCT resolution. For that reason, their sample mainly consisted 362 of particles larger than 5 μ m, which could be easily detected by MicroCT. However, the 363 natural basalt samples used by de Groot et al. (2021) and Kosters et al. (2023) undoubt-364 edly contain undetected ghost grains. Both articles study the magnetic moment of 1,646 365 iron-oxides detected by MicroCT in a 30 μ m thin section of Hawaiian basaltic rock, which 366 is a sister-specimen as subjected to our study. The resolution of the MicroCT used in 367 those studies is $\sim 1 \ \mu m$ and the mineral surface area is $\sim 4 \ mm^2$. According to Fig. 6, 368 this implies the existence of 4.8×10^6 undetected particles, which results in ~3,000 ghost 369 grains per detected iron-oxide. 370

All of these ghost grains might scramble the signal of the detected iron-oxides lead-371 ing to erroneous magnetic moments for the detected grains produced by the MMT in-372 version. Remarkably, both studies do not present unrealistic magnetic moments for the 373 detected iron-oxide particles, and their results are (somewhat) interpretable in terms of 374 magnetic moment and direction. Also, reported trends as function of grain-size and mag-375 netic treatment follow expectations based on earlier work (e.g. Dunlop & Özdemir, 1997). 376 It is therefore likely that, although high in number, the small ghost grains do not have 377 a large enough magnetic moment to have a major contribution to magnetic signal on the 378 surface of the sample. Another possibility is that the spatial resolution of the surface mag-379 netometry data might not be sufficient to locate the weak signal of these ghost particles. 380 As its magnetic moment is highly dependent on the size of a grain (Dunlop & Ozdemir, 381 1997) it is likely that the magnetic moments of larger grains are less affected by ghost 382 grains in an MMT inversion, compared to grains that are just above the detection thresh-383 old of the MicroCT analysis used. To characterise the reliability of MMT results as func-384

-18-

tion of grain size, a future study that models the impact of the undetected ghost grains on MMT results is critical.

387 4 Conclusions

Here we produced a grain-size distribution of all iron-oxides larger than the super-388 paramagnetic threshold in an Hawaiian basalt by combining MicroCT and FIB-SEM data 389 of sister-specimens. The MicroCT and FIB-SEM datasets could not be combined one-390 to-one due to volume difference, so they were scaled using the interface area of large, non-391 iron-oxide, minerals present in the sample (i.e. honeycomb structure), because iron-oxides 392 tend to cluster on the edges of large minerals. Through this scaling procedure we have 393 created a lognormal-like distribution between 20 nm to 10 μ m that spans the range of 394 SD, PSD, and MD iron-oxides. This distribution is probably a slight overestimation of 395 the concentration of iron-oxide grains, caused by the choice of FIB-SEM study area with 396 relatively many iron-oxides. This grain-size distribution can be used to populate real-397 istic models of the physical distribution of iron-oxides in a sample; and makes it possi-398 ble to assess the impact of undetected ghost grains on micromagnetic and MMT stud-399 ies. Due to the high demands of both MicroCT and FIB-SEM analyses we could only 400 process one sample from a well-characterised Hawaiian lava flow in this study. Future 401 studies can build on our theory of combining FIB-SEM and MicroCT data through the 402 surface area of the honeycomb structure of larger minerals produced by the MicroCT anal-403 ysis to study other sample material as well. 404

405 Data availability statement

The FIB-SEM, MicroCT, and SEM data used for obtaining and verifying, and the 406 jupyter notebooks employed for generating the grain-size distribution and honeycomb 407 structure in this study are available at YODA via https://doi.org/10.24416/UU01-QNASXB 408 with a CC-BY-NC-4.0 license (Out et al., 2023). Numerical calculations were executed 409 with support of NumPy (Harris et al., 2020), SciPy (Virtanen et al., 2020), OpenCV (Bradski, 410 2000), and Matplotlib (Hunter, 2007) Python libraries. The MicroCT analysis for this 411 study was generated using Dragonfly software, Version 2022.1.0.1259 for Windows. Ob-412 ject Research Systems (ORS) Inc, Montreal, Canada, 2022. 413

-19-

414 Acknowledgments

- ⁴¹⁵ This project has received funding from the European Research Council (ERC) under the
- ⁴¹⁶ European Union's Horizon 2020 research and innovation program (Grant agreement No.
- ⁴¹⁷ 851460 to L.V. de Groot). This publication results from work carried out under transna-
- ⁴¹⁸ tional access action under the support of EXCITE EC- HORIZON 2020 INFRAIA
- ⁴¹⁹ 2020 Integrating Activities for Starting Communities (Grant agreement No. 101005611).

420 **References**

445

- Berndt, T., Muxworthy, A. R., & Fabian, K. (2016). Does size matter? statistical
 limits of paleomagnetic field reconstruction from small rock specimens. Journal
 of Geophysical Research: Solid Earth, 121(1), 15–26.
- ⁴²⁴ Bradski, G. (2000). The OpenCV Library. Dr. Dobb's Journal of Software Tools.
- de Groot, L. V., Biggin, A. J., Dekkers, M. J., Langereis, C. G., & Herrero-Bervera,
- E. (2013, 12). Rapid regional perturbations to the recent global geomagnetic decay revealed by a new Hawaiian record. Nature Communications, 428 4(1), 2727. Retrieved from www.nature.com/naturecommunicationshttp:// 429 www.nature.com/articles/ncomms3727 doi: 10.1038/ncomms3727
- ⁴³⁰ De Groot, L. V., Fabian, K., Bakelaar, I. A., & Dekkers, M. J. (2014). Magnetic
 ⁴³¹ force microscopy reveals meta-stable magnetic domain states that prevent re⁴³² liable absolute palaeointensity experiments. Nature Communications, 5(1),
 ⁴³³ 4548.
- 434 de Groot, L. V., Fabian, K., Béguin, A., Kosters, M. E., Cortés-Ortuño, D., Fu,
- R. R., ... Barnhoorn, A. (2021). Micromagnetic tomography for paleomagnetism and rock-magnetism. Journal of Geophysical Research: Solid Earth, 126(10), e2021JB022364.
- de Groot, L. V., Fabian, K., Béguin, A., Reith, P., Barnhoorn, A., & Hilgenkamp,
 H. (2018). Determining individual particle magnetizations in assemblages of
 micrograins. *Geophysical Research Letters*, 45(7), 2995–3000.
- Dunlop, D. J., & Özdemir, Ö. (1997). Rock Magnetism. Cambridge University Press.
 Retrieved from https://www.cambridge.org/core/product/identifier/
 9780511612794/type/book doi: 10.1017/cbo9780511612794
- Einsle, J. F., Harrison, R. J., Kasama, T., Conbhuí, P. Ó., Fabian, K., Williams,
 - W., ... Midgley, P. A. (2016). Multi-scale three-dimensional charac-

446	terization of iron particles in dusty olivine: Implications for paleomag-
447	netism of chondritic meteorites. American Mineralogist, 101(9), 2070–
448	2084. Retrieved from https://doi.org/10.2138/am-2016-5738CCBY doi:
449	doi:10.2138/am-2016-5738CCBY
450	Fabian, K., & De Groot, L. V. (2019, 2). A uniqueness theorem for tomography-
451	assisted potential-field inversion. $Geophysical Journal International, 216(2),$
452	760-766. Retrieved from http://arxiv.org/abs/1712.06136https://
453	academic.oup.com/gji/article/216/2/760/5151336 doi: 10.1093/gji/
454	ggy455
455	Glenn, D. R., Fu, R. R., Kehayias, P., Le Sage, D., Lima, E. A., Weiss, B. P., &
456	Walsworth, R. L. (2017, 8). Micrometer-scale magnetic imaging of geological
457	samples using a quantum diamond microscope. Geochemistry, Geophysics,
458	Geosystems, 18(8), 3254–3267. Retrieved from https://onlinelibrary
459	.wiley.com/doi/abs/10.1002/2017GC006946 doi: 10.1002/2017GC006946
460	Harris, C. R., Millman, K. J., van der Walt, S. J., Gommers, R., Virtanen, P., Cour-
461	napeau, D., Oliphant, T. E. (2020, 9). Array programming with NumPy.
462	Nature, 585(7825), 357-362. Retrieved from http://www.nature.com/
463	articles/s41586-020-2649-2 doi: 10.1038/s41586-020-2649-2
464	Hunter, J. D. (2007, 5). Matplotlib: A 2D Graphics Environment. Com-
465	puting in Science & Engineering, $9(3)$, 90–95. Retrieved from http://
466	<pre>nipy.scipy.http://ieeexplore.ieee.org/document/4160265/ doi:</pre>
467	10.1109/MCSE.2007.55
468	Kosters, M. E., de Boer, R. A., Out, F., Cortés-Ortuño, D. I., & de Groot, L. V.
469	(2023). Unraveling the magnetic signal of individual grains in a hawaiian lava
470	using micromagnetic tomography. Geochemistry, Geophysics, Geosystems,
471	24(4), e2022GC010462. Retrieved from https://agupubs.onlinelibrary
472	.wiley.com/doi/abs/10.1029/2022GC010462 (e2022GC010462
473	2022GC010462) doi: https://doi.org/10.1029/2022GC010462
474	Lindblad, J. (2005). Surface area estimation of digitized 3d objects using weighted
475	local configurations. Image and Vision Computing, $23(2)$, $111-122$.
476	Nagy, L., Williams, W., Muxworthy, A. R., Fabian, K., Almeida, T. P., Conbhui,
477	P. O., & Shcherbakov, V. P. (2017). Stability of equidimensional pseudo-single-
478	domain magnetite over billion-year timescales. Proceedings of the National

manuscript submitted to *G*-cubed

479	Academy of Sciences of the United States of America, 114(39), 10356-10360.
480	doi: 10.1073/pnas.1708344114
481	Néel, L. (1955, 4). Some theoretical aspects of rock-magnetism. Advances in
482	Physics, 4(14), 191-243. Retrieved from https://www.tandfonline.com/
483	doi/abs/10.1080/00018735500101204http://www.tandfonline.com/doi/
484	abs/10.1080/00018735500101204 doi: $10.1080/00018735500101204$
485	Nikolaisen, E. S., Harrison, R. J., Fabian, K., & McEnroe, S. A. (2020). Hysteresis
486	of natural magnetite ensembles: Micromagnetics of silicate-hosted magnetite
487	inclusions based on focused-ion-beam nanotomography. Geochemistry, Geo-
488	physics, Geosystems, $21(11)$, e2020GC009389.
489	Out, F., Cortés-Ortuñno, D., Fabian, K., van Leeuwen, T., & de Groot, L. V.
490	(2022). A first-order statistical exploration of the mathematical limits of
491	micromagnetic tomography. $Geochemistry, Geophysics, Geosystems.$ doi:
492	e2021GC010184
493	Out, F., de Boer, R. A., Walmsley, J., & de Groot, L. V. (2023). Replication
494	Data for: Modeling the Distribution of Iron-oxides in Basalt by combining
495	FIB-SEM and MicroCT Measurements. YODA. Retrieved from https://
496	$\verb"public.yoda.uu.nl/geo/UU01/QNASXB.html" doi: 10.24416/UU01-QNASXB$
497	Smirnov, A. (2006). Memory of the magnetic field applied during cooling in the low- $% \left(2000,100,100,100,100,100,100,100,100,100$
498	temperature phase of magnetite: Grain size dependence. Journal of Geophysi-
499	cal Research: Solid Earth, 111(B12).
500	Tauxe, L. (2010). Essentials of paleomagnetism. Univ of California Press.
501	Tauxe, L., & Yamazaki, T. (2015). Paleointensities. In <i>Treatise on geophysics</i> (pp.
502	461–509). Elsevier. doi: 10.1016/B978-0-444-53802-4.00107-X
503	ter Maat, G. W., Pennock, G. M., & de Groot, L. V. (2018). Data descrip-
504	tor: A chemical, crystallographic and magnetic characterisation of indi-
505	vidual iron-oxide grains in Hawaiian lavas. Scientific Data, 5, 1–9. doi:
506	10.1038/sdata.2018.162
507	Virtanen, P., Gommers, R., Oliphant, T. E., Haberland, M., Reddy, T., Cournapeau,
508	D., van Mulbregt, P. (2020, 3). SciPy 1.0: fundamental algorithms for sci-
509	entific computing in Python. Nature Methods, 17(3), 261–272. Retrieved from
510	https://doi.org/10.1038/s41592-019-0686-2http://www.nature.com/
511	articles/s41592-019-0686-2 doi: $10.1038/s41592-019-0686-2$

- Yu, Y., Dunlop, D. J., & Özdemir, Ö. (2002). Partial anhysteretic remanent mag netization in magnetite 1. additivity. Journal of Geophysical Research: Solid
- ⁵¹⁴ Earth, 107(B10), EPM-7.

Modeling the Distribution of Iron-oxides in Basalt by combining FIB-SEM and MicroCT Measurements 2

Frenk Out¹, Rosa A. de Boer¹, John Walmsley², Lennart V. de Groot¹

4	$^1\mathrm{Paleomagnetic}$ laboratory Fort Hoofddijk, Department of Earth Sciences, Utrecht University,
5	Budapestlaan 17, 3584 CD Utrecht, The Netherlands.
6	$^{2} \mathrm{Department} \text{ of Materials Science \& Metallurgy, University of Cambridge, 27 Charles Babbage Road, CB3}$
7	0FS Cambridge, United Kingdom.

Key Points: 8

1

3

9	• We produced a grain-size distribution of iron-oxides in a typical Hawaiian basalt
10	from the superparamagnetic threshold (~40 nm) to 10 micron
11	• We combined FIB-SEM and MicroCT data from sister specimens by normalizing
12	them to the mineral surface area of the non-magnetic minerals
13	• Our grain-size distribution can be used to populate realistic models of iron-oxides
14	in a Hawaiian basalt

Corresponding author: Frenk Out, f.out@uu.nl

15 Abstract

Micromagnetic tomography (MMT) aims to go beyond paleomagnetic measurements on 16 bulk samples by obtaining magnetic moments for individual iron-oxide grains present in 17 a sample. To obtain accurate MMT results all magnetic sources and all their magnetic 18 signals should be known. Small particles ($<<1 \mu m$) are often not detected by MicroCT 19 analyses, but do have a magnetic signal, and therefore hamper obtaining reliable MMT 20 results. Currently it is unknown how many of these small 'ghost grains' are present in 21 basaltic samples. Here we aim to obtain a realistic grain-size distribution for iron-oxides 22 in a typical Hawaiian basalt. We characterize the entire grain-size range of interest to 23 paleomagnetism, from the superparamagnetic threshold of ~ 40 nm to multidomain grains 24 with sizes up to 10 μ m. This requires a combination of FIB-SEM slice-and-view and Mi-25 croCT techniques: FIB-SEM characterizes the grains between 20 nm and 1 μ m and Mi-26 croCT detects iron-oxides >750 nm. The FIB-SEM and MicroCT data are combined through 27 normalizing the grain-size distribution using the surface area of non-magnetic minerals 28 that are characterised in both datasets. Then, a lognormal-like grain-size distribution 29 is acquired for the entire grain-size range. Our dataset enables future studies to popu-30 late (MMT) models with a realistic distribution of even the smallest iron-oxide grains, 31 which ultimately may shed light on the confounding influence of such ghost grains on 32 MMT results. 33

34

Plain Language Summary

Micromagnetic tomography (MMT) is a promising new technique that extracts high-35 quality magnetic information from lavas. Magnetic particles in lavas obtain a magnetic 36 signal while cooling in presence of Earth's magnetic field. However, not all particles store 37 the signal well, meaning that both good and bad recorders are present. Classical pale-38 omagnetic techniques measure the magnetic signal of all recorders together, i.e. the bulk 39 signal. MMT, however, acquires the signal from individual recorders in the lava, enabling 40 selecting only the good recorders and rejecting the signal of bad recorders. MMT needs 41 two pieces of information: (1) the surface magnetic field produced by recorders and (2) 42 the location of all recorders in the lava. Unfortunately, the position of recorders smaller 43 than $\sim 1 \ \mu m$ are often unknown due to measurement limitations. If these small parti-44 cles are not detected, they may scramble the calculated signal of larger recorders. To in-45 vestigate how many disruptive undetected grains are present, we scanned a small vol-46

-2-

- 47 ume of lava on high resolution to extract all magnetic recorders and obtain a grain-size
- 48 distribution. With this distribution we may estimate if these small recorders distort the
- ⁴⁹ signal of larger detected grains.

50 1 Introduction

Volcanic rocks are often perceived to be excellent recorders of paleomagnetic fields, although only a small portion of all minerals in a volcanic rock has magnetic properties. These minerals, iron-oxide grains, present in a lava acquire a magnetization by cooling in the Earth's magnetic field. Their magnetization is thus a thermo-remanent magnetization (TRM), and these magnetic signals are regarded to be stable, often over millions or billions of years (e.g. Dunlop & Özdemir, 1997; Tauxe & Yamazaki, 2015).

Small particles, or single domain (SD) grains, have two configurations for storing 57 their magnetic signal: either parallel or anti-parallel to the grain's easy axis. One of these 58 two options is selected at random but with a slight preference to align with the prevail-59 ing magnetic field (e.g. Tauxe, 2010; Berndt et al., 2016). In slightly larger particles, which 60 are often referred to as pseudo-single domain (PSD) grains, the magnetic signal is stored 61 in a more complex vortex structure. Fortunately, this signal usually represents the orig-62 inal imparting magnetic field accurately (Nagy et al., 2017). The last group of particles 63 are iron-oxides with diameters $>1 \ \mu m$ and belong to the multi-domain (MD) class. Al-64 though these MD grains may store magnetic signals for millions of years, their expres-65 sion of the past field is often disturbed, overprinted, unstable, or lost (e.g. De Groot et 66 al., 2014). This instability is caused by magnetic reassembling of the original magnetic 67 structure of a grain over time to achieve energy minimization (Néel, 1955). MD grains 68 are therefore mostly perceived as bad paleomagnetic recorders, while SD and PSD grains 69 are considered to be more reliable recorders of paleomagnetic fields. To make it even more 70 difficult, most natural rock samples consist of a mixture of SD, PSD and MD grains. The 71 combination of these good and bad recorders in rocks is the main reason that generally 72 only 10 to 20% of all paleointensity experiments pass sufficient selection criteria to yield 73 usable results (e.g. Tauxe & Yamazaki, 2015; Nagy et al., 2017). 74

To circumvent the problem of acquiring signals simultaneously from good and bad recorders, de Groot et al. (2018) proposed a new paleomagnetic technique to obtain magnetic moments of individual iron-oxide grains in a sample: micromagnetic tomography (MMT). MMT infers magnetic moments by first acquiring a magnetic surface image of a (cut-out of a) thin-section of a sample using a surface magnetometry technique (e.g. a Quantum Diamond Microscope: Glenn et al., 2017). Then, the position of each ironoxide grain in that sample is obtained with X-Ray micro computed tomography (MicroCT).

-4-

These two pieces of information are combined in a mathematical inversion to retrieve individual magnetic moments by minimizing residuals in the magnetic surface map. Both de Groot et al. (2021) and Kosters et al. (2023) have shown the capabilities of MMT on Hawaiian rock samples. The mathematical uncertainties of the technique have been modeled in Out et al. (2022).

Theoretically, MMT is able to retrieve the magnetic moment of even the smallest 87 SD iron-oxide particles. These particles are unfortunately often not detected by present-88 day MicroCT equipment, because these machines have resolutions down to 350-500 nm. 89 As a grain can only be reliably resolved if it consists of a couple of voxels, it is only pos-90 sible to detect particles with a diameter of >1 μ m. Consequently, SD and PSD parti-91 cles may produce a detectable magnetic signal in the magnetic surface scan, while they 92 do not show up in the MicroCT data. This problem challenges MMT, because the method 93 can only deliver unique and reliable results if all measured magnetic signals originate from 94 detected iron-oxides in the sample (Fabian & De Groot, 2019). Currently it is unknown 95 how many small iron-oxides that go undetected by the MicroCT analyses (i.e. ghost par-96 ticles) are present in (Hawaiian) basaltic lavas. Here we aim to image and characterize 97 the distribution of these ghost particles using Focused Ion Beam Scanning Electron Mi-98 croscopy (FIB-SEM), following the slice-and-view approach of Nikolaisen et al. (2020). 99 With FIB-SEM it is possible to image a sample of micrometer size with a resolution of 100 ~ 20 nm by slicing (or milling) through the sample and imaging the cleared surface. In 101 this way, a 3D reconstruction of the sample volume with these previous undetected iron-102 oxide particles can be reconstructed. Since iron-oxide grains reflect electrons well, the 103 particles are simply detected through backscattered electron images obtained with scan-104 ning electron microscopy. Nikolaisen et al. (2020) used the grain shapes to model mi-105 cromagnetic properties of the individual grains that were characterized by the FIB-SEM 106 analyses. Here we will use the three dimensional characterisation of the small iron-oxides 107 present in our sample to reconstruct a grain-size distribution. We combine the grain-size 108 distribution obtained using the FIB-SEM with the MicroCT analyses on the same sam-109 ple material to obtain a grain-size distribution spanning both the nanometer and microm-110 eter scale. Linking these two datasets, however, is not straightforward because iron-oxides 111 are not uniformly present in a rock sample but cluster around large mineral interfaces. 112 We therefore use the surface area of relatively large, non-magnetic, grains in both the 113 FIB-SEM and MicroCT data to normalize and combine the grain-size distributions ob-114

-5-



Figure 1. Overview of the study area within a Hawaiian basalt examined by slice-and-view through FIB-SEM. After the ion beam removed a 20.18 nm slice from the bottom of the study area, a backscattered and secondary reflected electron image is recorded. Instrumental drift is constrained by calibrating the FIB-SEM every 100 slices using the cross-correlation mark in the top left corner.

tained from both analyses. Lastly, we will demonstrate how these particles can be placed
in numerical models to simulate a realistic basalt sample. With these models, it might
ultimately be possible to investigate the effect of resolution limitations of MicroCT leading to the presence of ghost particles in MMT analyses.

¹¹⁹ 2 Grain-size distribution

Pivotal in obtaining a realistic grain-size distribution is acquiring dimensional properties of all iron-oxide grains present in basaltic rocks. The sample we have used in this

study was taken from the 1907 lava flow from Hawaii (HW-03; de Groot et al., 2013; ter 122 Maat et al., 2018)). The MicroCT data of this sample had already been studied by de 123 Groot et al. (2021) and Kosters et al. (2023), facilitating the detection of grains larger 124 than $\sim 1 \ \mu m$, but leaving most of the SD and PSD particles undetected. Fortunately, with 125 the help of FIB-SEM (e.g. Einsle et al., 2016; Nikolaisen et al., 2020) these remaining 126 SD and PSD grains could finally be made visible. FIB-SEM allows a pixel size of 10 nm 127 and a field of view close to 20 by 20 μ m, which enables the detection of both the small-128 est SD grains above the superparamagnetic threshold of 40 nm (Dunlop & Özdemir, 1997) 129 and large MD grains of $>1 \ \mu m$, which are normally at the lower detection limit of Mi-130 croCT. 131

132

2.1 Image acquisition with FIB-SEM

We applied FIB-SEM on a sample of HW-03, following the slice-and-view proce-133 dure described in Einsle et al. (2016) and Nikolaisen et al. (2020). A 60 μ m thick slice 134 with a diameter of 2 mm was extracted from the specimen and coated with a nanome-135 ter thick layer of gold. Around the edges of the sample, silver was applied with a tooth-136 pick so that electrons could easily pass through the sample. The sample was placed in 137 a Zeiss Crossbeam 540 after which we searched the top of the sample for a 20 x 20 μ m 138 area containing many small iron-oxides. This area was prepared for the slice-and-view 139 procedure by first applying a 1 μ m thick layer of platinum while maintaining an ion beam 140 current of 700 pA and an accelerating voltage of 30 kV. This accelerating voltage did 141 not change unless explicitly stated otherwise. 142

Then, five 0.5 µm deep fiducials were engraved with a current of 50 pA (Fig. 1). 143 The three central fiducials were created parallel to each other to enable image alignment 144 during data processing. The outer two fiducials were created under an angle of 25 de-145 grees with respect to the three central fiducials. These two fiducials ensured that the real 146 thickness of the individual image slices could later be checked. To make these fiducials 147 visible, a 0.6 μ m thick carbon layer was deposited on top of the platinum layer with an 148 ion beam current of 700 pA and a dwell-time of 400 seconds. Then, three trenches with 149 a depth of 25 μ m and an area of 27 by 46 μ m² were eroded at three sides of the selected 150 area. For this erosion procedure, we used an ion beam current kept at 15 nA with a back-151 and-forth milling pattern (bidirectional). Using the same bidirectional routine, the cur-152 rent was reduced from 15 nA to 7 nA to create smooth trenches next to the studied area. 153

Instrumental drift during measurements was reduced to a minimum by applying a cross correlation mark outside the studied area (Fig. 1).

The data acquisition process was initiated after a 20.18 nm slice of the sample area 156 was removed with FIB (beam current at 1.5 nA with an accelerating voltage between 157 1 to 6 kV). Backscattered and secondary reflected electrons images where obtained us-158 ing SEM, after which a 20.18 nm layer of sample was removed using FIB. After this pro-159 cedure was repeated a hundred times instrumental drift was corrected by checking align-160 ment with the cross-correlation mark. The entire process was finished in one day, in which 161 601 grayscale images of both backscattered and secondary reflected electrons were saved 162 in TIFF format. Each file encompassed a pixel area of 3072 x 2104 pixels with a pixel 163 size of 10.09 nm and a 8-bit grayscale. The total volume of sample analysed was there-164 fore 31.0 x 21.2 x 12.1 μm^3 . 165

166

2.2 Dataprocessing

167

2.2.1 Extracting particles from FIB-SEM data

After image acquisition, a grain-size distribution could be extracted from the data. Every data processing step was executed with Python 3 and the OpenCV library (Bradski, 2000). We initiated our data processing by loading the 601 backscattered electron images and aligning them using the three parallel fiducials and the matchTemplate function of OpenCV. After alignment, we removed all duplicate images. To remove the effects of both curtaining and the platina and carbon deposited layers, every image was cropped to 2600×1000 pixels.

Through denoising and thresholding, SD and PSD iron-oxides could be extracted 175 from FIB-SEM images as shown by Nikolaisen et al. (2020). However, we designed the 176 extraction process in such a way that both the surface and the centre of each iron-oxide 177 grain were properly retrieved. We initiated denoising by 2×2 binning of the 549 remain-178 ing images, which resulted into 1300×500 pixels per image. A non-local means denois-179 ing filter from OpenCV (fastNlMeansDenoisingMulti) was used with a search window-180 size of 20 pixels, a filter strength of 20, and template windowsize of 3 pixels applied in 181 the x, y, and z-direction. Each image was then thresholded at a pixelvalue of 220. Af-182 ter thresholding, all images are stacked together into one 3D-volume. In this 3D-volume, 183

-8-

we grouped connected pixels into grains. With this high threshold, we have only obtained
the brightest central iron-oxide pixels.

However, the iron-oxide grains are not properly imaged yet, because the sides of 186 the grains frequently have a lower pixel value than their center and thus have dropped 187 out in the previous thresholding step. To correctly image the boundary of the iron-oxides 188 we applied a lower threshold value of 160 instead of 220 to the previous set of images. 189 Then we applied the canny edge algorithm of OpenCV (Canny). This algorithm detects 190 spatial changes in pixel value and draws, based on that change, boundaries between pixel 191 values. After drawing boundaries, small gaps in the boundaries were closed using the mor-192 phology close operation in order to create enclosed areas. Again, all images were com-193 bined into a 3D-volume and connected pixels were grouped into grains. However, the low 194 threshold outputted both well defined grain shapes of iron-oxides and poorly defined shapes 195 of other unwanted minerals. To get rid of these unwanted minerals, we mapped the grains 196 obtained with high thresholding (poorly defined boundaries, most certainly iron-oxides) 197 on top of the grains obtained with low thresholding (well defined boundaries, possible 198 iron-oxides) to only remain with actual iron-oxides. Lastly, we removed all iron-oxides 199 that extend beyond the edges of the FIB-SEM domain. We calculated the volume of 1,558 200 remaining iron-oxides, and transformed these volumes into an equivalent diameter as-201 suming spherical grains. 202

203

2.2.2 Scaling FIB-SEM and MicroCT data

To obtain a grain-size distribution spanning the entire range of naturally occur-204 ring iron-oxides, we combined our FIB-SEM data with MicroCT data that was obtained 205 on a sister specimen of HW-03 by Kosters et al. (2023). Since iron-oxides are not homo-206 geneously distributed throughout the rock, it is not straightforward to link MicroCT data 207 one-to-one to FIB-SEM data based on volume. Most iron-oxides are positioned on the 208 interface between larger grains (mainly plagioclase and pyroxene); especially the small-209 est iron-oxides attach themselves to this 'honeycomb' structure of other minerals (Fig. 210 2). Because this honeycomb structure of relatively large minerals can be imaged prop-211 erly using MicroCT, it allows for extrapolating the number of undetected small iron-oxides 212 from the FIB-SEM analysis to the MicroCT data and hence produce a continuous grain-213 size distribution over both analyses. 214



Figure 2. Overview of the $26.2 \times 10.1 \times 10.5 \ \mu m^3$ volume exposed to the slice-and-view procedure with FIB-SEM. The large grains are indicated by a grayish color. The iron-oxides residing in this volume are indicated in yellow-brown. Note that most iron-oxides reside on the large mineral interfaces.

Characterizing the mineral surface area of the larger minerals, mainly plagioclase 215 and pyroxene, required reprocessing the FIB-SEM data, with this specific goal in mind. 216 First, the FIB-SEM dataset was denoised using a non-local means filter; then we applied 217 a median filter with a spherical footprint of 13 pixels in diameter to smooth out irreg-218 ularities at the boundaries. After this pre-processing we applied a K-means algorithm 219 to create five groups of pixels. For each individual image, we removed the first group be-220 cause it corresponded to drilling artifacts. Additionally, we removed all areas smaller than 221 3000 pixels, because these areas do not represent grain interfaces of the larger minerals 222 present in the sample that are typically occupied by iron-oxides. Finally, the mineral sur-223 face area was calculated for the remaining pixel values using a 'Lindblad' algorithm (Lindblad, 224 2005). The mineral surface area consisted of about 3.8×10^6 squared pixels that cor-225 respond to an area of $1.53 \times 10^{-3} \text{ mm}^2$. Fig. 2 shows the position of the mineral sur-226 face area together with the iron-oxides. 227

To calculate the mineral surface area in the MicroCT dataset, we started with raw MicroCT data and processed it using Dragonfly software. After loading the DICOM files, we cropped about 100 μ m on both sides of the sample to remove scanning residuals. To align the sample with the coordinate axes, we rotated the sample 0.85° clockwise around

-10-

the z-axis, 0.50° anticlockwise around the x-axis, and resampled the image with 0.75 μ m 232 pixel size. Then the image was cropped further, which resulted into a region of 1979 \times 233 44×1901 voxels. We denoised the image by applying a non-local means filter with a 234 3D spherical kernel size of 11 pixels. To create a larger contrast we applied a median fil-235 ter with a 3D spherical kernel size of 13 pixels. The different minerals were separated 236 by applying a K-means filter of 3 groups, of which the group with the highest pixel value 237 corresponds to iron-oxides. The other two groups chiefly consisted of plagioclase and py-238 roxene. Since we observed that SD and PSD iron-oxides tend to cluster on the interfaces 239 of these minerals, we applied a 'sobel edge detection' filter to these groups. The result 240 of this filter enabled us to create a honeycomb structure of the sample, on which we could 241 populate SD and PSD iron-oxides to create realistic models of basalts. This mineral sur-242 face area of 4.00 mm^2 is approximately 2,620 times larger than the surface area found 243 in the FIB-SEM dataset. This means that we need to multiply the number of FIB-SEM 244 grains by 2,620 to scale correctly to the MicroCT dataset. This resulted into a combined 245 dataset of 4.7×10^6 FIB-SEM and 1.6×10^3 MicroCT iron-oxide grains. The FIB-SEM 246 and MicroCT datasets can now be combined to produce a continuous grain-size distri-247 bution of iron-oxides for the entire range of interest (Fig. 3). The median grain-size has 248 a diameter of 70 nm. From 30 to 70 nm the occurrence rapidly increases, after which the 249 occurrence of grain diameters between 70 nm and 10 μ m follows are more gradually de-250 caying trend. 251

252

2.3 Constructing the grain-size distribution

The trend in grain-size distribution can be described using a lognormal-like distri-253 bution between 20 nm and 10 μ m. We have chosen the lognormal distribution, because 254 this distribution is frequently used to describe grain-sizes for e.g. magnetite powders (Smirnov, 255 2006; Yu et al., 2002). Nevertheless, the lognormal trend line was originally fitted to the 256 10th-logarithm of the data. This logarithm ensures a better fit to grain-sizes larger than 257 70 nm. Through an iterative procedure reducing the least-squares error of our fit to the 258 data, we found the best fitting distribution after which we transformed it back to lin-259 ear space: 260

$$p = 3.51 * 10^{-6 + \frac{5.61}{d\sqrt{2\pi}}e^{-\frac{\ln^2(0.60d)}{6.29}}} \text{ if } d : [0, 10]$$
(1)



Figure 3. Overview of the non-normalized grain-size distribution. The histograms show the number of iron-oxides obtained from FIB-SEM or MicroCT data. The datasets are scaled to each other using mineral surface area (see main text). The logarithmic trend line shows the fit to these scaled distributions.

with d the diameter of the grain in μ m and p the occurrence or probability of finding that grain-size. Note that this probability density function, or pdf, is no longer a lognormal distribution, yet still accurately describes the data (Fig. 3). The function is only valid for grain-sizes between 0 and 10 μ m, although it could be extended to larger grain-sizes through rescaling.

To obtain a function from which grain-sizes can be sampled, the probability den-266 sity function is integrated into a cumulative probability density function (cdf). This func-267 tion takes a grain-size as input and returns the probability of finding that grain-size or 268 smaller as output. Then to actually create a sampling function, this cdf has to be inverted. 269 By inverting the function, a grain-size could be obtained as function of inputted cumu-270 lative probability. However, this inverted function has no analytical expression, so we 271 created a look-up table to be able to sample the function. To create this table we inputted 272 100,000 diameters between 0 and 10 μ m into the cdf, which returned 100,000 cumula-273 tive probabilities. Each grain-size-probability pair was then put into the table. To sam-274 ple grain-sizes from this look-up table we would generate a pool of random numbers, ob-275

-12-

tained from a uniform distribution between 0 and 1 (U(0,1)), equal to the number of samples requested. For each random number we would then extract the corresponding grain-size from the look-up table. Grain-sizes or probabilities not present in the table would be obtained through cubic interpolation of the nearest values. All Python codes to create and sample the grain-size distribution, and the look-up table are available in our repository (Out et al., 2023).

²⁸² 3 Discussion

283

3.1 Dataset limitations and validation

For this study only one FIB-SEM sample was used from a specific, but typical Hawai-284 ian lava flow. How well this sample represents the sample material of site HW-03, let alone 285 material from other sites, flows, and/or locations is a big unknown. Obtaining data from 286 rock samples using FIB-SEM, however, is not a straightforward task; sample prepara-287 tion, machine handling, and availability of resources complicate the scanning process. 288 To get the most out of our data, we decided during initial phase of the FIB-SEM anal-289 yses to target an area with the highest amount of iron-oxides in view. This implies that 290 relatively many small grains will be present in our FIB-SEM study, leading to an over-291 representation of small grains in our final grain-size distribution. 292

The overestimation, however, might be somewhat damped through how we con-203 structed the scaling factor. Our scaling factor is dependent on the amount of mineral 294 surface area which probably scales with the number of grains. This might mean that our 295 scaling factor is less sensitive to grain density variations in the studied FIB-SEM area, 296 than, for example, a scaling factor based on volume. In case of volume scaling, the FIB-297 SEM area would be scaled by a factor of 24,000, independent of the number of grains 298 or amount of mineral surface area present in the studied area. This would most defini-299 tively result in a severe overestimation of particles imaged by FIB-SEM compared to par-300 ticles imaged by MicroCT, because this volume scaling factor is one order of magnitude 301 larger than our scaling factor. We, therefore, consider scaling by grains per mineral in-302 terface area a stable protocol to ensure comparability between FIB-SEM and MicroCT 303 results as it dampens local variations of the studied area. 304

To obtain an independent verification for our method of combining FIB-SEM and MicroCT analyses, we validated our dataset through a scanning electron microscopy (SEM)

-13-

image on a sister specimen from site HW03 (de Groot et al., 2013; ter Maat et al., 2018). Two areas of $48.2 \times 71.3 \ \mu\text{m}^2$ and $33.2 \times 49.0 \ \mu\text{m}^2$ within the sample were imaged with a resolution of 80 and 55 nm per 8-bit pixel respectively. After noise was removed, the images were thresholded at a pixel value of 153 to only retain iron-oxides. Then the major axis *m* and minor axis *n* of each grain were determined and grain diameters, *d*, were calculated using (Yu et al., 2002):

$$d = 2\sqrt{m\frac{n}{\pi}} \tag{2}$$

The diameters of the grains were sorted to produce a normalized histogram with bin-313 sizes of 0.1 μ m ranging between 0.2 and 3.0 μ m - straddling the transition from the FIB-314 SEM realm to the MicroCT range. We discarded grains smaller than 0.2 μ m from the 315 SEM data because they would be constructed from two pixels or less. On top of the SEM 316 grain-size data, we overlaid the FIB-SEM and MicroCT normalized data (Fig. 4). The 317 SEM data misses some of the smaller grains, as expected. The trends in the FIB-SEM+MicroCT 318 and SEM datasets for grain-sizes $>0.3 \mu m$, however, correspond well. Taking into ac-319 count that SEM and FIB-SEM are very different scanning techniques compared to Mi-320 croCT, we are confident that our constructed grain-size distribution is a proper descrip-321 tion of the actual grain-size distribution of iron-oxides in our samples, although it may 322 emphasize small, sub-micron, grains. 323

324

3.2 Building models of the physical distribution of iron-oxides

MMT struggles with the smallest particles as they are not detected by MicroCT acquisitions. Although these grains are invisible for MicroCTs, they might still produce a signal in a magnetic surface scan, possibly leading to incorrect magnetic moments of other detected particles (Fabian & De Groot, 2019; Out et al., 2022). However, with the grain-size distribution presented here, it is now possible to study the effect of these ghost grains through generating realistic grain models of basalt and simulating their effect on MMT results in future studies.

To generate a realistic model of the physical distribution of iron-oxide grains in a basaltic sample, we aim to populate a honeycomb structure of the larger minerals obtained from MicroCT analyses with a simulated distribution of small iron-oxides that the MicroCT would miss (Fig. 5). The first step is to extract a honeycomb structure and

-14-



Figure 4. Verification of MicroCT and FIB-SEM data using separate SEM imaged iron-oxides using normalized histograms. Note that the distribution is similar for the two data-series except for the smallest grain-sizes.

the large iron-oxides from a scanned basaltic sample and determine the resolution of the 336 MicroCT. The MicroCT resolution determines how many ghost grains are missing per 337 mm^2 honeycomb structure. By multiplying this ghost grain density (Figure 6) with the 338 honeycomb structure surface area extracted from MicroCT images, we can calculate the 339 total number of missing grains. Then we use the grain-size distribution obtained in this 340 study to make a list of the missing ghost grains and their diameters by feeding the num-341 ber of missing ghost grains together with the MicroCT resolution into the grain-size dis-342 tribution sampling function (available in the repository of Out et al., 2023). Finally, we 343 simulate grains from the list of ghost grains and place them on the honeycomb struc-344 ture obtained from the MicroCT analysis besides the larger iron-oxide particles that were 345 already detected. Hence, by extracting the honeycomb structure and MicroCT resolu-346 tion, we are able to simulate a realistic distribution of all iron-oxides in a basaltic sam-347 ple, with diameters down to the superparamagnetic range. 348

As an example we consider a hypothetical MicroCT analysis with a resolution of $0.1 \ \mu m$ that produced a honeycomb structure with a surface area of 4 mm². Then we would need to populate this honeycomb structure with 9×10^5 particles per mm² (dashed line in Fig. 6). This implies that we miss 3.6×10^6 ghost grains in the entire sample with

-15-



Figure 5. Overview of the protocol to build realistic models of the physical distribution of iron-oxides in a basalt. By segmenting the MicroCT image, we only retain the honeycomb structure and the visible iron-oxides. The MicroCT resolution is required as input for the relation to sample the grain-size distribution of iron-oxides (Fig. 6). Multiplying the ghost grain density with the honeycomb surface area extracted from the MicroCT image, we are able to determine how many grains are left undetected by MicroCT. Given the MicroCT resolution and the number of ghost grains, we sample the grain-size distribution function (see Out et al., 2023) to obtain a list of grain diameters. These grain diameters are used to construct particles that are randomly placed on the extracted honeycomb structure. A graphical overview can be found in the supplementary information (Figure S1).



Figure 6. Maximum probability and particles per mm² mineral interface area as function of MicroCT resolution. Based on this figure we can determine how many grains have to be sampled from the provided grain-size distribution. By providing the maximum cumulative probability to the grain-size distribution, we can select the maximum grain-size that should be sampled. For example, for 4 mm² mineral interface area and a MicroCT resolution of 0.1 μ m, we would sample 900.000 particles/mm² × 4 mm² = 3.6 million particles with a maximum cumulative probability of 0.57. For a resolution larger than 1.0 μ m, we would sample 1.20 million grains per mm² mineral interface area.

a surface area of 4 mm². To obtain the list of grains that we need to add to the honeycomb structure we sample the grain-size distribution 3.6 million times using a maximum cumulative probability of 0.57 (dashed line in Fig. 6) for grains with a diameter between 0 and 0.1 μ m. These grains are then randomly placed on the honeycomb structure together next to the larger iron-oxide particles that were already detected with MicroCT.

358

3.3 Implications for Micromagnetic Tomography

The results of previous MMT based on synthetic samples in which the grain size 359 distribution was controlled (e.g. (most) results in de Groot et al., 2018, 2021)) are most 360 likely not disturbed by ghost grains, because the iron-oxide grains have been sieved for 361 a size larger than the MicroCT resolution. For that reason, their sample mainly consisted 362 of particles larger than 5 μ m, which could be easily detected by MicroCT. However, the 363 natural basalt samples used by de Groot et al. (2021) and Kosters et al. (2023) undoubt-364 edly contain undetected ghost grains. Both articles study the magnetic moment of 1,646 365 iron-oxides detected by MicroCT in a 30 μ m thin section of Hawaiian basaltic rock, which 366 is a sister-specimen as subjected to our study. The resolution of the MicroCT used in 367 those studies is $\sim 1 \ \mu m$ and the mineral surface area is $\sim 4 \ mm^2$. According to Fig. 6, 368 this implies the existence of 4.8×10^6 undetected particles, which results in ~3,000 ghost 369 grains per detected iron-oxide. 370

All of these ghost grains might scramble the signal of the detected iron-oxides lead-371 ing to erroneous magnetic moments for the detected grains produced by the MMT in-372 version. Remarkably, both studies do not present unrealistic magnetic moments for the 373 detected iron-oxide particles, and their results are (somewhat) interpretable in terms of 374 magnetic moment and direction. Also, reported trends as function of grain-size and mag-375 netic treatment follow expectations based on earlier work (e.g. Dunlop & Özdemir, 1997). 376 It is therefore likely that, although high in number, the small ghost grains do not have 377 a large enough magnetic moment to have a major contribution to magnetic signal on the 378 surface of the sample. Another possibility is that the spatial resolution of the surface mag-379 netometry data might not be sufficient to locate the weak signal of these ghost particles. 380 As its magnetic moment is highly dependent on the size of a grain (Dunlop & Ozdemir, 381 1997) it is likely that the magnetic moments of larger grains are less affected by ghost 382 grains in an MMT inversion, compared to grains that are just above the detection thresh-383 old of the MicroCT analysis used. To characterise the reliability of MMT results as func-384

-18-

tion of grain size, a future study that models the impact of the undetected ghost grains on MMT results is critical.

387 4 Conclusions

Here we produced a grain-size distribution of all iron-oxides larger than the super-388 paramagnetic threshold in an Hawaiian basalt by combining MicroCT and FIB-SEM data 389 of sister-specimens. The MicroCT and FIB-SEM datasets could not be combined one-390 to-one due to volume difference, so they were scaled using the interface area of large, non-391 iron-oxide, minerals present in the sample (i.e. honeycomb structure), because iron-oxides 392 tend to cluster on the edges of large minerals. Through this scaling procedure we have 393 created a lognormal-like distribution between 20 nm to 10 μ m that spans the range of 394 SD, PSD, and MD iron-oxides. This distribution is probably a slight overestimation of 395 the concentration of iron-oxide grains, caused by the choice of FIB-SEM study area with 396 relatively many iron-oxides. This grain-size distribution can be used to populate real-397 istic models of the physical distribution of iron-oxides in a sample; and makes it possi-398 ble to assess the impact of undetected ghost grains on micromagnetic and MMT stud-399 ies. Due to the high demands of both MicroCT and FIB-SEM analyses we could only 400 process one sample from a well-characterised Hawaiian lava flow in this study. Future 401 studies can build on our theory of combining FIB-SEM and MicroCT data through the 402 surface area of the honeycomb structure of larger minerals produced by the MicroCT anal-403 ysis to study other sample material as well. 404

405 Data availability statement

The FIB-SEM, MicroCT, and SEM data used for obtaining and verifying, and the 406 jupyter notebooks employed for generating the grain-size distribution and honeycomb 407 structure in this study are available at YODA via https://doi.org/10.24416/UU01-QNASXB 408 with a CC-BY-NC-4.0 license (Out et al., 2023). Numerical calculations were executed 409 with support of NumPy (Harris et al., 2020), SciPy (Virtanen et al., 2020), OpenCV (Bradski, 410 2000), and Matplotlib (Hunter, 2007) Python libraries. The MicroCT analysis for this 411 study was generated using Dragonfly software, Version 2022.1.0.1259 for Windows. Ob-412 ject Research Systems (ORS) Inc, Montreal, Canada, 2022. 413

-19-

414 Acknowledgments

- ⁴¹⁵ This project has received funding from the European Research Council (ERC) under the
- ⁴¹⁶ European Union's Horizon 2020 research and innovation program (Grant agreement No.
- ⁴¹⁷ 851460 to L.V. de Groot). This publication results from work carried out under transna-
- ⁴¹⁸ tional access action under the support of EXCITE EC- HORIZON 2020 INFRAIA
- ⁴¹⁹ 2020 Integrating Activities for Starting Communities (Grant agreement No. 101005611).

420 **References**

445

- Berndt, T., Muxworthy, A. R., & Fabian, K. (2016). Does size matter? statistical
 limits of paleomagnetic field reconstruction from small rock specimens. Journal
 of Geophysical Research: Solid Earth, 121(1), 15–26.
- ⁴²⁴ Bradski, G. (2000). The OpenCV Library. Dr. Dobb's Journal of Software Tools.
- de Groot, L. V., Biggin, A. J., Dekkers, M. J., Langereis, C. G., & Herrero-Bervera,
- E. (2013, 12). Rapid regional perturbations to the recent global geomagnetic decay revealed by a new Hawaiian record. Nature Communications, 428 4(1), 2727. Retrieved from www.nature.com/naturecommunicationshttp:// 429 www.nature.com/articles/ncomms3727 doi: 10.1038/ncomms3727
- ⁴³⁰ De Groot, L. V., Fabian, K., Bakelaar, I. A., & Dekkers, M. J. (2014). Magnetic
 ⁴³¹ force microscopy reveals meta-stable magnetic domain states that prevent re⁴³² liable absolute palaeointensity experiments. Nature Communications, 5(1),
 ⁴³³ 4548.
- 434 de Groot, L. V., Fabian, K., Béguin, A., Kosters, M. E., Cortés-Ortuño, D., Fu,
- R. R., ... Barnhoorn, A. (2021). Micromagnetic tomography for paleomagnetism and rock-magnetism. Journal of Geophysical Research: Solid Earth, 126(10), e2021JB022364.
- de Groot, L. V., Fabian, K., Béguin, A., Reith, P., Barnhoorn, A., & Hilgenkamp,
 H. (2018). Determining individual particle magnetizations in assemblages of
 micrograins. *Geophysical Research Letters*, 45(7), 2995–3000.
- Dunlop, D. J., & Özdemir, Ö. (1997). Rock Magnetism. Cambridge University Press.
 Retrieved from https://www.cambridge.org/core/product/identifier/
 9780511612794/type/book doi: 10.1017/cbo9780511612794
- Einsle, J. F., Harrison, R. J., Kasama, T., Conbhuí, P. Ó., Fabian, K., Williams,
 - W., ... Midgley, P. A. (2016). Multi-scale three-dimensional charac-

446	terization of iron particles in dusty olivine: Implications for paleomag-
447	netism of chondritic meteorites. American Mineralogist, 101(9), 2070–
448	2084. Retrieved from https://doi.org/10.2138/am-2016-5738CCBY doi:
449	doi:10.2138/am-2016-5738CCBY
450	Fabian, K., & De Groot, L. V. (2019, 2). A uniqueness theorem for tomography-
451	assisted potential-field inversion. $Geophysical Journal International, 216(2),$
452	760-766. Retrieved from http://arxiv.org/abs/1712.06136https://
453	academic.oup.com/gji/article/216/2/760/5151336 doi: 10.1093/gji/
454	ggy455
455	Glenn, D. R., Fu, R. R., Kehayias, P., Le Sage, D., Lima, E. A., Weiss, B. P., &
456	Walsworth, R. L. (2017, 8). Micrometer-scale magnetic imaging of geological
457	samples using a quantum diamond microscope. Geochemistry, Geophysics,
458	Geosystems, 18(8), 3254–3267. Retrieved from https://onlinelibrary
459	.wiley.com/doi/abs/10.1002/2017GC006946 doi: 10.1002/2017GC006946
460	Harris, C. R., Millman, K. J., van der Walt, S. J., Gommers, R., Virtanen, P., Cour-
461	napeau, D., Oliphant, T. E. (2020, 9). Array programming with NumPy.
462	Nature, 585(7825), 357-362. Retrieved from http://www.nature.com/
463	articles/s41586-020-2649-2 doi: 10.1038/s41586-020-2649-2
464	Hunter, J. D. (2007, 5). Matplotlib: A 2D Graphics Environment. Com-
465	puting in Science & Engineering, $9(3)$, 90–95. Retrieved from http://
466	<pre>nipy.scipy.http://ieeexplore.ieee.org/document/4160265/ doi:</pre>
467	10.1109/MCSE.2007.55
468	Kosters, M. E., de Boer, R. A., Out, F., Cortés-Ortuño, D. I., & de Groot, L. V.
469	(2023). Unraveling the magnetic signal of individual grains in a hawaiian lava
470	using micromagnetic tomography. Geochemistry, Geophysics, Geosystems,
471	24(4), e2022GC010462. Retrieved from https://agupubs.onlinelibrary
472	.wiley.com/doi/abs/10.1029/2022GC010462 (e2022GC010462
473	2022GC010462) doi: https://doi.org/10.1029/2022GC010462
474	Lindblad, J. (2005). Surface area estimation of digitized 3d objects using weighted
475	local configurations. Image and Vision Computing, $23(2)$, $111-122$.
476	Nagy, L., Williams, W., Muxworthy, A. R., Fabian, K., Almeida, T. P., Conbhui,
477	P. O., & Shcherbakov, V. P. (2017). Stability of equidimensional pseudo-single-
478	domain magnetite over billion-year timescales. Proceedings of the National

manuscript submitted to *G*-cubed

479	Academy of Sciences of the United States of America, 114(39), 10356-10360.
480	doi: 10.1073/pnas.1708344114
481	Néel, L. (1955, 4). Some theoretical aspects of rock-magnetism. Advances in
482	Physics, 4(14), 191-243. Retrieved from https://www.tandfonline.com/
483	doi/abs/10.1080/00018735500101204http://www.tandfonline.com/doi/
484	abs/10.1080/00018735500101204 doi: $10.1080/00018735500101204$
485	Nikolaisen, E. S., Harrison, R. J., Fabian, K., & McEnroe, S. A. (2020). Hysteresis
486	of natural magnetite ensembles: Micromagnetics of silicate-hosted magnetite
487	inclusions based on focused-ion-beam nanotomography. Geochemistry, Geo-
488	physics, Geosystems, $21(11)$, e2020GC009389.
489	Out, F., Cortés-Ortuñno, D., Fabian, K., van Leeuwen, T., & de Groot, L. V.
490	(2022). A first-order statistical exploration of the mathematical limits of
491	micromagnetic tomography. $Geochemistry, Geophysics, Geosystems.$ doi:
492	e2021GC010184
493	Out, F., de Boer, R. A., Walmsley, J., & de Groot, L. V. (2023). Replication
494	Data for: Modeling the Distribution of Iron-oxides in Basalt by combining
495	FIB-SEM and MicroCT Measurements. YODA. Retrieved from https://
496	public.yoda.uu.nl/geo/UU01/QNASXB.html doi: 10.24416/UU01-QNASXB
497	Smirnov, A. (2006). Memory of the magnetic field applied during cooling in the low- $% \left(2006\right) \left(1-2000\right) \left(1-20000\right) \left(1-2000\right) \left(1-2000\right) \left(1-20000\right) \left(1-2000\right) \left($
498	temperature phase of magnetite: Grain size dependence. Journal of Geophysi-
499	cal Research: Solid Earth, 111(B12).
500	Tauxe, L. (2010). Essentials of paleomagnetism. Univ of California Press.
501	Tauxe, L., & Yamazaki, T. (2015). Paleointensities. In <i>Treatise on geophysics</i> (pp.
502	461–509). Elsevier. doi: 10.1016/B978-0-444-53802-4.00107-X
503	ter Maat, G. W., Pennock, G. M., & de Groot, L. V. (2018). Data descrip-
504	tor: A chemical, crystallographic and magnetic characterisation of indi-
505	vidual iron-oxide grains in Hawaiian lavas. Scientific Data, 5, 1–9. doi:
506	10.1038/sdata.2018.162
507	Virtanen, P., Gommers, R., Oliphant, T. E., Haberland, M., Reddy, T., Cournapeau,
508	D., van Mulbregt, P. (2020, 3). Sci Py 1.0: fundamental algorithms for sci-
509	entific computing in Python. Nature Methods, 17(3), 261–272. Retrieved from
510	https://doi.org/10.1038/s41592-019-0686-2http://www.nature.com/
511	articles/s41592-019-0686-2 doi: $10.1038/s41592-019-0686-2$

- Yu, Y., Dunlop, D. J., & Özdemir, Ö. (2002). Partial anhysteretic remanent mag netization in magnetite 1. additivity. Journal of Geophysical Research: Solid
- ⁵¹⁴ Earth, 107(B10), EPM-7.

Supporting Information for "Modeling the Distribution of Iron-oxides in Basalt by combining FIB-SEM and MicroCT Measurements"

Frenk Out¹, Rosa A. de Boer¹, John Walmsley², Lennart V. de Groot¹

¹Paleomagnetic laboratory Fort Hoofddijk, Department of Earth Sciences, Utrecht University, Budapestlaan 17, 3584 CD Utrecht,

The Netherlands.

²Department of Materials Science & Metallurgy, University of Cambridge, 27 Charles Babbage Road, CB3 0FS Cambridge, United

Kingdom.

Contents of this file

1. Figures S1

Corresponding author: F. Out, Paleomagnetic laboratory Fort Hoofddijk, Department of Earth Sciences, Utrecht University, Budapestlaan 17, 3584 CD Utrecht, The Netherlands. (f.out@uu.nl) OUT ET AL.: MODELING THE DISTRIBUTION OF IRON-OXIDES IN BASALTS





Figure S1. Graphical overview of the protocol to build a realistic iron-oxide model. See section3.2 of the main text for an explanation of the protocol.