

# Shallow Tectonic Stress Magnitudes at the Hikurangi Subduction Margin, New Zealand

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December 22, 2022

## Abstract

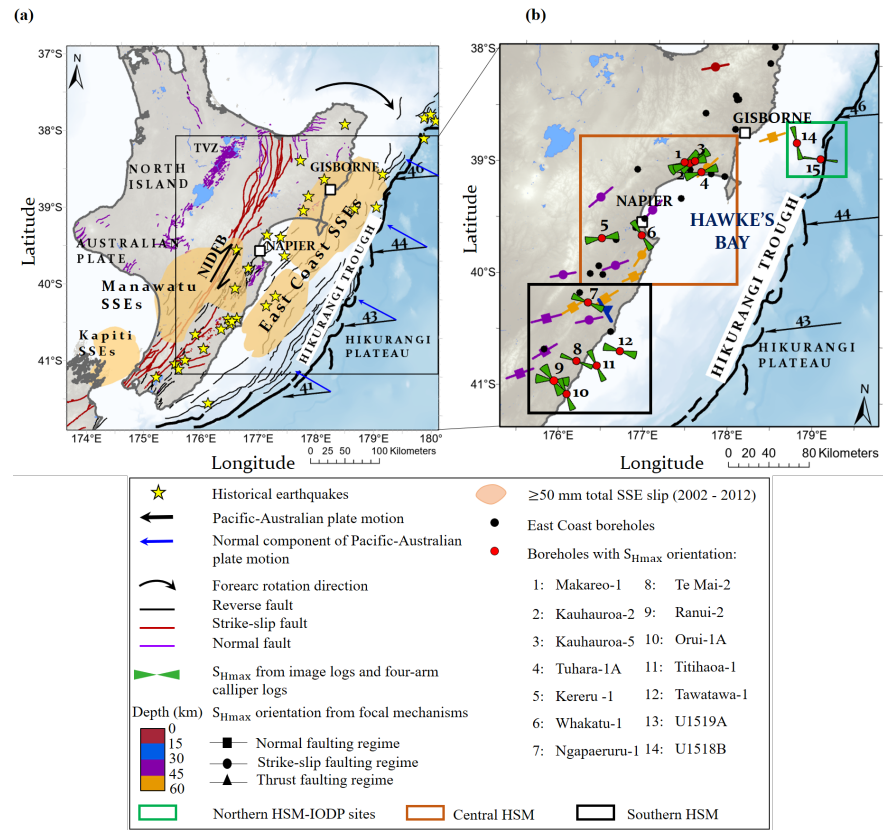
Quantifying tectonic stress magnitudes is crucial in understanding crustal deformation processes, fault geomechanics, and variable plate interface slip behaviors in subduction zones. The Hikurangi Subduction Margin (HSM), New Zealand is characterized by along-strike variation in interface slip behavior, which may be linked to tectonic stress variations within the overriding plate. This study constrains in-situ stress magnitudes of the shallow (<3km) overriding plate of the HSM to better understand its tectonics and how they relate to larger scale subduction dynamics. Results reveal  $\sigma_3$ :  $S_v$  ratios of 0.6-1 at depths above 650-700 m TVD and 0.92-1 below this depth interval along the HSM and  $SH_{max}$ :  $S_v$  ratios of 0.95-1.81 in the central HSM, and 0.95-3.12 in the southern HSM. These stress ratios suggest a prevalent thrust to strike-slip ( $\sigma_1=SH_{max}$ ) faulting regime across the central and southern HSM. In the central HSM, the presence of NNE-NE striking reverse faults co-existing with a modern  $\sigma_1$  aligned ENE-WSW ( $SH_{max}$ ) suggests that overtime the stress state here evolved from a contractional to a strike-slip state, where the compressional direction changes from perpendicular (NW-SE) to subparallel (ENE-WSW) to the Hikurangi margin. This temporal change in stress state may be explained by forearc rotation, likely combined with development of upper plate overpressures. In the southern HSM, the modern WNW-ESE/ NW-SE  $\sigma_1$  ( $SH_{max}$ ) and pre-existing NNE-NE striking reverse faults indicate that stress state remains contractional and subparallel (NW-SE) to the Hikurangi margin overtime. This may reflect the interseismic locked nature of the plate interface.

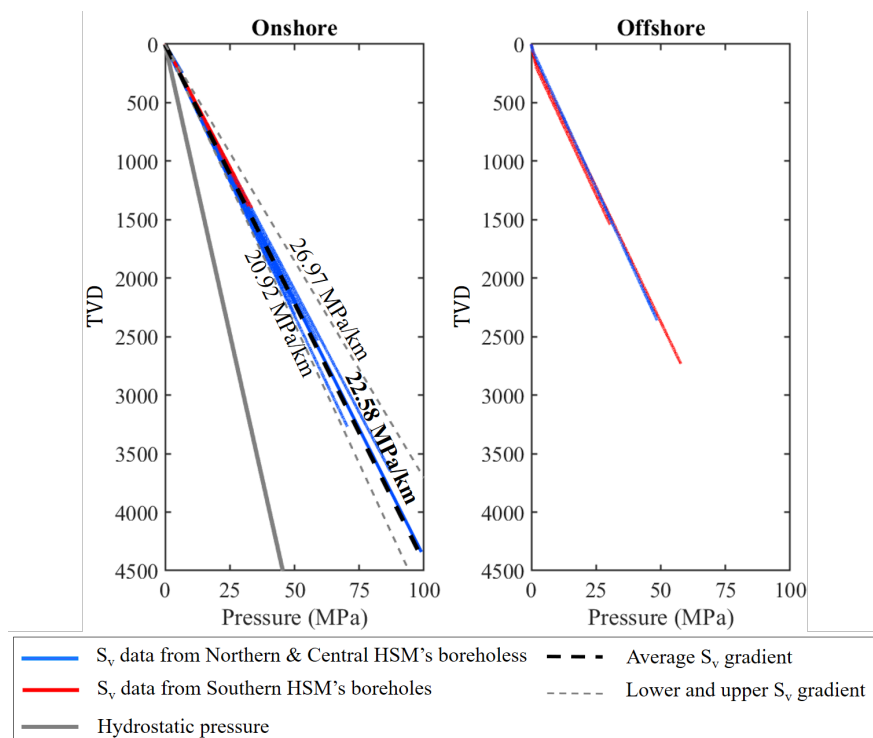
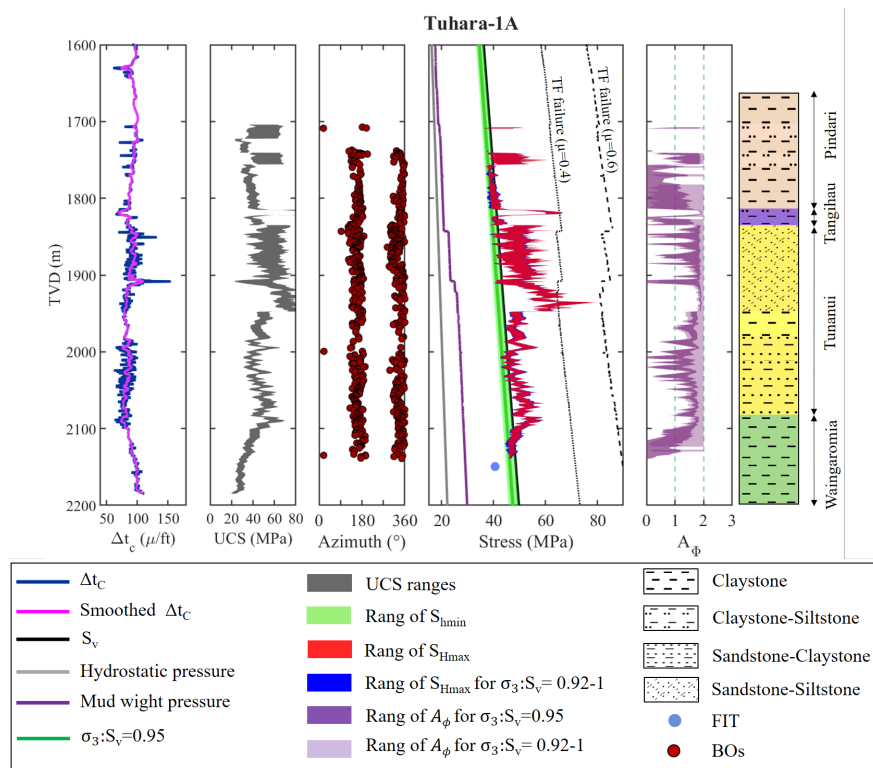
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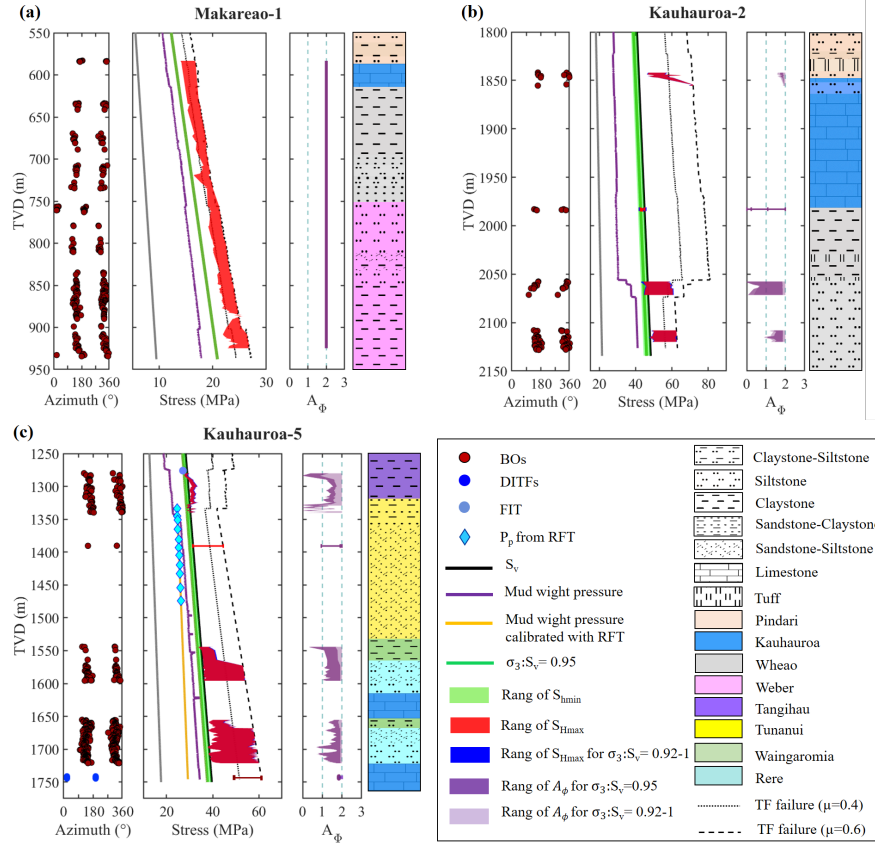
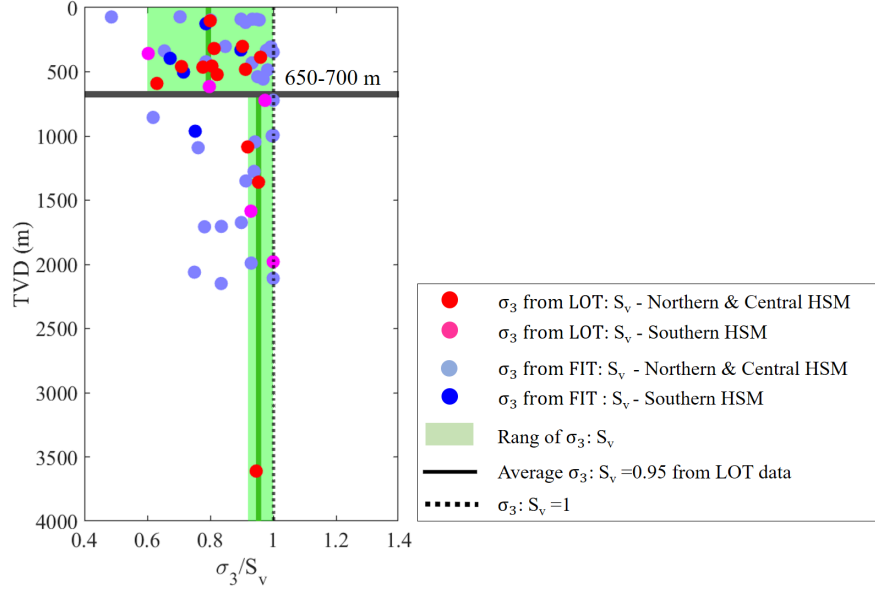
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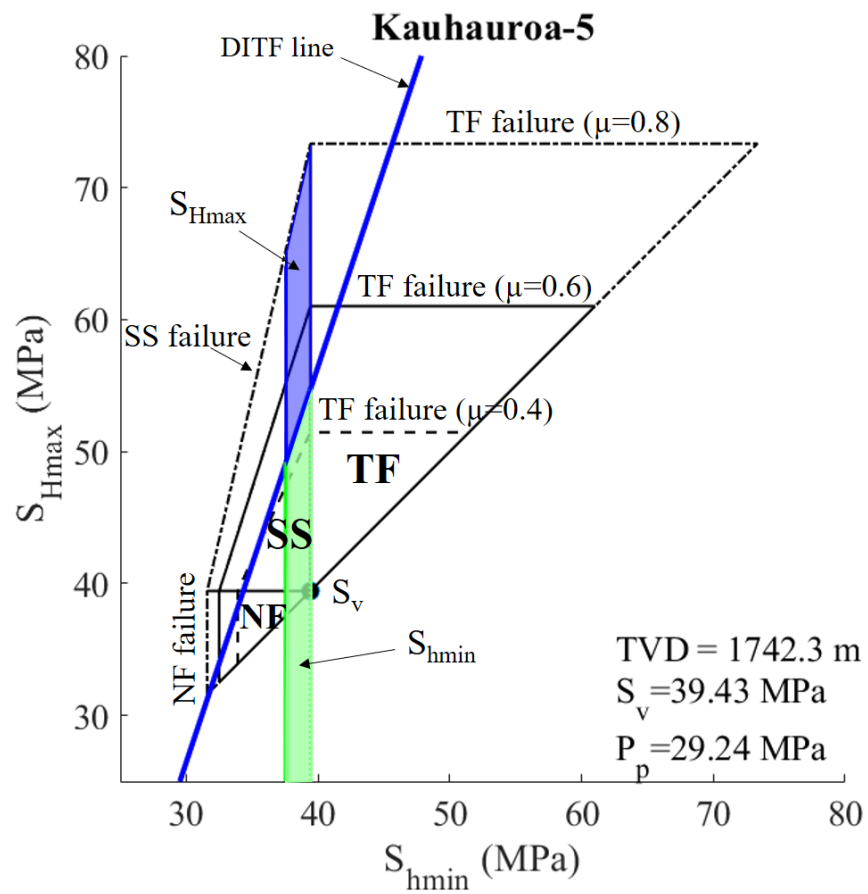
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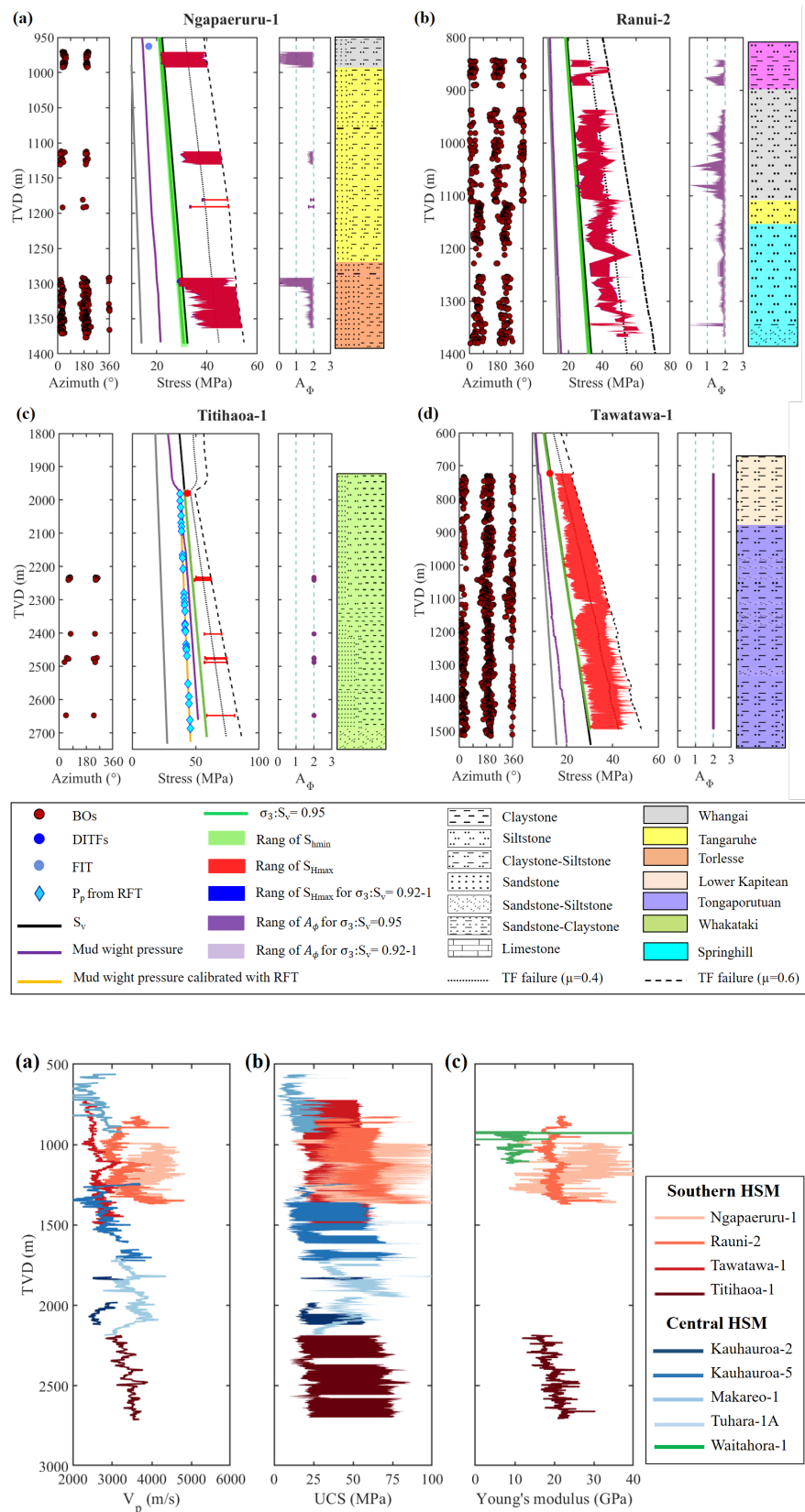


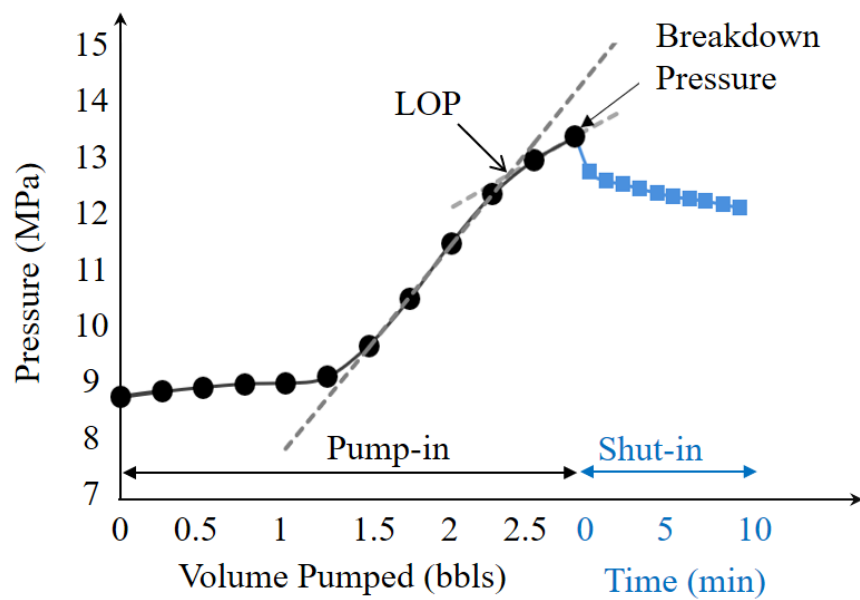


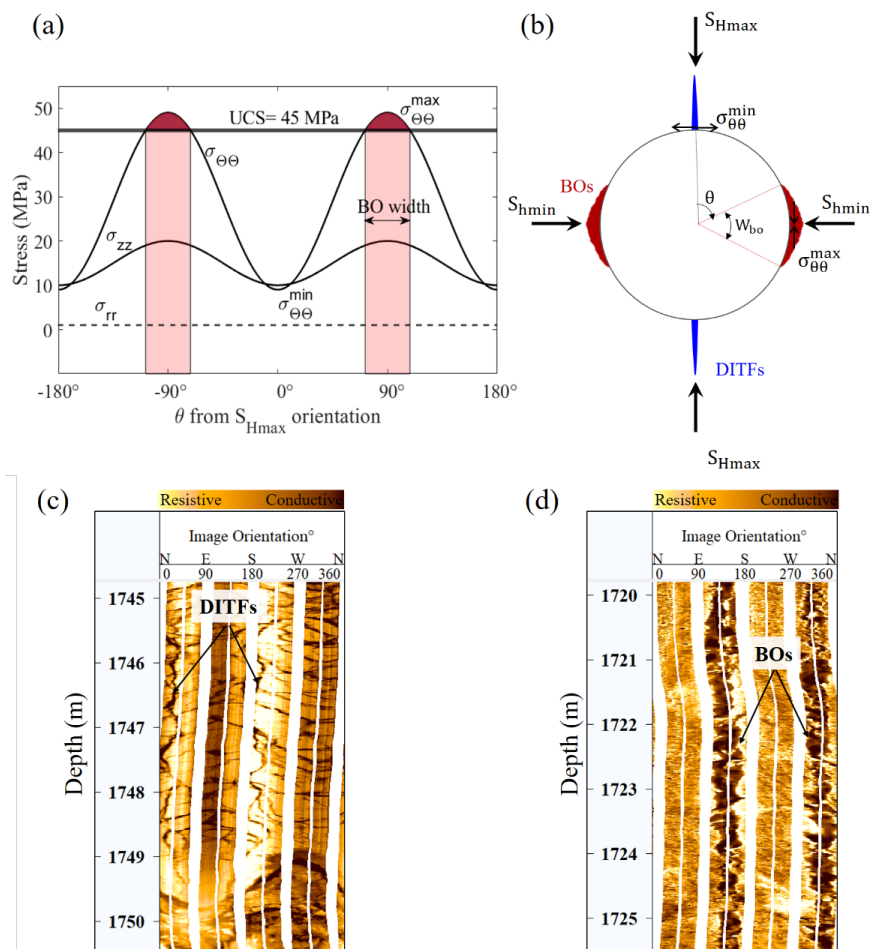












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2                                     **Margin, New Zealand**

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## Abstract

Quantifying tectonic stress magnitudes is crucial in understanding crustal deformation processes, fault geomechanics, and variable plate interface slip behaviors in subduction zones. The Hikurangi Subduction Margin (HSM), New Zealand is characterized by along-strike variation in interface slip behavior, which may be linked to tectonic stress variations within the overriding plate. This study constrains *in-situ* stress magnitudes of the shallow (<3km) overriding plate of the HSM to better understand its tectonics and how they relate to larger scale subduction dynamics. Results reveal  $\sigma_3$ :  $S_v$  ratios of 0.6-1 at depths above 650-700 m TVD and 0.92-1 below this depth interval along the HSM and  $S_{Hmax}$ :  $S_v$  ratios of 0.95-1.81 in the central HSM, and 0.95-3.12 in the southern HSM. These stress ratios suggest a prevalent thrust to strike-slip ( $\sigma_1=S_{Hmax}$ ) faulting regime across the central and southern HSM. In the central HSM, the presence of NNE-NE striking reverse faults co-existing with a modern  $\sigma_1$  aligned ENE-WSW ( $S_{Hmax}$ ) suggests that overtime the stress state here evolved from a contractional to a strike-slip state, where the compressional direction changes from perpendicular (NW-SE) to subparallel (ENE-WSW) to the Hikurangi margin. This temporal change in stress state may be explained by forearc rotation, likely combined with development of upper plate overpressures. In the southern HSM, the modern WNW-ESE/NW-SE  $\sigma_1$  ( $S_{Hmax}$ ) and pre-existing NNE-NE striking reverse faults indicate that stress state remains contractional and subparallel (NW-SE) to the Hikurangi margin overtime. This may reflect the interseismic locked nature of the plate interface.

## Plain Language Summary

The type of geological faults and their movement are partially controlled by forces generated from plate movement, known as in-situ stress. This stress state can also be changed overtime due to the occurrence of earthquakes on such faults. The HSM is New Zealand's largest and most hazardous plate boundary fault and experiences different types of earthquakes that may be related to variations in in-situ stress of the plates involved in this subduction boundary. This study quantifies for the first time the stresses associated with the modern HSM, and finds that they and their resulting tectonic behavior have changed with geological time in the central regions. This change is likely related to the effects of other nearby tectonic processes further inland and to the development of high pore pressures in the overriding plate in this region.

## Key Points

- For the shallow crust (upper 3 km) of the Hikurangi Subduction Margin,  $\sigma_1 = S_{Hmax}$ .
- $\sigma_1$  rotates from margin-parallel (NW-SE) to margin-perpendicular (WNW-ESE) in the central Hikurangi Subduction Margin overtime.
- The shift in the stress state overtime in the central HSM may be driven by forearc rotation and shallow overpressures in this region.
- $\sigma_1$  remains perpendicular (NW-SE/WNW-ESE) to the margin overtime in the southern HSM, may reflect the interseismic locked nature of the plate interface.

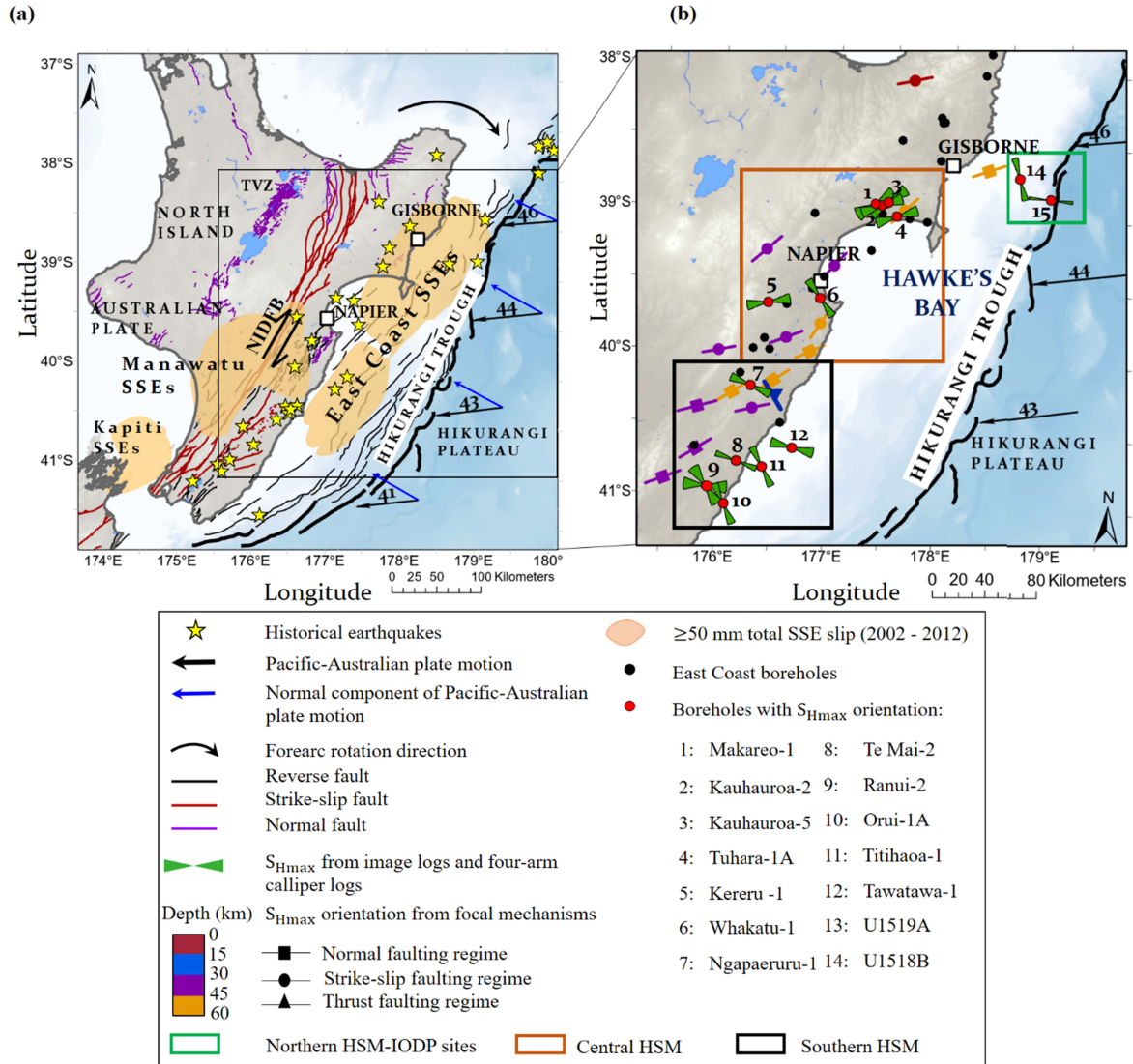
## 1 Introduction

Large magnitude, tsunamigenic earthquakes commonly occur at subduction plate boundaries and are associated with a wide range of tectonic fault slip behaviors along the subduction interface including slow slip events (SSEs), low-frequency earthquakes (LFEs), very-low-frequency earthquakes (VLFEs), and episodic tremor and slip (ETS) (Audet et al., 2009; Ito & Obara, 2006; Kodaira et al., 2004; Liu & Rice, 2007; Ujiie & Kimura, 2014). Earthquake occurrence such as nucleation of earthquake ruptures and rupture propagations, and a variety of seismic slip behaviors are, in part, controlled by the interaction between *in-situ* stresses (their orientations and magnitudes), the mechanical and geometrical properties of crustal faults, and pore pressure (Jaeger et al., 2009; Schellart & Rawlinson, 2013; Vavrycuk, 2015). Furthermore, seismic cycling and slip on faults are known to drive temporal changes in the stress state on adjacent fault planes and surrounding rocks (Brodsky et al., 2017, 2020; Hardebeck & Okada, 2018; K. F. Ma et al., 2005; Seeber & Armbruster, 2000; Stein, 1999). For example, significant principal stress rotations followed the 2011  $M_w$  9.0 Tohoku earthquake in Japan, 2010  $M_w$  8.8 Maule earthquake in Chile; and 2004  $M_w$  9.2 earthquake in Sumatra-Andaman are suggested to be related to near-complete stress drops (Hardebeck, 2012). Therefore, quantitative knowledge of stress is an essential step to characterize and understand the nature and causes of earthquake processes, the mechanical behavior of plate boundary faults, the origin and controls of diverse fault slip patterns; and to better assess seismic and tsunamigenic hazards along subduction zones (Huffman & Saffer, 2016; Riedel et al., 2016; Wu et al., 2019).

The Hikurangi Subduction Margin (HSM), New Zealand displays along-strike variation in plate interface slip behavior, ranging from episodic SSEs and creep at the northern and

98 central HSM, to deep interseismic locking beneath the southern North Island (Wallace &  
99 Beavan, 2010) (Figure 1a). Creep and shallow (<15 km depth) SSEs, lasting for 2–3 weeks,  
100 recur every 1 to 2 years offshore the northern and central HSM (Wallace, Beavan, et al.,  
101 2012) (Figure 1a). Deep (>25 km), long-term (>1 year) SSEs occur approximately every ~5  
102 years at the southern HSM (Wallace & Beavan, 2010), down-dip from a portion of the plate  
103 interface that is locked and accumulating stress (Wallace et al., 2009). The physical processes  
104 controlling SSEs are currently debated, with studies suggesting they are linked to the  
105 frictional properties of fault zone materials (e.g., strength and coefficient of friction), low  
106 effective stress linked to high pore pressure, fault heterogeneity, and fault rheology (Ando et  
107 al., 2012; Kodaira et al., 2004; Kurzwski et al., 2018; Saffer & Wallace, 2015). More than  
108 80% of HSM historic earthquakes and  $M_w \geq 6$  earthquakes occur on upper plate ( $\leq 30$  km)  
109 faults or at the plate interface (Figure 1a) (Doser & Webb, 2003; Downes, 2006; Grapes &  
110 Downes, 1997; Webb & Anderson, 1998). Earthquakes located within the subducting slab or  
111 at the plate interface have also been known to trigger slope failures or series of smaller  
112 earthquakes hosted on upper plate faults, some of which can be tsunamigenic (Beetham et al.,  
113 2018; Lange & Moon, 2004; Power et al., 2008).





**Figure 1.** (a) Map of the tectonic structures and regions that have experienced cumulative slow slip of  $\geq 50$  mm between 2002 and 2012 in North Island, New Zealand (Wallace & Eberhart-Phillips, 2013). Fault traces from Barnes et al. (2010), Langridge et al. (2016), Mountjoy and Barnes (2011), and Pedley et al. (Pedley et al., 2010). Yellow stars are historic earthquakes (Doser & Webb, 2003; Downes, 2006; Grapes & Downes, 1997; Webb & Anderson, 1998) and  $M_w \geq 6$  earthquakes from August 2000 to 2022 (<https://www.geonet.org.nz/>). Black arrows indicate long-term relative motion between Pacific and Australian plates (Beavan et al., 2002). Blue arrows show motion of the Pacific Plate relative to overriding plate (or normal component of Pacific-Australian plate motion). (b) Map showing borehole-derived  $S_{Hmax}$  orientations (Behboudi et al., 2022; Mcnamara et al., 2021), and focal mechanisms derived  $S_{Hmax}$  orientations (Townend et al., 2012). Abbreviations: NIDFB = North Island Dextral Fault Belt; TVZ = Taupo Volcanic Zone.

Shallow horizontal stress orientations within the HSM have recently been constrained via borehole data analyses (Behboudi et al., 2022; Griffin, 2019; Griffin et al., 2021; Heidbach et

al., 2018; Lawrence, 2018; McNamara et al., 2021). Behboudi et al. (2022) provides a comprehensive overview of the along-strike and depth related variability in HSM stress orientations. Borehole-derived  $S_{Hmax}$  orientations rotate from ENE-WSW ( $065^{\circ}/245^{\circ} \pm 10^{\circ}$ ) in the central HSM to WNW- ESE ( $112^{\circ}/292^{\circ} \pm 20^{\circ}$ ) and NW- SE ( $140^{\circ}/320^{\circ} \pm 22^{\circ}$ ) in the southern HSM (Figure 1b). Deep stress orientations are defined by focal mechanism inversions (Townend et al., 2012), shear wave anisotropy (Illsley-Kemp et al., 2019), and gravitational stresses (Evanzia et al., 2017). Earthquake focal mechanism solutions ( $\leq 60$  km depth) indicate a regional  $S_{Hmax}$  orientation of  $060^{\circ}/240^{\circ} \pm 17^{\circ}$  and  $066^{\circ}/246^{\circ} \pm 22^{\circ}$  in the central and southern HSM, respectively (Figure 1b) (Behboudi et al., 2022; Townend et al., 2012).

The characterization of stress magnitudes at the HSM are currently limited to relative stress magnitudes derived from earthquake focal mechanisms at seismogenic depths (Townend et al., 2012), and direct measurements of the minimum principal stress magnitudes ( $\sigma_3$ ), vertical stress magnitudes ( $S_v$ ), pore pressures ( $P_p$ ) at shallow depths ( $< 3$  km) (Burgreen-Chan et al., 2016; D. Darby & Ellis, 2001; D. Darby & Funnell, 2001), and stress regime in one borehole (Tuhara-1A) in the central HSM (HRT, 2000). Observations of relative stress magnitudes ( $\leq 60$  km) by Townend et al. (2012) indicate a predominantly strike-slip and normal faulting regime along the HSM.  $P_p$  measured from repeat formation tests (RFTs) and modular dynamic tests (MDTs), and inferred from drilling mud weights reveal shallow ( $< 3$  km) overpressures within the upper plate of the central HSM (Burgreen-Chan et al., 2016; D. Darby & Funnell, 2001). High pore pressure in central and northern HSM are attributed to disequilibrium compaction of Miocene sediments and porosity reduction due to high horizontal compressive stresses associated with subduction of Hikurangi Plateau beneath the continental crust of North Island, New Zealand (Burgreen-Chan et al., 2016; David Darby & Funnell, 2001).  $\sigma_3$  magnitudes determined from leak-off tests are less than or close to  $S_v$  magnitudes (Burgreen-Chan et al., 2016), suggesting variable normal, strike-slip, and a reverse faulting regimes along the HSM.

In this study, we apply an indirect approach to constrain the three principal stress magnitudes along the shallow HSM crust using openly available borehole data. We discuss our findings in the context of understanding the upper plate tectonics within the HSM forearc. This study, in combination with stress orientation studies already completed for the HSM, provides a deeper insight into the variable tectonic behaviors associated with subduction margins, and will serve as crucial information to assist in future hazard assessments of this region.

## 2 Geologic setting and background

The HSM at the east coast of North Island, New Zealand is a site of recent significant scientific investigation into the complexity of subduction dynamics. The HSM is formed by westward subduction of the oceanic crust of the Hikurangi Plateau beneath the continental crust of the North Island of New Zealand (Davy, 1992; Davy et al., 2008). The oblique relative motion of the Australian-Pacific plate increases from ~31 mm/year in the southern to ~48 mm/year in the northern North Island (Figure 1a) (Wallace et al., 2004). Tectonic deformation across the HSM ranges from subduction-related shortening at the Hikurangi Trough, strike-slip faulting along the North Island Dextral Fault Belt (NIDFB), and back-arc extensional tectonics in the Taupo Volcanic Zone (TVZ) at the center of North Island (Wallace et al., 2004; Figure 1a). The East Coast forearc has rotated at rate of 3°–4°/Myr relative to the Australian plate, resulting in the TVZ back-arc rifting, strike-slip and/or normal faulting in the onshore portion of the northern and central HSM, transpressional faulting in the southern HSM, and a large along-strike variation in convergence rate at the Hikurangi Trough (Figure 1a) (Fagereng & Ellis, 2009; Nicol et al., 2007; Wallace et al., 2004; Wallace, Fagereng, et al., 2012). The oblique motion of the Australian-Pacific plate is partitioned into a margin-perpendicular component and a margin-parallel component. The margin-perpendicular component occurs along the Hikurangi subduction interface and provides NW-SE shortening mostly accommodated by slip on the subduction interface (>80%) and active frontal thrusts in the overriding plate (Nicol & Beavan, 2003). The margin-parallel component is largely accommodated by a combination of right-lateral strike-slip on the North Island Dextral Fault Belt (NIDFB) and clockwise rotation of the North Island forearc (Beanland & Haines, 1998; Nicol et al., 2007; Wallace et al., 2004).

## 3 Methodology and Data

### 3.1 Data Sources and Limits

Data used in this study is sourced from 44 boreholes along the HSM (Figure 1), 41 of them are located within the onshore forearc and 3 are located offshore the east coast of NZ but west of the Hikurangi Trough. Data utilised includes wireline logging acquired over the period 1967-2013 from 0 to a maximum depth of 4350 m below ground level. Wireline data includes density logs from 26 boreholes, sonic velocity logs from 24 boreholes, and borehole image logs from 10 boreholes. Data presented here include the analysis of 21 leak-off tests and 39 formation integrity tests from 30 boreholes spanning a depth range of 371.5 to 3610.6

m, mud weight logs from 44 boreholes, and repeat formation test results from 2 boreholes spanning a depth range of 1335-2700 m. How each of these data are utilised in determining aspects of the in situ-stress magnitudes across the HSM is detailed below. All depths in this study is referenced to ground level for onshore boreholes and sea level for offshore boreholes.

### 3.2 Vertical stress magnitude ( $S_v$ )

Assuming the vertical stress ( $S_v$ ) is aligned to one of the principal stresses, the  $S_v$  magnitude at any specific subsurface depth can be determined by the integration of rock densities from the surface to the depth of interest (equation 1):

$$S_v = \rho_w g Z_w + \int_{Z_w}^Z \rho(Z) g dZ \approx \rho_w g Z_w + \bar{\rho} g (Z - Z_w) \quad 1$$

where  $\rho_w$  is the average seawater density (1.03 g/cm<sup>3</sup>),  $g$  is the gravitational acceleration constant ( $\sim 9.81 \text{ m/s}^2$ ),  $Z_w$  is the depth of the water column (m),  $Z$  is the depth of interest (m),  $\rho(Z)$  (g/cm<sup>3</sup>) is bulk density of the rock as a function of depth, and  $\bar{\rho}$  (g/cm<sup>3</sup>) is the average density of the rock column above  $Z$ . For onshore boreholes,  $Z_w$  is equal to zero.

We utilise 26 density wireline logs to estimate  $S_v$  profiles. At times wireline density logs are not acquired within the top depth intervals of drilled boreholes, the rock density is extrapolated from the top of a density log to the surface (seafloor for offshore boreholes) to more accurately determine a complete  $S_v$  profile. This study uses several extrapolation methods: 1) using wireline sonic logs to convert compressional velocity to density values in boreholes where checkshot data or vertical seismic profile (VSP) surveys are available (Kereru-1, Hawke Bay-1, Opoutama-1, Whakatu-1, Ngapaeruru-1, Tawatawa-1, and Titihaoa-1), 2) using average densities from nearby boreholes with similar stratigraphy (e.g. boreholes Kauhauroa-1, Kauhauroa-2, Kauhauroa-5, Makareao-1, and Tuhara-1A are all within <20 km of each other), or 3) using standard Gardner's relationship (Gardner et al., 1974) and/or regional Gardner's relationship (Table S5 in Supporting Information S1) to convert compressional velocity data from sonic wireline logs to density data logs (e.g. Hawke Bay-1, Rere-1). All density logs used in this study, supplied by the New Zealand Petroleum and Minerals group (NZPM), have been undergone borehole environmental corrections.

### 3.3 Minimum principal stress magnitude ( $\sigma_3$ )

$\sigma_3$  can be measured directly from pressure-time plots produced during leak-off tests (LOTs), extended leak-off tests (XLOTs), or mini-frac tests (Addis et al., 1998; Bell, 2003; White et

al., 2002; Zoback et al., 2003). In the HSM, LOTs are the most common tests available to calculate *in-situ*  $\sigma_3$  magnitudes. LOTs are pumping pressure tests conducted in a borehole a few meters below recently set casing shoes. During constant fluid volume pumping, the recorded fluid pressure increase stops behaving linearly with time as the injected fluid pressure surpasses the  $\sigma_3$  confining stress around the borehole and fluid starts to penetrate into the formation around the borehole (Addis et al., 1998; Bell, 1996). The point when the fluid pressure-time curve becomes non-linear (leak-off pressure (LOP)) can be read as an approximation of  $\sigma_3$  magnitude. If a LOT is stopped at any point before the LOP is reached the test is called a formation integrity test (FIT) and fluid pressure has not exceeded  $\sigma_3$  magnitude. In this case, the final fluid pressure value recorded during the FIT can be used as an estimate of the lower boundary of the  $\sigma_3$  magnitude (e.g. Makareao-1, Zoback et al., 2003).

In the majority of boreholes studied here the validity and accuracy of LOTs cannot be assessed as the pressure–time record data is not fully reported, with only the final LOP being provided in the text reports by drilling companies. Furthermore, pressure–time records are sometimes estimated by only a few distinct data points, obtained from pressure measurements on fluctuating gauges or flow rate estimations from counting pump strokes, making it impossible to determine the specific and accurate LOP values (Zoback, 2007). It is therefore possible for  $\sigma_3$  to be reported slightly higher or extremely close to  $S_v$  when the measurements are not carefully taken or reported. Further consideration for subduction margins is provided by Couzens-Schultz and Chan (2010), who demonstrate that in active compressional settings and seismically active regions, LOTs cause shear failure along pre-existing fractures rather than generating new tensile fractures, leading to an underestimation of the  $\sigma_3$  magnitude.

We first calculate  $\sigma_3:S_v$  for all boreholes for which LOP measurements are available and then use the average of these data to extrapolate the  $\sigma_3$  values beyond the depth of measurements. The FIT: $S_v$  and  $\sigma_3:S_v=1$  are used to define the lower and upper limit of the  $\sigma_3$  profile, respectively.

### **3.4 Maximum horizontal stress magnitude ( $S_{Hmax}$ )**

#### **3.3.1 $S_{Hmax}$ estimation from borehole failure analysis**

When a vertical borehole is drilled into a homogeneous, isotropic, and elastic medium parallel to one of the three principal stress orientations, the stress at the borehole wall is

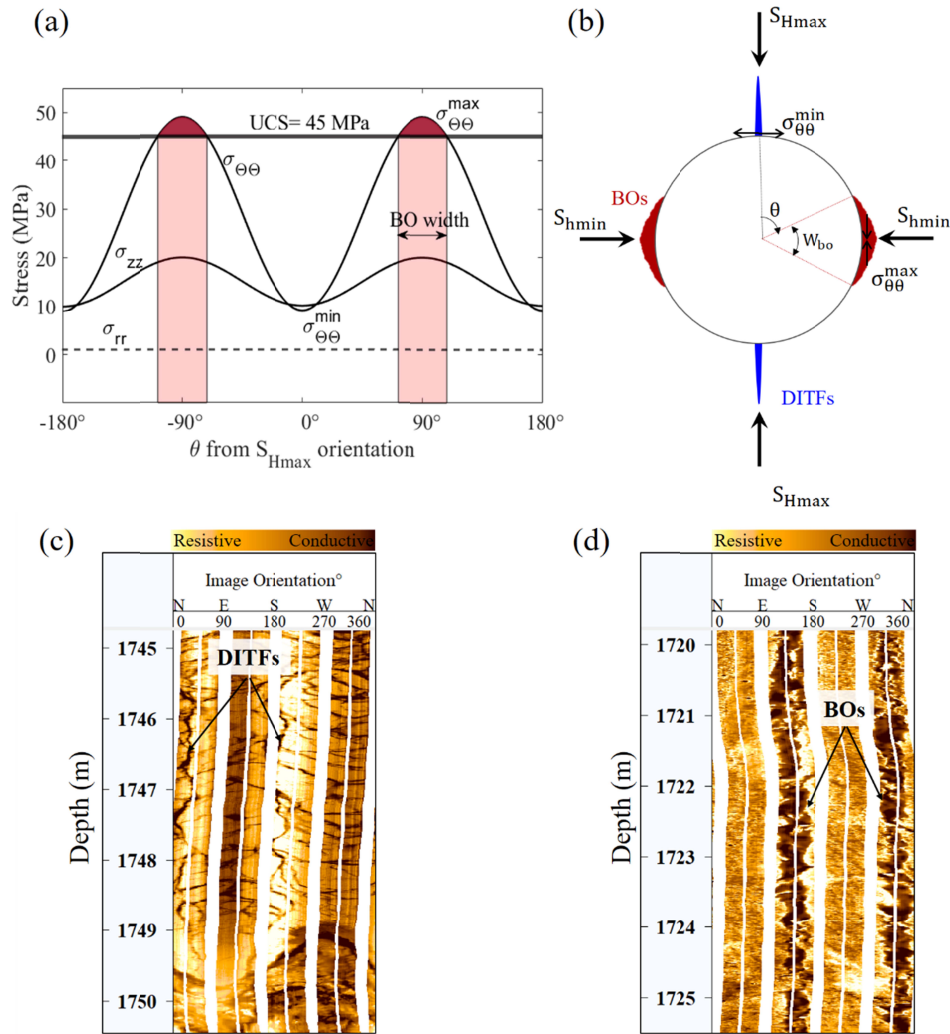
254 redistributed regarding to non-uniform, far-field principal stresses (Jaeger et al., 2009;  
 255 Zoback, 2007). Assuming far field principal stresses are vertical and horizontal, the local  
 256 principal effective stresses at a vertical borehole wall can be defined (Moos & Zoback, 1990;  
 257 Zoback, 2007):

$$\sigma_{\theta\theta} = S_{Hmax} + S_{hmin} - 2\cos 2\theta (S_{Hmax} - S_{hmin}) - P_p - APRS \quad 2a$$

$$\sigma_{ZZ} = S_v - 2\theta \cos 2\theta (S_{Hmax} - S_{hmin}) - P_p \quad 2b$$

$$\sigma_{rr} = APRS - P_p \quad 2c$$

258 where  $\sigma_{\theta\theta}$  is the effective hoop stress (acting parallel to the borehole wall),  $\sigma_{ZZ}$  is the  
 259 effective vertical stress,  $\sigma_{rr}$  is the effective radial stress (acting perpendicular to the borehole  
 260 wall),  $S_{Hmax}$  and  $S_{hmin}$  are the maximum and minimum horizontal principal stress magnitudes,  
 261  $\theta$  is Poisson's ratio, APRS is the annulus pressure at the time of borehole failure (or mud  
 262 weight pressure),  $P_p$  is pore pressure, and  $\theta$  is the angle between the edge of borehole  
 263 breakout and the  $S_{Hmax}$  orientation (Figure 2a & 2b).



**Figure 2.** (a) Borehole schematic showing local principal stresses ( $\sigma_{\theta\theta}$ ,  $\sigma_{zz}$ , and  $\sigma_{rr}$ ) at the borehole wall as a function of azimuth ( $\theta$ ) measured relative to  $S_{Hmax}$  orientation and presence of breakouts for an example in which  $S_{Hmax} = 50$  MPa,  $S_v = 45$  MPa,  $S_{hmin} = 40$  MPa, and UCS=45 MPa. The red shaded region shows schematically the circumference where  $\sigma_{\theta\theta}$  is large enough to exceed the compressional strength of the formation and induce BOs. (b) Diagram of a borehole cross-section showing the relationship between BOs, DITFs, and the horizontal principal stress orientations. (c) Example of DITFs as they appear on a resistivity image log. (d) Examples of BOs as they appear on a resistivity image log. Figures 2c-d are from resistivity image logs of borehole Kauhauroa-5. Abbreviations: UCS = unconfined compressive strength; BO = borehole breakout; DITF: drilling induced tensile fracture.

Where local effective stresses exceed the tensile or compressive rock strength of the formation around the borehole, borehole failures such as drilling induced tensile fractures (DITFs) and borehole breakouts (BOs) can form, respectively (Figure 2a). Measurements of the properties of these borehole failures, e.g. the azimuth angle of BOs and/or DITFs and the

angular width of BOs can be used to determine *in-situ* principal stress orientations and to calculate in situ stress magnitudes present at the time of drilling.

DITFs form on the borehole wall where local effective stress concentrations around the borehole wall lead to a minimum  $\sigma_{\theta\theta}$  less than the tensile strength of the rock ( $\sigma_{\theta\theta}^{\min} \leq 0$ ) (Aadnoy, 1990), at a borehole azimuth parallel to  $S_{Hmax}$  ( $\theta=0^\circ/180^\circ$ ) (Figure 2b) (Aadnoy, 1990; Bell, 2003; Bell & Gough, 1979; Brudy & Zoback, 1999). DITFs typically appear as narrow, conductive (on resistivity image logs) or low amplitude and slower travel time (on acoustic image logs) pairs,  $\sim 180^\circ$  from each other around the borehole wall circumference (Figure 2c). DITFs are generally parallel or slightly inclined to the borehole axis in vertical to semi-vertical boreholes (Brudy & Zoback, 1999; Zoback, 2007). Where DITFs are observed the magnitude of the far-field  $S_{Hmax}$  can be constrained using Equation 4 (Zoback, 2007):

$$3S_{hmin} - T_0 - P_p - APRS - \sigma^{\Delta T} \leq S_{Hmax} \quad 3$$

where  $S_{Hmax}$  and  $S_{hmin}$  are maximum and minimum horizontal principal stresses respectively,  $T_0$  is the formation tensile strength,  $P_p$  is pore pressure, APRS is annulus pressure (or mud weight), and  $\sigma^{\Delta T}$  is thermal stress arising from the difference between the drilling mud temperature and formation temperature.  $\sigma^{\Delta T}$  is applied where there is a noticeable difference between mud and rock temperature, such as geothermal boreholes. The tensile rock strength in sedimentary rocks is often quite small (a few MPa) and can be assumed to be zero in the analysis of DITFs (Brudy & Zoback, 1999). In this study,  $\sigma^{\Delta T}$  is considered negligible.

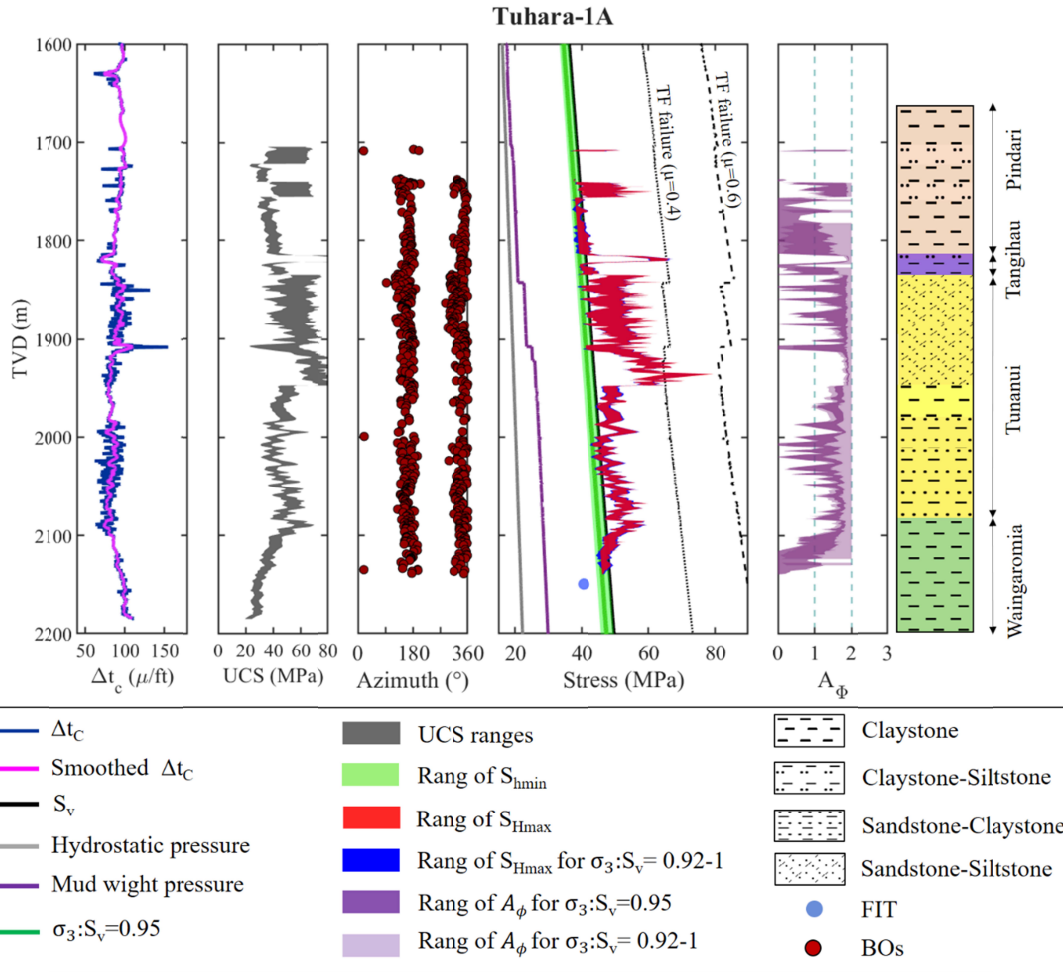
BOs form as enlargements of the borehole diameter on opposite sides of the borehole wall where  $\sigma_{\theta\theta}$  is large enough to exceed the formations compressional strength (Figure 2a) (Bell & Gough, 1979; Zoback, 2007). The  $\sigma_{\theta\theta}$  magnitude reaches a maximum at  $\theta=\pm 90^\circ$  (Figure 2a), which occurs at a borehole azimuth oriented perpendicular to the  $S_{Hmax}$  direction (Figure 2b). BOs typically appear as a pair of wide, out-of-focus, conductive (in water-based mud; Figure 2d) or resistive (in oil-based mud) zones on resistivity image logs, or as zones of low acoustic amplitude and slower travel time on acoustic image logs. BOs are located  $\sim 180^\circ$  from each other around the circumference of the borehole wall (Figure 2b & 2d).  $S_{Hmax}$  magnitudes can be estimated by measuring BO widths ( $W_{bo}$ ) from borehole image logs using Equation 5 (Barton et al., 1988; Vernik & Zoback, 1992):

$$S_{Hmax} = \frac{(UCS + P_p + APRS + \sigma^{\Delta T}) - S_{hmin}(1 + 2 \cos(\pi - W_{bo}))}{1 - 2 \cos(\pi - W_{bo})} \quad 4$$



where  $W_{bo}$  is the angular width of the BO; UCS is unconfined compressive strength of the formation,  $P_p$  is pore pressure, APRS is annulus pressure or mud weight,  $S_{hmin}$  is the minimum horizontal principal stress magnitude, and  $\sigma^{\Delta T}$  is the thermal stress effect resulting from the difference between the drilling mud temperature and formation temperature. In this study,  $\sigma^{\Delta T}$  is considered negligible.

UCS is a key parameter in estimating  $S_{Hmax}$  magnitude (Equation 4), and can either be directly measured from laboratory strength tests on core samples, or estimated using empirical relationships between UCS and other rock properties (Chang et al., 2006). Direct measurements of rock strength are rare for the HSM. Borehole Waingaromia-2 in the northern HSM is the only borehole where a laboratory strength test was conducted on calcareous claystone and mudstone core samples (acquired from 132 and 362 m measured depth, respectively), providing UCS values of 1.1-1.2 MPa and friction angles of 20.5°-32.1° (friction coefficient 0.37-0.64) (Indo-Pacific Energy (NZ) Ltd., 2002). However, no relationship between P-wave slowness ( $\Delta t_c$ ) and UCS was established because no geophysical logs were obtained and velocity measurements on core samples are unavailable. Therefore, in this study, UCS values are indirectly estimated by using empirical relationships between rock strength and  $\Delta t_c$ . Empirical equations have been developed for different rock types, relating various rock properties to UCS across the world. In this study we utilize a variety of empirical relationships between UCS and sonic velocity by matching appropriate equations to dominant lithologies encountered along each studied borehole in an effort to reduce uncertainty in UCS values and thus  $S_{Hmax}$  magnitude values. Upper and lower bounds of the UCS are determined using various published empirical relationships (Chang et al., 2006) to provide a range of possible  $S_{Hmax}$  magnitudes (Figure 3). Details on the equations used in individual boreholes to determine the lower and upper limits of UCS can be found in Table S1, Table S2, and S3 in Supporting Information S1.



**Figure 3.** Calculated far-field in situ stress magnitudes, referenced to the sea level in borehole Tuhara-1A. (a) P-wave slowness (blue line) de-spiked and smoothed over 3m intervals (pink line). (b) Range of UCS values derived from P-wave slowness using relations in Table S2 and S3 (c) Azimuth of borehole BOs. (d) Calculated  $S_v$  (solid black line),  $S_{hmin}$  (green field), and  $S_{Hmax}$  (red field) magnitudes by considering that pore pressure is equal to mud weight. The hydrostatic pressure (grey line) is computed assuming a sea water density of 1.03 g/cc. The  $\sigma_3$  and the range of  $S_{hmin}$  is determined from the average  $\sigma_3:S_v = 0.95$  and  $\sigma_3:S_v = 0.92-1$ , respectively. Abbreviations:  $\Delta t_c$ : P-wave slowness; UCS = uniaxial compressive strength; BO = borehole breakout; FIT = formation integrity test; TF failure: thrust faulting failure;  $\mu$ : friction coefficient;  $S_v$ : vertical stress;  $S_{hmin}$ : minimum horizontal stress;  $S_{Hmax}$ : maximum horizontal stress;  $\sigma_3$ : minimum principal stress;  $A_\phi$ : Tectonic stress regime index.

A further important parameter required to calculate  $S_{Hmax}$  magnitudes and effective stresses is  $P_p$ . Direct  $P_p$  measurements tests such as RFTs and MDTs are the most reliable measurements (Gunter & Moore, 1986; Zoback, 2007). However, these direct  $P_p$  measurements are difficult to acquire, particularly in low permeability formations, and are often only conducted at

depths where possible overpressures may exist (Dutta et al., 2021; Lee et al., 2022; Y. Z. Ma & Holditch, 2015; Zoback, 2007). Drilling mud weight logs can provide indirect, continuous approximations of the  $P_p$  along a borehole, and can be used as a proxy of  $P_p$  assuming the mud weights have been chosen to stabilize the borehole during drilling, and if no significant mud losses or kicks are reported (Van Ruth et al., 2002) . Mud Losses of greater than 25 bbl/hr for water-based mud (Zhang & Yin, 2017) may indicate that annulus pressure exceeded  $P_p$  or/and  $\sigma_3$  value, resulting in the loss of fluids into the formation. While kicks and high fluid influx indicate that  $P_p$  is greater than the annulus pressure. In both cases, the  $P_p$  derived from drilling mud weight logs should be corrected to generate a good estimation of  $P_p$ . In this study we use mud weight logs from 44 boreholes to calculate  $P_p$ . Minor seepage (mud losses <22 bbl/hr) is reported for boreholes Kauhauroa-2/5, Makareao-1, Tuhara-1A, Ngapaeruru-1, Tawatawa-1, and Titihaoa-1, and Ranui-2 in the intervals where BOs are observed, providing confidence in the use of mud weight logs in those intervals for  $P_p$  determination. A minor mud loss of 28 bbl/hr has been observed at severely fractured depth interval of 1030-1225 m TVD in borehole Ngapaeruru-1 which was treated by remedial techniques and procedures easily. Moreover, minor background gas and fluid influx are reported in boreholes Kahauuroa-5, Makareao-1, Tuhara-1A, Tawatawa-1, and Titihaoa-1, which were controlled by mud weight such that they never flowed. Since no significant mud losses or kicks are reported in the depth intervals where BOs are observed, we consider the annulus pressure records a good proxy of  $P_p$  in those depth intervals.

The  $P_p$  calculated from mud weight logs in Kauhauroa-5 and Titihaoa-1 boreholes are further calibrated using direct  $P_p$  measurements obtained from RFTs. Formation tests conducted in 17 further HSM boreholes (Awatere-1, Hukarere-1, Kauhauroa-2/3/4B, Kiakia-1A, Makareao-1, Mangaone-1, Morere-1, Opoutama-1, Ruakituri-1, Takapau-1, Te Hoe-1, Tuhara-1A/1B, Waitahora-1, and Waitaria-2) are not included in this study due to incomplete pressure build ups during testing in low-permeability formations, test seal failures, or tests conducted in formation intervals supercharged to hydrostatic pressure.

### 3.3.2 $S_{Hmax}$ magnitude estimation from frictional limit theory

To constrain  $S_{Hmax}$  magnitudes that result in the observed BO and DITF occurrences, the stress state is assumed to be limited by Coulomb frictional sliding on an optimally oriented and pre-existing fault plane (Zoback, 2007). This means that the maximum effective principal

stress cannot exceed the stress value required to cause slip, defined by the friction coefficient ( $\mu$ ) of adjacent faults, on a critically oriented fault plane (Jaeger et al., 2009; Sibson, 1974):

$$\frac{(\sigma_1 - P_p)}{(\sigma_3 - P_p)} \leq ((1 + \mu^2)^{0.5} + \mu)^2 \quad 5$$

where  $\sigma_1$  is the maximum principal stress,  $\sigma_3$  minimum principal stress,  $P_p$  is pore pressure, and  $\mu$  is coefficient of friction on an optimally oriented, cohesionless, pre-existing fault.

This constraint is typically displayed as a stress polygon, which shows the permissible values of horizontal principal stress magnitudes for a specific depth,  $S_v$ ,  $\mu$ , and  $P_p$  for normal, strike-slip, and thrust faulting tectonics (Zoback, 2007). Although this method only provides the upper and lower limits for the  $S_{Hmax}$  magnitude, it can yield more accurate ranges of permissible  $S_{Hmax}$  magnitudes when combined with  $S_{Hmax}$  magnitude estimates from borehole failure analysis (Chang et al., 2010; Huffman & Saffer, 2016).

### 3.5 Tectonic stress regime index ( $A_\phi$ )

In order to characterize a stress regime or faulting style with stress magnitude data, we use the stress regime index ( $A_\phi$ , Equation 6a and 6b) described by Simpson (Delvaux et al., 1997):

$$A_\phi = (n + 0.5) + (-1)^n (R - 0.5) \quad 6a$$

$$R = (\sigma_2 - \sigma_3) / (\sigma_1 - \sigma_3) \quad 6b$$

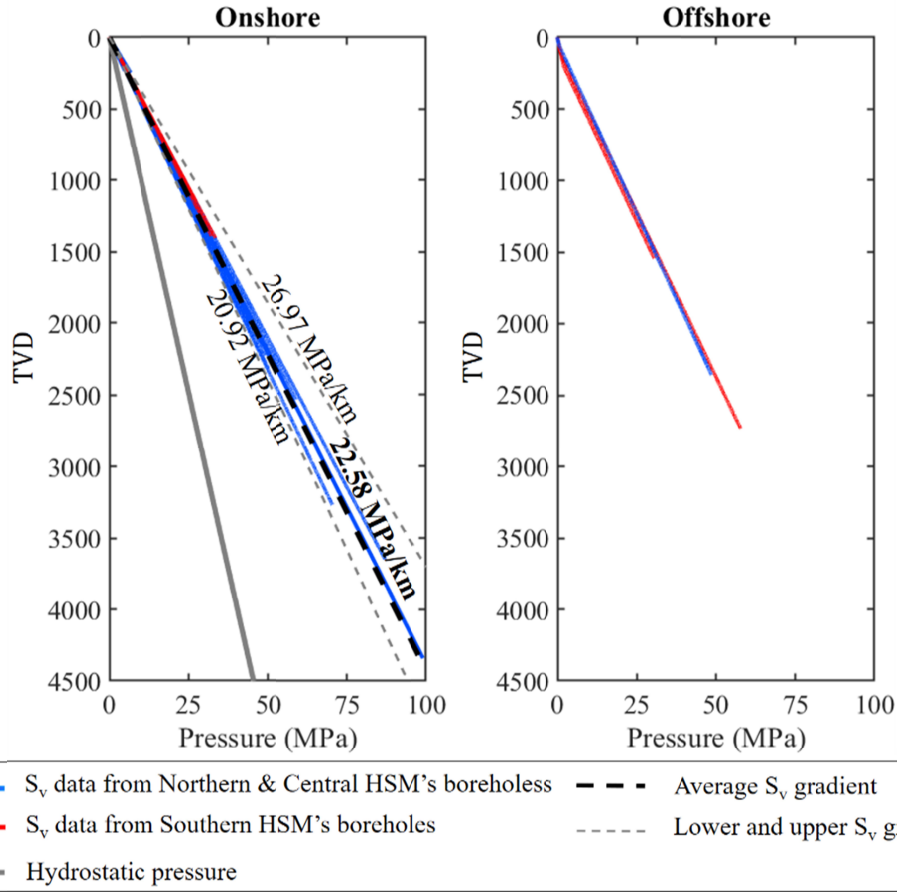
where  $n$  is the number of principal stress components greater than the principal stress whose axis is closest to the vertical,  $R$  is the stress ratio, and  $\sigma_1$ ,  $\sigma_2$ ,  $\sigma_3$  are the maximum, medium, and minimum principal stress magnitudes, respectively.

$A_\phi$  values range from 0 to 1 in normal faulting regimes, 1 to 2 in strike-slip regimes, and 2 to 3 in thrust faulting regimes.

## 4 Results

### 4.1 Vertical Stress Magnitudes

$S_v$  magnitudes determined from 24 onshore boreholes provide overburden stress gradients ranging from 20.92 to 26.97 MPa/km, with a mean value of  $22.58 \pm 1.23$  MPa/km (Figure 4a; Table S5).  $S_v$  magnitudes measured within the 3 offshore boreholes range from 20.9 to 21.7 MPa/km with a mean value of  $21.26 \pm 0.4$  MPa/km.



**Figure 4.**  $S_v$  gradient profiles from boreholes in the a) HSM onshore and b) HSM offshore.  $S_v$ : vertical stress.

## 4.2 Minimum Principal Stress Magnitudes

$\sigma_3$  magnitudes calculated from LOT data range from 1.9 MPa at 102.3 m TVD (borehole Waitaria-2) to 77.9 MPa at 3610.6 m TVD (borehole Rere-1) (Table 1). For all examined LOT data (with the exception of one test in borehole Titihoa-1 where  $\sigma_3$  derived from LOP is greater than  $S_v$ ), the normalized effective  $\sigma_3$  ratio ranges from 0.23 to 1 (Table 1).

Table 1.  $\sigma_3$  magnitudes calculated from LOP measurements.

	Borehole	Depth <sup>a</sup> (m)	$\sigma_3^b$ (MPa)	$S_v^c$ (MPa)	$\sigma_3/S_v$
Central HSM	Awatere-1	301.9	6.1	6.5	0.94
		1085.3	22.4	24.0	0.93
	Hawke Bay-1	386.6	7.2	7.4	0.97
		1359.6	26.4	27.7	0.95
	Kauhauroa-1	455.8	7.9	9.5	0.83
	Kauhauroa-2	463.8	8.2	10.1	0.81
	Kauhauroa-5	459.2	7.4	10.2	0.73
	Kereru-1	481.1	9.4	10.2	0.92
	Kiakia-1/1A	318.6	5.7	6.9	0.83

<b>Southern HSM</b>	Opoho-1	521.3	9.3	11.7	0.79
	Rere-1	3610.6	77.9	82.3	0.95
	Tuhara-1/1A	590.8	8.4	13.3	0.63
	Waitaria-2	102.3	1.9	2.4	0.79
	Ranui-1	357.3	5.2	8.6	0.6
	Tawatawa-1	722.5	12.7	12.7	1
	Titihaoa-1	614.9 1585.7 1979.8	9.5 30.6 43.4 <sup>a</sup>	11 32.4 41.3	0.86 0.94 1.05

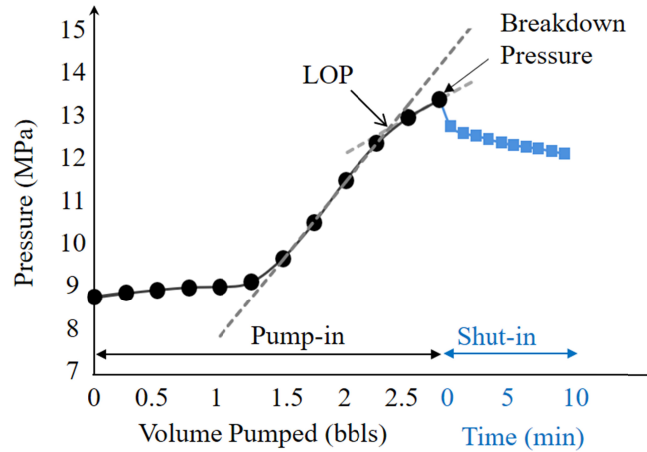
<sup>a</sup> True vertical depth from ground level for onshore boreholes and sea level for offshore boreholes

<sup>b</sup>  $\sigma_3$  derived from Leak-off pressure

<sup>c</sup>  $S_v$  vertical stress

\*  $\sigma_3$  derived from LOP is greater than  $S_v$

LOP values measured in boreholes Tuhara-1/1A and Ranui-1 (11 m west of Ranui-2) are remarkably low, such that LOP values are less than the  $\sigma_3$  values estimated by normal faulting failure with friction coefficients less than 0.6. In borehole Titihaoa-1, the  $\sigma_3$  value (43.4 MPa) derived from an LOT performed at ~1979.8 m TVD is greater than the  $S_v$  for this depth (41.3 MPa). In this case,  $\sigma_3$  is considered to be vertical, indicating a thrust/reverse faulting regime (Zoback, 2007). In borehole Tawatawa-1 two LOTs were performed at 722.5 m TVD. The initial test yielded a LOP of 13.23 MPa, while the second LOT yielded a LOP of 13.36 MPa (Tap Oil Limited, 2004). Our reassessment of pressure-time curve of the second LOT (which had more data defining the time-pressure plot) reveals that the formation breakdown pressure (FBP) was reported rather than LOP, resulting in an overestimation of  $\sigma_3$  making it appear greater than the  $S_v$  for this depth. We determine the LOP of the second test by intersecting the straight line of the linear section with the tangent line of the ascending section on the pressure-volume curve (Figure 5), and report a  $\sigma_3$  magnitude of 12.7 MPa, almost equal to  $S_v$  (13 MPa). Assuming  $\sigma_3$  measurements made from the reported LOT data in the study boreholes are a proxy of  $\sigma_3$  (after correcting  $\sigma_3$  derived from LOP  $> S_v$  to  $\sigma_3 = S_v$ ), a HSM average minimum normalized effective stress ratio of  $0.66 \pm 0.2$  and  $0.7 \pm 0.3$  are derived for the central and southern HSM respectively (Figure 6).



**Figure 5.** Results of the leak-off test run at 722.5 m TVD in borehole Tawatawa-1. Pressure versus volume of mud pumped to the formation curve reveals that the leak-off pressure (LOP) is 12.7 MPa.

FIT data show the lower limit of  $\sigma_3$  magnitudes are typically below  $S_v$  in most of boreholes in this study. However, in some boreholes (Table 2), FIT results are approximately equal to or greater than  $S_v$ . The entire FIT dataset for all boreholes in this region can be found in Table S3 in Supporting Information S1, and are used to constrain  $\sigma_3$  profiles within boreholes where LOP measurements are not available.

Table 2. The lower limit of  $\sigma_3$  values calculated from FIT data.

	Borehole	Depth <sup>a</sup> (m)	The lower limit of $\sigma_3$ <sup>b</sup> from FIT <sup>c</sup> (MPa)	$S_v$ <sup>d</sup> (MPa)
Central HSM	Hukarere-1	89.9	1.81	1.69
		430.2	8.64	8.55
	Kauhauroa-3	332.7	7.33	7.42
		999.2	22.44	22.56
	Kauhauroa-4	346.1	7.63	7.36
	Kauhauroa-4B	91	1.95	1.97
		538.3	11.56	11.75
	Makareao-1	306.2	6.76	6.55
		484.8	10.6	10.59
	Rere-1	115	2.39	2.65
		2109.4	49.56	48.17
	Waitahora-1	95.8	2	2.01
		722	16.24	15.53
		994.5	22.93	21.76
	Waitaria-2	556.1	12.53	12.85

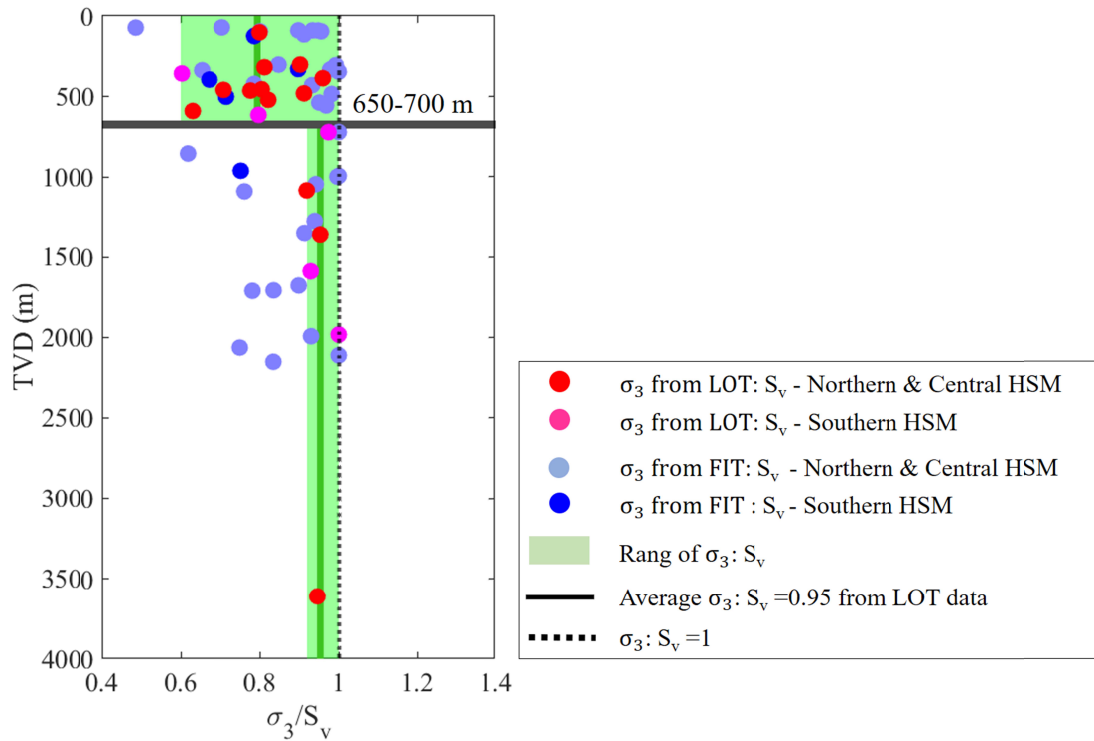
<sup>a</sup> True vertical depth from ground level for onshore boreholes and sea level for offshore boreholes

<sup>b</sup> Minimum principal stress magnitudes

<sup>c</sup> Formation integrity test

<sup>d</sup>  $S_v$  vertical stress

We calculated  $\sigma_3:S_v$  (after correcting  $\sigma_3$  derived from LOP and FIT  $>S_v$  to  $\sigma_3=S_v$ ) for all 21 LOTs and 39 FIT measurements in the study area and found that  $\sigma_3:S_v$  varies significantly above and below 650-700 m TVD (Figure 6). The  $\sigma_3:S_v$  ranges from 0.6-1 at depths above 650-700 m TVD, while it ranges from 0.92-1 below this depth interval (Figure 6). Above and below 650-700 m TVD, the average  $\sigma_3:S_v$  values derived from LOT measurements are 0.79 and 0.95, respectively.



**Figure 6.**  $\sigma_3: S_v$  along the HSM. Abbreviations: LOT: leak of test; FIT = formation integrity test;  $S_v$ : vertical stress;  $\sigma_3$ : minimum principal stress.

In this study, to create the  $\sigma_3$  profile in boreholes where LOT measurements are not available at the depth of interest, we only consider LOT and FIT measurements recorded at depths below 650-700 m TVD for two main reasons: 1) to exclude the influence of topographic effects and shallow processes such as gravitational collapse, erosion, and subsidence in the  $\sigma_3$  calculation, and 2) because our borehole breakouts data to estimate  $S_{Hmax}$  magnitudes in 7 out of 8 boreholes are located below 700m TVD.

### 4.3 Stress magnitude from borehole data

#### 4.3.1 Central HSM (Hawke's Bay region)

A range of potential  $S_{Hmax}$  magnitudes are estimated along individual boreholes at depths intervals where BO widths and DITFs are measured from image logs. Further the lower and

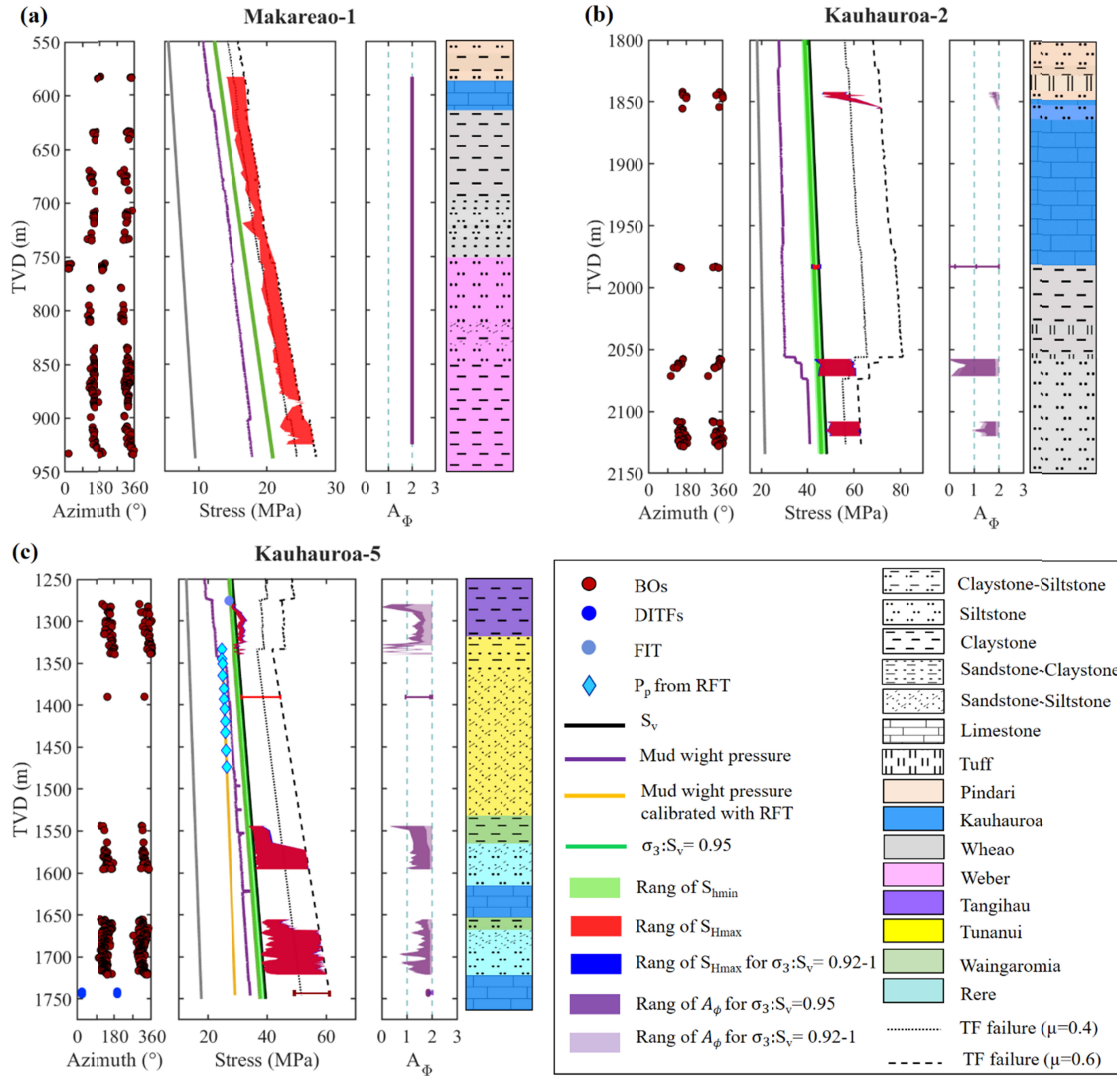


upper limit of  $S_{Hmax}$  magnitudes are constrained by theoretical limits provided by slip on pre-existing faults with a friction coefficient of 0.6 (described in 3.3.2).

### Makareao-1 borehole

The  $\sigma_3:S_v$  ratios of 1.03 and 1 are determined from  $\sigma_3$  values calculated using FIT data at depths 306.2 and 484.8 m TVD respectively (Table 2). The  $\sigma_3:S_v \geq 1$  indicates that  $\sigma_3 = S_v$  in this borehole. The  $S_{Hmax}:S_v$  ratio of 1.03-1.31 is determined for borehole Makareao-1 using the  $S_{Hmax}$  values calculated from the lower and upper value of UCS.

The  $S_{Hmax}:S_v$  ratio of 1.03-1.31 and a  $\sigma_3:S_v = 1$  along this borehole indicates a stress regime such that  $\sigma_3 = S_v < S_{Hmax}$  (Figure 7a).  $A_\phi = 2$  is determined from calculated stress magnitude data in this borehole.



**Figure 7.** Calculated far-field in situ stress magnitudes, referenced to the sea level in borehole (a) Makareao-1 (b) Kauhauroa-2 (c) Kauhauroa-5 in the central HSM. Abbreviations:  $\Delta t_C$ : P-wave slowness; UCS = uniaxial compressive strength; BO = breakout; DITF: drilling induced fracture; FIT = formation integrity test; TF failure: thrust faulting failure;  $\mu$ : friction coefficient; RFT: repeat formation test;  $S_v$ : vertical stress;  $S_{hmin}$ : minimum horizontal stress;  $S_{Hmax}$ : maximum horizontal stress;  $\sigma_3$ : minimum principal stress;  $A_\phi$ : tectonic stress regime index.

#### **Kauhauroa-2 borehole**

A  $\sigma_3:S_v$  ratio of 0.81 is determined from  $\sigma_3$  value calculated using LOT data at 463.8 m TVD (Table 1). The  $\sigma_3$  values in the deeper part of the borehole are calculated from the average  $\sigma_3:S_v$  ratio of 0.95 along the HSM and are further constrained by the lower limit of  $\sigma_3$  value determined from an FIT= 30.23 MPa at 1707.3 m TVD. The  $S_{Hmax}:S_v$  ratio of 0.95-1.71 is determined for borehole Kauhauroa-2 using the  $S_{Hmax}$  values calculated from the lower and upper value of UCS.

The  $S_{Hmax}:S_v$  ratio of 0.95-1.71 and the  $\sigma_3:S_v$  ratio of 0.95 indicate a dominant stress regime such that  $S_{hmin} \leq S_v \leq S_{Hmax}$  (Figure 7b). A  $0 \leq A_\phi \leq 1.94$  is determined from calculated stress magnitude data in this borehole.  $S_v$ ,  $S_{hmin}$ , and the lower limit of  $S_{Hmax}$  are nearly equal below 1980 m TVD such that  $S_{hmin} \approx S_{Hmax} \approx S_v$ .

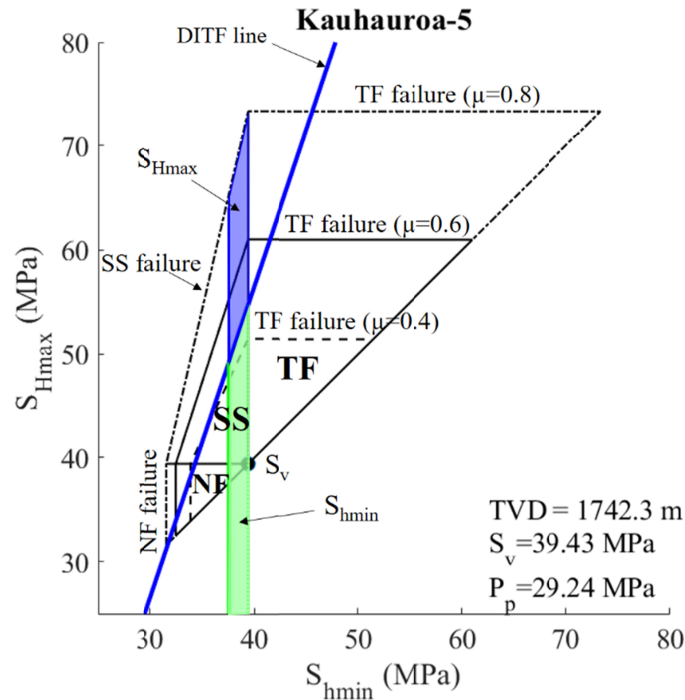
#### **Kauhauroa-5 borehole**

A  $\sigma_3:S_v$  ratio of 0.73 is determined from  $\sigma_3$  value calculated using LOT data at 459.2 m TVD (Table 1). The  $\sigma_3$  values in the deeper part of the borehole are calculated from the average  $\sigma_3:S_v$  ratio of 0.95 along the HSM and are further constrained by the lower limit of  $\sigma_3$  value calculated from an FIT value of 27.13 MPa at 1276.1 m TVD. The  $S_{Hmax}:S_v$  ratios of 0.95-1.13 in 1280-1350 m TVD and 0.97-1.54 in 1390-1750 m TVD are determined in this borehole using the  $S_{Hmax}$  values calculated from the lower and upper value of UCS.

The analysis of  $S_{Hmax}$  magnitudes and the  $\sigma_3:S_v$  ratio of 0.95 indicate a dominant  $S_{hmin} \approx S_v \approx S_{Hmax}$  ( $S_{Hmax} - S_{hmin} < 5$  MPa) and  $0 \leq A_\phi \leq 1.13$  in the depth interval of 1280-1350 m TVD (Figure 7c). Moving along the depth to 1390-1750 m TVD,  $S_{hmin} \leq S_v \leq S_{Hmax}$  and  $0.44 \leq A_\phi \leq 1.92$  are observed.

Further constraints on stress magnitudes are made in this borehole using the presence of DITFs between 1741-1745 m on FMI borehole image logs (Figure 7c). The presence of DITFs at 1742.3 m suggests that the  $S_{Hmax}$  should be above the DITF line (Figure 8), where the local hoop stress can be tensile (Equation 3), but also inside the stress polygon with  $\mu =$

0.6. The possible range of  $S_{Hmax}$  and  $S_{hmin}$  constrained using this information lie inside the blue shaded area (Figure 8) and suggest a stress state such that  $S_{hmin} \leq S_v < S_{Hmax}$ .



**Figure 8.** Analysis of stress magnitudes using stress polygon defined by Coulomb friction law with a friction coefficient ( $\mu$ ) of 0.4 and 0.6, and 0.8 in borehole Kauhauroa-5 at depth of 1742.3 m where DITFs are observed. The green shaded area represents  $\sigma_3$  range estimated from the average  $\sigma_3 : S_v$  ratio of 0.95 along the HSM. The blue shaded area represents  $S_{Hmax}$  range which local hoop stress is tensile and DITFs are formed. NF: normal faulting, SS: strike-slip faulting, TF: thrust faulting; DITF: drilling induced tensile fracture.

#### Tuhara-1A borehole

A  $\sigma_3 : S_v$  ratio of 0.63 is determined from  $\sigma_3$  value calculated using LOT data at 590.8 m TVD (Table 1). The  $\sigma_3$  values in the deeper part of the borehole are calculated from the average  $\sigma_3 : S_v$  ratio of 0.95 along the HSM, and are further constrained by the lower limit of  $\sigma_3$  value determined from FIT value of 40.6 MPa at 2149.5 m TVD. The  $S_{Hmax} : S_v$  ratio of 0.95-1.81 is determined for borehole Tuhara-1A using the  $S_{Hmax}$  values calculated from the lower and upper value of UCS.

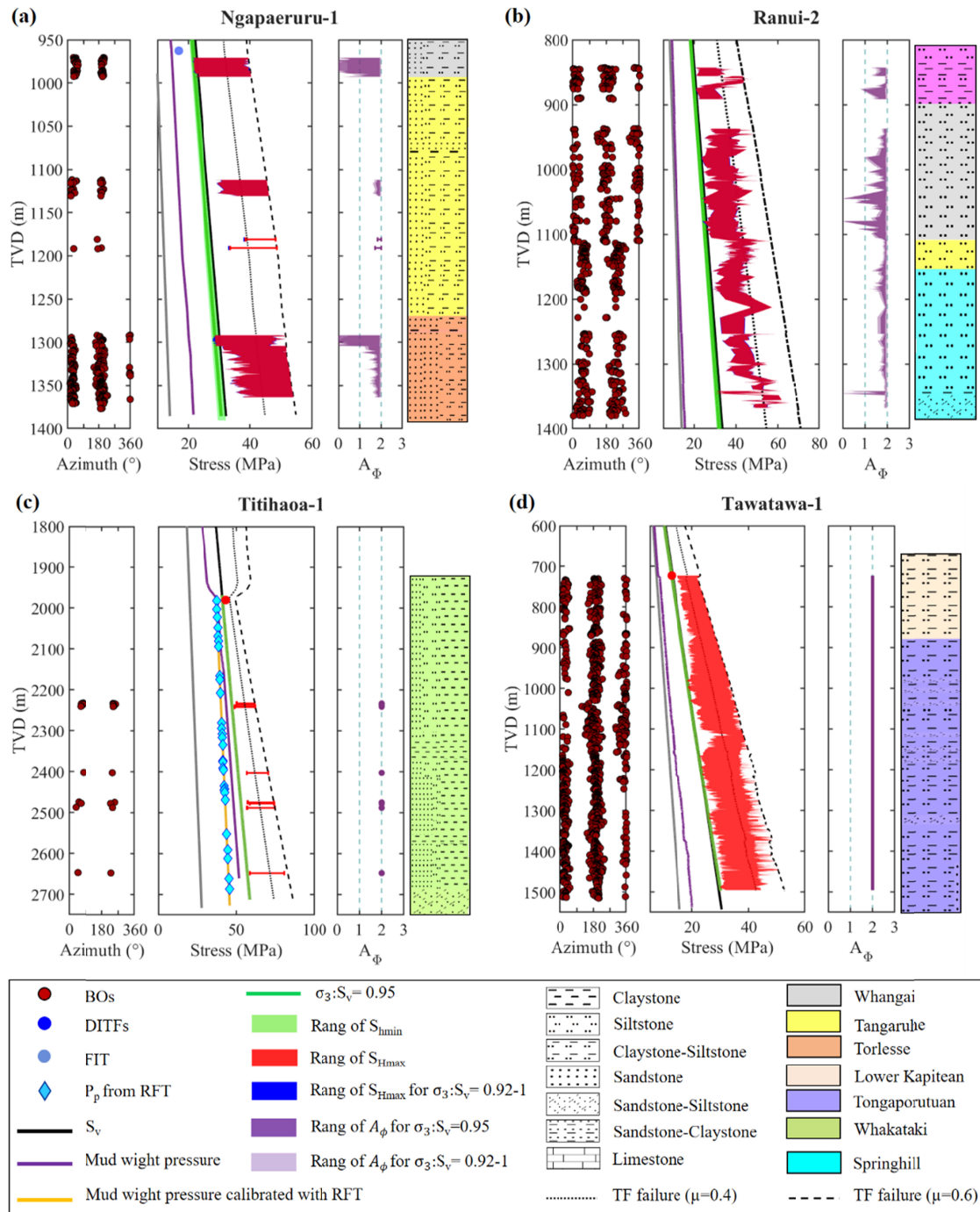
The  $S_{Hmax} : S_v$  ratio of 0.95-1.81 and the  $\sigma_3 : S_v$  ratio of 0.95 indicate a dominant stress regime such that  $S_{hmin} \leq S_v \leq S_{Hmax}$  (Figure 3). A  $0 \leq A_\phi \leq 1.95$  is determined from calculated stress magnitude data in this borehole.

#### 4.2.2 Southern HSM

##### Ngapaeruru-1 borehole

The  $\sigma_3$  values in this borehole are calculated from the average  $\sigma_3 : S_v$  ratio of 0.95 along the HSM and are further constrained by the lower limit of  $\sigma_3$  value determined from FIT values of 8.35 and 16.86 MPa at 501.9 and 962.7 m TVD, respectively. The  $S_{Hmax} : S_v$  ratio of 0.95-1.75 is determined for borehole Ngapaeruru-1 using the  $S_{Hmax}$  values calculated from the lower and upper value of UCS.

The  $S_{Hmax} : S_v$  ratio of 0.95-1.75 and the  $\sigma_3 : S_v$  ratio of 0.95 indicate a dominant stress regime such that  $S_{hmin} \leq S_v \leq S_{Hmax}$  (Figure 9a). A  $0 \leq A_\phi \leq 1.94$  is determined from calculated stress magnitude data in this borehole (Figure 9a). The upper limit of  $S_{Hmax}$  magnitudes from the upper values of UCS are constrained by the limits provided by slip on pre-existing faults with  $\mu=0.6$ .



**Figure 9.** The constrained in situ stress profile with depth in (a) Ngapaeruru-1 (b) Tawatawa-1 (c) Titihaoa-1 in the southern HSM. Abbreviations: BO = breakout; FIT = formation integrity test; LOT: leak of test; TF failure: thrust faulting failure;  $\mu$ : friction coefficient; RFT: repeat formation test;  $S_v$ : vertical stress;  $S_{hmin}$ : minimum horizontal stress;  $S_{Hmax}$ : maximum horizontal stress;  $\sigma_3$ : minimum principal stress;  $A_\phi$ : tectonic stress regime index.

### **Tawatawa-1 borehole**

A  $\sigma_3 : S_v$  ratio of 1 is determined from  $\sigma_3$  value calculated using LOT data at 722.5 m TVD (Table 1). The  $S_{Hmax} : S_v$  ratio of 1-1.82 is determined for borehole Tawatawa-1 using the  $S_{Hmax}$  values calculated from the lower and upper value of UCS.

The  $S_{Hmax} : S_v$  ratio of 1-1.82 and the  $\sigma_3 : S_v = 1$  indicate a dominant stress regime such that  $\sigma_3 = S_v \leq S_{Hmax}$  (Figure 9d). The upper limit of  $S_{Hmax}$  magnitudes from the upper values of UCS are constrained by the limits provided by slip on pre-existing faults with  $\mu=0.6$ .  $A_\phi = 2$  is determined from calculated stress magnitude data in this borehole.

### **Titihaoa-1 borehole**

The  $\sigma_3 : S_v$  ratios of 0.86, 0.94, and 1.05 are determined from  $\sigma_3$  values calculated using LOT data at 614, 1585.7, and 1979.8 m TVD in this borehole (Table 1). The  $\sigma_3 : S_v$  ratio of 1.05 at 1979.8 m TVD indicate that  $\sigma_3 = S_v$  at depth intervals of 2200-2700 m TVD. The  $S_{Hmax} : S_v$  ratio of 1.02-1.41 are determined for borehole Titihaoa-1 using the  $S_{Hmax}$  values calculated from the lower and upper value of UCS.

The analysis of  $S_{Hmax}$  magnitudes and  $\sigma_3 : S_v = 1$  at depth intervals of 2200-2700 m TVD indicate a stress regime such that  $\sigma_3 = S_v \leq S_{Hmax}$  (Figure 9c). The upper limit of  $S_{Hmax}$  magnitudes from the upper values of UCS are constrained by the limits provided by slip on pre-existing faults with  $\mu=0.6$ .  $A_\phi = 2$  is determined from calculated stress magnitude data in this borehole.

### **Ranui-2 borehole**

The  $\sigma_3$  profile in this borehole is calculated from the average HSM  $\sigma_3 : S_v$  ratio of 0.95 are further constrained by the lower limit of  $\sigma_3$  value determined from FIT value of 6.35 MPa at 395 m TVD. The  $S_{Hmax} : S_v$  ratio of 0.95-3.12 are determined for borehole Ranui-2 using the  $S_{Hmax}$  values calculated from the lower and upper value of UCS.

The  $S_{Hmax} : S_v$  ratio of 0.95-3.12 and the  $\sigma_3 : S_v$  ratio of 0.95 indicate a dominant stress regime such that  $S_{hmin} \leq S_v \leq S_{Hmax}$  (Figure 9b). A  $0 \leq A_\phi \leq 1.96$  is determined from calculated stress magnitude data in this borehole.

## **5 Discussion**

### **5.1 Shallow HSM tectonics**

Stress magnitudes calculated from borehole data indicate that the  $S_{Hmax} : S_v$  ratios ranging from 0.95-1.81 in the central HSM and 0.95-3.12 in the the southern HSM. Additionally,

585  $\sigma_3:S_v$  ratios of 0.6-1 are measured at depths above 650-700 m TVD, while 0.92-1 are  
586 measured below this depth interval along the HSM. These stress magnitude results reveal that  
587 across the central and southern HSM,  $S_{Hmax}$  is dominantly  $\sigma_1$ , indicating a thrust to strike-slip  
588 faulting regime. The observed dominant thrust to strike-slip faulting regime is consistent with  
589 observed contractional tectonics in the HSM developed by the subduction of the Hikurangi  
590 Plateau beneath the North Island (Barnes et al., 1998; Nicol & Beavan, 2003), and the strike-  
591 slip faulting generated by forearc rotation of the East Coast (Beanland & Haines, 1998;  
592 Litchfield et al., 2014; Nicol et al., 2007; Wallace et al., 2004).

593 Behboudi et al. (2022) report a dominant ENE-WSW shallow crust  $S_{Hmax}$  orientation within  
594 the central HSM, and WNW-ESE or NW-SE  $S_{Hmax}$  orientations for the southern HSM  
595 (Figure 1b). Considering  $\sigma_1=S_{Hmax}$  along the HSM, observed  $S_{Hmax}$  orientations suggest the  
596 contemporary maximum compressional stress switches from subparallel (ENE-WSW) to the  
597 Hikurangi margin in the north and central HSM, to roughly perpendicular (WNW-ESE or  
598 NW-SE) to the Hikurangi margin in the southern HSM. Based on our confirmation here that  
599  $\sigma_1=S_{Hmax}$  along the HSM, it is likely that contemporary tectonics in the central HSM are  
600 dominantly strike-slip, while in the southern HSM, more contractional tectonics may be  
601 expected.

602 The NNE/NE striking faults in the central HSM, while currently inactive, express reverse  
603 dip-slip components to them based on seismic survey data (Western Energy New Zealand,  
604 2001). This tectonic slip is at odds with the contemporary fault strike-parallel  $\sigma_1$  ( $S_{Hmax}$ ). We  
605 suggest here that these central HSM faults formed in an initially contractional stress state  
606 such that  $\sigma_3=S_v$ ,  $\sigma_1=S_{Hmax}$  oriented NW-SE which would have been consistent with the NW-  
607 SE component of Pacific-Australian plate motion. Overtime, this stress state changed from  
608 this contractional state to the modern strike-slip/contractional/contractional-oblique stress  
609 state ( $\sigma_3:S_v=0.92-1$ ,  $\sigma_1=S_{Hmax}$  oriented ENE-WSW).

610 This switch in  $\sigma_1$  orientation overtime and along HSM strike may be explained by (a) long-  
611 term clockwise rotation of the Hikurangi forearc (b) clockwise rotation of the Hikurangi  
612 forearc in conjunction with high shallow crust overpressures and/or mechanical property  
613 variations, and/or (c) along-strike variation in slip behavior in the HSM.

614 Clockwise rotation of the forearc, which accommodates the margin-parallel component of  
615 oblique Pacific-Australian plate motion, drives strike-slip and/or normal faulting within the  
616 onshore portion of the northern and central HSM, and transpressional faulting in the southern

HSM (Figure 5, Fagereng & Ellis, 2009; Nicol et al., 2007; Wallace et al., 2004; Wallace, Fagereng, & Ellis, 2012). Behboudi et al. (2022) suggest that this forearc rotation is likely responsible for generating strike-slip stress state with ENE-WSW  $S_{Hmax} = \sigma_1$  in the central HSM, and contemporary contractional stress state with WNW-ESE/ NW-SE  $S_{Hmax} = \sigma_1$  in the southern HSM. However, our stress magnitude results of  $\sigma_3 : S_v = 0.92-1$  and  $\sigma_1 = S_{Hmax}$  leave a possibility for both strike-slip and contractional stress states to occur across both the central and southern HSM due to poorly constrained UCS values used in this study, a limitation of the study that could be restricted by laboratory rock strength testing of both onshore and offshore HSM lithologies.

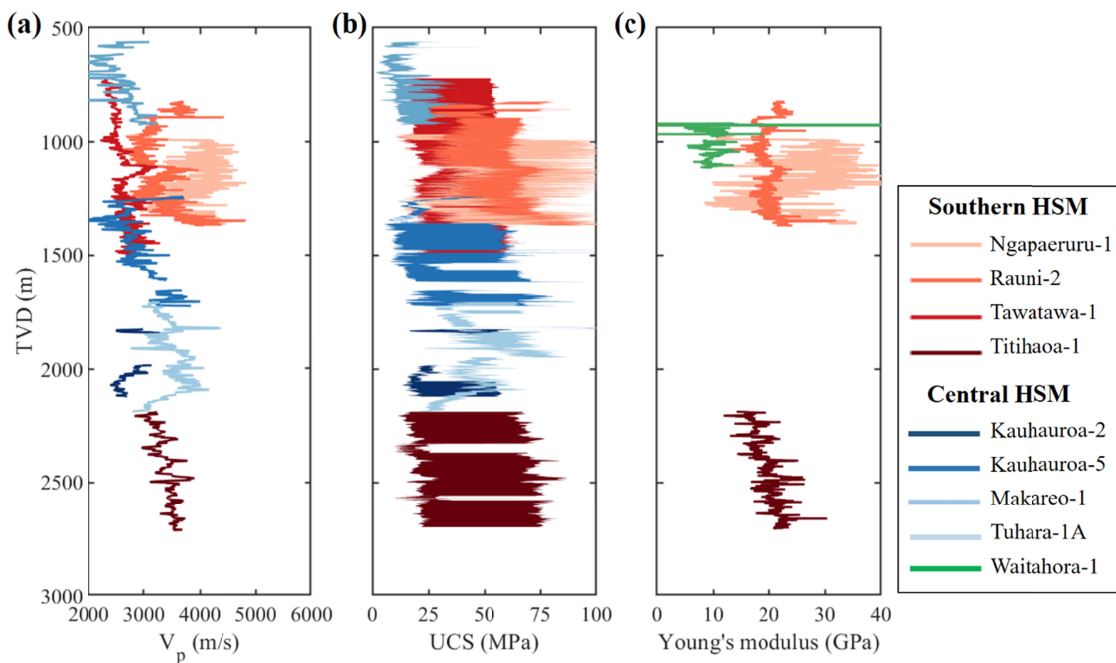
The northern and central HSM have high  $P_p$  based on borehole data (Burgreen-Chan et al., 2016; D. Darby & Funnell, 2001), magnetotellurics (Heise et al., 2019), and seismic tomography (Bassett et al., 2014; Eberhart-Phillips et al., 2017). Overpressure reduces the effective normal stress on fault planes, meaning that the existing NNE/NE striking faults in this region will be able to slip at lower shear stresses. Therefore, as the result of this overpressure, these faults could be less stable, allowing the hangingwall of upper plate faults to move more easily in response to NE-SW forces raised from forearc rotation. In this scenario, forces raised from forearc rotation were able to alter stress state overtime from  $\sigma_3 : S_v = 1$  and  $\sigma_1 = S_{Hmax}$  with NW-SE  $S_{Hmax}$  orientation, compatible with NW-SE component of Pacific-Australian plate motion and old geological structures, to  $\sigma_3 : S_v = 0.92-1$  and  $\sigma_1 = S_{Hmax}$  with ENE-WSW  $S_{Hmax}$  orientation. Similar shallow, high overpressures are not observed in the hangingwalls of upper plate faults in the onshore of the southern HSM. Therefore it is possible that the NE-SW forces resulted from forearc rotation alone are insufficient to exceed the fault shear resistance and change the orientation of  $\sigma_1$  away from the NW-SE component of Pacific-Australian plate motion, however they may have been high enough to play a role in reducing  $\sigma_3$  magnitudes to the point that they become  $\leq S_v$ , resulting in a more transtensional tectonic regime overtime.

The mechanical properties of fault gauges and formations hosting faults (friction coefficient and rock strength) can play a role in controlling upper plate tectonic stresses (Mantovani et al., 2000; Marotta et al., 2002). Reiter (2021) investigated the impact of physical and elastic parameter contrasts on  $S_{Hmax}$  orientation and proposed that contrasts in Young's modulus can introduce  $S_{Hmax}$  rotations up to  $78^\circ$ , with larger stress rotations occurring within the softer lithologies. Behboudi et al. (2022) proposed that basement uplift in the southern HSM may introduce lateral geomechanical heterogeneities and variations in rock and sediment physical



650 properties along the HSM which may influence  $S_{Hmax}$  orientations. Such that clay and sand-  
 651 siltstone sediments (Miocene to present), where our stress data are calculated, in the upper  
 652 plate of the central HSM may geomechanically differ from clay and sand-siltstone sediments  
 653 (Miocene to Cretaceous) in the onshore of the southern HSM. Therefore, we analyzed  
 654 physical properties of aforementioned sediments and discovered that P-wave velocity ( $V_p$ ),  
 655 UCS ranges, and Young's modulus in the central HSM are lower than the onshore of southern  
 656 HSM (Figure 10). In this scenario,  $S_{Hmax}$  orientations in the sediments of the central HSM,  
 657 which have lower UCS and Young's modulus, could be easily reoriented in response to long-  
 658 term forces such as forearc rotation compared to southern HSM. This theory, however, does  
 659 not explain why the offshore boreholes in the southern HSM have not reoriented in response  
 660 to long-term forearc rotation forces, while having comparable  $V_p$ , UCS range, and Young's  
 661 Modulus to the boreholes in the central HSM.

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**Figure 10.** Graph shows (a) p-wave velocity, (b) rock strength (UCS), and (c) Young's modulus in  
 clay and sand-siltstone sediments as a function of depth across the central and southern HSM.

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This along-strike variation in contemporary stress state is spatially consistent with north to  
 south variation in slip behavior along the Hikurangi subduction interface (**Figure 1a**). In the  
 northern and central HSM, the subduction interface is largely creeping and experiences  
 shallow (<15 km), episodic slow slip events that extend offshore and possibly to the trench.  
 At the southern HSM the plate interface is strongly interseismically locked to ~30 km depth,  
 and is currently accumulating elastic strain in the surrounding crust (Wallace, 2020). Some

studies suggest that SSEs can release the amount of energy equivalent to a  $M_w$  6.5–8 earthquakes (Dixon et al., 2014; Wallace, Beavan, et al., 2012). In the central HSM the recurring SSEs and frequent earthquakes may release energy overtime such that the normal to shear stress ratio on pre-existing faults has changed in a way that make it easier to slip in response to forces deriving from long-term forearc rotation. While stress accumulation due to locked nature of the southern HSM, don't allow the normal to shear stress ratio change considerably on the existing NNE/NE striking compressional faults and make it difficult for the hanging wall of these faults to slip in response to forearc rotation forces; therefore stress state has not changed overtime in the southern HSM. However, the static stress drop of SSEs is estimated to range 0.01–1.0 MPa (Gao et al., 2012). Given that the contemporary  $\sigma_3$ :  $S_v \approx 0.95$  and  $S_v - \sigma_3$  ranges between 0–3 MPa (for depths less than 3 km), these SSEs should have existed in the central HSM for more than 20 years such that they were able to release energy in order of 3 MPa (for depths less than 3 km) to change the initial  $\sigma_3$ :  $S_v = 1$  to the contemporary  $\sigma_3$ :  $S_v = 0.92$ –1 and reorient the  $S_{Hmax}$  orientation from NW-SE to ENE-WSW in this region. However, further research and modeling are required to determine and quantify the initial stress state and whether the amount of stress released during SSEs in the central HSM was sufficient to support such a theory.

## **5.2 Extensional tectonics within the HSM forearc**

There are locales in the central and southern HSM where stress magnitude determination suggests a normal faulting regime ( $\sigma_3:S_v < 1$  and  $0 \leq A_\phi \leq 1$ ). Also  $\sigma_3:S_v < 1$  where  $\sigma_3$  calculated from LOT data is observed for 13 tests conducted at depth intervals anywhere from  $\approx 102$  to 3611 m TVD in northern and central HSM boreholes (Table 1), and from 3 tests at depth intervals of  $\approx 357$ –1586 m TVD in southern HSM boreholes (Table 1). Several factors can result in localized normal faulting regime at subduction margins including 1) uncertainties in calculated UCS values and/or  $\sigma_3$  magnitudes used to determine stress states in this study, 2) the presence of local, active normal faults, and 3) fluctuations in stress magnitudes modulated by seismic cycles.

### **5.2.1. Uncertainties in calculated UCS values and $\sigma_3$ magnitudes**

Estimations of  $S_{Hmax}$  magnitudes are highly sensitive to the UCS values used, particularly when UCS is determined from empirical relationships not constrained by laboratory testing (Zoback, 2007). Due to lack of direct UCS data in this region, and a lack of empirical relationships for the formations of this region to determine UCS from other rock properties,

this study relied on the use of a range of empirical relationships developed elsewhere to generate a low and high limit for UCS at the HSM. These UCS ranges were then used to generate the lowest and highest limits of  $S_{Hmax}$  magnitude. When the lowest limit of UCS is used it can result in a potentially extensional stress state such that  $S_{hmin} \approx S_{Hmax} \leq S_v$ . As such, the uncertainty in calculated UCS values, and the resulting potential errors it can introduce into a stress model for the HSM, highlight the importance of dedicated laboratory tests for developing robust empirical relationships for UCS in the HSM region, and subduction regions like this, where stress is a critical geological consideration for hazard and resource management.

Inaccuracy involved in LOT measurements (section 3.3) along with the lack of detail reported on LOT results introduces an unknown level of uncertainty on estimated  $\sigma_3$  magnitudes, and hence on estimated  $\sigma_3:S_v$  ratios using this data. Additionally, lack of LOT data along each borehole necessitates the estimation of  $\sigma_3$  profiles from the average  $\sigma_3:S_v = 0.95$ , which also carries uncertainty. As a result, we recognize the potential impact this has on calculations of  $S_{Hmax}$  magnitudes here, as well as on any interpretations of regional stress state and tectonics. To investigate the potential effect of  $\sigma_3$  uncertainties on  $S_{Hmax}$  calculations, we use both the lower and upper limits of  $\sigma_3$  values calculated from  $\sigma_3:S_v = 0.92-1$ , BO widths, and the lower and upper boundary of UCS values. This analysis reveals that the  $\sigma_3$  magnitude uncertainties at the scale explored here have little influence on  $S_{Hmax}$  magnitude calculations ( $\pm 3.5$  Mpa) and hence do not change our findings about the stress regime and tectonics within the HSM (blue areas in Figure 3; Figure 7b,c; Figure 9a,b).

### 5.2.2 Presence of active normal faults

Extensional structures are common within the overriding plate of many subduction margins (Loveless et al., 2010; Moore et al., 2013). Normal faults in subduction zones are often attributed to gravitational instabilities associated with subduction erosions and subsidence, density imbalances produced by forearc uplifts, strain releases during earthquake cycles, and flexural rigidity of the subduction interface (Barnes & Nicol, 2004; Collot et al., 1996; Loveless et al., 2005; Park et al., 2002; Sacks et al., 2013). Within the HSM, localized extensional stresses within the overriding plate are suggested to result from processes such as slab rollback, forearc rotation (Nicol et al., 2007; Wallace et al., 2004), subduction erosion and related subsidence, gravitational collapse due to forearc uplift, and growth of bending-

moment faults (Barnes & Nicol, 2004; Chanier et al., 1999; Upton et al., 2003; Walcott, 1987; Wallace, Fagereng, et al., 2012).

The  $\sigma_3$  magnitude of 8.4 MPa measured from LOP in borehole Tuhara-1/1A (590.8 m TVD; Table 1) is lower than  $\sigma_3$  values of 8.95 MPa estimated from normal faulting failure with a friction coefficient of 0.6 (Equation 7). This lower  $\sigma_3$  magnitude may indicate there are active normal faults at this depth along this borehole. In addition, borehole Tuhara-1A is located within the Tuhara anticline structure, formed by contractional stresses resulting from two blind thrust faults beneath the structure (Western Energy New Zealand, 1999). Our stress magnitudes and HRT's (2000) analysis suggests that the Tuhara structure currently experiences a dominant strike-slip faulting regime ( $S_{hmin} \leq S_v \leq S_{Hmax}$ ;  $1 \leq A_\phi \leq 2$ ) along the majority of the borehole, interspersed with intervals of normal faulting regime ( $S_{hmin} \leq S_{Hmax} \leq S_v$ ;  $0 \leq A_\phi < 1$ ) mainly within the 1700-1820 m and 2100-2145 m TVD depth interval (Figure 3). A prominent feature of the Tuhara structure, as indicated by seismic reflection profiles, is observation of relatively short steep east- and west-dipping normal faults throughout Pliocene and Miocene successions (Western Energy New Zealand, 1999; Barnes et al., 2002). Accordingly, we relate the appearance of normal stress states in our data to the normal structures that develop as part of the larger compressional structural architecture of this borehole site, and not due to the previously discussed uncertainties in the calculated UCS and/or  $\sigma_3$  magnitude values. This could particularly be the case where both the calculated lower and upper limit of  $S_{Hmax}$  magnitudes are less than  $S_v$  (for example at 1700-1820 mTVD in Tuhara-1A; Figure 3).

### 5.2.3. Stress field fluctuations modulated by seismic cycling

Fluctuations in stress magnitudes can be caused by seismic cycling. It has been reported that earthquake events generate stress drops of 0.01 to 100 MPa, depending on the rheology, roughness of fault, geometry of slip area, and heterogeneous stress fields (Allmann & Shearer, 2009; Baltay et al., 2011; Candela et al., 2011; Cocco et al., 2016; Oth et al., 2010). The observation of localized normal faulting regimes in the HSM may be related to seismic cycling in the region. The normal faulting regimes observed along central HSM boreholes Kauhauroa-2 (1980-2075 m TVD), Kauhauroa-5 (1330-1345 m TVD), and Tuhara-1A (1700-1820 m TVD) occur where  $S_{Hmax}$  and  $S_v$  are very similar and are greater than  $S_{hmin}$  (Figure 3, Figure 7b & 7c). In such stress state scenarios, a post-seismic stress drop of only a few MPa after great earthquakes or frequent moderate earthquakes in the HSM region could perturb the

767 delicately balanced stress magnitudes surrounding these boreholes, switching  $\sigma_1 = S_{Hmax}$  to  
768  $\sigma_1 = S_v$  i.e. from a reverse/strike-slip to a normal stress state, accompanied by small rotations  
769 in the  $S_{Hmax}$  orientation.

## 770 **6 Conclusions**

771 This work represents the first comprehensive determination of the *in-situ* stress state of the  
772 HSM margin using available borehole data. We found a  $\sigma_3:S_v = 0.6-1$  at depths above 650-  
773 700 m TVD, while  $\sigma_3:S_v = 0.92-1$  below this depth interval along the HSM. Stress  
774 magnitudes calculated from borehole data indicate that the  $S_{Hmax}:S_v$  ratios ranging from 0.95-  
775 1.81 in the central HSM and 0.95-3.12 in the the southern HSM. These principal stress  
776 magnitude results indicate a  $\sigma_1 = S_{Hmax}$  and a thrust to strike-slip faulting regime across the  
777 both central and southern HSM. The pre-existing NNE/NE striking reverse faults along the  
778 both central and southern HSM infer that stress regime was initially in a contractional state  
779 such that  $\sigma_3:S_v = 1$ ,  $\sigma_1 = S_{Hmax}$ , and a dominant NW-SE  $S_{Hmax}$ , consistent with NW-SE  
780 component of Pacific-Australian plate motion. Taking contemporary stress state of  $\sigma_1 = S_{Hmax}$   
781 and ENE-WSW  $S_{Hmax}$  orientation and initial stress state into account in the central HSM,  
782 these observations suggest that the compressional regime has shifted from subparallel to  
783 perpendicular to the NW-SE Hikurangi convergence direction overtime in this region.  
784 Variation of the central HSM stress state overtime may result from forces arising from  
785 Hikurangi forearc rotation either by itself or facilitated by the upper plate, shallow, high  
786 overpressures in the central HSM. Along-strike variation in slip behavior may also play a role  
787 by releasing stress overtime due to SSEs and frequent earthquakes, hence changing the stress  
788 state in the central HSM, while in the southern HSM, the modern WNW-ESE/ NW-SE  $\sigma_1$   
789 ( $S_{Hmax}$ ) remains subparallel to NW-SE Hikurangi convergence direction overtime, may reflect  
790 the interseismic locked nature of the plate interface. Finally, stress determination highlights  
791 localized normal stress states within the HSM forearc interpreted to be due to processes such  
792 as the presence of localized active normal faults or fluctuations in stress magnitudes  
793 modulated by seismic cycles. The determination of HSM *in-situ* stresses in this study will  
794 provide an invaluable tool for improving our understanding of the stability of upper plate  
795 faults and will facilitate more quantitative efforts to assess the seismic hazard potential of the  
796 HSM that will support of disaster risk reduction plans.

## 797 **Acknowledgments**

This publication has emanated from research conducted with the financial support of Science Foundation Ireland (SFI) under Grant Number 17/RC-PhD/3481 and co funded by Geological survey of Ireland (GSI). The authors would like to thank the New Zealand Petroleum and Minerals group (NZPM) within the Ministry for Business, Innovation and Employment (MBIE) for providing access to borehole data and supporting materials for this study. We also thank Schlumberger for providing academic licenses for Techlog 2018.1 to University College Dublin and University of Liverpool. We thank MathWorks for providing academic licenses for MATLAB to University College Dublin. We thank Esri for providing the academic license of ArcGIS Pro to University College Dublin. For the purpose of Open Access, the author has applied a CC BY public copyright license to any Author Accepted Manuscript version arising from this submission.

### Data Availability Statement

This research used data provided by the New Zealand Petroleum and Minerals group (NZPM) within the Ministry for Business, Innovation and Employment (MBIE). The borehole image logs used in this paper can accessed through MBIE's online free database (<https://data.nzpam.govt.nz/GOLD/system/mainframe.asp>). Borehole breakout measurements presented in this study can be accessed at [https://github.com/BehboudiEffatGeo/StressCharacterization\\_HSM.git](https://github.com/BehboudiEffatGeo/StressCharacterization_HSM.git) and <https://doi.org/10.5281/zenodo.7450966>.

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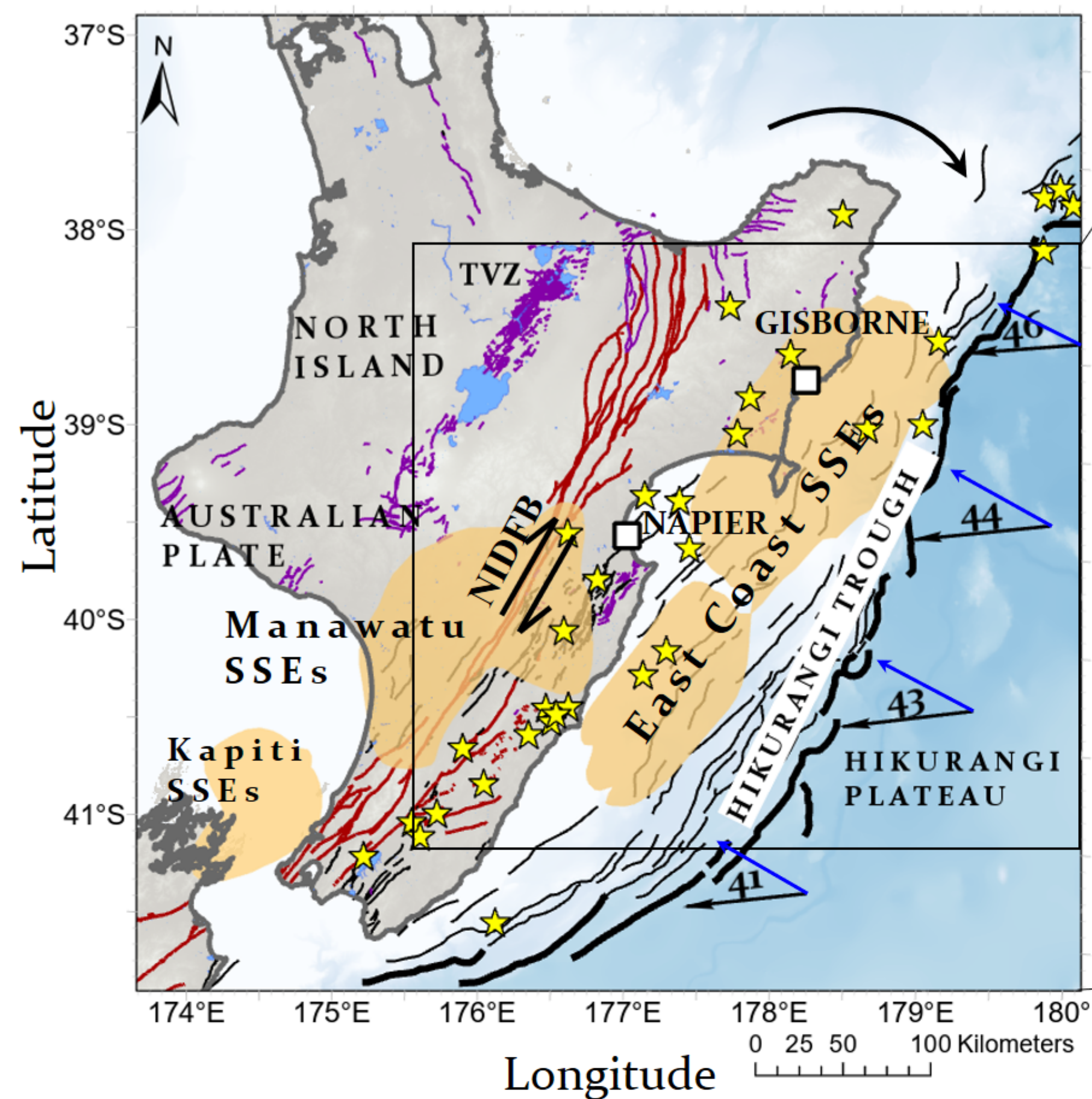
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Figure 1.



(a)



(b)

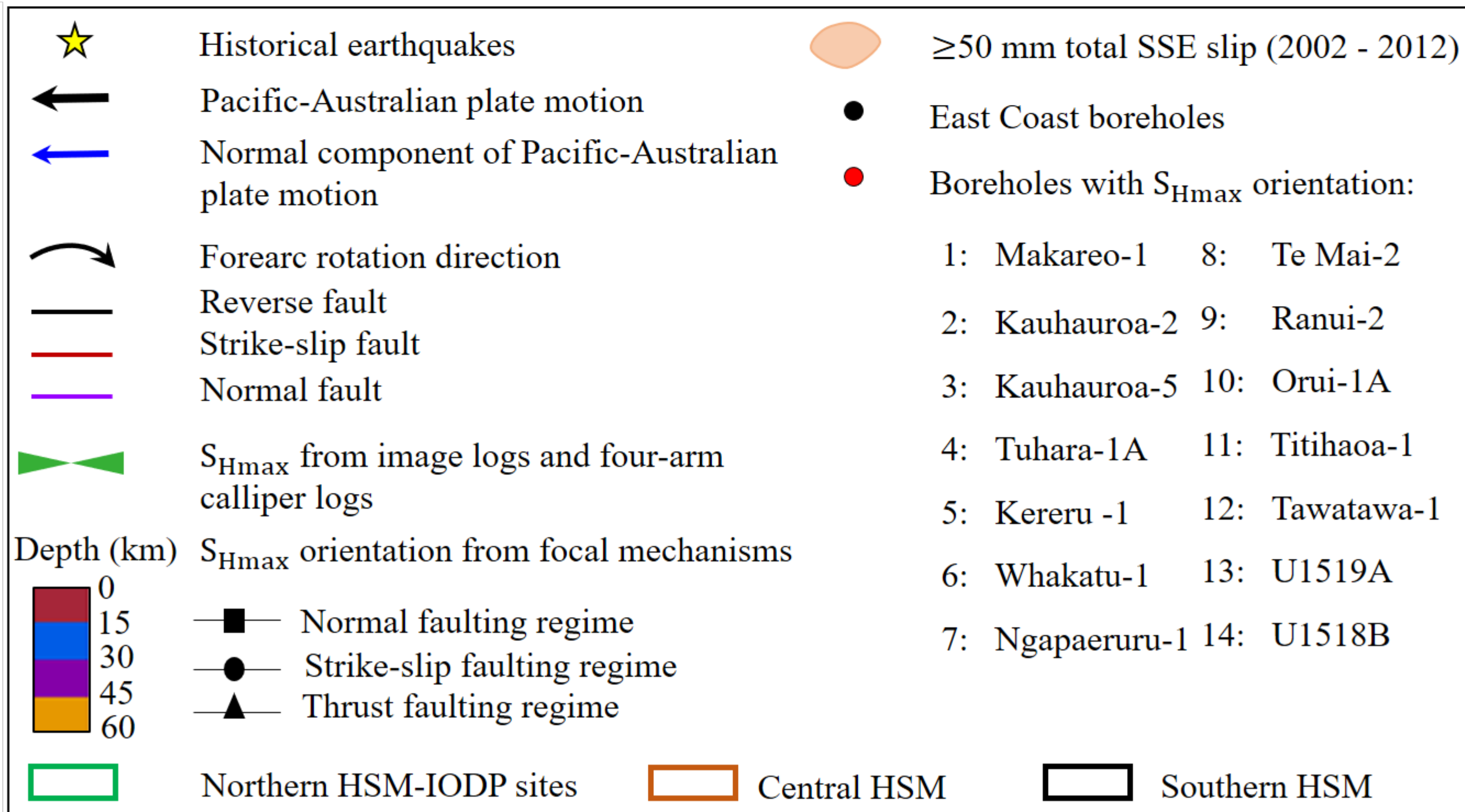
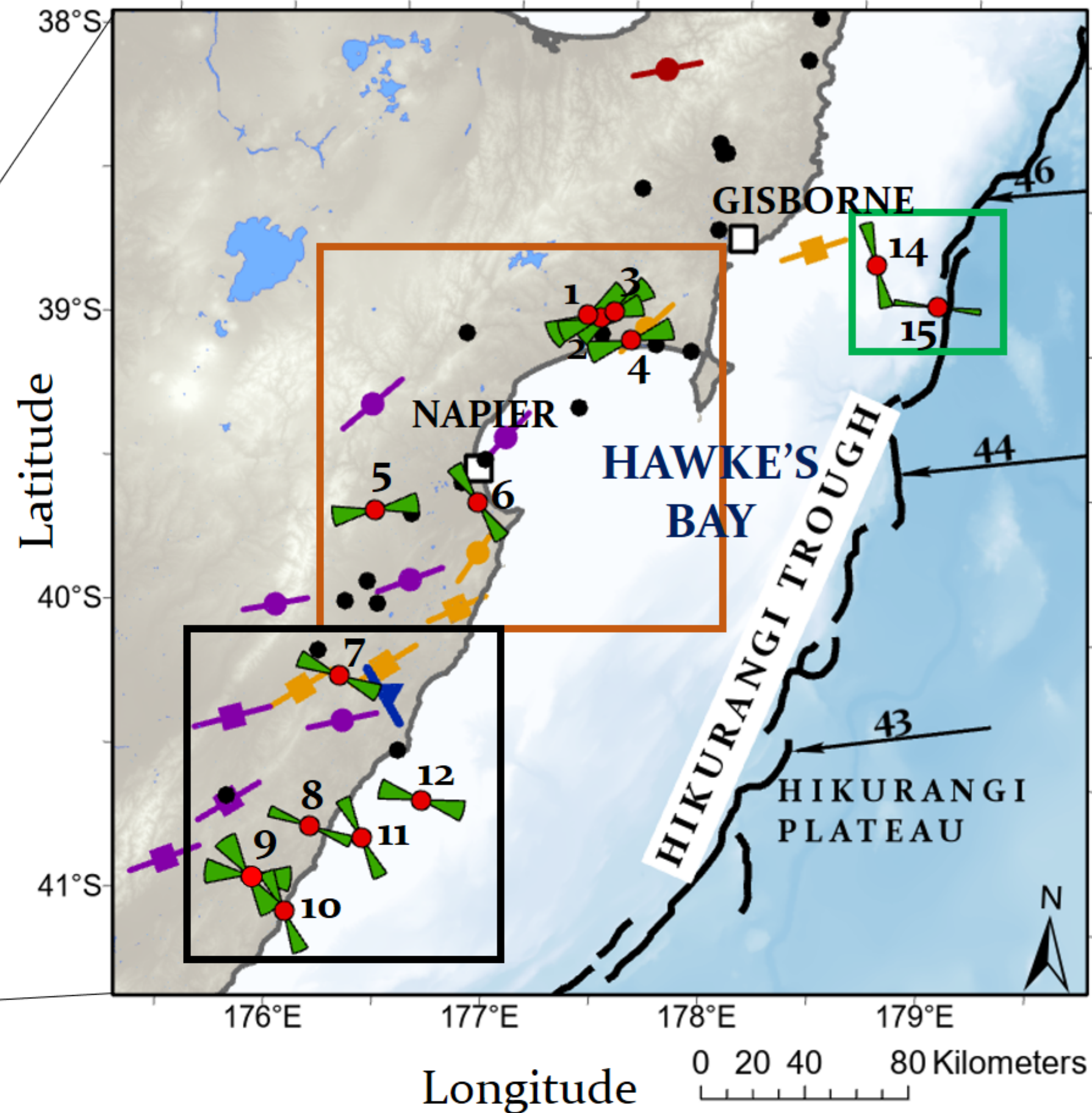




Figure 2.

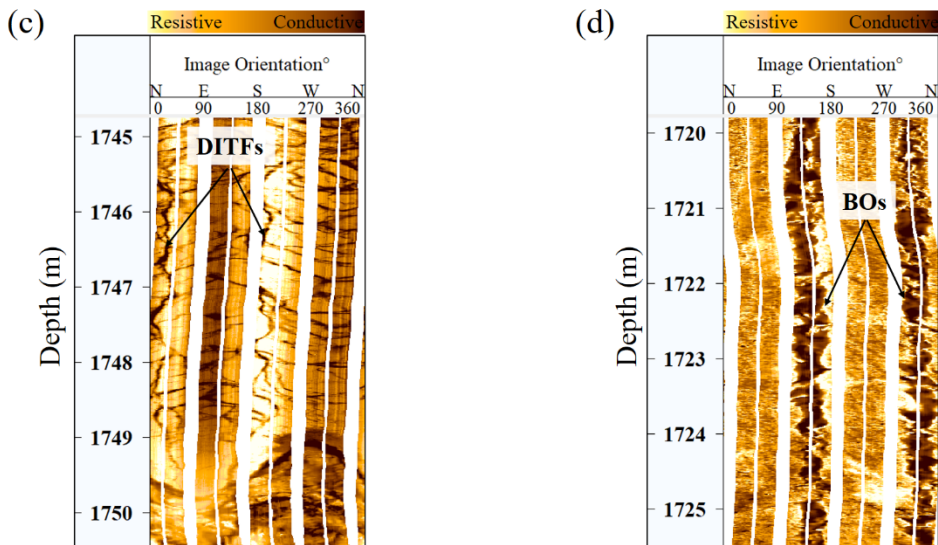
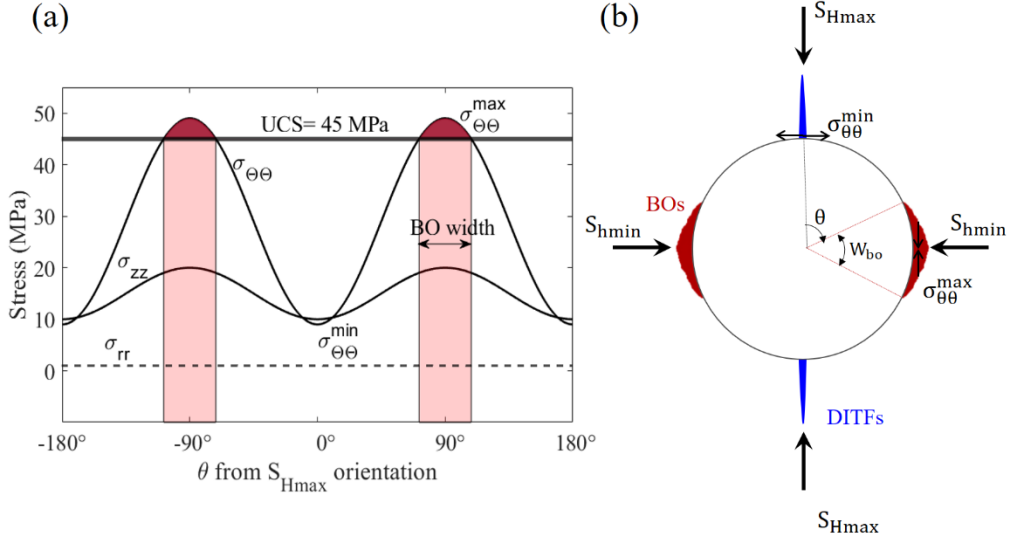


Figure 3.

# Tuhara-1A

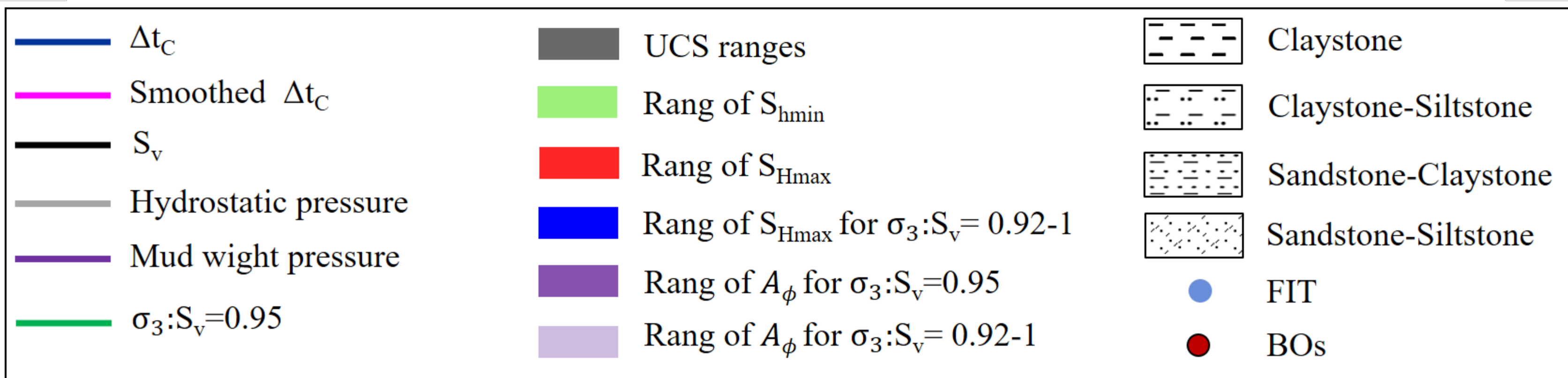
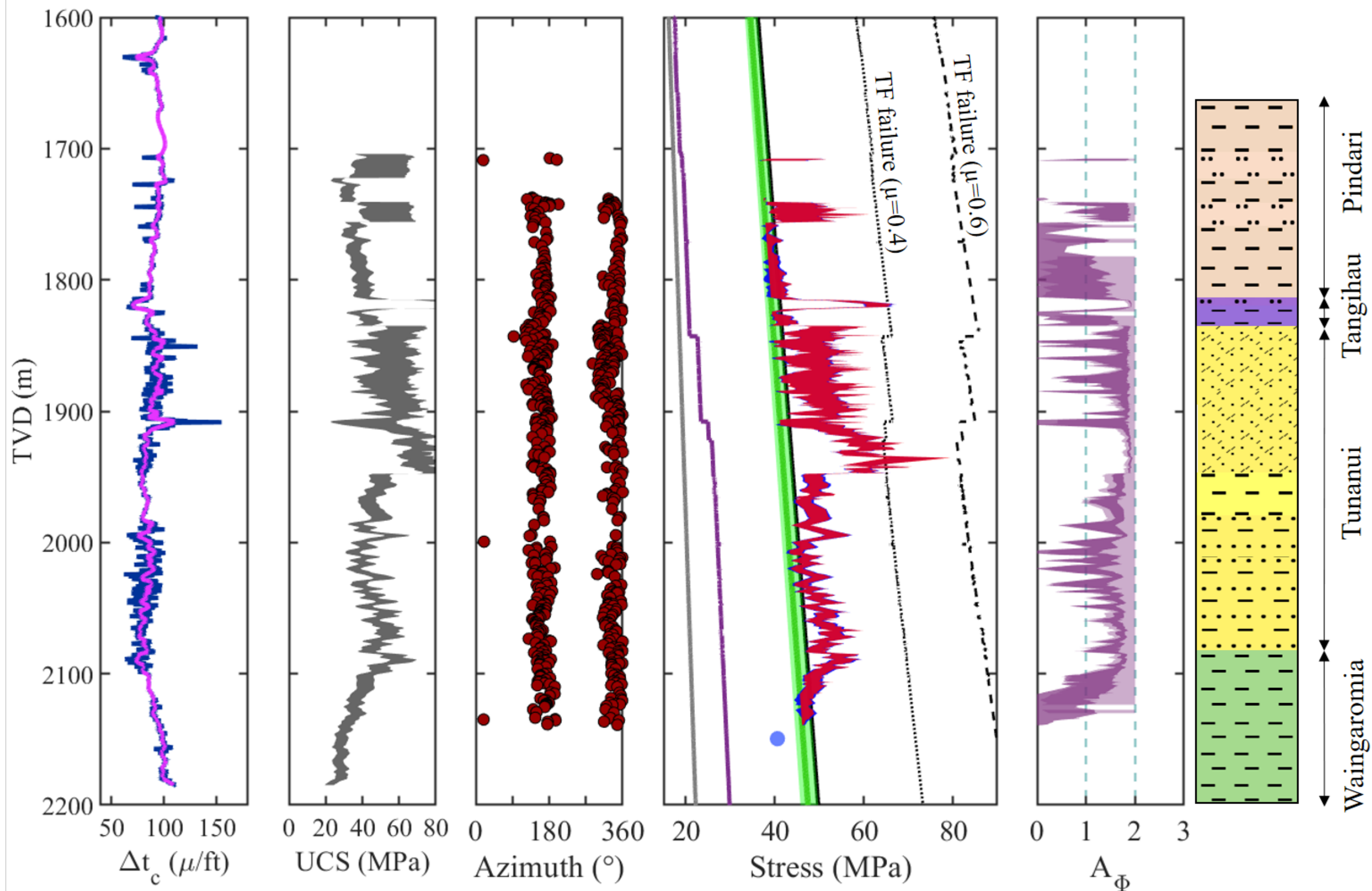


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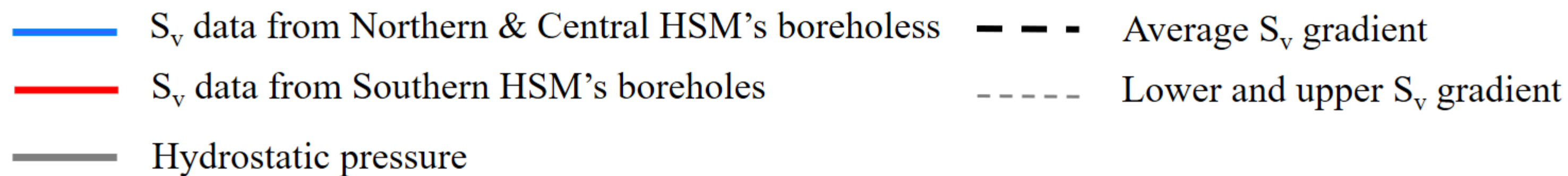
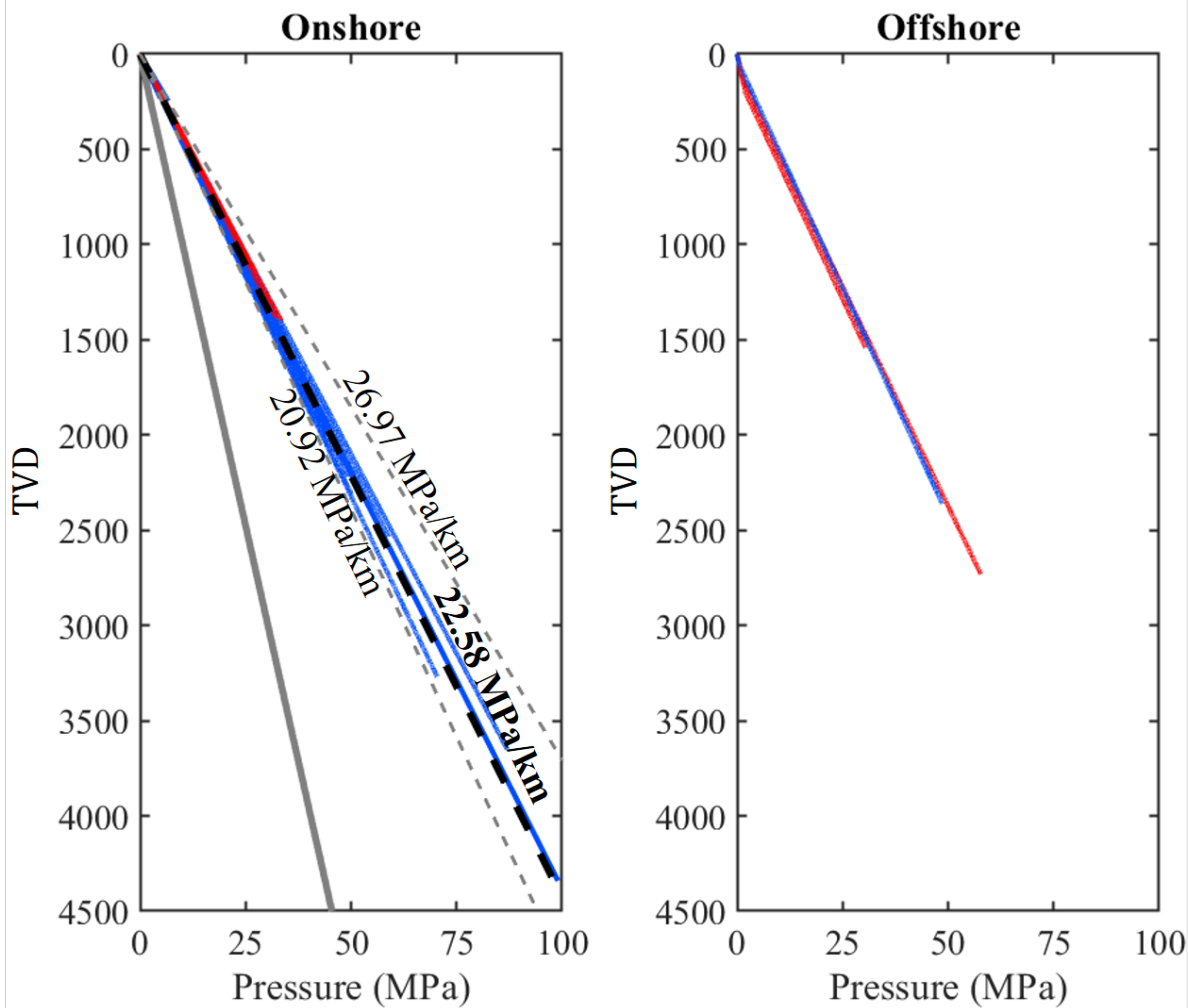
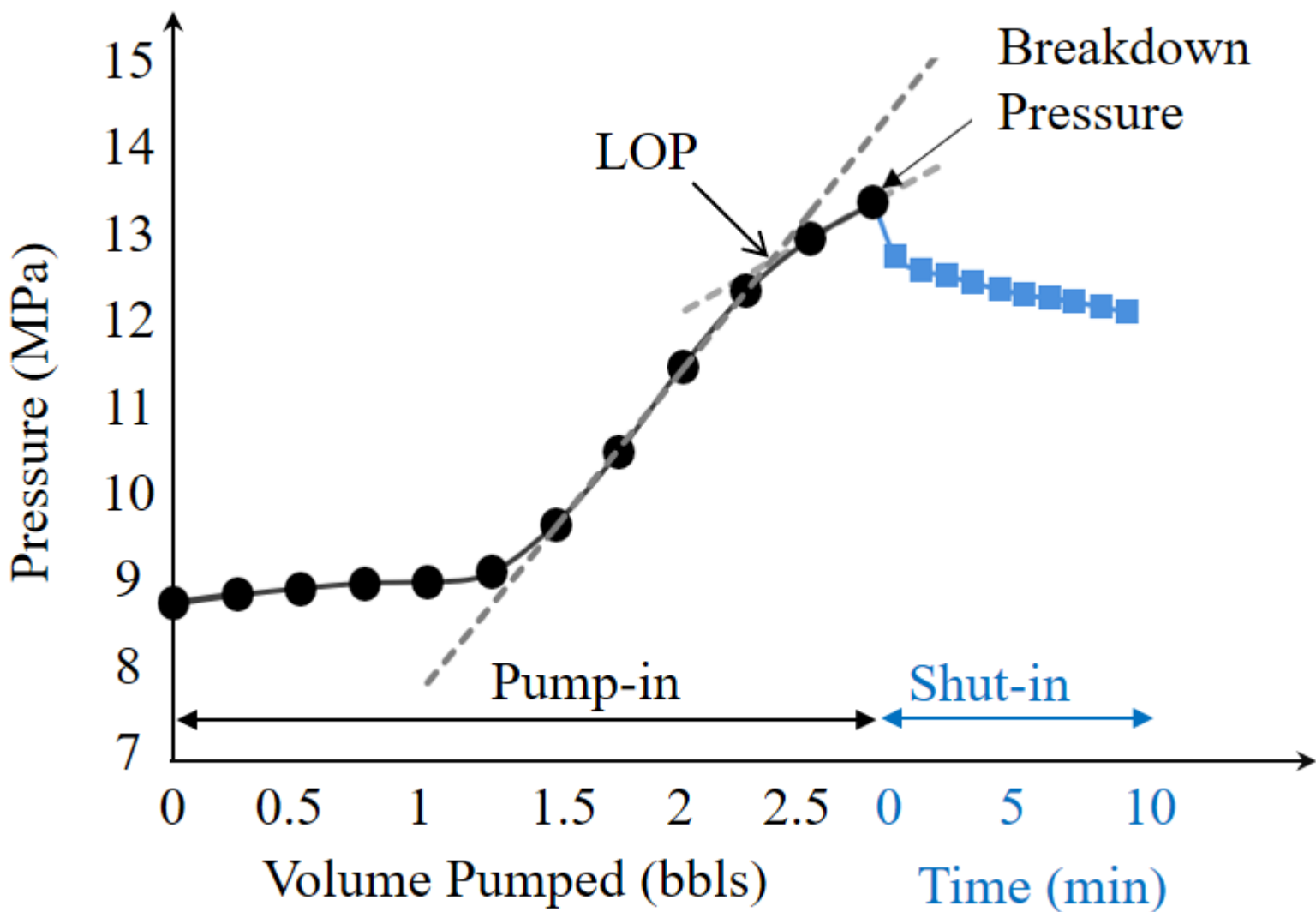


Figure 5.







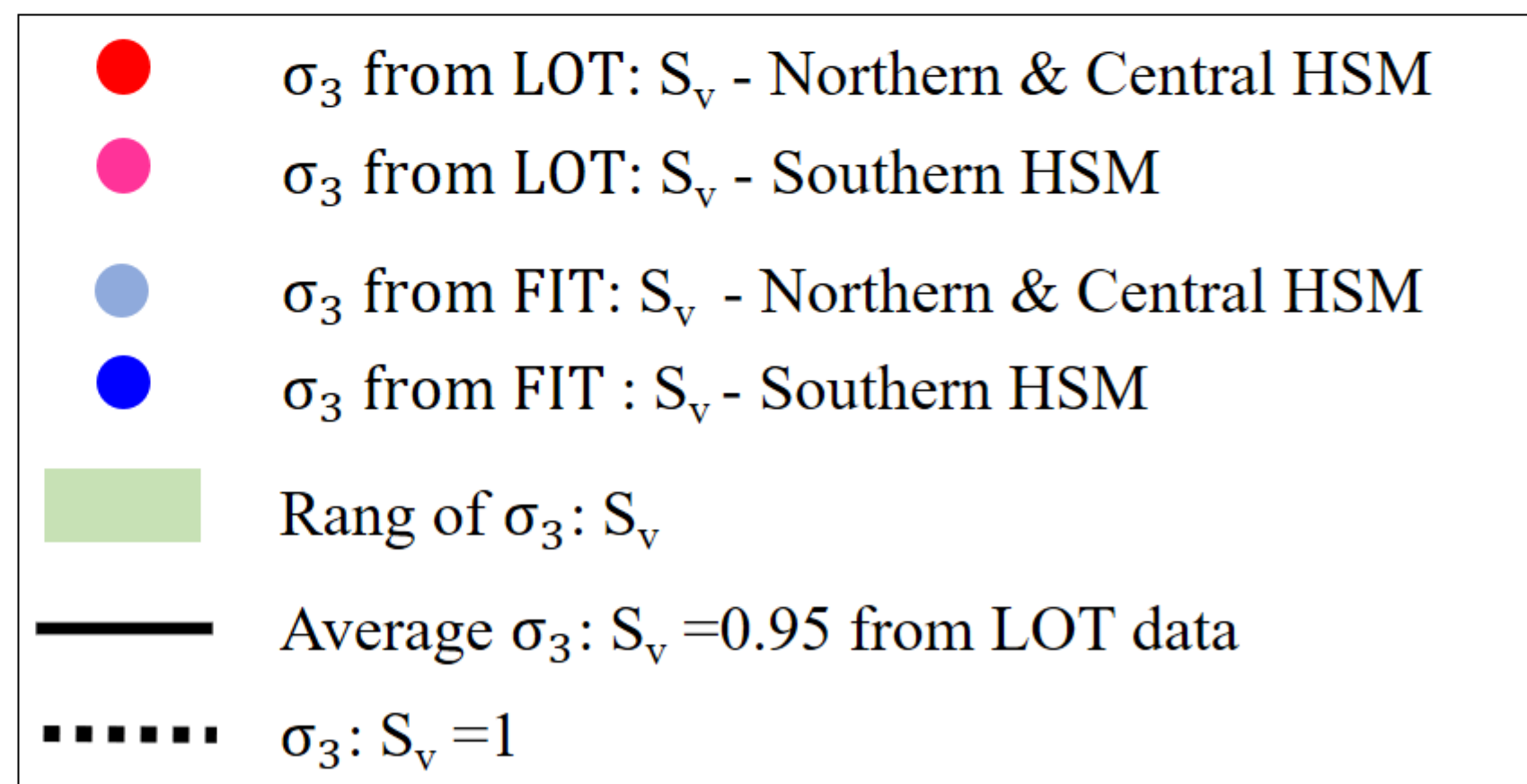
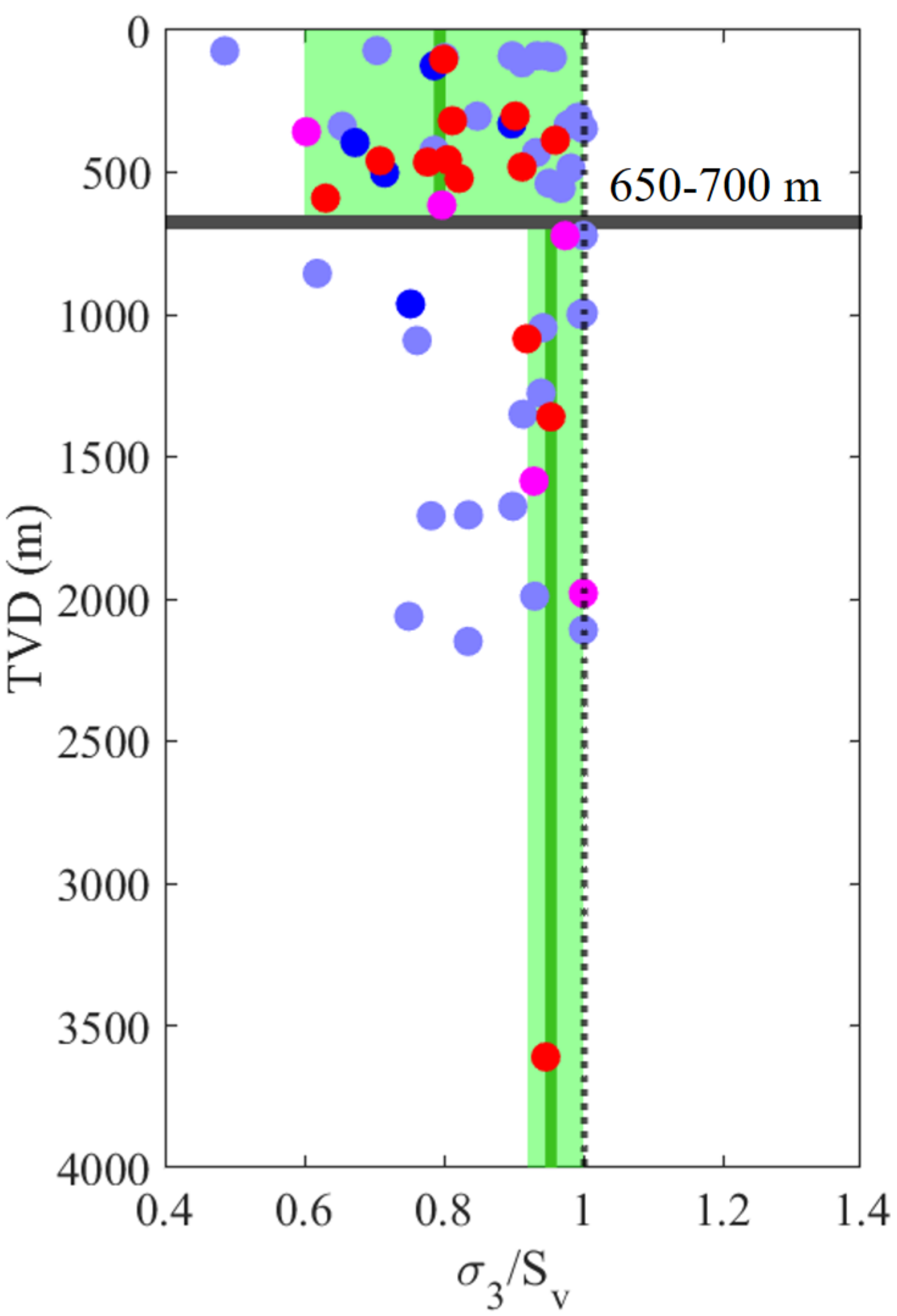


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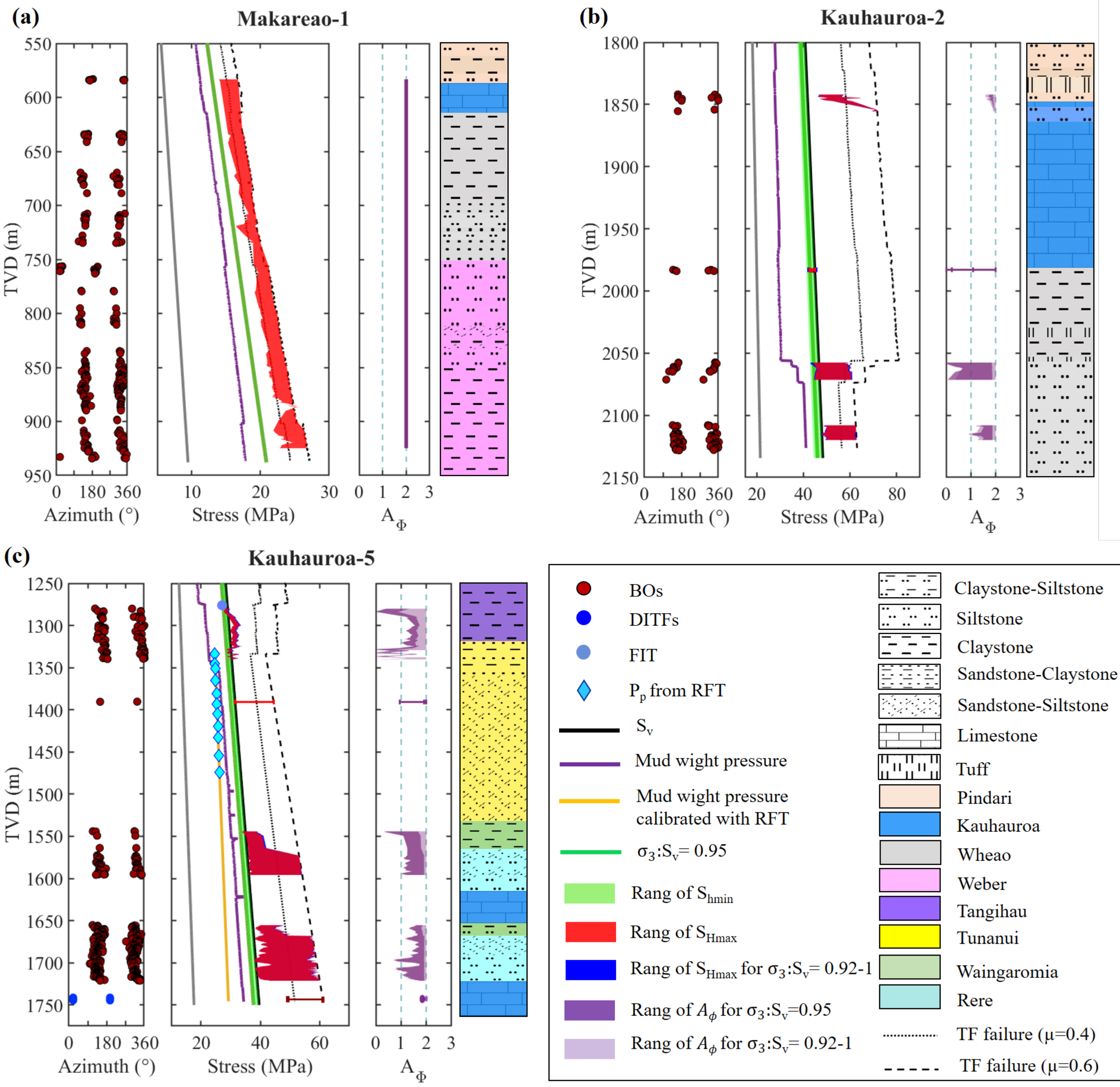


Figure 8.

# Kauhauroa-5

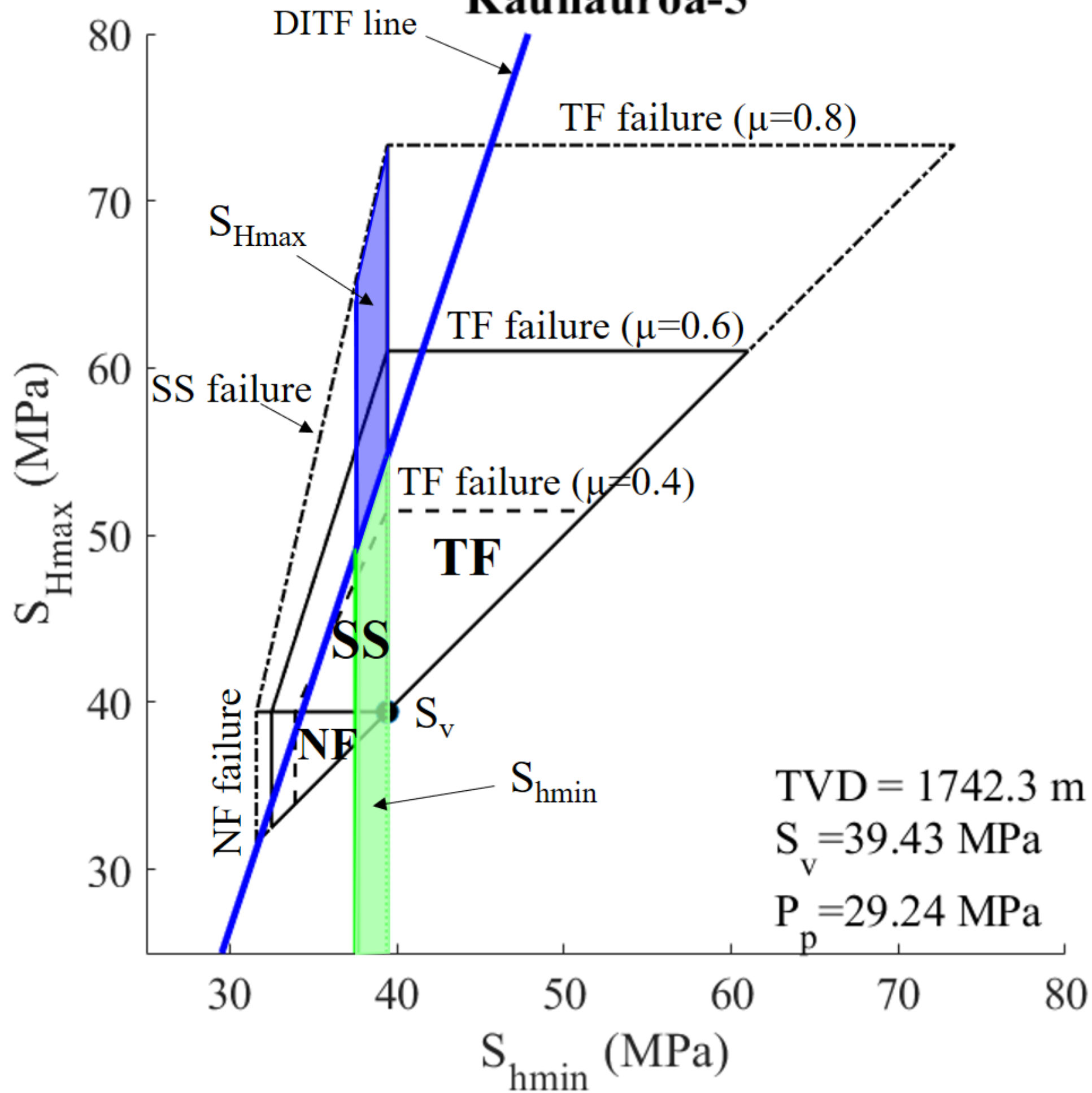


Figure 9.



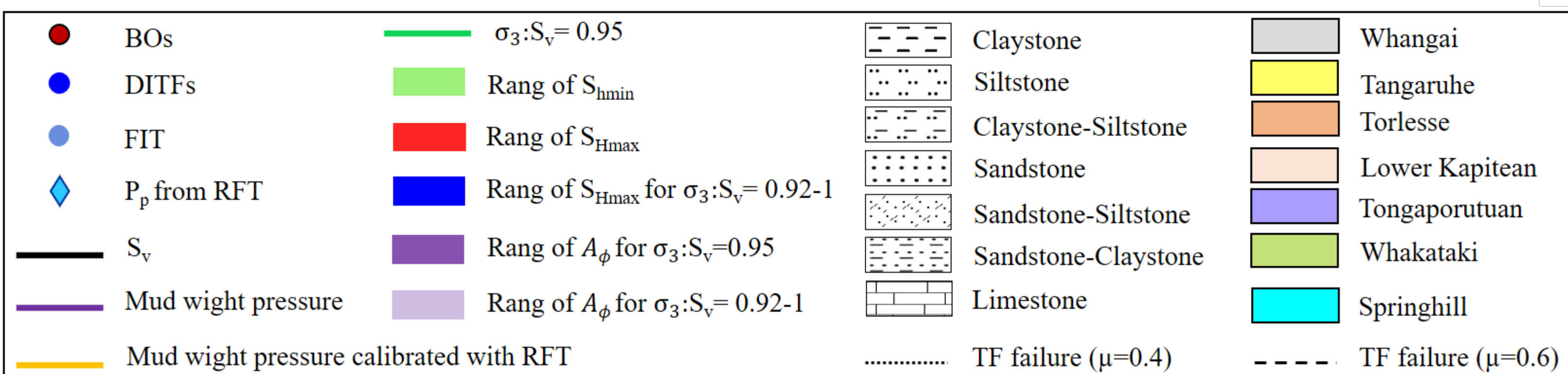
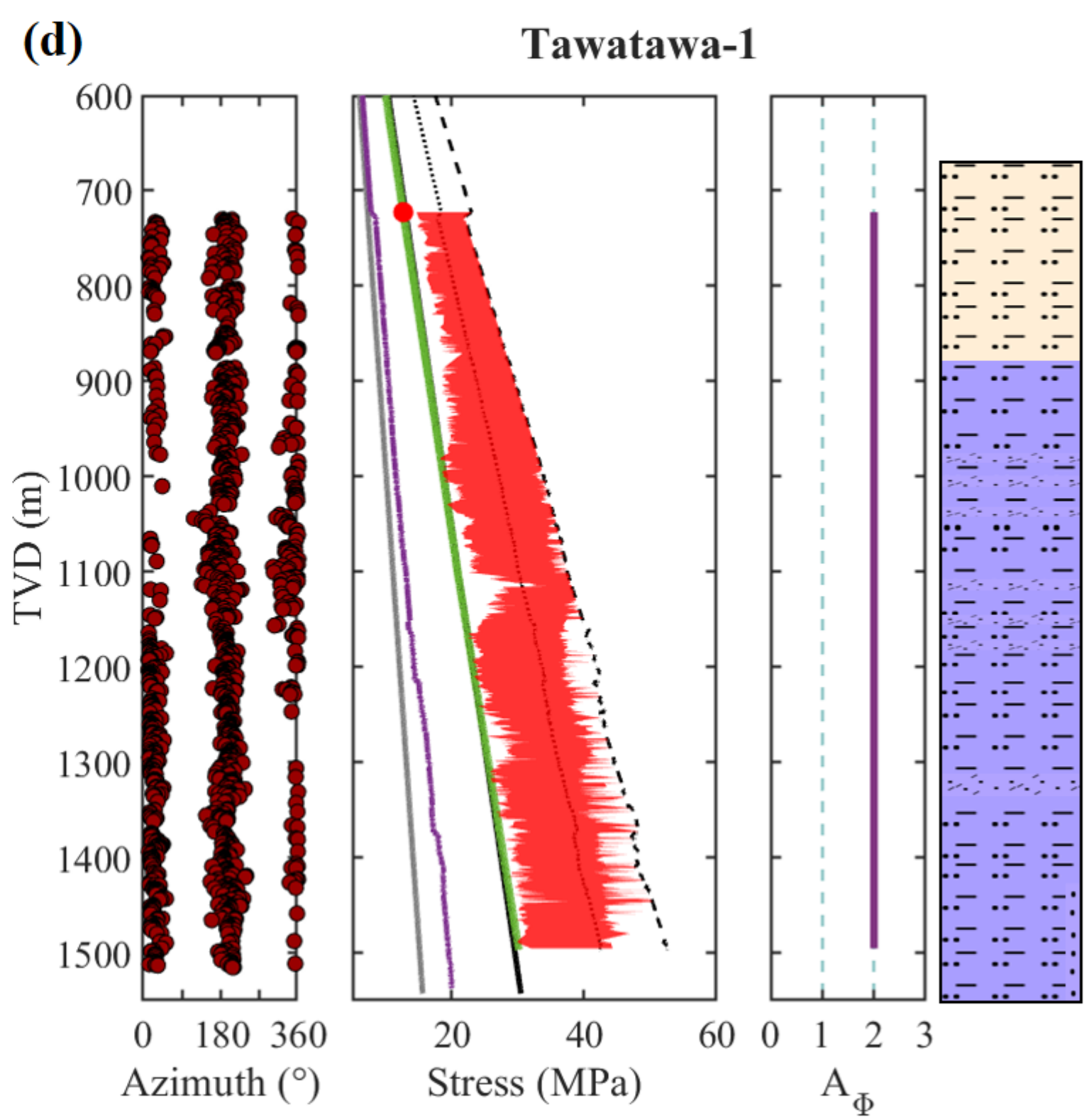
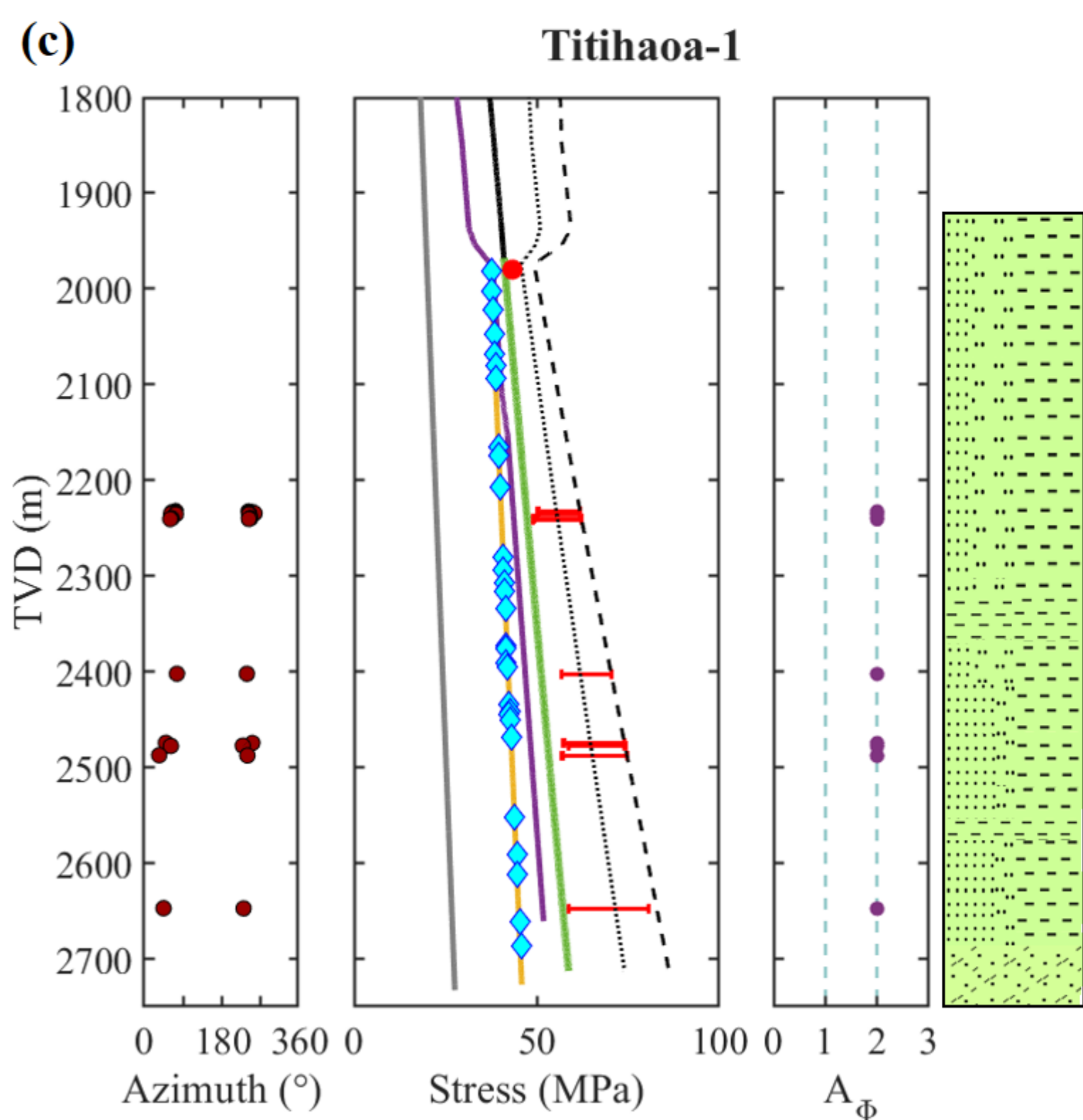
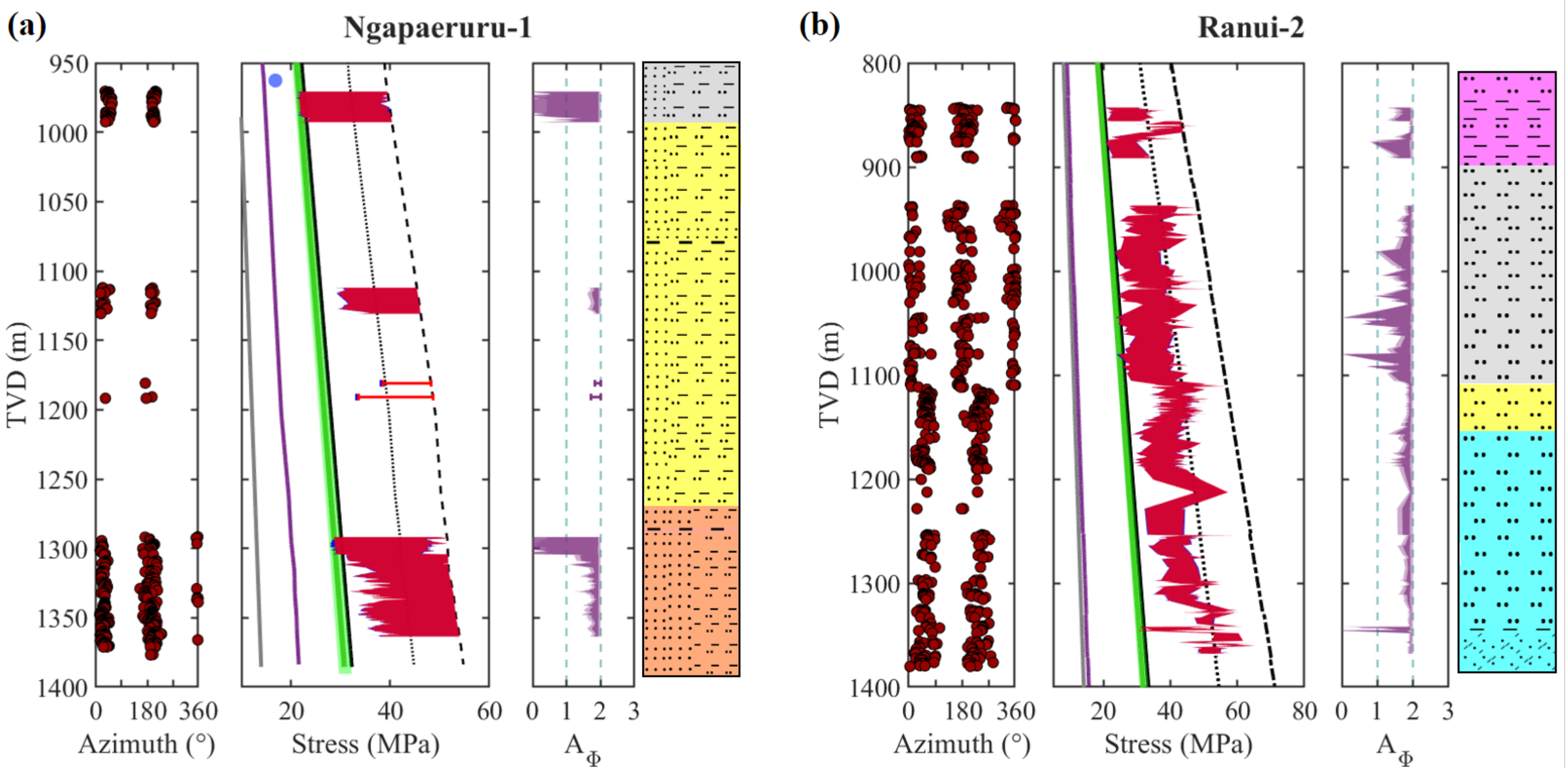




Figure 10.

