Nonlinear earthquake response of marine sediments with distributed acoustic sensing

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November 22, 2022

Abstract

Seismic waves can be significantly amplified by soft sediment layers. Large dynamic strains can trigger a nonlinear response of shallow soils having low strength, which is characterized by a shift of the resonance frequencies, ground motion deamplification, and in some cases, soil liquefaction. We investigate the response of marine sediments during earthquake ground motions recorded along a fiber-optic cable offshore the Tohoku region, Japan, with Distributed Acoustic Sensing (DAS). We compute AutoCorrelation Functions (ACFs) of the ground motions from 103 earthquakes in different frequency bands. We detect time delays in the ACF waveforms that are converted to relative velocity changes (dv/v). dv/v drops, which are characteristic of soil nonlinearity, are observed during the strongest ground motions. Moreover, the dv/v values show a strong variability along the cable. This study demonstrates that DAS can be used to infer the dynamic properties of the shallow Earth with an unprecedented spatial resolution.

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Key Points:

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10	• AutoCorrelation Functions (ACFs) of earthquakes recorded by Distributed Acous-
11	tic Sensing (DAS) exhibit phase delays during ground motions
12	• ACF time delays are converted to relative velocity drops in the medium, which
13	characterize soil non-linearity
14	• DAS is used to infer the nonlinear behavior of soils with an unprecedented spa-
15	tial resolution

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²⁹ Plain Language Summary

Seismic waves from earthquakes are amplified by shallow and soft sediment layers 30 of the Earth. This amplification is linear for weak seismic waves, but can become highly 31 nonlinear during strong ground motions. Nonlinear soil response, which can lead to a 32 complete failure of the ground through soil liquefaction, threatens the safety of human-33 made constructions and needs to be accurately characterized. We study the response of 34 marine sediments offshore the Tohoku region in Japan using earthquake data recorded 35 along 43.3 km of a fiber-optic cable with Distributed Acoustic Sensing (DAS). We use 36 an autocorrelation approach to analyze the ground motions from 103 earthquakes recorded 37 by thousands of DAS channels. We detect a clear nonlinear behavior of shallow sediments 38 during the strongest ground motions. Moreover, we show that soil nonlinearity signif-39 icantly varies along the cable. Our methodology could easily be applied to earthquake 40 DAS data recorded in populated and seismically active regions to help better understand-41 ing the dynamic behavior of shallow soils. 42

43 **1** Introduction

Local geological conditions can significantly impact the propagation of incoming 44 seismic waves from earthquakes. In particular, shallow, soft, and unconsolidated sedi-45 ment layers are well known to amplify earthquake ground motions (Sanchez-Sesma, 1987), 46 which can lead to catastrophic events such as during the 1985 moment magnitude (M_w) 47 8.0 Michoacán earthquake in Mexico (Anderson et al., 1986; Campillo et al., 1989). When 48 subjected to weak dynamic strains (i.e., less than 10^{-4} and 10^{-8} for field observations 49 and laboratory experiments, respectively; Ishihara, 1996; TenCate et al., 2004), shallow 50 soils linearly amplify seismic waves. During large dynamic strains, however, soft sedi-51 ments can behave nonlinearly (e.g., Field et al., 1997; Ostrovsky & Johnson, 2001). Soil 52 nonlinearity is generally characterized by a relative reduction of the high-frequency ground-53 motion amplification, which is related to an increase of damping in the medium, and a 54 shift of the resonance frequency to lower frequencies due to a reduction of the shear mod-55 ulus (Beresnev & Wen, 1996; Brunet et al., 2008; Bonilla et al., 2011; Lyakhovsky et al., 56 2009; Zaitsev et al., 2005). In some cases, large dynamic strains can trigger a complete 57 failure of cohesionless and saturated shallow sediments through soil liquefaction (Kramer, 58 1996), which can have disastrous consequences for human infrastructures as observed dur-59 ing the 1964 Niigata (Japan, Ohsaki, 1966) and 2010–2011 Christchurch (New Zealand, 60 Quigley et al., 2013) earthquakes. Therefore, characterizing the nonlinear response of 61 shallow sediments to earthquake ground motions is critical for better mitigating seismic 62 risk. 63

Several empirical methods have been developed to assess the response of soils to
 ground motions. A classical approach relies on computing the spectral ratio of earthquakes

recorded at a soft-soil site and at a nearby reference rock site (Borcherdt, 1970; Field 66 & Jacob, 1995; Bonilla et al., 1997). However, this method suffers from the fact that a 67 reference site may not always be available in the vicinity of the site of interest. Another 68 approach consists in using pairs of surface-borehole stations to detect potential soil nonlinear elastic behavior between the two sensors (Bonilla et al., 2011; Minato et al., 2012; 70 Nakata & Snieder, 2011; Régnier et al., 2013; Sawazaki et al., 2006; Takagi et al., 2012; 71 Wen et al., 1995). While this technique allows us to isolate the shallow subsurface re-72 sponse from the earthquake source and path effects, pairs of surface-borehole instruments 73 are expensive to install and their low spatial coverage prevents us from capturing small-74 scale lateral variations. 75

AutoCorrelation Functions (ACFs) calculated from data recorded by surface seis-76 mometers yield the reflectivity response of the underlying elastic structure (Claerbout, 77 1968; Wapenaar, 2003). This technique has been primarily applied to image interfaces 78 with strong seismic impedance contrasts using earthquake (Delph et al., 2019; Pham & 79 Tkalčić, 2017; Tork Qashqai et al., 2019; Viens, Jiang, & Denolle, 2022) and ambient seis-80 mic field (ASF; Gorbatov et al., 2013; Ito et al., 2012; Kennett, 2015; Saygin et al., 2017; 81 Spica et al., 2020; Viens, Jiang, & Denolle, 2022) datasets. Repeated ACF computations 82 through time from continuous ASF time series have also been used to monitor tempo-83 ral seismic velocity changes in the subsurface in different environments, such as volcanic 84 (De Plaen et al., 2016; Sens-Schönfelder & Wegler, 2006; Yates et al., 2019) and earth-85 quake source (Hobiger et al., 2014; Ohmi et al., 2008; Wegler et al., 2009) regions. How-86 ever, the partitioning of surface and body waves in ACFs computed from the ASF is gen-87 erally unknown and hinders the interpretation of the measured velocity changes (Nakahara, 88 2015). To ease the interpretation, ACFs have also been computed from earthquake P-89 , S-, or coda-wave windows (Bonilla et al., 2019; Bonilla & Ben-Zion, 2020; Nakahara, 90 2015; Qin et al., 2020). Bonilla and Ben-Zion (2020) showed that the first negative peak 91 of ACFs calculated during earthquake ground motions corresponds to the seismic-wave 92 two-way travel time between the sensor and the first major interface below the station, 93 and captures the soil non-linear response. Moreover, Bonilla et al. (2019) and Qin et al. 94 (2020) showed that the response of the shallow subsurface obtained from ACFs at sur-95 face stations yields a similar estimation of the soil nonlinear behavior as that from a surface-96 borehole station configuration. In other words, ACFs can isolate the site response term 97 from the earthquake source and path effects, which makes single-component stations a 98 powerful tool to analyze shallow sediment nonlinear behavior. 99

Mapping local site effects with data-driven techniques remains challenging due to 100 the large density of seismometers needed to capture complex spatial variations of the seis-101 mic wavefield. In some cases, a large station coverage can be nearly impossible to attain 102 due to environmental or physical constrains, especially in urban and underwater areas. 103 Nevertheless, recent technological advances in Distributed Acoustic Sensing (DAS) of-104 fer an unprecedented opportunity to measure the Earth's vibrations over tens of kilo-105 meters with a dense spatial resolution ($\sim 1-10$ m) by turning ground-coupled fiber-optic 106 cables into arrays of sensors (Hartog, 2017). DAS uses an optoelectrical interrogator to 107 probe fibers with a laser sending thousands of short pulses of light every second. As each 108 pulse of light travels down the fiber, some of the light is reflected back to the interroga-109 tor in a process known as Rayleigh backscatter. External forcing, such as seismic waves, 110 generate phase shifts of the back-scattered Rayleigh light, which are measured by the 111 interrogator. The measured phase shifts are finally linearly converted to longitudinal strain 112 (or strain-rate) along the cable over a sliding spatial distance (i.e., the gauge length). Both 113 fit-to-purpose and existing telecommunication fiber-optic cables have been used to record 114 high-fidelity earthquake wavefields (Lellouch et al., 2019; Spica et al., 2022; Wang et al., 115 2018; Zeng et al., 2017). One great advantage of telecommunication fibers is that they 116 have been widely deployed, from the oceans' bottom to nearly every street in large de-117 veloped cities, to sustain our modern telecommunication network. Therefore, DAS could 118

complement expensive urban and offshore seismic array deployments by probing exist ing telecommunication cables to capture the full extent of earthquake wavefields.

In this study, we analyze the response of shallow marine sediments to 103 earth-121 quakes recorded along a telecommunication cable offshore the Sanriku coast in Japan 122 by a DAS experiment (Figures 1a-b). We calculate ACFs from the earthquake ground 123 motions after filtering the data into different frequency bands to infer the soil response 124 at different depths. We detect changes in the ACF time series that are converted to rel-125 ative velocity changes to characterize the soil linear and nonlinear regimes below each 126 127 DAS channel. We show that the relative velocity changes exhibit spatial variations along the cable which are strongly influenced by the ACF frequency ranges. We finally discuss 128 our results and the potential of DAS for extracting soil parameters with an unprecedented 129 spatial resolution. 130

¹³¹ 2 Data and Methods

2.1 DAS data

The Earthquake Research Institute, The University of Tokyo, operates an ocean-133 bottom observatory composed of three 3-component accelerometers and two tsunami me-134 ters offshore the Sanriku Coast (Figure 1b; Kanazawa & Hasegawa, 1997; Shinohara et 135 al., 2021, 2022). The data recorded by the instruments are streamed in real-time to the 136 landing station located in the city of Kamaishi through a submarine telecommunication 137 cable. The cable contains six dark (unused) dispersion-shifted single-mode optical fibers 138 with a wavelength of 1,550 nm, which are suitable for DAS measurements. Moreover, 139 the first 47.7 km of the cable are relatively straight and are buried under 0.6-0.7 m of 140 sediments, which guarantees a good coupling of the fiber. 141

An AP Sensing N5200A DAS interrogator unit (Cedilnik et al., 2019) probed one 142 of the dark fibers between November 18 and December 2, 2019, and recorded continu-143 ous data over the first 70-km of the cable with a sampling rate of 500 Hz. The gauge length 144 and spatial sampling are set to 40 m and 5.1 m, respectively. During the two weeks of 145 measurement, hundreds of earthquakes were recorded by the DAS system. We first con-146 vert the raw DAS data to strain (Shinohara et al., 2022) and focus on the ground mo-147 tions from 103 earthquakes that were clearly recorded by all the DAS channels (Figure 148 1a-b). The velocity magnitude (M_V) of the earthquakes ranges between 1.0 and 6.3, and 149 we show the strain waveforms of a M_V 2.5 earthquake in Figure 1c. This event occurred 150 on November 28, 2019 at 14:17:32UTC at a depth of 30 km. Clear P- and S-wave arrivals 151 can be observed at most channels as well as locally generated surface waves which sig-152 nificantly extend the ground motion duration. 153

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2.2 Autocorrelation functions and relative velocity changes

For each earthquake and each DAS channel, we compute the time derivative of the 155 strain data to retrieve strain-rate waveforms, which are proportional to acceleration time 156 series. We then bandpass filter the strain-rate data between 2 and 30 Hz (all filters are 157 two-pass four-pole Butterworth bandpass filters) and select a fixed 15-s window start-158 ing 5 s before the earthquake absolute maximum amplitude. We then further bandpass 159 filter the strain-rate waveforms into 19 frequencies bands (e.g., 2-4, 3-6,..., 20-40 Hz) and 160 compute ACFs over the fixed 15-s window using the phase correlation method in the fre-161 quency domain (Schimmel & Paulssen, 1997; Ventosa et al., 2019). We show the band-162 pass filtered strain-rate waveforms of the M_V 2.5 earthquake together with their corre-163 sponding ACFs at channel 5000 in Figures 1d and 1e, respectively. The ACFs are cal-164 culated around the S-wave direct arrival and we therefore expect their first negative peak 165 to capture the S-wave two-way travel time (Bonilla & Ben-Zion, 2020). Moreover, the 166 different frequency bands allow us to sample different depth of the media, with low-frequency 167



Figure 1. (a) Topographic map of the Japanese Islands and their surroundings including the 103 earthquakes used in this study. The red rectangle denotes the region near the cable shown in (b). (b) Bathymetric map offshore the Sanriku coast including the location of the seafloor cable observation system. The orange line denotes the buried section of the cable used in this study (i.e., channels 500 to 9000) and the location of channel 5000 is indicated by the red cross. The white circles and purple inverted triangles show the positions of the accelerometers and tsunamimeters, respectively. The location of the velocity magnitude (M_V) 2.5 event (red circle) shown in (c) and that of other nearby earthquakes (gray circles) are highlighted. The magnitude scale is the same as in (a). (c) Strain waveforms of the M_V 2.5 event bandpass filtered between 2 and 30 Hz between channels 500 and 9000. The waveform amplitudes are clipped for visibility. (d) Strain-rate waveforms of the M_V 2.5 earthquake bandpass filtered in different frequency bands at channel 5000. The gray area denotes the time period over which ACFs are calculated. (e) Amplitude normalized ACFs computed from the waveforms shown in (d).

bandpass filtered ACFs displaying later arrivals as they sample deeper media compared
 to high-frequency ACFs.

The 103 earthquake waveforms analyzed in this study generated various levels of 170 dynamic strain along the cable. In Figures 2a-d, we show the ACFs calculated for all the 171 earthquakes after bandpass filtering the strain-rate data in the 10-20 Hz and 15-30 Hz 172 frequency bands at channels 5000 and 7000. We also show the dynamic peak strains com-173 puted as the maximum absolute amplitude of the bandpass filtered strain data in Fig-174 ures 2e-f. For both frequency bands, the ACF first negative peaks exhibit similar lag-175 times for weak dynamic strains (e.g., less than $\sim 5 \times 10^{-10}$), but clear delays can be ob-176 served for larger dynamic peak strains. 177

Soil nonlinear behavior during ground motions delays the ACF first negative peak and can therefore be interpreted as a velocity reduction of the medium (Bonilla & Ben-Zion, 2020). Under the assumption that the changes in the medium are uniformly distributed, we can estimate the relative velocity changes (dv/v) of each ACF with respect to a reference ACF with the stretching method (Lobkis & Weaver, 2003; Sens-Schönfelder & Wegler, 2006) as

 $\tau = \frac{dt}{t} = -\frac{dv}{v},\tag{1}$

where τ , dt/t, and dv/v are the stretching coefficient, the relative time shift, and the rel-185 ative velocity change, respectively. For each channel and each frequency band, we first 186 compute a reference ACF by stacking the ACFs from the earthquake waveforms that gen-187 erated the ten weakest dynamic peak strains. For each frequency band, we then select 188 a time window that corresponds to 75% of the inverse of the lower cutoff frequency (e.g., 189 the first 0.375 s of the ACF for the 2-4 Hz frequency band) to focus on the first nega-190 tive peak of the ACFs. We then stretch and compress the selected window of the ref-191 erence ACF to find the stretching coefficient that maximizes the fit between the refer-192 ence and each ACF waveform, and therefore infer relative velocity changes. The stretch-193 ing is performed in two steps; we first use ten values uniformly distributed between -50 194 and 50% of stretching to find an initial guess of the stretching coefficient, and then re-195 fine the measurement by interpolating the stretched waveforms 500 times between the 196 neighboring values (similar to Viens et al., 2018). 197

198 **3 Results**

We show the relative velocity changes for all the frequency bands and earthquakes 199 for two ranges of channels in Figures 3a-b. While the soil nonlinear response can rapidly 200 evolve spatially, we display the combined results at 10 neighboring channels (i.e., over 201 51 m) between channels 5000–5010 and 7000–7010 for visibility. For dynamic peak strains 202 smaller than $\sim 5 \times 10^{-10}$, dv/v measurements are generally equal to zero for all the fre-203 quency bands at both locations, which indicates that there is no change in the medium. 204 However, clear dv/v drops can be observed in different frequency bands at the two lo-205 cations with increasing dynamic peak strains. For example, we primarily observe dv/v206 reductions between central frequencies (i.e., the central frequency of the bandpass filter; 207 15 Hz for the 10–20 Hz bandpass filter) of 15–24 Hz for channels 5000–5010 and between 208 12–28 Hz for channels 7000–7010. Moreover, we also note that the intensity of the dv/v209 changes varies, with larger changes observed at channels 5000–5010 compared to those 210 at channels 7000–7010. 211

²¹² Spatial variations of the relative velocity changes can also be tracked along the ca-²¹³ ble thanks to the high density of DAS channels. In Figures 3c-d, we show the relative ²¹⁴ velocity changes along the cable in two frequency bands (e.g., 5–10 and 10–20 Hz). Clear ²¹⁵ differences can be observed between the two frequency ranges. In the 5–10 Hz frequency ²¹⁶ band, almost no dv/v changes can be observe between channels 500 and 6900, even dur-²¹⁷ ing the strongest dynamic peak strains. However, we detect clear dv/v drops for dynamic



Figure 2. (a) ACFs computed from the 103 earthquakes bandpass filtered between 10 and 20 Hz at channels (a) 5000 and (b) 7000. The amplitude of the data is clipped for visibility. (c–d) Same as (a–b) for the data bandpass filtered between 15 and 30 Hz. In (a–d), the ACFs are sorted by increasing dynamic peak strain values, which are computed after bandpass filtering the strain waveforms in their respective frequency bands. (e–f) Dynamic peak strains after bandpass filtering the earthquake waveforms between 10–20 Hz and 15–30 Hz at channels 5000 and 7000, respectively.



Figure 3. (a) dv/v measurements at channels 5000–5010 for the 19 frequency bands and the 103 earthquakes. Dynamic peak strains are computed for each event and each station after bandpass filtering the strain data. The central frequency corresponds to the central frequency of the bandpass filter (e.g., 15 Hz for the 10–20 Hz bandpass filter). (b) Same as (a) at channels 7000–7010. (c) dv/v measurements from the ACFs computed from the 103 earthquakes bandpass filtered between 5 to 10 Hz between channels 500 and 9000 as a function of the dynamic peak strain. (d) Same as (c) for the 10–20 Hz frequency band. The dv/v color-scale shown in (b) is the same for all panels.

peak strains above 10^{-9} between channels 6900 and 8200. In the 10–20 Hz frequency band, almost no changes are found between channels 500–3500, but large dv/v drops are observed after channels 3500 for dynamic peak strains larger than ~ 10^{-9} .

To isolate and investigate the average sediment response during strong ground mo-221 tions, we also compute dv/v measurements between each reference ACF (i.e., the stack 222 of the ACFs computed during the 10 weakest dynamic peak strains) and a stack of the 223 ACFs computed during the five largest dynamic peak strains. dv/v changes between weak 224 and strong ground motion ACFs exhibit clear spatial and frequency variations (Figure 225 4a). Between channels 500 and 2000, we do not observe any large dv/v changes in any 226 frequency band. However, we observe spatial variations of the dv/v reductions at cen-227 tral frequencies above 15 Hz between channels 2000 and 9000. We also observe clear dv/v228 changes at frequencies below 15 Hz between channels 6300 and 9000. Such coherent spa-229 tial changes across frequency bands highlight the sensitivity of DAS ACFs to local site 230 conditions as well as their depth sensitivities. 231

The amplitude of dv/v reductions is expected to increase with increasing dynamic 232 peak strains. In Figure 4b-d, we show the dv/v measurements calculated between the 233 weak and strong ground motion ACFs as a function of the dynamic peak strains in three 234 frequency bands. We only show the results at 650 locations between channels 1700-8200 235 as we average the dv/v and dynamic peak strain values over 10 neighboring channels (e.g., 236 channels 1995-2005 for channel 2000). This step is perform to compare our results with 237 local site condition data from a velocity model of the region, as discussed below. In the 238 2-4 and 20-40 Hz, the largest dv/v drops correlate with the channels where the largest 239 dynamic peak strains are recorded, typically near the beginning of the cable. In the 15-240 30 Hz frequency band, however, dv/v changes are almost constant between dynamic peak 241 strains of 10^{-9} and 2×10^{-8} with an average value of -20%. This suggests that the non-242 linearity threshold in this frequency range is lower than dynamic peak strains of 10^{-9} . 243

In Figures 4e–g, we show the dv/v measurements at the same 650 locations along 244 the cable as a function of the average S-wave velocity in the first 30 m of the ground (V_{S30}) 245 obtained from the velocity model derived by Viens, Perton, et al. (2022). The Viens, Per-246 ton, et al. (2022) model was obtained by inverting Rayleigh-wave phase velocity disper-247 sion curves calculated by seismic interferometry using virtual sources located every 10 248 channels (e.g., 51 m). We compute V_{S30} from the 650 locations of the velocity model and 249 apply a smoothing of the V_{S30} values over 5 locations. We observe a decrease of the dv/v250 values with decreasing V_{S30} values in the 2-4 Hz and 20-40 Hz frequency band. However, 251 we do not observe any correlation between V_{S30} and the dv/v results in the 10-20 Hz fre-252 quency band. Nevertheless, the correlation between dv/v values and V_{S30} is relatively 253 weak, which suggests that V_{S30} is not the best parameter to characterize the nature of 254 soil nonlinearity as also shown by Bonilla et al. (2021). 255

256 4 Discussion

While larger dynamic peak strains generally correlate with larger dv/v drops, the 257 correlation with V_{S30} is weaker or even nonexistent. Three hypotheses can explain this 258 behavior. First, the velocity model, which was obtained from ASF cross-correlation func-259 tions spanning over 2 km (i.e., 400 channels), only captures a smoothed representation 260 of the shallow Earth structure. Therefore, the V_{S30} parameter extracted from the veloc-261 ity model may not fully capture the structural changes that can rapidly occur at shal-262 low depth. Secondly, we expect the ACFs to have different depth sensitivities based of 263 their frequency ranges. Therefore, a single parameter, namely V_{S30} , does not account 264 for such depth sensitivity variations. Thirdly, while an accurate value of V_{S30} can be use-265 ful for some geotechnical engineering purposes, it may not be the best parameter to ex-266 plain the intensity of the dv/v drops. For example, in a 30 m profile composed of a very 267 shallow and soft sediment layer overlaying a stiffer material, nonlinearity is expected to 268



Figure 4. (a) dv/v measurements computed between a reference ACF, which represents the soil linear response, and average ACF obtained from the earthquakes that generated the five largest peak strains, which captures the nonlinear behavior of sediments, at each channel and each frequency band. (b) Relative velocity changes as a functions of the filtered dynamic peak strain in the 2-4 Hz frequency bands. (c-d) Same as (b) for the 10-20 and 20-40 Hz frequency bands. (e) Relative velocity changes as a function of the average S-wave velocity within the first 30 m of the ground (V_{S30}) for the 2-4 Hz frequency bands. (f-g) Same as (g) for the 10-20 and 20-40 Hz frequency bands. In (b-g), the color-bar corresponds to the channel number.

only occur in the first layer. Therefore, a complete velocity profile of each site is likely to be more informative than a summarizing parameter such as V_{S30} (Bonilla et al., 2021), and future work should focus on refining the shallow structure of the velocity model.

While the recorded dynamic strains from the earthquakes considered in this study 272 are relatively weak, we observe significant relative velocity changes from the ACFs, which 273 indicate a nonlinear response of marine sediments. Due to the weak levels of shaking, 274 the soil nonlinear behavior only occurs during the passing of seismic waves, and no long-275 term effects, as those observed at land stations after the 2011 M_w 9.0 Tohoku-Oki earth-276 277 quake (Bonilla et al., 2021), could be detected. Nevertheless, the nonlinearity thresholds of strain levels obtained along the cable are consistent with those from laboratory ex-278 periments (Pasqualini et al., 2007; Remillieux et al., 2017; TenCate et al., 2004) and from 279 ACFs computed from a seismic array in California (Bonilla & Ben-Zion, 2020). To fur-280 ther validate our approach, we also compute ACFs from earthquake data recorded by 281 the horizontal accelerometer along the axis of the cable from the SOB3 station (Figure 282 1b). ACFs are computed for the same earthquakes as for the DAS dataset as well as 138 283 nearby M_w 5+ earthquakes which occurred between 2015 and 2021 (Figure S1). The ACFs from the SOB3 station exhibit similar features, with a nonlinearity threshold of the same 285 order as that obtained with the DAS data, which validates our approach. 286

The dynamic peak strain recorded by the DAS channels are in the direction of the 287 cable. However, DAS has different theoretical sensitivities depending on the type of seis-288 mic waves and their incidence angles (Martin et al., 2021). For example, DAS records 289 from earthquakes occurring at a 90 degree angle from the direction of the cable are ex-290 pected to exhibit less energy than events happening along the axis of the cable. More-291 over, the relatively long gauge length (e.g., 40 m) used to record the DAS data could po-292 tentially create notches in the frequency spectrum between 2 and 40 Hz (Dean et al., 2017). 293 Nevertheless, the steep subduction zone in the Tohoku region (Haves et al., 2018) com-294 bined to shallow and slow sediment layers interfere with the propagation of seismic waves, 295 which likely arrive with almost vertical angles to the cable. This translates into high apparent velocities of all earthquake wavefields recorded by the cable (e.g., Figure 1c), which 297 limits both the azimuthal and gauge length effects on the recorded data. We further con-298 firm this point by comparing the maximum amplitudes of DAS and SOB3 data during 299 the 103 earthquakes considered in this study in Figure S2. Both datasets exhibit sim-300 ilar maximum amplitudes with respect to azimuth angles to the earthquake epicenters, 301 which confirms that there is no noticeable azimuthal effect for the DAS data. 302

The largest ground motions during the two week experiment occurred during a M_V 303 5.6, which occurred 50 km east of the SOB3 station (Figure S3). Unfortunately, the data 304 recorded by the DAS system clipped and are therefore not usable in our analysis. The 305 clipping of the data is caused by rapid phase changes that occurred during strong ground 306 motions, which wraps the signal's phase. To reduce clipping effects and improve the dy-307 namic range of DAS experiments, one can increase the laser's pulse rate frequency, which 308 would limit the maximum distance that can be sampled by the DAS system, and/or re-309 duce the gauge length, which could result in a decrease of the SNR of the recorded wave-310 field (Mellors et al., 2022). Despite these drawbacks, a better tuning of the DAS param-311 eters could allow us to record strong ground motions that are likely to trigger stronger 312 nonlinear soil responses. 313

314 5 Conclusions

We analyzed the ground motions of 103 earthquakes recorded along a fiber-optic cable during a two-week DAS campaign offshore the Tohoku region, Japan. We computed ACFs of earthquake ground motions and detected relative velocity changes in the marine sediments surrounding the cable from the ACFs. Large drops of dv/v are observed along the cable and are typical of a nonlinear behavior of the medium. Moreover, the dv/v changes are frequency and spatially dependent, which highlights the sensitivity of DAS ACFs to the shallow Earth structure.

This study demonstrates that earthquakes recorded by DAS can be used to char-322 acterize the nonlinear behavior of soils during ground motions. This characterization could 323 be of critical importance for fiber-optic cables used for earthquake early warning pur-324 poses as soil nonlinearity impacts the amplitude and frequency content of the recorded 325 wavefield, and could bias rapid magnitude estimations. Nevertheless, the ACF approach 326 could easily be applied to other DAS datasets recorded in populated regions located on 327 top of sedimentary basins, such as Mexico City and Los Angeles, to better characterize 328 seismic hazard. 329

330 Acknowledgments

We thank Takeshi Akuhara for providing useful information about DAS measurements. 331 We thank *Fujitsu* for cooperating with the Earthquake Research Institute (ERI), The 332 University of Tokyo, for the DAS measurement campaigns. All the Figures are plotted 333 with Matplotlib (Hunter, 2007). Some of the data processing steps have been performed 334 using ObsPy (Beyreuther et al., 2010). Funding: This project was partly supported by 335 the discretionary budget of the director of ERI. The observations were carried out as part 336 of the Earthquake and Volcano Hazards Observation and Research Program by the Min-337 istry of Education, Culture, Sports, Science, and Technology of Japan. L.V. is supported 338 by NSF award EAR2022716. Z.J.S acknowledges support from the Air Force Research 339 Laboratory grant FA9453-21-2-0018. Competing interests: The authors declare that 340 they have no competing interests. Data availability: The codes developed to perform 341 the technical analysis and to reproduce most figures are available at https://doi.org/ 342 10.5281/zenodo.6672479. 343

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¹ Supporting Information for 'Nonlinear earthquake

² response of marine sediments with distributed

acoustic sensing'

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¹¹ Text S1.

¹² To validate the dv/v measurements obtained from the DAS cable, we also analyze the ¹³ earthquake ground motions recorded by the SOB3 station. We first downloaded the data ¹⁴ (100 Hz sampling rate) recorded by the horizontal accelerometer sensor in the axis of the ¹⁵ cable from 159 earthquakes, which occurred within the two week time period of the DAS ¹⁶ experiment. We are able to add 56 events to the 103 earthquakes analyzed in the main

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manuscript as the signal-to-noise ratio (SNR) of the accelerometer data is higher than 17 that of the DAS data. Moreover, the largest ground motions are recorded during the M_V 18 5.6 $(M_w 5.4)$ earthquake, which occurred close to the DAS cable, but is not usable in our 19 study as the DAS data clipped (see Text and Figure S3). The SOB3 accelerometion data 20 are processed similarly to the DAS data and we show the dv/v results for the 159 events 21 in Figure S1a. We compute a metric as the peak ground velocity (PGV) of the bandpass 22 filtered data divided by an estimate of V_{S30} (300 m/s) to obtain results with the unit of 23 strain that are comparable to the results from the DAS data. We observe that the largest 24 dv/v drops occur at frequencies of 14–20 Hz at strain values slightly above 10⁻⁹, which is 25 similar to the DAS data. 26

To confirm that the reliability of the analysis performed with relatively weak ground motions, we also investigate the response of the sediments to 138 earthquakes that generated stronger ground motions at the SOB3 station. The 138 earthquakes are M_w 5+ events and occurred within 300 km from the cable between 2015 and May 2021. The results are shown in Figure S1b and confirm the dv/v measurements obtained for the 159 earthquakes of the two-week time period of the DAS experiment. We observe strong dv/vdrops at most frequencies above 5 Hz during the strongest ground motions and different non-linearity thresholds based on the frequency band of the ACFs.

$_{35}$ Text S2.

The azimuthal and gauge length effects described in the main manuscript are likely to be weak for the DAS data offshore the Sanriku Coast. Shinohara, Yamada, Akuhara, Mochizuki, and Sakai (2022) converted strain data of a magnitude 3.0 earthquake recorded

at channel 10265 to acceleration data and compared the resulting waveforms to the acceleration waveform recorded by the SOB3 station. They found a high similarity between
the OBS and converted DAS waveforms for this event, which confirms the fidelity of the
wavefield recorded by DAS.

To further confirm this, we convert the strain data of the 103 earthquakes of our dataset 43 to acceleration by considering a constant apparent phase velocity of 2800 m/s. We then 44 select the peak ground acceleration of the converted DAS and OBS data and display 45 them as a function of the azimuth angle between each earthquake epicenter and the OBS 46 station in Figure S2. First, we observe similar PGAs between the two datasets, which 47 confirms that the reliability of the DAS data. Secondly, the azimuthal variations are also 48 very similar, which indicates that gauge length and seismic wave azimuth angles do not impact our results. Finally, some of the differences between the two datasets are likely 50 to be caused by the constant apparent velocity used to convert strain data to velocity. 51 Nevertheless, this first order correction allows us to validate our approach. 52

53 Text S3.

The November 29, 2019 M_w 5.4 earthquake, which occurred in the vicinity of the cable (Figure S3a), generated the strongest ground motions during the two-week experiment. The raw DAS data are recorded as 2-byte integers. Therefore, the amplitude of the raw data varies between -32768 to 32767, which corresponds to $-\pi$ and π of phase. To retrieve strain data, the cumulative sum of the raw data needs to be computed before scaling the amplitude with a constant (Equation 1 in Shinohara et al., 2022). Unfortunately, the amplitude of the raw DAS data clipped during this event (Figure S3b), which makes

⁶¹ impossible the conversion to strain data for this earthquake. The clipping of the data ⁶² only occurs during the strongest part of the ground motion.

References

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Figure S1. (a) dv/v measurements computed at the SOB3 station for 159 earthquakes which occurred during the two-week experiment. The x-axis has the unit of strain as the Peak Ground Velocity (PGV) is divided by an estimate of V_{S30} of 300 m/s. (b) dv/v measurements computed for the 159 earthquakes from November 2019 and 155 M_w 5+ events, which occurred within 300 km from the cable between January 2014 and May 2021.



Figure S2. Peak Ground Acceleration (PGA) from the 103 earthquakes recorded at the (orange) SOB3 station and at the (blue) DAS channel 10265 after conversion to acceleration as a function of the azimuth from the epicenter. The PGAs are obtained after filtering the waveforms between 1 and 5 Hz. Zero azimuth is north.



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Figure S3. (b) Bathymetric map near the seafloor cable observation system. The orange line denotes the buried section of the cable used in this study, and the locations of channels 2000, 4000, and 6000 are highlighted by red crosses. The white circles and orange inverted triangles show the positions of the accelerometers and tsunami-meters, respectively. The location of the M_w 5.4 event together with its focal mechanism is also shown. (b) Raw DAS data of the M_w 5.4 event recorded at channels 2000, 4000, and 6000.