Mechanisms driving the dispersal of hydrothermal iron from the northern Mid Atlantic Ridge

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Abstract

The dispersal of dissolved iron (DFe) from hydrothermal vents is poorly constrained. Combining field observations and a hierarchy of models, we show that the dispersal of DFe from the Trans-Atlantic-Geotraverse vent site occurs predominantly in the colloidal phase and is controlled by multiple physical processes. Enhanced mixing near the seafloor and transport through fracture zones at fine-scales interacts with the wider ocean circulation to drive predominant westward DFe dispersal away from the Mid-Atlantic ridge at the 100km scale. In contrast, diapycnal mixing predominantly drives northward DFe transport within the ridge axial valley. The observed DFe dispersal is not reproduced by the coarse resolution ocean models typically used to assess ocean iron cycling due to their omission of local topography and mixing. Unless biogeochemical models include high-resolution nested grids, they will inaccurately represent DFe dispersal from axial valley ridge systems, which make up half of the global ocean ridge crest.

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- 16 constrained. Combining field observations and a hierarchy of models, we show that
- 17 the dispersal of DFe from the Trans-Atlantic-Geotraverse vent site occurs
- 18 predominantly in the colloidal phase and is controlled by multiple physical processes.
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- models typically used to assess ocean iron cycling due to their omission of local 25 topography and mixing. Unless biogeochemical models include high-resolution
- 26 nested grids, they will inaccurately represent DFe dispersal from axial valley ridge
- 27 systems, which make up half of the global ocean ridge crest.
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- 29 149/150wds 30

31 Plain Language Summary

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33 Hydrothermal venting along mid ocean ridges supplies large quantities of the trace 34 metal iron to the ocean. Once it mixed with oxygenated seawater, precipitation leads 35 to iron being lost from the dissolved phase to generate seafloor metal deposits. 36 However, a small fraction of iron supplied escapes precipitation and remains in the 37 dissolved phase. The processes that control the retention and ocean transport of 38 hydrothermal dissolved iron is important as it has a disproportionate influence on the 39 global carbon cycle. In this work we examined the processes driving the dispersal of 40 dissolved iron from a major site of hydrothermal venting on the northern mid Atlantic 41 ridge. We found that the complex topography of the mid Atlantic ridge was crucial in steering the escape of dissolved iron in the colloidal size range out of the immediate 42 43 mid ocean ridge system. This raises challenges for the large scale ocean models 44 used to represent the global ocean iron cycle as they are typically not parameterised 45 at high enough spatial resolution. The use of multiple grids, with higher resolution 46 nests, may offer a solution to the challenge of representing the interactions of tracer 47 dispersal with complex topography.

- 48
- 49 Key points:
- 50

- 51 1. Iron is dispersed from TAG predominantly northward within the axial valley and
- 52 westward off axis, dominated by the colloidal size fraction
- 53 2. A combination of fine-scale processes are necessary to explain the dispersal both54 within and outside the axial valley
- 55 3. Coarse resolution models are impaired in their ability to constrain the broader
- 56 influence of iron supplied from axial valley ridge systems
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59 **1. Introduction**

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61 Dissolved iron (DFe) supply from hydrothermal vents has emerged as an important 62 component of the ocean iron cycle [Tagliabue et al., 2017]. Moreover, as 63 hydrothermally sourced iron is ventilated in the iron-limited Southern Ocean, there is an important link to the ocean carbon cycle [Resing et al., 2015; Tagliabue et al., 64 65 2010; Tagliabue and Resing, 2016]. Consequently, there is a need to include hydrothermal DFe supply in ocean biogeochemical models to accurately represent 66 67 the supply and cycling of this key micronutrient. Elevated iron signals have been 68 observed in plumes above most mid ocean ridge systems [Baker et al., 2002: Gamo et al., 1996; Hahm et al., 2015; Massoth et al., 1994; Rudnicki and Elderfield, 1993]. 69 70 More recently, as part of the GEOTRACES programme, iron has been shown to 71 persist as DFe above and beyond the global ridge crest system [Hatta et al., 2015; 72 Klunder et al., 2011; Nishioka et al., 2013; Resing et al., 2015; Tagliabue et al., 73 2022]. Crucial in this growing role for hydrothermalism in shaping basin scale 74 distributions is the question of how DFe is transported away from hydrothermal vent 75 sites at the >100km scale [Tagliabue and Resing, 2016].

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77 DFe (<0.2 μ m) is an operational definition that encompasses a complex array of 78 chemical species. In particular, contributions of biogenic and non-biogenic phases 79 will play an important role in the colloidal size fraction (>0.02 μ m, but <0.2 μ m) 80 [Tagliabue et al., 2017]. This is particularly true in hydrothermal settings, where large 81 fluxes of reduced soluble forms of iron interact with oxygenated seawater to drive 82 rapid changes in physico-chemical speciation [Field and Sherrell, 2000; Rudnicki and 83 Elderfield, 1993]. In the deep ocean, colloidal and soluble forms of DFe have been 84 observed to exist in a 1:1 ratio, but closer to iron sources or in the upper ocean the colloidal contribution can fluctuate notably [Bergquist et al., 2007; Fitzsimmons and 85 Boyle, 2014; Kunde et al., 2019; Nishioka et al., 2001]. Colloidal iron is typically 86 87 made up of iron (oxy)hydroxide phases and small lithogenic particles, as well as 88 biomolecules and small bacteria or viruses that interact with organics [Lough et al., 89 2019; Tagliabue et al., 2017].

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91 The global mid ocean ridge crest displays variable spreading rates and associated 92 topographic settings, with potential implications for DFe supply and transport. Inert 93 passive tracers of hydrothermal inputs, like mantle helium-3 (³He), are elevated in 94 basins with fast spreading ridges, like the Pacific, and depressed where ridge 95 spreading rates are lower, like the Atlantic [Jenkins et al., 2019]. This is important as 96 although hydrothermal dFe inputs are parameterised in global models with constant 97 DFe:³He ratios following *Tagliabue et al.* [2010], hydrothermal DFe anomalies are 98 higher than would be expected along the slow spreading mid Atlantic ridge based on 99 ³He values (e.g. [*Hatta et al.*, 2015; *Saito et al.*, 2013]). Adding further complexity is 100 the fact that the slow spreading ridges like the mid-Atlantic ridge are typified by large 101 axial valleys with topographic relief varying by kilometres across relatively small

102 spatial scales. It is not known how resolving these scales of variability affects the

transport of DFe into the wider basin and the implications for coarse resolution ocean

biogeochemical models that are typically used to test hypothesis about hydrothermal
 DFe supply and cycling [*Roshan et al.*, 2020; *Somes et al.*, 2021; *Tagliabue et al.*,

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108 Here we present new observations of DFe and colloidal Fe from the trans Atlantic

109 geotraverse (TAG) hydrothermal site northern mid-Atlantic ridge as part of the UK

110 GEOTRACES GA13 section. Our data document transport of DFe northwards within 111 the axial valley and westward off axis into the wider basin at a range of spatial

112 scales. Using a suite of model experiments at a range of resolutions, we diagnose

113 the candidate physical processes that drive this behaviour and demonstrate that they

are absent in coarse resolution models. This raises important questions about

115 whether coarse resolution models are appropriate tools to explore iron cycle

116 pathways associated with DFe supply from slow spreading ridge systems.

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118 **2. Methods** 119

2022].

The UK GEOTRACES GA13 voyage sailed between Southampton and Guadeloupe in 2017/8 and as part of the sampling a detailed process study was conducted around the TAG hydrothermal vent system. A number of stations were sampled north and south of TAG within the axial valley, as well as east and west off axis into the Atlantic basin (Figure 1, Supplementary Figure 1). Station spacing ranged from 10-30km close to TAG and up to 100-200km for the farthest stations.

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127 All sampling protocols followed those established by the GEOTRACES program 128 [Cutter et al., 2010]. Water samples were collected using Teflon coated Niskin-X 129 bottles (Ocean Test Equipment) on a kevlar coated conducting wire. Water samples were filtered (0.2 µm, Sartorius) into acid clean low-density polyethylene bottles for 130 131 DFe. A separate aliquot seawater was filtered through 0.02 µm filters (Anotop, 132 Whatman) for soluble Fe (SFe) (Ussher et al., 2010). All filtration was done in a class 133 100 clean laboratory on board the ship. Samples were acidified onboard to 0.024 M 134 (UpA HCI, Romil). Samples were analysed for Fe concentrations using flow injection 135 chemiluminescence and inductively coupled plasma mass spectrometry. In both 136 cases GEOTRACES reference materials (D2, GSC, GSP) were analysed and there 137 was a maximum difference of 0.14 nM between measured and consensus values 138 (Measured GSC 2.04 ± 0.03 nM, consensus GSC 2.18 ± 0.08 nM). For further 139 details see [Lough et al., 2022].

140

141 Two Lagrangian dispersion experiments were carried out using the 3D velocity

field of the GIGATL3 simulation. GIGATL3 is a regional simulation of the

143 ocean physical state in the Atlantic Ocean based on the primitive-equation

144 CROCO model, developed using the Regional Oceanic Modeling System

145 [Shchepetkin and McWilliams, 2005]. The GIGATL3 simulation has a nominal

horizontal resolution of 3 km and features 100 terrain-following vertical levels, with

147 stretching near the surface and seafloor (supplementary figure 2). The

148 GIGATL3 bathymetry is taken from the global 30 arc second SRTM30plus data set

- 149 [Becker et al., 2009]. The initial state and lateral boundary conditions for velocity, sea
- surface height, temperature, and salinity are supplied by the Simple Ocean Data

151 Assimilation data set [Carton and Giese, 2008]. Atmospheric forcing was supplied at 152 hourly resolution by the Climate Forecast System Reanalysis [Saha et al., 2010]. 153 Tidal forcing derived from TPXO7 is included. The 3D velocity field is saved hourly 154 and linearly interpolated in space and time to perform two Lagrangian experiments 155 using the Pyticles software [Gula et al., 2014]. We used these two experiments to 156 improve the robustness of the dispersion patterns and diagnostics are calculated 157 using all particles across both experiments. The two experiments are strictly 158 identical in terms of the particle seeding set up and integration time, only differing by 159 starting point. The first experiment starts on 2008-08-29 whereas the second starts 160 on 2010-06-05. Each experiment consists of releasing 25 particles every 6 h at the 161 TAG vent site plume depth for 8 months. Diagnostics on particle spreading were 162 performed for particles with ages between 10 and 180 days. As particles are 163 continuously released, approximately 22000, 14000 or 5800 particles have an overall 164 lifespan of 10, 90 or 180 days, respectively, for each experiment. 165 166 Modelling experiments are also conducted using a global scale ocean general 167 circulation model. We conducted passive release experiments using two 168 configurations of the Nucleus for European Modelling of the Ocean (NEMO) model. 169 The first was the standard global configuration using the ORCA2 configuration at a 170 horizontal resolution of 2°x cos(latitude) curvilinear grid, with an enhancement to 171 0.5° around the equator and 31 irregularly spaced vertical levels. This NEMO-172 ORCA2 configuration is typical of those models coupled to biogeochemical models 173 to address guestions regarding biogeochemical cycling (e.g. Tagliabue et al., 2022). 174 We applied the default settings and boundary conditions of the reference 175 configuration ORCA2 ICE PISCES [NEMO-Consortium, 2019]. NEMO-ORCA2 is 176 forced with CORE-II normal year atmospheric forcing, with the NCAR bulk formulae 177 [Large and Yeager, 2008]. After the initial spin-up, we conducted an idealised 178 passive tracer release. The passive tracer concentration is continuously set to one at 179 the deepest grid cell closest to the TAG site. The passive tracer fluxes at the surface, 180 at the lateral boundaries and at the bottom are set to zero and the surface passive 181 tracer concentration is restored to zero. The model is run for another 30 years (year 182 130 to 160) and the tracer spread is monitored. This scenario is repeated with a 183 NEMO-AGRIF configuration with the addition of two nested regions, covering the 184 TAG site, via adaptive mesh refinement package [Debreu et al., 2008]. Two level, 185 two-way nesting is used: the first level covers a region in the subtropical North 186 Atlantic (dashed rectangle in supplementary figure 3) with refinement ratio of 4 in 187 both latitude and longitude (to give a horizontal resolution of $1/2^{\circ}$). The second level 188 of nesting is applied over a region with the TAG site in the centre, with further 189 refinement ratio of 4 (solid line rectangle in supplementary figure 3) reaching a 190 horizontal resolution of 1/8° or 12.48km. The passive tracer is released only in this 191 high-resolution region. The model bathymetry of the two nested regions is 192 constructed from the 5 arc minute resolution global bathymetry from ETOPO5. The 193 initial conditions and the surface forcing functions of the nested regions in NEMO-194 AGRIF are interpolated from NEMO-ORCA2 fields using the NEMO nesting tools. 195

196 **3. Results**

198 **3.1 Dispersion of DFe from the TAG hydrothermal vent field**

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200 The TAG site is a well-studied hydrothermal vent system, sited within the mid-201 Atlantic ridge axial valley (Figure 1). As part of the GA03 GEOTRACES section, DFe measurements were taken at TAG, but the station spacing for full-depth profiles 202 203 exceeded 500km [Hatta et al., 2015]. Within the axial valley DFe is predominantly 204 dispersed northwards in the colloidal phase. The hydrothermal DFe anomaly of around 80 nM we observed at TAG matches that observed during the GA03 voyage 205 206 [Hatta et al., 2015] and persists at 3-4 nM at the stations 19 and 30km north (Figure 207 1). In contrast, DFe drops below concentrations of 2 nM for the station 30km south, indicating greater dilution and/or removal from the dissolved phase (Figure 1). At 208 209 TAG, the highest concentrations of DFe are associated with very low soluble Fe 210 fraction (<10%), indicating the dominance of colloidal Fe. At the depths of the 211 greatest hydrothermal DFe signals, the soluble dFe fraction within the axial valley 212 also remains low (<25%) within 30km of TAG, again indicating dominance of Fe 213 colloids in the hydrothermal DFe signal throughout the valley.

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215 There is a contrasting DFe signal east and west from TAG off axis from the mid-216 Atlantic ridge. At the largest spatial scales, strong hydrothermal DFe anomalies 217 persist 140-250km west of TAG (stations 28 and 29), but are absent at stations 140-218 250km to the east (stations 33 and 32). A marked dFe anomaly between 2,200-219 3,400m (centered on 3000m) of 2.64nM 140km from TAG (station 29) declines to a 220 more localised anomaly of 0.95nM 250km west of TAG (station 28). Both of these 221 DFe signals are above the concentrations observed at this depth and latitude at the 222 eastern stations. Notably, the elevated DFe concentrations 140-250km west of TAG 223 remain associated with low soluble DFe fractions, indicating the importance of Fe 224 colloids. Consistent with the absence of hydrothermal DFe input or transport, the 225 soluble iron fraction 140-250km to the east is closer to the 50% typically observed in 226 the deep ocean [Kunde et al., 2019]. 227

3.2 Processes shaping the dispersion from TAG over different space and time scales

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The particle release experiments conducted with the high-resolution particle tracking
model reveal the role of different processes shaping dispersion from TAG over
different space and timescales (Figure 2). We particularly note three stages of
physical dispersal.

235

236 10-30 days: In the immediate period following their release, particles preferentially 237 spread within the axial valley in a northeast-southwest direction, largely under 238 topographic control (Figure 2ab). The impact of topography is illustrated by the 239 topostrophy parameter, τ [Holloway, 2008], which is elevated and positive within the 240 valley and much reduced off axis (Supplementary Figure 4). That τ is >0 indicates a 241 cyclonic circulation within the axial valley basins, consistent with observed and 242 modelled circulations within the MAR valley [Lahaye et al., 2019]. The topography of 243 the mid Atlantic ridge axial valley also allows particles to rapidly escape the axial 244 valley along isopycnal surfaces to the southwest via fracture zones within only a few 245 days (Figure 2a).

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247 *60-90 days:* Particles that escape the valley spread isotropically along density

- surfaces (Figures 2c,d). The topographic control is now very weak on average as
- topographic slopes are weaker and particles are now well above the seafloor

(Supplementary Figure 4). Consequently, particles spread mostly along isopycnals
due to the lesser influence of diapycnal mixing processes that were occurring on and
within the axial valley system.

254 120-180 days: After 4-6 months, particles preferentially spread westward outside the 255 axial valley (Figure 2ef) due to the combination of the large-scale mean sub-tropical 256 gyre circulation and the planetary beta effect that constrains mesoscale vortices to 257 travel westward [Killworth, 1983; Nof, 1981]. By this stage, transport within the axial 258 valley also demonstrates a predominantly northward signal as particles fill the 259 northern axial valley basin within the mid-Atlantic ridge north of TAG. Notably, 260 dispersal westward off axis into the wider Atlantic Ocean basin and northward within 261 the axial valley matches the observed DFe concentration anomalies closely (Figure 262 1), as well as understanding from prior work [Thurnherr et al., 2002; Vic et al., 2018; 263 Yearsley et al., 2020]. 264

3.3 The importance of bottom topography in representing hydrothermal Fe supply in global ocean models

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268 The set of simulations conducted with two NEMO allow us to explore how the 269 dispersion of hydrothermal tracers from TAG are represented by coarse resolution 270 NEMO-ORCA2 global models and with the NEMO-AGRIF configuration. The NEMO-271 AGRIF configuration has a 1/8° (or 12.48km) resolution regional nested grid around TAG (supplementary figure 3), but the vertical resolution of NEMO-AGRIF and 272 273 NEMO-ORCA2 are around three-fold lower than the particle tracking model. This set 274 of model configurations were designed to link the very high-resolution regional 275 modelling with the types of models used for larger scale biogeochemical modelling 276 that tend to have horizontal resolutions of between 1-5° [Roshan et al., 2020; Somes 277 et al., 2021; Tagliabue et al., 2022].

278

279 In general, the NEMO-AGRIF model shows very similar dispersal patterns to the 280 high-resolution particle model, with strong signals remaining localised within the axial 281 valley and spreading preferentially north within the axial valley and westward off-axis 282 (Figure 3). In contrast, the NEMO-ORCA2 configuration typically used for 283 assessments of iron biogeochemistry show two major deficiencies, relative to the 284 nested NEMO-AGRIF and high-resolution particle models (Figure 2), as well as the 285 observations (Figure 1). First, high concentrations of tracer do not remain trapped in 286 the ridge system. Second, large scale dispersal operates equally away from the 287 ridge both east and west, rather than predominantly to the west. These deficiencies 288 in the zonal dispersal can be illustrated by a section taken along 26N, with the 289 NEMO-AGRIF nested model displaying dispersal was (i) more restricted and (ii) 290 predominantly westward as compared to the coarse resolution NEMO-ORCA2 291 model. Notably, both NEMO-ORCA2 and NEMO-AGRIF share the same number 292 and arrangement of vertical levels, highlighting the importance of the horizontal 293 resolution in improving the agreement with both high-resolution particle tracking 294 models and inferences regarding dispersal gleaned from direct observations of DFe. 295 To compare the model configurations more directly against observations we merged 296 the GA13 data from this study with DFe observations from the GA03 GEOTRACES 297 section that followed the same cruise track, which displays a similar westward propagation off axis from TAG (Figure 3e). 298 299

300 4. Discussion

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302 4.1 Processes driving the dispersal of iron from the mid-Atlantic ridge 303

304 Dispersal of DFe from TAG is controlled by a combination of local mixing across 305 density surfaces and the specific geometry of the mid Atlantic ridge. Using our 306 particle tracking simulations, we tracked the cumulative changes in density during 307 the particle lifetimes to identify the importance of across isopycnal mixing in 308 explaining whether DFe dispersal west of TAG (Figure 3e) is over or around 309 topography. Enhanced mixing is associated with small scale internal tides and 310 mesoscale currents interacting with topographic features, such as mid ocean ridges 311 [Vic et al., 2019]. The average cumulative change in density across all particles 312 highlights strong transfer to lighter density surfaces on the ridge crest and within the 313 axial valley (Figure 4a). Small average changes outside the axial valley are 314 associated with large variability (Figure 4b). Focussing on an example site outside 315 the valley, we can see that while average cumulative density changes are close to 316 zero, 64% of all particles experience lightening (Figure 4c, Supplementary Figure 5), 317 leading to around 200m elevation in absolute depth (Figure 4d). Within the valley, 318 the change in cumulative density is much more striking (Figure 4c), with changes in 319 absolute elevation of closer to 100m (Figure 4d). The larger changes in absolute 320 depth, despite smaller changes in cumulative density outside the valley are 321 associated with slumping and heaving of isopycnal layers. Overall, the significant 322 population of particles outside the valley without strong cumulative density changes 323 demonstrates the transfer of particles through fracture zones and highlights the role 324 of local geometry of the mid Atlantic ridge system (see Sec 3.2 and 3.3, 325 Supplementary Figure 5). Within the valley, across density mixing associated with 326 topography is much more important.

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328 By examining the lifetime of particles reaching the local particle maxima at any given 329 location, we can quantify timescales of dispersion across the two Lagrangian 330 experiments (Figure 4e). Dispersion from TAG takes at least 100 days to reach 331 140km west of TAG (station 29) or exceeds 150 days to reach station 28 at 250km 332 from TAG (Figure 4). Since most hydrothermal iron is associated with maxima in 333 colloidal Fe, this indicates relatively strong stability of colloidal iron phases during 334 transport (for at least 6 months). We note that these estimates should be seen as 335 minimum estimates for transport to the different locations. If we integrated the 336 particle model longer, we would potentially also observe the arrival of older particles 337 and find evidence for longer term stability of colloidal iron. The extended lifetime of 338 colloidal iron may be driven by slower oxidation of colloidal iron [Gartman and 339 Luther, 2014] or its association with organic phases [Fitzsimmons et al., 2017]. 340

341

4.2 Using model-observation studies to quantify iron cycle mechanisms 342

343 Combining global ocean biogeochemical modelling experiments with GEOTRACES 344 datasets have played a key role in identifying significant processes shaping the 345 ocean iron cycle and their wider impacts. However, their results may be 346 compromised when the model resolution is insufficient for the system of interest. Our 347 results show that only model experiments at the 10s of km horizontal resolution can 348 accurately represent the dispersal of DFe from the TAG site on the mid-Atlantic 349 ridge, far exceeding the typical resolution of global ocean biogeochemical models

350 (100-500km). This is due to the complex topography of the axial valley setting of the 351 TAG site that is not resolved in global ocean biogeochemical models. Where ridge 352 topography is less variable, for instance at faster ridge spreading sites, e.g. the East Pacific Rise, coarser resolutions may be sufficient [Resing et al., 2015]. Our results 353 354 suggest that new solutions, accounting for high-resolution sub grids, are required to properly represent DFe dispersal at scales exceeding 100km from the axial valley 355 356 settings that make up around half of the mid ocean ridge crest. Improved vertical 357 resolution may also be important, but this was not assessed directly in this study as 358 both NEMO model configurations had the same number and arrangement of vertical 359 levels. 360

361 The role of representing DFe input and wider transport around local topography in 362 coarse resolution models may be more generically relevant. For instance, continental 363 shelves can also be associated with complex topographic geometry and coarse 364 resolution models may face similar challenges in properly representing the dispersal 365 of tracers supplied, such as DFe. For instance, in the North Pacific, the DFe inputs 366 from sediment resuspension disperses from the sea of Okhotsk into the wider North 367 Pacific basin via North Pacific Intermediate Water (NPIW) [Nishioka et al., 2021; 368 Nishioka et al., 2020]. An important component of this dispersal is the strong 369 topographically induced diapycnal mixing that occurs over the Kuril straits [Yagi and 370 Yasuda, 2012], transporting DFe onto the NPIW isopycnals to then spread 371 throughout the North Pacific basin [Nishioka et al., 2020]. Tidal mixing along the 372 shelf break may also be an important component of DFe supply in the southern 373 Bering Sea [Tanaka et al., 2012], as well as seamounts [Lavelle et al., 2004]. 374 Therefore, it is important to consider any biases in the DFe dispersal from both 375 hydrothermal and sediment inputs that may arise from insufficient resolution in 376 process-based models.

377

378 **5. Conclusions**

379

380 Using a novel sampling strategy at a scale of 10s and 100s of km around the TAG 381 vent site, we document the predominant transport pathways of hydrothermal DFe. 382 Our observations indicated DFe was transported northward within the axial valley 383 and westward off axis in the colloidal size fraction. Dispersal within the valley arose 384 due to the diapycnal mixing that resulted from topographic interaction. Transport off 385 axis resulted from both diapycnal mixing and the fine scales of axial valley geometry, 386 especially fracture zones. The dispersal patterns of DFe were reproduced with a 387 high-resolution particle tracking model and a global model with a nested ~10km 388 horizontal grid, but not with a global ocean configuration. This raises challenges for 389 correctly representing DFe supply and the associated biogeochemical impacts from 390 axial valley hydrothermal vent systems, as well as other supply mechanisms around 391 local topographical features.

392

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394

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- 422 https://www.geotraces.org/geotraces-intermediate-data-product-2021/
- 423
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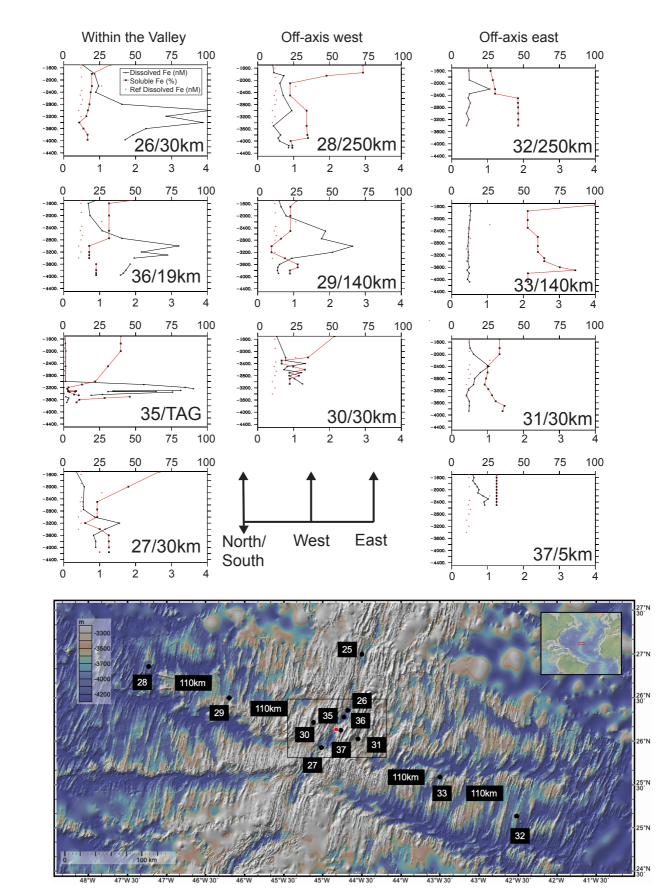
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595 Figure 1. Vertical profiles of dissolved Fe (nM, black line, bottom x-axis), percentage 596 of dFe present as soluble Fe (%, red line, top x-axis) and the iron profile from a reference station (station 32, red dots, nM) for the rage of stations within and outside 597 598 the axial valley. The left-hand column shows stations within the valley, the central column shows stations from the west and the rightmost column shows stations to the 599 600 east. Consult the map and supplementary figure 1 for more information on the 601 stations and their spacing. Distances noted on each panel are the km from TAG. 602

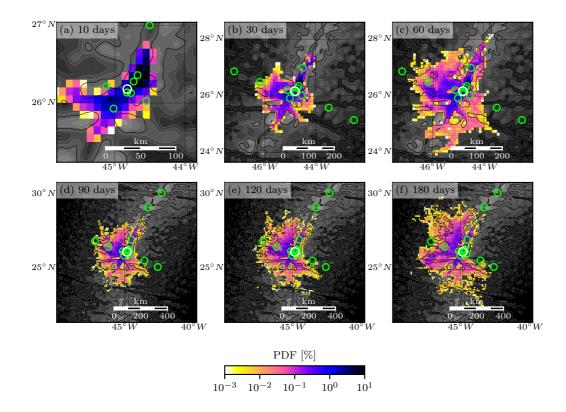
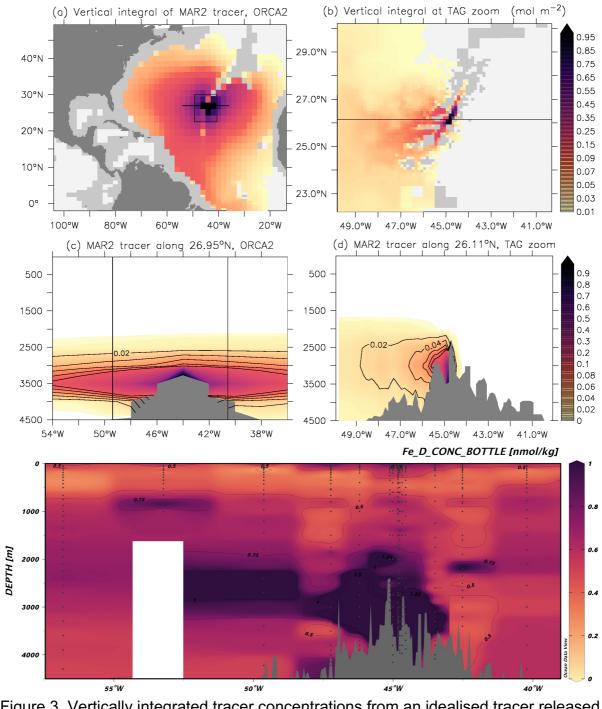
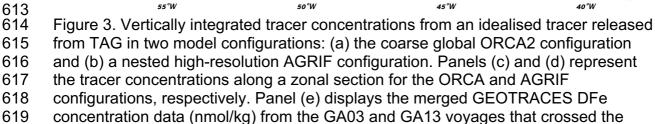
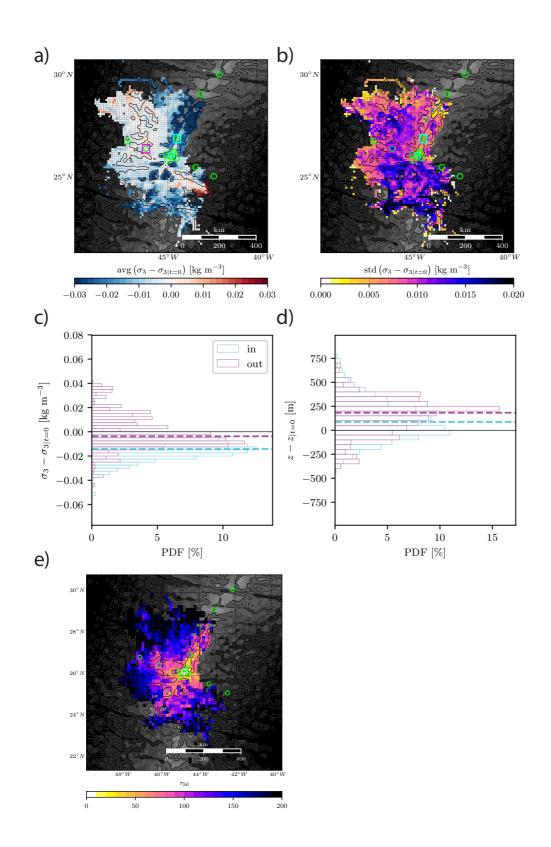


Figure 2. Probability density function of particle presence after (a) 10, (b) 30, (c) 60, (d) 90, (e) 120 and (f) 180 days of dispersion from TAG (white circle). Individual positions are binned onto a 0.1-degree resolution grid. Green circles are the Fridge stations.





- 620 same region.

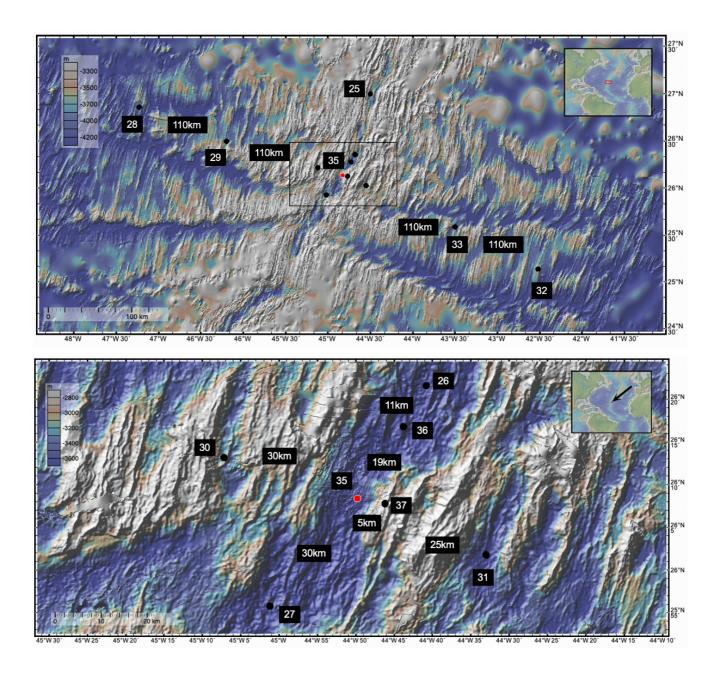


626 Figure 4. a) mean and b) standard deviation of the cumulative density change relative to initial density. Panels c) and d) represent histograms of density and

absolute depth for particles inside and outside the valley (see squares on map on

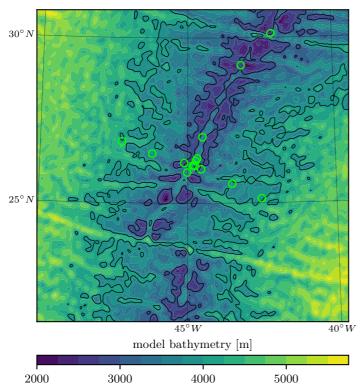
- panel a). e) Median ages (in days) of particles at the depth of maximum particle density after 6 months of model simulation

634 Supplementary Figures

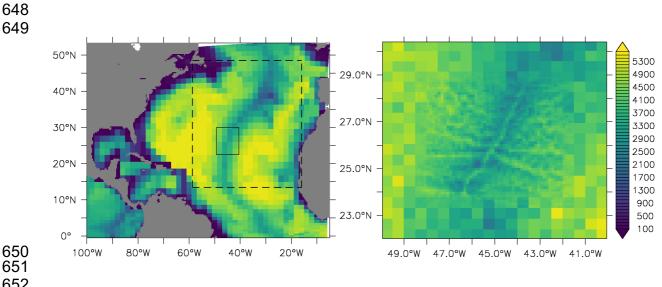


Supp Fig 1. JC156 Cruise stations. Red circle marks TAG at station 35 and labels
represent the spacing between stations. The closest station spacing is between 530km close to TAG and extends to 110km further off axis. Upper panel shows the full
domain around the TAG site and the lower panel zooms in on stations immediately

- 642 adjacent to TAG enclosed by the black square.

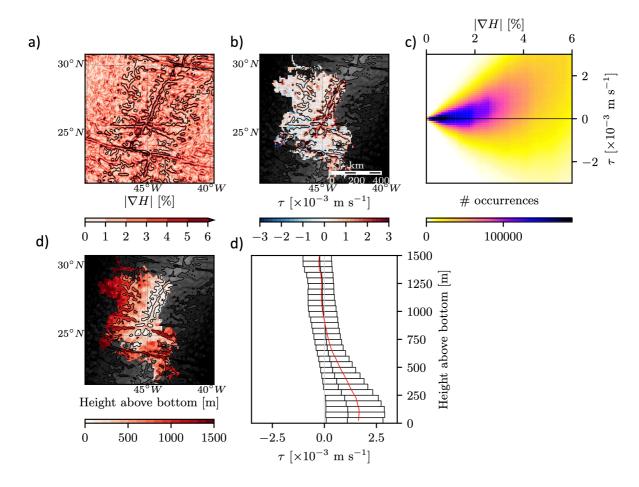


Supp Fig 2. Bathymetry in the GIGATL3 model with the GA13 sampling locations as green circles.



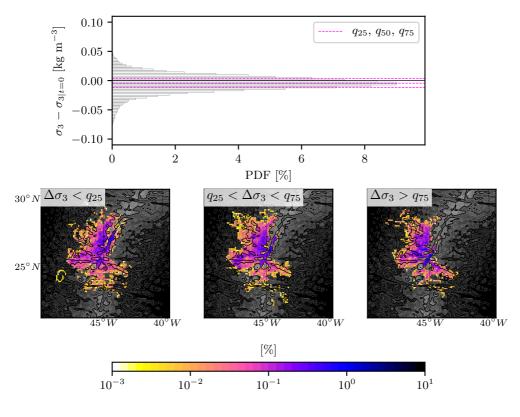
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Supp Fig 3. Bathymetry in the ORCA2 and the AGRIF nested model configuration. The AGRIF nesting is at 0.5x0.5 degrees (dashed box) and at 1/8 x1/8 degrees (black box)





Supp Fig 4. Diagnostics related to topostrophy. (a) slope of bathymetry (grad(H)) (b) topostrophy (τ) computed along particle trajectories and bin-averaged (c) Histogram of τ vs grad(H) (d) Bin-averaged height above bottom of particles and (e) quartiles and mean of in height-above-bottom coordinates.



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667 Supp Fig 5. (Top) histogram of density change and (bottom) histogram of position 668 discriminated on density change for all particles. Lower panels show maps of the 669 different distributions according to the density change quartile: (Left, lower quartile) 670 the 25% of particles that have lightened the most, (middle, middle quartile) 50% of 671 the particles that have shown the least change in density (i.e. in the middle of the 672 histogram) and (right, upper quartile) the 25% of the particles that have become 673 most dense.