# A Three Fractional Order Jerk EquationWith Anti Periodic Conditions 

Zoubir Dahmani ${ }^{1}$, Meriem Mansouria BELHAMITI ${ }^{1}$, and Mehmet Zeki Sarikaya ${ }^{2}$<br>${ }^{1}$ Université Abdelhamid Ibn Badis de Mostaganem<br>${ }^{2}$ Duzce Universitesi

March 07, 2024


#### Abstract

We study a new Jerk equation involving three fractional derivatives and anti periodic conditions. By Banach contraction principle, we present an existence and uniqueness result for the considered problem. Then, by applications of Krasnoselskii fixed point theorem, another result for the existence of at least one solution is established. Also, An illustrative example is discussed. At the end, an approximation for Caputo derivaitive is proposed and some chaotic behaviours are discussed by means of the Runge Kutta 4th order method.

























# A Three Fractional Order Jerk Equation With Anti Periodic Conditions 

Zoubir DAHMANI ${ }^{1}$, Meriem Mansouria BELHAMITI ${ }^{2}$ and Mehmed Zeki SARIKAYA ${ }^{3}$<br>${ }^{1}$ LMPA, Faculty SEI, UMAB Mostaganem, Algeria.<br>e-mail: zzdahmani@yahoo.fr<br>${ }^{2}$ LMPA, Faculty SEI, UMAB Mostaganem, Algeria.<br>e-mail: meriembelhamiti@gmail.com<br>${ }^{3}$ Department of Mathematics, Faculty of Science and Arts, Düzce University, Düzce, Turkey.<br>e-mail: sarikayamz@gmail.com


#### Abstract

We study a new Jerk equation involving three fractional derivatives and anti periodic conditions. By Banach contraction principle, we present an existence and uniqueness result for the considered problem. Then, by applications of Krasnoselskii fixed point theorem, another result for the existence of at least one solution is established. Also, An illustrative example is discussed. At the end, an approximation for Caputo derivaitive is proposed and some chaotic behaviours are discussed by means of the Runge Kutta 4th order method.


2010 Mathematics Subject Classification: 30C45, 39B72, 39B82.
Keywords: Caputo derivative, anti-periodic conditions, Jerk equation, chaotic behaviour, existence and uniqueness, Runge-Kutta method.

## 1 Introduction

Fractional differential equations theory has been a powerful tool for modeling several phenomena in applied sciences and engineering, such as, visco-elasticity, chemistry, fluid flow, electrical networks, electrical circuits, optics, chaotic phenomena in dynamical systems, so on. For more details, one can consult the papers $[2,3,5,6,4,8,9,16,21]$.
In this paper, we are concerned with some applications of dynamical systems of chaotic behaviours. To do this, we begin by recalling some papers that have motivated the present work. We begin by the references [18], where J.C. Sportt discovered some new chaotic systems through an extensive computer search on this phenomenon, with five terms and two quadratic nonlinearities or six terms and a single quadratic nonlinearity.
Then, H.P.W. Gottlieb [11] studied the following problem of chaotic type:
$x^{\prime \prime \prime}=-x^{\prime} s+x^{\prime \prime}\left(x+x^{\prime \prime}\right) / x^{\prime}$,
which he called a Jerk equation.
Then, J.C. Sportt [19] discovered the following particular case of Jerk equation with chaotic behaviours:
$x^{\prime \prime \prime}+a x^{\prime \prime} \pm x^{\prime 2}+x=0$,
which has only three terms in its Jerk representation or five terms in its dynamical system representation with a single quadratic nonlinearity and a single parameter $a$.
We cite also the work of Z. Fu and J. Heidel [10], where the authors proved that there can be no simpler system with a quadratic nonlinearity.
Then, B. Munmuangsaen et al. [14] studied several simple chaotic systems of the form:
$x^{\prime \prime \prime}+x^{\prime \prime}+x=h\left(x^{\prime}\right)$.
Other papers dealing with Jerk equations and systems for chaotic behaviours can be found in $[1,7,12,13,15,17,20]$.
In this work, we try to find a suitable fractional presentation for a simple Jerk circuit that allows us to study some chaotic behaviours.

So, we consider the following problem:

$$
\left\{\begin{array}{l}
D^{\alpha}\left[D^{\beta}\left(D^{\gamma}+\lambda\right)\right] y(t)=f\left(t, y(t), D^{\gamma} y(t)\right), t \in J,  \tag{1}\\
y(0)+y(T)=0,\left(D^{\gamma}\right) y(0)+\left(D^{\gamma}\right) y(T)=0, \\
D^{\gamma} D^{\gamma} y(0)+D^{\gamma} D^{\gamma} y(T)=0 .
\end{array}\right.
$$

In (1), the derivatives $D^{\alpha}, D^{\beta}, D^{\gamma}$ are taken in the sense of Caputo, with $0 \leq \gamma \leq \beta \leq \alpha \leq$ $1, J:=[0, T]$ and $f: J \times \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{R}$ is a given function.
It is to important to note that in the present work:
1-We take the derivatives of Caputo in both sides of the problem.
2-We consider three parameters of derivation $\alpha, \beta$ and $\gamma$ which allow us to be concerned with a sequential Jerk problem without commutativity and semi group properties.
3-The above sequential processes with the anti-periodic conditions on the problem allow us to consider a new type of Jerk problem of three fractional order.
4-Also, it is important to note that Eq. (1) is general enough to describe many problems that arise in mathematical physics, and depending on the values of the constants and functions involved in (1), there are several particular types of equations with important practical applications. For example, Eq. (1) includes the standard Jerk equation of Gottlieb [11] as a particular case. Also, it includes the above two models of Sportt [19] and Munmuangsaen et al. [14].
To the best of our knowledge, this is the first time in the literature where such problem is considered.
The organization of the paper is as follows: In the second section, we recall some fundamental results about fractional calculus. In the third section, the theoretical main results are obtained by using Banach contraction mapping principle and Krasnoselskii fixed point theorem. An illustrative example is discussed in the fourth section. Moreover, some applied results are discussed in section five; an approximation for Caputo derivative is proposed and some numerical simulations and some chaotic behaviours for problem (1) are discussed in this section. At the last section, a conclusion follows.

## 2 Preliminaries on Fractional Calculus

We recall some definitions and lemmas that will be used later. For more details, we refer to [8].

Definition 1 Let $\alpha>0$ and $f:[0, T] \longmapsto \mathbb{R}$ be a continuous function. The Riemann-Liouville integral of order $\alpha$ is defined by:

$$
J^{\alpha} f(t)=\frac{1}{\Gamma(\alpha)} \int_{0}^{t}(t-\tau)^{\alpha-1} f(\tau) d \tau, \alpha>0,0<t<T .
$$

where $\Gamma(\alpha):=\int_{0}^{\infty} e^{-u} u^{\alpha-1} d u$.
Definition 2 For a function $h \in C^{n}([0, T], \mathbb{R})$ and $n-1<\alpha \leq n$, the Caputo fractional derivative is defined by:

$$
\begin{aligned}
D^{\alpha} h(t) & =J^{n-\alpha} \frac{d^{n}}{d t^{n}}(h(t)) \\
& =\frac{1}{\Gamma(n-\alpha)} \int_{0}^{t}(t-s)^{n-\alpha-1} h^{(n)}(s) d s .
\end{aligned}
$$

In order to study the problem (1), we need the following lemmas:
Lemma 3 Let $n \in \mathbb{N}^{*}$, and $n-1<\alpha<n$. The general solution of $D^{\alpha} y(t)=0$ is given by

$$
y(t)=c_{0}+c_{1} t+c_{2} t^{2}+\ldots+c_{n-1} t^{n-1},
$$

where $c_{i} \in \mathbb{R}, i=0,1,2, \ldots, n-1$.
Lemma 4 Let $n \in \mathbb{N}^{*}, n-1<\alpha<n$. Then

$$
I^{\alpha} D^{\alpha} y(t)=y(t)+c_{0}+c_{1} t+c_{2} t^{2}+\ldots+c_{n-1} t^{n-1},
$$

for some $c_{i} \in \mathbb{R}, i=0,1,2, \ldots, n-1$.
Lemma 5 (Krasnoselskii Theorem) Let $\Omega$ be a closed convex and nonempty subset of Banach space $\mathcal{H}$. Let $\mathcal{H}_{1}$ and $\mathcal{H}_{2}$ be two operators such that:

1. $\mathcal{H}_{1} y_{1}+\mathcal{H}_{2} y_{2} \in \Omega, \forall y_{1}, y_{2} \in \Omega$.
2. $\mathcal{H}_{1}$ is compact and continuous.
3. $\mathcal{H}_{2}$ is a contraction mapping.

Then there exists $y_{3} \in \Omega$ such that $\mathcal{H}_{1} y_{3}+\mathcal{H}_{2} y_{3}=y_{3}$.
We need also to prove the following integral representation for (1):
Lemma 6 Let $G \in C(J)$. Then, the problem

$$
\left\{\begin{array}{l}
D^{\alpha}\left[D^{\beta}\left(D^{\gamma}+\lambda\right)\right] y(t)=G(t), t \in J, \\
y(0)+y(T)=0,\left(D^{\gamma}\right) y(0)+\left(D^{\gamma}\right) y(T)=0, \\
D^{\gamma} D^{\gamma} y(0)+D^{\gamma} D^{\gamma} y(T)=0,
\end{array}\right.
$$

is equivalent to the following integral representation:

$$
\begin{aligned}
y(t) & =\frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+1)} J^{\alpha+\beta-\gamma} G(T)\left(\frac{T^{\gamma} t^{\gamma}}{2 \Gamma(\gamma+1)}-\frac{T^{2 \gamma}}{4 \Gamma(\gamma+1)}\right) \\
& +\frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+\gamma+1)} J^{\alpha+\beta-\gamma} G(T)\left(\frac{t^{\beta+\gamma}}{T^{\beta-\gamma}}-\frac{T^{2 \gamma}}{2}\right) \\
& +J^{\alpha+\beta} G(T)\left(\frac{T^{\gamma}}{4 \Gamma(\gamma+1)}-\frac{t^{\gamma}}{2 \Gamma(\gamma+1)}\right) \\
& +J^{\alpha+\beta+\gamma} G(t)-\frac{1}{2} J^{\alpha+\beta+\gamma} G(T) \\
& -\lambda J^{\gamma} y(t)+\frac{\lambda}{2} J^{\gamma} y(T) .
\end{aligned}
$$

Proof. Thanks to lemma 4, we can written:

$$
\left[D^{\beta}\left(D^{\gamma}+\lambda\right)\right] y(t)=J^{\alpha} G(t)-c_{0} .
$$

With the the same idea, we can write

$$
\left(D^{\gamma}+\lambda\right) y(t)=J^{\alpha+\beta} G(t)-c_{0} J^{\beta}(T)-c_{1} .
$$

Hence, it yields that

$$
y(t)=J^{\alpha+\beta+\gamma} G(t)-c_{0} \frac{t^{\beta+\gamma}}{\Gamma(\beta+\gamma+1)}-c_{1} \frac{t^{\gamma}}{\Gamma(\gamma+1)}-\lambda J^{\gamma} y(t)-c_{2} .
$$

Thanks to the anti-periodic conditions, we observe that

$$
\left\{\begin{array}{l}
-2 c_{2}+J^{\alpha+\beta+\gamma} G(T)-c_{0} J^{\beta+\gamma}(T)-c_{1} J^{\gamma}(T)-\lambda J^{\gamma} y(T)=0, \\
-2 c_{1}+J^{\alpha+\beta} G(T)-c_{0} J^{\beta}(T)=0, \\
J^{\alpha+\beta-\gamma} G(T)-c_{0} J^{\beta-\gamma}(T)=0 .
\end{array}\right.
$$

Solving the above system, we get:

$$
\left\{\begin{aligned}
c_{0} & =\Gamma(\beta-\gamma+1) J^{\alpha+\beta-\gamma} G(T), \\
2 c_{1} & =J^{\alpha+\beta} G(T)-\Gamma(\beta-\gamma+1) J^{\alpha+\beta-\gamma} G(t) J^{\beta}(T), \\
2 c_{2} & =J^{\alpha+\beta+\gamma} G(T)-\Gamma(\beta-\gamma+1) J^{\alpha+\beta-\gamma} G(T) J^{\beta+\gamma}(T) \\
& -\frac{1}{2}\left(J^{\alpha+\beta} G(T)-\Gamma(\beta-\gamma+1) J^{\alpha+\beta-\gamma} G(T) J^{\beta}(1)\right) J^{\gamma}(T)-\lambda J^{\gamma} y(T) .
\end{aligned}\right.
$$

The lemma is thus proved.

## 3 Main Results

Now, we introduce the Banach space:

$$
X:=\left\{x \in C(J, \mathbb{R}), D^{\gamma} x \in C(J, \mathbb{R})\right\},
$$

endowed with the norm:

$$
\|x\|_{X}=\|x\|_{\infty}+\left\|D^{\gamma} x\right\|_{\infty},
$$

where,

$$
\|x\|_{\infty}=\sup _{t \in J}|x(t)|,\left\|D^{\gamma} x\right\|_{\infty}=\sup _{t \in J}\left|D^{\gamma} x(t)\right| .
$$

Then, we consider the notations:

$$
\alpha_{1}=\alpha+\beta+\gamma, \alpha_{2}=\alpha+\beta-\gamma, \alpha_{3}=\alpha+\beta, \alpha_{4}=\beta+\gamma
$$

Over the above Banach space, we define the nonlinear operator $\mathcal{H}: X \rightarrow X$ by:

$$
\mathcal{H} y(t)=\mathcal{H}_{1} y(t)+\mathcal{H}_{2} y(t),
$$

where,

$$
\begin{aligned}
\mathcal{H}_{1} y(t) & =\frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+1)} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{2}-1}}{\Gamma\left(\alpha_{2}\right)} f\left(\tau, y(\tau), D^{\gamma} y(\tau)\right) d \tau\left(\frac{T^{\gamma} t^{\gamma}}{2 \Gamma(\gamma+1)}-\frac{T^{2 \gamma}}{4 \Gamma(\gamma+1)}\right) \\
& +\frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+\gamma+1)} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{2}-1}}{\Gamma\left(\alpha_{2}\right)} f\left(\tau, y(\tau), D^{\gamma} y(\tau)\right) d \tau\left(\frac{t^{\beta+\gamma}}{T^{\beta-\gamma}}-\frac{T^{2 \gamma}}{2}\right) \\
& +\int_{0}^{T} \frac{(T-\tau)^{\alpha_{3}-1}}{\Gamma\left(\alpha_{3}\right)} f\left(\tau, y(\tau), D^{\gamma} y(\tau)\right) d \tau\left(\frac{T^{\gamma}}{4 \Gamma(\gamma+1)}-\frac{t^{\gamma}}{2 \Gamma(\gamma+1)}\right) \\
& +\int_{0}^{t} \frac{(t-\tau)^{\alpha_{1}-1}}{\Gamma\left(\alpha_{1}\right)} f\left(\tau, y(\tau), D^{\gamma} y(\tau)\right) d \tau-\frac{1}{2} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{1}-1}}{\Gamma\left(\alpha_{1}\right)} f\left(\tau, y(\tau), D^{\gamma} y(\tau)\right) d \tau
\end{aligned}
$$

and

$$
\mathcal{H}_{2} y(t)=-\lambda \int_{0}^{t} \frac{(t-\tau)^{\gamma-1}}{\Gamma(\gamma)} y(t) d \tau+\frac{\lambda}{2} \int_{0}^{T} \frac{(T-\tau)^{\gamma-1}}{\Gamma(\gamma)} y(t) d \tau .
$$

We need also to consider the following hypotheses:
(H1): There exist $K_{1}, K_{2}>0$, such that for all $t \in[0, T]$ and $u_{i}, v_{i} \in \mathbb{R}, i=1,2$, we have

$$
\left.\left|f\left(t, u_{1}, u_{2}\right)-f\left(t, v_{1}, v_{1}\right)\right| \leq K_{1}\left|u_{1}-v_{1}\right|+K_{2}\left|u_{2}-v_{2}\right|\right), K:=\max \left(K_{1}, K_{2}\right) .
$$

(H2): Let $f:[0, T] \times \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{R}$ be a jointly continuous function.
(H3): There exists a continuous function $\psi$ defined over $[0, T]$; such that for all $t \in[0, T]$ and $u_{i} \in \mathbb{R}, i=1$, 2 , we have

$$
\left|f\left(t, u_{1}, u_{2}\right)\right| \leq \psi(t) ; \sup _{t \in[0, T]}=\|\psi\| .
$$

For computation convenience, we define the function $F_{y}:[0, T] \rightarrow \mathbb{R}$ by:

$$
F_{y}(t):=f\left(t, y(t), D^{\alpha} y(t)\right),
$$

and

$$
\begin{gathered}
\mathcal{I}_{1}:=\frac{3 T^{\alpha_{2}+2 \gamma}}{2 \Gamma\left(\alpha_{2}+1\right)} \frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+\gamma+1)}+\frac{3 T^{\alpha_{2}+2 \gamma}}{4 \Gamma(\gamma+1) \Gamma\left(\alpha_{2}+1\right)} \frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+1)}+\frac{3 T^{\alpha_{3}+\gamma}}{4 \Gamma\left(\alpha_{3}+1\right) \Gamma(\gamma+1)}+\frac{3 T^{\alpha_{1}}}{2 \Gamma\left(\alpha_{1}+1\right)} . \\
\mathcal{L}_{1}:=\frac{T^{\alpha_{2}+\gamma}}{\Gamma\left(\alpha_{2}+1\right)} \frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+1)}+\frac{T^{\alpha_{2}+\gamma}}{2 \Gamma\left(\alpha_{2}+1\right)} \frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+1)}+\frac{T^{\alpha_{3}}}{2 \Gamma\left(\alpha_{3}+1\right)}+\frac{T^{\alpha_{1}-\gamma}}{\Gamma\left(\alpha_{1}-\gamma+1\right)} . \\
\mathcal{I}_{2}:=\frac{3|\lambda|}{2 \Gamma(\gamma+1)}, \mathcal{L}_{2}:=|\lambda| .
\end{gathered}
$$

Also, we take:

$$
\Delta:=K\left(\mathcal{I}_{1}+\mathcal{L}_{1}\right)+\left(\mathcal{I}_{2}+\mathcal{L}_{2}\right) .
$$

Now, we are ready to study the above problem by means of the fixed point theory.

### 3.1 A Unique Solution Via Banach Contraction

Theorem 7 Assume that (H1) is valid and $\Delta<1$. Then, the problem (1) has a unique solution on $[0, T]$.

Proof. We show that the operator $\mathcal{H}$ is contractive. Let $y_{i} \in X, i=1,2$. Then, for each $t \in[0, T]$, we have

$$
\begin{aligned}
\left|\mathcal{H}_{1} y_{1}(t)-\mathcal{H}_{1} y_{2}(t)\right| & =\frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+1)} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{2}-1}}{\Gamma\left(\alpha_{2}\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau\left|\frac{T^{\gamma} t^{\gamma}}{2 \Gamma(\gamma+1)}+\frac{T^{2 \gamma}}{4 \Gamma(\gamma+1)}\right| \\
& +\frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+\gamma+1)} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{2}-1}}{\Gamma\left(\alpha_{2}\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau\left|\frac{t^{\beta+\gamma}}{T^{\beta-\gamma}}+\frac{T^{2 \gamma}}{2}\right| \\
& +\int_{0}^{T} \frac{(T-\tau)^{\alpha_{3}-1}}{\Gamma\left(\alpha_{3}\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau\left|\frac{T^{\gamma}}{4 \Gamma(\gamma+1)}+\frac{t^{\gamma}}{2 \Gamma(\gamma+1)}\right| \\
& +\int_{0}^{t} \frac{(T-\tau)^{\alpha_{1}-1}}{\Gamma\left(\alpha_{1}\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau+\frac{1}{2} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{1}-1}}{\Gamma\left(\alpha_{1}\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau .
\end{aligned}
$$

Thanks to (H1), yields

$$
\begin{aligned}
\left|\mathcal{H}_{1} y_{1}(t)-\mathcal{H}_{1} y_{2}(t)\right| & \leq \frac{K\left\|y_{1}-y_{2}\right\|_{X} T^{\alpha_{2}}}{\Gamma\left(\alpha_{2}+1\right)}\left|\frac{T^{\gamma} t^{\gamma}}{2 \Gamma(\gamma+1)}+\frac{T^{2 \gamma}}{4 \Gamma(\gamma+1)}\right| \frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+1)} \\
& +\frac{K\left\|y_{1}-y_{2}\right\|_{X} T^{\alpha_{2}}}{\Gamma\left(\alpha_{2}+1\right)}\left|\frac{t^{\beta+\gamma}}{T^{\beta-\gamma}}+\frac{T^{2 \gamma}}{2}\right| \frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+\gamma+1)} \\
& +\frac{K\left\|y_{1}-y_{2}\right\|_{X} T^{\alpha_{3}}}{\Gamma\left(\alpha_{3}+1\right)}\left|\frac{T^{\gamma}}{4 \Gamma(\gamma+1)}+\frac{T^{\gamma}}{2 \Gamma(\gamma+1)}\right| \\
& +\frac{K\left\|y_{1}-y_{2}\right\|_{X} T^{\alpha_{1}}}{\Gamma\left(\alpha_{1}+1\right)}+\frac{K\left\|y_{1}-y_{2}\right\|_{X}}{2} \frac{T^{\alpha_{1}}}{\Gamma\left(\alpha_{1}+1\right)} .
\end{aligned}
$$

Also, we have

$$
\begin{align*}
\left|\mathcal{H}_{2} y_{1}(t)-\mathcal{H}_{2} y_{2}(t)\right| \leq & |\lambda| \int_{0}^{t} \frac{(t-\tau)^{\gamma-1}}{\Gamma(\gamma)}\left|y_{1}(\tau)-y_{2}(\tau)\right| d \tau \\
& +\frac{|\lambda|}{2} \int_{0}^{T} \frac{(T-\tau)^{\gamma-1}}{\Gamma(\gamma)}\left|y_{1}(\tau)-y_{2}(\tau)\right| d \tau . \\
\left\|\mathcal{H}_{2} y_{1}-\mathcal{H}_{2} y_{2}\right\|_{\infty} \leq & \mathcal{I}_{2}\left\|y_{1}-y_{2}\right\|_{\infty} \leq \mathcal{I}_{2}\left\|y_{1}-y_{2}\right\|_{X} . \tag{2}
\end{align*}
$$

Therefore,

$$
\begin{equation*}
\left\|\mathcal{H} y_{1}-\mathcal{H} y_{2}\right\|_{\infty} \leq\left(\mathcal{I}_{1} K+\mathcal{I}_{2}\right)\left\|y_{1}-y_{2}\right\|_{X} \tag{3}
\end{equation*}
$$

In the same manner, we can write

$$
\begin{aligned}
\left|D \mathcal{H}_{1} y_{1}(t)-D \mathcal{H}_{1} y_{2}(t)\right| & \leq \frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+1)} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{2}-1}}{\Gamma\left(\alpha_{2}\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau\left|\frac{\gamma^{\gamma} t^{\gamma-1}}{2 \Gamma(\gamma+1)}\right| \\
& +\frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+\gamma+1)} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{2}-1}}{\Gamma\left(\alpha_{2}\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau\left|\frac{(\beta+\gamma) t^{\beta+\gamma-1}}{T^{\beta-\gamma}}\right| \\
& +\int_{0}^{T} \frac{(T-\tau)^{\alpha_{3}-1}}{\Gamma\left(\alpha_{3}\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau\left|\frac{\gamma t^{\gamma-1}}{2 \Gamma(\gamma+1)}\right| \\
& +\int_{0}^{t} \frac{(T-\tau)^{\alpha_{1}-2}}{\Gamma\left(\alpha_{1}-1\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau .
\end{aligned}
$$

Hence, we have

$$
\begin{aligned}
\left|D^{\gamma} \mathcal{H}_{1} y_{1}(t)-D^{\gamma} \mathcal{H}_{1} y_{2}(t)\right| & \leq \frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+\gamma+1)} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{2}-1}}{\Gamma\left(\alpha_{2}\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau\left|\frac{\Gamma(\beta+\gamma+1) t^{\beta}}{T^{\beta-\gamma} \Gamma(\beta+1)}\right| \\
& +\frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+1)} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{2}-1}}{\Gamma\left(\alpha_{2}\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau\left|\frac{T^{\gamma}}{2}\right| \\
& +\frac{1}{2} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{3}-1}}{\Gamma\left(\alpha_{3}\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau \\
& +\int_{0}^{t} \frac{(T-\tau)^{\alpha_{1}-\gamma-1}}{\Gamma\left(\alpha_{1}-\gamma\right)}\left|F_{y_{1}}(\tau)-F_{y_{2}}(\tau)\right| d \tau .
\end{aligned}
$$

Consequently,

$$
\left|D^{\gamma} \mathcal{H}_{1} y_{1}(t)-D^{\gamma} \mathcal{H}_{1} y_{2}(t)\right| \leq \mathcal{L}_{1}\left\|y_{1}-y_{2}\right\|_{X},
$$

and

$$
\begin{equation*}
\left\|D^{\gamma} \mathcal{H}_{2} y_{1}-D^{\gamma} \mathcal{H}_{2} y_{2}\right\|_{\infty} \leq \mathcal{L}_{2}\left\|y_{1}-y_{2}\right\|_{X} . \tag{4}
\end{equation*}
$$

It yields then that

$$
\begin{equation*}
\left\|D^{\gamma} \mathcal{H} y_{1}-D^{\gamma} \mathcal{H} y_{2}\right\|_{\infty} \leq\left(\mathcal{L}_{1} K+\mathcal{L}_{2}\right)\left\|y_{1}-y_{2}\right\|_{X} . \tag{5}
\end{equation*}
$$

Thanks to (3)-(5), we get

$$
\left\|\mathcal{H} y_{1}-\mathcal{H} y_{2}\right\|_{X} \leq\left(\left(\mathcal{I}_{1}+\mathcal{L}_{1}\right) K+\left(\mathcal{I}_{2}+\mathcal{L}_{2}\right)\right)\left\|y_{1}-y_{2}\right\|_{X}
$$

By Theorem 7, we deduce that $\mathcal{H}$ is contractive. As a consequence of Banach contraction principle, we conclude that $\mathcal{H}$ has a unique fixed point which is the solution of (1).

### 3.2 Existence via Krasnoselskii Theorem

Theorem 8 Assume that (H2) and (H3) are valid and $\Delta<1$.
Then, the problem (1) has at least one solution $y(t), t \in[0, T]$ that satisfies:
$\|y\|_{X} \leq \epsilon$,
where $\epsilon$ is an infinitesimal that satisfies $\epsilon \geq \frac{\|\psi\|\left(\mathcal{I}_{1}+\mathcal{L}_{1}\right)}{1-\left(\mathcal{I}_{2}+\mathcal{L}_{2}\right)}$.

Proof. Let us consider the following closed convex subset: $B(\epsilon):=\left\{y \in X,\|y\|_{X} \leq \epsilon\right\} \subset X$, with

$$
\epsilon \geq \frac{\|\psi\|\left(\mathcal{I}_{1}+\mathcal{L}_{1}\right)}{1-\left(\mathcal{I}_{2}+\mathcal{L}_{2}\right)}
$$

Our first claim is to prove that for any $y_{1}, y_{2} \in B(\epsilon)$, we have $\mathcal{H} B(\epsilon) \subset B(\epsilon)$.
Let $y_{1}, y_{2} \in B(\epsilon)$. Then, by (H3), it follows that

$$
\left\|\mathcal{H}_{1} y_{1}+\mathcal{H}_{2} y_{2}\right\|_{\infty} \leq \mathcal{I}_{1}\|\psi\|+\mathcal{I}_{2} \epsilon .
$$

On the other hand, we have

$$
\left\|D^{\gamma} \mathcal{H}_{1} y_{1}+D^{\gamma} \mathcal{H}_{2} y_{2}\right\|_{\infty} \leq \mathcal{L}_{1}\|\psi\|+\mathcal{L}_{2} \epsilon .
$$

Therefore,

$$
\left\|\mathcal{H}_{1} y_{1}+\mathcal{H}_{2} y_{2}\right\|_{X} \leq \epsilon .
$$

Furthermore, $\mathcal{H}_{2}$ satisfies the Banach contraction principle (as it has been shown in the inequalities (2)-(4)). In fact, we have

$$
\left\|\mathcal{H}_{2} y_{1}-\mathcal{H}_{2} y_{2}\right\|_{X} \leq\left(\mathcal{I}_{2}+\mathcal{L}_{2}\right)\left\|y_{1}-y_{2}\right\|_{X} .
$$

The last step in this proof is to show that $\mathcal{H}_{1}$ is compact and continuous.
The continuity of the operator $\mathcal{H}_{1}$ is due to the continuity of $f($ see (H2))
Moreover, $\mathcal{H}_{1}$ is bounded on $B(\epsilon)$. Indeed, for any $y \in B(\epsilon)$, we have:

$$
\left\|\mathcal{H}_{1} y\right\|_{X} \leq\left(\mathcal{I}_{1}+\mathcal{L}_{1}\right)\|\psi\| .
$$

Next, we shall prove that $\mathcal{H}_{1}$ is equicontinuous.
Let $t_{1}, t_{2} \in[0, T]$ with $t_{1}<t_{2}$. Then, we have:

$$
\begin{aligned}
\left|\mathcal{H}_{1} y\left(t_{1}\right)-\mathcal{H}_{1} y\left(t_{2}\right)\right| & \leq \frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+1)} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{2}-1}}{\Gamma\left(\alpha_{2}\right)}\left|F_{y}(\tau)\right| d \tau\left|\frac{T^{\gamma}\left(t_{1}^{\gamma}-t_{2}^{\gamma}\right)}{2 \Gamma(\gamma+1)}\right| \\
& +\frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+\gamma+1)} \int_{0}^{T} \frac{(T-\tau)^{\alpha_{2}-1}}{\Gamma\left(\alpha_{2}\right)}\left|F_{y}(\tau)\right| d \tau\left|\frac{t_{1}^{\beta+\gamma}-t_{2}^{\beta+\gamma}}{T^{\beta-\gamma}}\right| \\
& +\int_{0}^{T} \frac{(T-\tau)^{\alpha_{3}-1}}{\Gamma\left(\alpha_{3}\right)}\left|F_{y}(\tau)\right| d \tau\left|\frac{\left(t_{1}^{\gamma}-t_{2}^{\gamma}\right)}{2 \Gamma(\gamma+1)}\right| \\
& +\int_{0}^{t_{1}} \frac{\left|\left(t_{1}-\tau\right)^{\alpha_{1}-1}-\left(t_{2}-\tau\right)^{\alpha_{1}-1}\right|}{\Gamma\left(\alpha_{1}\right)}\left|F_{y}(\tau)\right| d \tau+\int_{t_{1}}^{t_{2}} \frac{\left(t_{2}-\tau\right)^{\alpha_{1}-1}}{\Gamma\left(\alpha_{1}\right)}\left|F_{y}(\tau)\right| d \tau .
\end{aligned}
$$

Thanks to (H3), we can write

$$
\begin{align*}
\left|\mathcal{H}_{1} y\left(t_{1}\right)-\mathcal{H}_{1} y\left(t_{2}\right)\right| & \leq \int_{0}^{t_{1}} \frac{\left|\left(t_{1}-\tau\right)^{\alpha_{1}-1}-\left(t_{2}-\tau\right)^{\alpha_{1}-1}\right|}{\Gamma\left(\alpha_{1}\right)}\left|F_{y}(\tau)\right| d \tau+\int_{t_{1}}^{t_{2}} \frac{\left(t_{2}-\tau\right)^{\alpha_{1}-1}}{\Gamma\left(\alpha_{1}\right)}\left|F_{y}(\tau)\right| d \tau \\
& +\frac{\|\psi\|(T)^{\alpha_{2}+\gamma}}{\Gamma\left(\alpha_{2}+1\right)} \frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+1)}\left|\frac{\left(t_{1}^{\gamma}-t_{2}^{\gamma}\right)}{2 \Gamma(\gamma+1)}\right|+\frac{\|\psi\|(T)^{\alpha_{3}}}{\Gamma\left(\alpha_{3}+1\right)}\left|\frac{\left(t_{1}^{\gamma}-t_{2}^{\gamma}\right)}{2 \Gamma(\gamma+1)}\right| \\
& +\frac{\|\psi\|(T)^{\alpha_{2}}}{\Gamma\left(\alpha_{2}+1\right)} \frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+\gamma+1)}\left|\frac{t_{1}^{\beta+\gamma}-t_{2}^{\beta+\gamma}}{T^{\beta-\gamma}}\right| . \tag{6}
\end{align*}
$$

With the same arguments as before, we have

$$
\begin{align*}
\left|D^{\gamma} \mathcal{H}_{1} y\left(t_{1}\right)-D^{\gamma} \mathcal{H}_{1} y\left(t_{2}\right)\right| & \leq \int_{0}^{t_{1}} \frac{\left|\left(t_{1}-\tau\right)^{\alpha_{1}-1}-\left(t_{2}-\tau\right)^{\alpha_{1}-\gamma-1}\right|}{\Gamma\left(\alpha_{1}\right)}\left|F_{y}(\tau)\right| d \tau \\
& +\int_{t_{1}}^{t_{2}} \frac{\left(t_{2}-\tau\right)^{\alpha_{1}-\gamma-1}}{\Gamma\left(\alpha_{1}\right)}\left|F_{y}(\tau)\right| d \tau+\frac{\|\psi\|(T)^{\alpha_{2}}}{\Gamma\left(\alpha_{2}+1\right)} \frac{\Gamma(\beta-\gamma+1)}{\Gamma(\beta+\gamma+1)}\left|\frac{t_{1}^{\beta}-t_{2}^{\beta}}{T^{\beta-\gamma}}\right| . \tag{7}
\end{align*}
$$

The right hand sides of (6) and (7) tend to zero as $t_{1} \rightarrow t_{2}$. Then, $\mathcal{H}_{1}$ is equicontinuous. Hence by the Arzela-Ascoli theorem, the operator $\mathcal{H}_{1}$ is compact on $B(\epsilon)$.
Thus, thanks to Krasnoselskii theorem, the problem (1) has at least one solution $y ;\|y\|_{X} \leq \epsilon$, with $\epsilon$ satisfies the condition in Theorem 8.

## 4 An Illustrative Example

We consider the following problem:

$$
\left\{\begin{array}{l}
D^{\alpha}\left[D^{\beta}\left(D^{\gamma}+\lambda\right)\right] y(t)=f\left(t, y(t), D^{\alpha} y(t)\right), t \in[0,1], 0<\alpha, \beta, \gamma \leq 1 .  \tag{8}\\
y(0)+y(1)=0,\left(D^{\gamma}\right) y(0)+\left(D^{\gamma}\right) y(1)=0, D^{\gamma} D^{\gamma} y(0)+D^{\gamma} D^{\gamma} y(1)=0 .
\end{array}\right.
$$

Here, we take

$$
\begin{gathered}
f\left(t, y(t), D^{\alpha} y(t)\right)= \pm \sinh \left(D^{\alpha} y(t)\right)-0.95 D^{\alpha} y(t)-0.08 y(t), \\
T=1, \alpha=0.9, \beta=0.75, \gamma=0.6, \lambda=0.1 .
\end{gathered}
$$

For all $\left(x_{1}, y_{1}\right),\left(x_{2}, y_{2}\right) \in \mathbb{R}^{2}$, and $t \in[0.1]$, we have:

$$
\left|f\left(t, x_{1}, x_{2}\right)-f\left(t, y_{1}, y_{2}\right)\right| \leq 0.08\left(\left|x_{1}-y_{1}\right|+\left|x_{2}-y_{2}\right|\right) .
$$

We have also

$$
\mathcal{I}_{1}=4.263, \mathcal{I}_{2}=0.167, \mathcal{L}_{1}=4.486, \mathcal{L}_{2}=0.1, \Delta=0.967
$$

Thanks to Theorem 7, we can state that the problem (8) has a unique solution on $[0,1]$.

## 5 Existence of Fractional Chaotic Behaviours

In the following two subsections, we present a numerical approach for the Caputo derivative. Then, in order to study dynamic behavior of the above fractional Jerk problem, we present a reduced fractional differential system that is equivalent to our studied problem. For this reduction, we will show sensitive attractors to initial anti-periodic conditions in phase space, which is one of the properties of certain chaotic behaviours.
The numerical simulation of the fractional system is done by the fourth-order Runge-Kutta applied to Caputo derivative.

### 5.1 Numerical Approach for Caputo Derivative

Since the Caputo derivative has wide applications, so we put our attention to the numerical approach of this derivative. To do this, we begin this section by recalling the following theorem [9] in which the authors presented an important numerical approach for the Riemann-Liouville fractional integral.

Theorem 9 Assume that $y \in \mathcal{C}^{1}([0, T], \mathbb{R})$. The numerical approach for fractional integral is given by:

$$
J^{\alpha} y\left(t_{i}\right) \simeq \frac{h^{\alpha}}{\Gamma(\alpha+2)} \sum_{j=0}^{i} K_{j} y_{j}, \quad i=0, \ldots, n+1, y_{0} \text { initial condition }
$$

where

$$
K_{j}:=\sigma_{j}(\alpha)=\left\{\begin{array}{l}
(n+2-j)^{(\alpha+1)}+(n-j)^{(\alpha+1)}-2(n-j+1)^{(\alpha+1)}, j=1 \ldots i-1 \\
(n)^{(\alpha+1)}-(n-\alpha)(n+1)^{\alpha}, j=0, \quad 1, j=i
\end{array}\right.
$$

Based on Theorem 9, we propose the following approximation for Caputo derivative:
Theorem 10 Assume that $y \in \mathcal{C}^{1}([0, T], \mathbb{R})$ and $0<\alpha \leq 1$. Then, we have:

$$
D^{\alpha} y\left(t_{i}\right) \simeq \frac{h^{1-\alpha}}{\Gamma(1-\alpha+2)} \sum_{j=0}^{i} L_{j} y^{(j)}\left(t_{j}\right), \quad i=0, \ldots, n
$$

Where,

$$
L_{j}:=\sigma_{j}(1-\alpha) .
$$

and

$$
y^{(j)}=\left\{\frac{y_{1}-y_{0}}{h}, j=0, \quad \frac{y_{j+1}-y_{j-1}}{2 h}, j=1 \ldots i-1, \quad \frac{y_{i}-y_{i-1}}{h}, j=i .\right.
$$

Proof. Let $y \in \mathcal{C}^{1}([0, T], \mathbb{R}), 0<\alpha \leq 1,0 \leq t \leq T$ and setting $h:=\frac{T}{n} ; n \in \mathbb{N}^{*}$.
Then, thanks to Theorem 9 and using the above definition of Caputo derivative, we obtain:

$$
D^{\alpha} y\left(t_{n}\right)=J^{1-\alpha} D y\left(t_{n}\right) \approx \frac{h^{1-\alpha}}{\Gamma(1-\alpha+2)} \sum_{j=0}^{n} L_{j} D y\left(t_{j}\right)
$$

By finite difference scheme, instead of using "central difference scheme" for $j=0, \ldots, n$, we use "forward difference scheme" for $j=0$, "backward difference scheme" for $j=n$, and " central difference scheme" for $j=1, \ldots n-1$.
By substitution in the above formula, we obtain:
$D^{\alpha} y\left(t_{n}\right) \approx \frac{h^{1-\alpha}}{\Gamma(1-\alpha+2)}\left(\sum_{j=1}^{n-1} L_{j}\left(\frac{y_{j+1}-y_{j-1}}{2 h}\right)+L_{0}\left(\frac{y_{1}-y_{0}}{h}\right)+L_{n}\left(\frac{y_{n}-y_{n-1}}{h}\right)\right)$.
Theorem 10 is thus proved and a Caputo derivative approximation is obtained.


Figure 1: Different phase portrait of incommensurate order of system(9) for (10)

### 5.2 Simulation for Chaotic Behaviours

We note that the problem (1) can reduced to the following system:

$$
\begin{aligned}
& D^{\gamma} y(t)=z(t)-\lambda y(t) \\
& D^{\beta} z(t)=w(t) \\
& D^{\alpha} w(t)=f\left(t, y(t), D^{\alpha} y(t)\right),
\end{aligned}
$$

that is

$$
\begin{align*}
& D y(t)=D^{1-\gamma}(z(t)-\lambda y(t)) \\
& D z(t)=D^{1-\beta} w(t)  \tag{9}\\
& D w(t)=D^{1-\alpha} f\left(t, y(t), D^{\alpha} y(t)\right) .
\end{align*}
$$

( a:) As a first simulation, we consider the case where $f$ is given by:

$$
\begin{equation*}
f\left(t, y(t), D^{\alpha} y(t)\right)= \pm\left(\left(D^{\alpha} y(t)\right)^{2} \pm 2\right)^{-2}-y(t) \tag{10}
\end{equation*}
$$

with initial conditions ( $0.0945,0,-0.0945$ ), $\lambda=0.7$ and $h=0.005$. The integration of system (9) is carried out by the fourth-order Runge-Kutta method and the Caputo approach.

- For incommensurate order $(\gamma, \beta, \alpha)=(0.8,0.85,0.9)$, we get
- For commensurate order $(\gamma, \beta, \alpha)=(0.9,0.9,0.9)$, so with the same data as above, we have
- For $\gamma=\beta=\alpha=1$, the incorporation of system (9) is carried out by the 4 th Runge-Kutta method, we obtain
(b:) As a second simulation, we consider the case where $f$ is given by:

$$
\begin{equation*}
f\left(t, y(t), D^{\alpha} y(t)\right)= \pm 0.1 \exp \left(\mp D^{\alpha} y(t)\right)-y(t) \tag{11}
\end{equation*}
$$

with initial conditions $(0.1021,0,-0.1021), \lambda=0.7$ and $h=0.005$, the integration of system (9) is carried out by the fourth-order Runge-Kutta method and Caputo approach.

- For incommensurate order $(\gamma, \beta, \alpha)=(0.85,0.9,0.9)$, we get graphical illustrations in Figure 1-6.


Figure 2: Different phase portrait of commensurate order of system(9) for (10)


Figure 3: Different phase portrait of $\gamma=\beta=\alpha=1$, of system (9) for (10)


Figure 4: Different phase portrait of incommensurate order of system(9) for (11)


Figure 5: Different phase portrait of commensurate order of system (9) for (11)


Figure 6: Different phase portrait of $\gamma=\beta=\alpha=1$, of system (9) for (11)

- For commensurate order $(\gamma, \beta, \alpha)=(0.9,0.9,0.9)$, with the same data as above, we obtain
- For $\gamma=\beta=\alpha=1$, the incorporation of system (9) is carried out only by the fourthorder Runge-Kutta method.


## Remark 11 .

- Comparison of numerical simulations showed great correlation for specific parameters. Unfortunately, it is not the same for all cases.
- Numerical simulations have displayed a strange attractors that can coexist in these fractional order systems.
- The incorporation have indicated a high qualitative agreement between chaotic systems for $\alpha, \beta, \gamma \longrightarrow 1$.
- The commensurate-order and incommensurate-order reflect the influence of fractional order on chaotic systems.


## 6 Conclusion

We have introduced a new problem for Jerk circuits of chaotic phenomena by means of fractional derivatives. An existence and uniqueness result has been established by means of Banach contraction principle. Then, using Krasnoselskii fixed point theorem, another main result for the existence of one solution has also been discussed. An illustrative example has been presented to show the applicability of our main result. At the end, an approximation for Caputo derivative has been proved and some numerical solutions having chaotic behaviours have been illustrated and discussed.

## References

[1] Khan NA., Hameed T., Qureshi MA, Abkar S., Alzahrani AK. Emulate the chaotic flows of fractional jerk system to scramble the sound and image memo with circuit execution, Phys. Scripta, IOP Publishing Ltd, Accepted Manuscript online 1 April 2020, 2020.
[2] Dahmani Z., Belarbi S. New results for fractional evolution equations using Banach fixed point theorem, International Journal of Nonlinear Analysis and Applications, 5 (2), 2014, 22-30
[3] Dahmani Z., Abdellaoui MA. On a three point boundary value problem of arbitrary order, Journal of Interdisciplinary Mathematics, Volume 19, 2016, Issue 5-6.
[4] Fazil H. Fractional Langevin equations with anti peridoic conditions, Chaos, Solitons and Fractals, 114, (2018), 332-337.
[5] Houas M., Dahmani Z. On existence of solutions for fractional differential equations with nonlocal multi-point boundary conditions, Lobachevskii Journal of Mathematics, 37(2), 2016, Pages 120-127.
[6] Houas M., Dahmani Z. New results for a system of two fractional differential equations involving n Caputo derivatives, Kragujevac Journal of Mathematics, 38(2), 2014,Pages 283-301.
[7] Kingni ST, Jafari S., Simo H., Woafo P. Three-dimensional chaotic autonomous system with only one stable equilibrium: Analysis, circuit design, parameter estimation, control, synchronization and its fractional-order form, Eur. Phys. J. Plus (2014) 129: 76.
[8] Kilbas AA, Srivastava HM, Trujillo JJ. Theory and Applications of Fractional Differential Equations, North-Holland Mathematics Studies 204-Elsevier, (2006).
[9] Li C., Chen A., Ye J. Numerical approaches to fractional calculus and fractional ordinary differential equation, Journal of Computational Physics, 2011 Elsevier.
[10] Fu Z., Heidel J. Nonlinearity 10 (1997) 1289.
[11] Gottlieb HPW. Am. J. Phys. 64 (1996) 525.
[12] Linz SJ, Sprott JC. Elementary chaotic flow, Physics Letters A, 1999 Elsevier.
[13] Mendoza J., Araque-Lameda L., Morles EC. Understanding chaos through a Jerk circuit, Conference Paper, Jun 2016, DOI: 10.1109/TAEE.2016.7528241.
[14] Munmuangsaen B., Srisuchinwong B., Sprott JC. Generalization of the simplest autonomous chaotic system, Physics Letters A, 2011 Elsevier.
[15] Salem A., Alzahrani F. Langevin equation involving one fractional order with threepoint boundary conditions, Mathematics, 2019, 402, pp: 1-13.
[16] Salem A., Alzahrani F., Almaghamsi L. Fractional Langevin Equations with Nonlocal Integral Boundary Conditions, Mathematics, 2019, 7, 402, 1-10.
[17] Sambas A., Mamat M., Salleh Z. Design, Numerical Simulation of Jerk Circuit and Its Circuit Implementation, Advanced Studies in Theoretical Physics, July 2015.
[18] Sprott JC. Phys. Rev. E 50 (1994) R647.
[19] Sprott JC. Phys. Lett. A 228 (1997) 271.
[20] Sprott JC. A New Chaotic Jerk Circuit, IEEE Transactions on Circuits and Systems II: Express Briefs, 58(4), 2011, pages 240-243.
[21] Zourmba K., Oumate AA., Gambo B., Effa JY., Mohamadou A. Chaos in the incommensurate fractional order system and circuit simulations, International Journal of Dynamics and Control, volume 7, 2019, 94-111.

