

Methane's Solar Radiative Forcing

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Key Points:

- The net effect of solar radiative forcing by CH₄ is assessed with spatially and temporally resolved calculations.
- CH₄ solar radiative forcing enhances CH₄ longwave-only radiative forcing by 7%.
- CH₄ solar radiative forcing is significantly sensitive to CH₄ profile specification, surface albedo and solar mid-infrared CH₄ absorption.

Abstract

Methane (CH_4) has significant absorption bands at wavelengths of 1.7, 2.3, 3.3 and 7.6 μm which absorb incoming solar ‘shortwave’ (SW) radiation and cause a radiative forcing (RF) of climate. This letter presents a comprehensive quantification of CH_4 SW RF using a narrow-band radiative transfer model to calculate spatially and monthly resolved estimates of the direct CH_4 SW effect and its impact on longwave (LW) RF via stratospheric temperature change. This significantly extends previous assessments. Calculations include satellite measurements of CH_4 distribution, spectrally-varying surface albedo and absorption of solar mid-infrared radiation by methane’s 7.6 μm band. These factors substantially influence the SW effect; all-sky CH_4 SW tropopause instantaneous RF is significantly smaller than previous estimates at 0.002 W m^{-2} ($\pm 25\%$). In total, the CH_4 SW effect enhances CH_4 LW-only RF by 7% from 0.574 W m^{-2} to 0.613 W m^{-2} for a 750-1800 ppb perturbation.

Plain Language Summary

Radiative forcing is a widely used metric to quantitatively compare the strength of different drivers of climate change. It measures the perturbation to Earth’s energy balance, such as that due to an increase in greenhouse gas concentrations. Most greenhouse gas forcings arise due to absorption of thermal infrared (“longwave”) radiation, but greenhouse gases also absorb solar radiation. However, methane’s solar radiative forcing has received relatively little attention. We present the most detailed calculation of methane’s solar radiative forcing to date, when concentrations increase from pre-industrial to present day levels (750 – 1800 parts per billion). Our calculations are highly sensitive to the spatial and seasonal variation in surface albedo and methane’s vertical profile, for which we exploit recent satellite measurements; we also quantify the effect of absorption of solar radiation at mid-infrared wavelengths. These factors significantly influence the magnitude of the forcing. Consequently, it is significantly smaller than previous estimates at 0.002 W m^{-2} ($\pm 25\%$). However, methane’s absorption of solar radiation heats the stratosphere; this heating significantly modifies methane’s longwave radiative forcing. When this effect is included we find that methane’s absorption of solar radiation enhances methane’s longwave radiative forcing by 7% from 0.574 W m^{-2} to 0.613 W m^{-2} .

1 Introduction

CH₄ is the second largest contributor to greenhouse gas RF (Forster et al., 2021). It predominantly interacts with thermal-infrared LW irradiances, but also absorbs incoming solar SW radiation, mainly at near-infrared (NIR) wavelengths (0.7 – 4μm). CH₄ SW absorption is not explicitly treated in radiative transfer codes in many climate models and its radiative effect has been relatively unexplored.

Recent studies have demonstrated the significance of CH₄ SW RF. Etminan et al., (2016) presented calculations using the Oslo line-by-line (LBL) code (OLBL; Myhre et al., 2006) for two atmospheric profiles (tropical/extratropical) to estimate global-mean tropopause CH₄ SW instantaneous RF (IRF) and its effect on LW stratospheric temperature-adjusted RF (SARF; see Myhre et al., 2013a for definitions). They demonstrated two key components of CH₄ SW RF. Firstly, CH₄ SW absorption bands exert a positive all-sky tropopause IRF of 0.03 W m⁻² following a 750-1800 ppb perturbation. Secondly, the consequent stratospheric heating enhances the magnitude of CH₄ SARF. They reported that SW RF increases the Intergovernmental Panel on Climate Change (IPCC) Fifth Assessment Report (Myhre et al., 2013a) 1750-2011 CH₄ SARF by 15%. This revision underscored the importance of incorporating CH₄ SW absorption in RF calculations and emission metrics. They estimated a ±25% uncertainty due to a range of factors including: specification of surface albedo (α_s), computation of day-averaged forcings and their simple method of spatial and temporal averaging. Collins et al., (2018) later presented spatially-resolved calculations of CH₄ SW tropopause IRF using a narrow-band radiative transfer code and background climatologies from two Earth System Models (ESMs). For an 806-1760 ppb perturbation, they reported global annual-mean all-sky CH₄ SW tropopause IRFs of $\approx 0.026 \text{ W m}^{-2}$ and demonstrated its distribution to depend significantly on NIR α_s . Collins et al., (2018) did not quantify the effect of SW absorption on SARF.

This letter presents the most detailed quantification of the CH₄ SW effect to date. A narrow-band radiative transfer model is used to calculate seasonally and spatially-resolved RF (Section 2). These calculations utilise satellite measurements of CH₄ and spectrally-resolved α_s to derive a best estimate of global annual-mean CH₄ SW IRF and SARF (Section 3). Comparisons with Etminan et al., (2016) and Collins et al., (2018) highlight key sensitivities, namely the effect of CH₄ SW

absorption at solar mid-infrared wavelengths, the vertical representation of CH₄ mole fraction and the specification of spectrally-resolved α_s .

2 Models, methods and data

RF calculations use the two-stream radiative transfer model, SOCRATES-RF (Checa-Garcia et al., 2018), which utilises the radiance core of the UK Met Office Suite Of Community RAdiative Transfer codes (SOCRATES, Edwards & Slingo, 1996; Manners et al., 2015) to compute IRF and SARF using the fixed dynamical heating (FDH) method (e.g. Ramaswamy et al., 2001).

SW IRFs are calculated using a 260-band spectral file (sp_sw_260_jm2) covering wavelengths from 0.17-10 μm . This file has been verified against LBL codes for a variety of scenarios by Walters et al., (2019) and is independently validated here against the LBL Reference Forward Model (RFM; Dudhia, 2017). It shows good agreement across a range of conditions. sp_sw_260_jm2 uses solar spectral data from the Naval Research Laboratory Solar Spectral Irradiance model (NRLSSI; Lean et al., 2005) averaged across 2000-2011. LW forcings use the standard broadband spectral file configuration (sp_lw_ga7) with 9 bands across 3.34-10000 μm (Walters et al., 2019). This lower resolution configuration was used due to the computational expense of performing spatially and seasonally resolved FDH calculations.

Gaseous absorption uses the correlated-k distribution method based on HITRAN2012 spectroscopic data (Rothman et al., 2013). SW and LW spectral gaseous overlap are treated using the equivalent extinction with correlated scaling method and the random overlap assumption, respectively. H₂O continuum absorption is based on MT_CKD 2.5 (Mlawer et al., 2012) with modifications in some NIR window regions using experimental data (Ptashnik et al., 2011, 2012). Calculations are conducted under clear-sky and all-sky conditions. Temperature, humidity and cloud fields are taken from ERA-interim reanalysis (Dee et al., 2011). Vertical cloud overlap is treated assuming random overlap; ice and water clouds are treated separately within three atmospheric layers. Aerosols are absent from calculations.

SOCRATES-RF adopts a tropopause pressure based on the parametric equation of Checa-Garcia et al., (2018) which describes a hemispherically-symmetric tropopause (see Equation S1). This prescribes annual-mean tropopause height following the climatological tropopause representation

by Hansen et al., (2005). A six-point Gaussian quadrature method is used to calculate the diurnal variation of solar radiation. All calculations are performed using monthly-mean input fields.

CH₄ mole fractions are derived from Michelson Interferometer for Passive Atmospheric Sounding (MIPAS) climatology (Hegglin et al., 2021), comprising zonal-mean values averaged between 2005-2012 from 300 hPa to 0.1 hPa at a 5° spatial resolution (see Text S2). This results in well-mixed tropospheric CH₄ and latitudinally-dependent, vertically-varying CH₄ at pressures below 300 hPa (Figure S3).

α_s is derived from 33 one-nm wavelength bands between 0.355-2.314 μm from the Scanning Imaging Absorption spectroMeter for Atmospheric CHartographY (SCIAMACHY) Lambertian-equivalent reflectivity (LER) database (version 2.6; Tilstra et al., 2017) averaged between 2002-2012 at 0.5° x 0.5° spatial resolution. LER is defined as the reflectance of an isotropic surface required to match observed TOA reflectance in a clear-sky, aerosol-free Rayleigh scattering atmosphere (e.g., Kleipool et al., 2008). This dataset implicitly includes the effect of spatial and temporal variation in snow cover and sea-ice on spectral α_s (see Text S3). Given the lack of alternative observationally-based, spectrally-resolved global α_s datasets at $0.355 > \lambda > 2.314 \mu\text{m}$, this approach constructs a best estimate of spectral surface reflectance across 0.17-10 μm . LER data exhibit unrealistically low NIR α_s (≈ 0.001 -0.008) over sea-surfaces, therefore these datapoints have been replaced with spectrally-varying α_s calculated using a SOCRATES internal subroutine (see Text S3). Consequently, α_s fields amalgamate SCIAMACHY and SOCRATES sea-surface α_s (henceforth referred to as SCIAMACHY-SSA) re-gridded to a 5° x 5° spatial resolution.

3.1 Clear-sky and all-sky CH₄ SW IRF

Table 1 shows global annual-mean SW IRF at top-of-atmosphere (TOA), tropopause and surface for a CH₄ perturbation from 750 to 1800 ppb, the approximate pre-industrial and present-day values used in Etminan et al., (2016). Corresponding LW and net IRFs are also shown to demonstrate the impact of SW absorption on total IRF.

	Clear-sky IRF (W m^{-2})			All-sky IRF (W m^{-2})		
	SW	LW	Net	SW	LW	Net
TOA	0.040	0.650	0.690	0.082	0.553	0.635
Tropopause	-0.038	0.674	0.636	0.002	0.578	0.580
Surface	-0.235	0.306	0.071	-0.163	0.180	0.017

Table 1: Global annual-mean clear-sky and all-sky SW, LW and net IRF at TOA, tropopause and surface following a 750-1800 ppb CH_4 perturbation. Derived using MIPAS CH_4 fields and SCIAMACHY-SSA α_s .

For clear-skies, SW IRF is positive at TOA (0.040 W m^{-2}) driven by enhanced absorption of surface-reflected radiation between the surface and TOA. At the tropopause and surface, SW IRF is negative due to a reduction in downward SW irradiance with enhanced CH_4 absorption in the stratosphere and troposphere. The surface SW IRF of -0.235 W m^{-2} counterbalances the surface LW IRF of 0.306 W m^{-2} , giving a net forcing of 0.071 W m^{-2} . Under all-sky conditions this offset is enhanced resulting in a net forcing of just 0.017 W m^{-2} . This demonstrates the importance of CH_4 SW absorption on surface energy budgets.

The all-sky TOA SW IRF increases from 0.040 W m^{-2} to 0.082 W m^{-2} due to absorption of increased reflected SW radiation. This mechanism is also responsible for the switch in sign of the tropopause SW IRF from -0.038 W m^{-2} to 0.002 W m^{-2} . This all-sky tropopause SW IRF is significantly smaller than Etminan et al., (2016) and Collins et al., (2018) (0.03 and $\approx 0.026 \text{ W m}^{-2}$, respectively). Section 3.2 addresses this disparity. Etminan et al., (2016) were the first to highlight the importance of cloud radiative effect (CRE) in determining the sign and magnitude of tropopause SW IRF; CRE increased their clear-sky SW tropopause IRF of -0.045 W m^{-2} by 0.075 W m^{-2} to a value of 0.03 W m^{-2} . The results presented here confirm the importance of CRE, albeit with a less pronounced increase (0.04 W m^{-2}).

Figure 1 shows the spatial distribution of annual-mean all-sky SW tropopause IRF (upper left) and the corresponding CRE (lower left). The distribution of SW IRF is comparable to Collins et al., (2018) exhibiting regions of maxima over the Arabian Peninsula and the Sahara (exceeding 0.20 W m^{-2}). Strong spatial gradients in the forcing are driven by spectrally-resolved NIR α_s for desert surfaces. As explained by Collins et al., (2018), such surfaces increase the reflection of

downwelling SW irradiance, enhancing the magnitude of tropospheric CH₄ SW absorption causing a more positive IRF. CRE is predominantly positive with a localised maximum of 0.1 W m⁻² off the west coast of South America associated with oceanic stratocumulus cloud decks. Regions with negative CRE dampen the influence of α_s on IRF magnitude, particularly across northern Africa due to high NIR α_s of desert surfaces.

Notably, Collins et al., (2018) calculate positive (rather than negative) clear-sky tropopause SW IRFs of 0.008 W m⁻² and 0.011 W m⁻² using their two ESM climatologies. They attribute positive clear-sky forcings to the use of realistic surface conditions in their spatially-resolved simulations. Consequently, CRE enhances their IRF by just 0.014 W m⁻² and 0.018 W m⁻². Collins et al., (2018) report a global-mean all-sky SW tropopause IRF of approximately 0.026 W m⁻², similar to the 0.03 W m⁻² reported by Etminan et al., (2016). However, disagreement in the sign of clear-sky IRF and the size of CRE indicate that this agreement is coincidental.

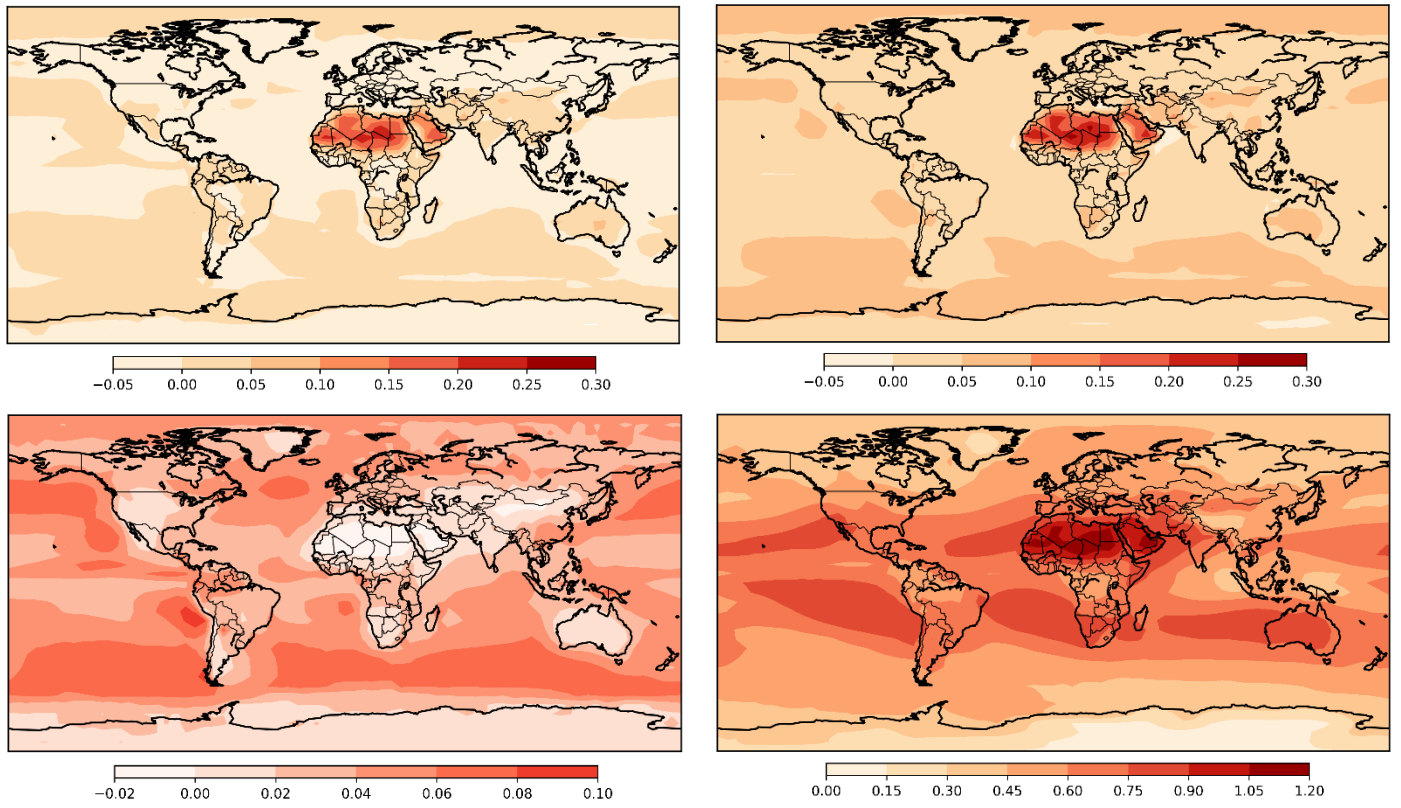


Figure 1: All annual-mean (W m^{-2}). Upper left: All-sky tropopause SW IRF. Lower left: Cloud radiative effect on tropopause SW IRF calculated as the difference between all-sky and clear-sky conditions. Upper right: Total all-sky SW RF at the tropopause derived as the difference between net SARF (i.e., LW SARF (including SW in FDH calculation) + SW IRF) and LW-only SARF. Lower right: All-sky net SARF at the tropopause. Derived using conditions stated in Table 1 caption. Global-mean forcings are given in Table 2.

171

172 3.2 Explanation of reduction in size of all-sky CH_4 SW tropopause IRF

173 The all-sky SW tropopause IRF of 0.002 W m^{-2} is significantly smaller than Etminan et al., (2016)
 174 and Collins et al., (2018) due to a combination of factors that alter the relative importance of
 175 downward and upward tropopause forcing components. As explained by Etminan et al., (2016)
 176 (see also Shine et al., 2022), the tropopause SW IRF is a residual of reduced downward irradiance
 177 (due to absorption in the stratosphere, which constitutes a negative IRF) and reduced upward
 178 irradiance (due to absorption of surface-troposphere reflected irradiance in the troposphere, which

constitutes a positive IRF). The following sub-sections detail these factors using a range of sensitivity tests.

3.2.1 The representation of CH₄ vertical profiles

CH₄ decreases with height above the tropopause due to chemical destruction. A realistic vertical distribution in forcing calculations is necessary because stratospheric SW absorption contributes significantly to the sign and magnitude of all-sky SW tropopause IRF (Figure S7).

Idealised sensitivity tests calculated using the RFM (not shown) and SOCRATES-RF demonstrate the need for a realistic description of the spatial and temporal variability in stratospheric CH₄. SOCRATES-RF tests were conducted using the Section 2 methodology, albeit with idealised zonal-mean CH₄ mole fraction scale-heights from the CAM3.0 model (Collins et al., 2004), which were used by Collins et al., (2018) (based on early 2-dimensional chemical model simulations – see Text S4). These do not represent hemispherical and seasonal differences in stratospheric CH₄.

Using ‘CAM3.0’ scale-heights results in an all-sky SW tropopause IRF of 0.009 W m⁻². The equivalent calculation using MIPAS mole fractions gives an IRF of 0.002 W m⁻² (Table 1). This difference originates from the size of the downward SW forcing component. With ‘MIPAS’, more CH₄ is present in the stratosphere compared to CAM3.0. MIPAS CH₄ mole fractions vary monthly and exhibit a variable rate of decrease with height at each latitude (see Figure S6) unlike CAM3.0’s smooth exponential decrease (see Figure S5). MIPAS values are at least 2.5 times larger than CAM3.0 values at 300 hPa at 0°. This disparity increases poleward to over 10 times larger at 300 hPa at 60° and 90° N/S. MIPAS values increase stratospheric CH₄ SW absorption at both pre-industrial and present-day levels, causing a more negative downward IRF at the tropopause. Hence, the positive upward forcing component (due to CRE) makes the MIPAS forcing only weakly positive, but decisively positive for CAM3.0.

3.2.2 Impact of CH₄ absorption between 5–10 μm

Solar spectral irradiance extends beyond the NIR; around 5.4 W m⁻² occurs in the 5–10 μm region. Although only about 0.5% of total solar irradiance, this spectral region contains the strong 7.6 μm CH₄ absorption band (primarily responsible for CH₄ LW RF). In exploratory calculations using

CAM3.0 CH₄ profiles and spectrally-constant α_s fields from ERA-interim reanalysis data (Dee et al., 2011), the 5-10 μm all-sky CH₄ SW tropopause IRF is found to be -0.004 W m^{-2} compared to the 0.17-10 μm value of 0.016 W m^{-2} . The negative forcing demonstrates that this spectral region mainly impacts stratospheric absorption. Neglecting this contribution would overestimate the all-sky SW tropopause IRF by 25%, for this configuration; hence the 5-10 μm region is included in calculations presented here. Both Etminan et al., (2016) and Collins et al., (2018) use an upper wavelength limit of 5 μm . The spectral nature of CH₄ SW IRF was considered in detail in Etminan et al., (2016); Figure S7 extends this to include methane's 7.6 μm band.

3.2.3 Sensitivity to spectrally-varying albedo

The spatial and spectral variation in α_s depends on factors including surface type, season and solar zenith angle. Tests conducted using the Reading LBL model RFMDISORT (see Table 7 of Collins et al., 2006) and SOCRATES-RF demonstrate the importance of spectrally-varying α_s in controlling the amount of surface-reflected SW radiation absorbed by CH₄. The use of spectrally-constant ERA-interim α_s in forcing calculations overestimates the reflected SW radiation absorbed by CH₄ predominantly in the polar regions, yielding a global-mean all-sky SW tropopause IRF of 0.016 W m^{-2} compared to 0.009 W m^{-2} using SCIAMACHY-SSA α_s (both calculated using SOCRATES-RF with CAM3.0 CH₄; see Figure S8).

Both Etminan et al., (2016) and Collins et al., (2018) incorporate spectrally-varying α_s . Etminan et al., (2016) use broadband values taken from OsloCTM2 simulations (Myhre et al., 2013b). Collins et al., (2018) use land and snow α_s “obtained from climate model grid spectral seven-band Moderate Resolution Imaging Spectroradiometer bi-directional reflectance distribution functions” covering the 0.620-2.155 μm wavelength range. Here, the use of SCIAMACHY α_s data (with 33 bands across 0.355-2.314 μm) provides a better resolved description of land and sea-ice α_s . The representation of ocean α_s here (Section 2) and in Collins et al., (2018) is of similar complexity, based on Cox and Munk (1954).

Summarising, the combined effect of the above sensitivities results in a significantly smaller all-sky CH₄ SW tropopause IRF compared to earlier estimates. MIPAS fields significantly improve the representation of the spatial and temporal variation in stratospheric CH₄, SCIAMACHY-SSA

provides a more highly spectrally-resolved description of α_s and SOCRATES includes solar CH₄ absorption between 5-10 μm .

3.3 Impact of stratospheric CH₄ SW absorption on SARF

The total SW effect is quantified by comparing the net SARF, calculated due to both SW and LW absorption, against the SARF calculated due to LW absorption alone. Table 2 shows all-sky SW IRF, LW SARF (where SW absorption is *included* in the FDH calculation), net SARF (SW IRF plus LW SARF) and LW-only SARF (where SW absorption is *excluded* from the FDH calculation). The SW forcing is only minimally affected by the temperature change; hence SW IRF and SW SARF are identical.

	All-sky CH ₄ RF (W m ⁻²)			
	SW IRF	LW SARF	Net SARF	LW-only SARF
TOA	0.082	0.530	0.613	0.574
Tropopause	0.002	0.611	0.613	0.574
Surface	-0.163	0.181	0.018	0.179

Table 2: Global annual-mean all-sky TOA, tropopause and surface SW IRF, LW SARF (including SW heating rates in the FDH calculation), net SARF (SW IRF + LW SARF) and LW-only SARF (excluding SW heating rates in the FDH calculation). Derived using conditions stated in Table 1 caption.

The tropopause SW IRF of 0.002 W m⁻² is only a small fraction (0.3%) of net SARF (0.613 W m⁻²). However, when SW absorption is included in the calculation of stratospheric temperature adjustment, net SARF increases by 7% to 0.613 W m⁻² from the LW-only SARF of 0.574 W m⁻². This difference (0.04 W m⁻²) is 20 times larger than the tropopause SW IRF itself, emphasising that the contribution of SW absorption to net SARF should not be assessed by considering the impact on SW fluxes alone.

Figure 2 compares the stratospheric temperature difference (ΔT) calculated using FDH for LW-only and LW plus SW heating rates. In the LW-only case (upper panel), increased CH₄ leads to a cooling across most of the stratosphere, exceeding 0.2 K in the upper stratosphere. A small

warming is evident above the tropical tropopause; however, this does not counterbalance the impact of the cooling on downward LW irradiance at the tropopause, which is found to slightly decrease the LW forcing from 0.578 W m^{-2} (Table 1) to 0.574 W m^{-2} (Table 2). When SW absorption is included (lower panel), the cooling is reduced across the entire stratosphere and reverses sign over much of the lower stratosphere, giving a strong warming reaching a maximum ΔT of 0.5 K above the tropical tropopause. As a result, the LW IRF of 0.578 W m^{-2} (Table 1) is increased to a SARF value of 0.611 W m^{-2} (Table 2). This increases further to 0.613 W m^{-2} when the SW IRF is included.

The temperature change is consistent with the idealised experiments of Modak et al., (2018) in which SW absorption causes stratospheric warming of around 0.8 K, peaking at 100 hPa for a 10x increase in CH_4 . The peak 0.5 K warming calculated here is equivalent to a warming of around 0.05 decade^{-1} , if it occurred at a steady rate.

Figure 1 (upper right) presents the first spatially-resolved calculation of the total CH_4 SW radiative effect (to our knowledge) including both direct all-sky tropopause SW IRF and the effect of SW absorption on stratospheric temperatures. The inclusion of SW absorption increases net SARF (Figure 1, lower right) largely across the entire globe. Figure 1 (upper right) shows the enhancement is largest across northern Africa and the Arabian Peninsula, but also in regions with extensive low cloud, such as the southern oceans.

The total SW effect of 7% is smaller than the 15% reported by Etminan et al., (2016). The difference is largely explained by the substantial difference in the magnitude of the SW tropopause IRF (0.002 W m^{-2} compared to 0.03 W m^{-2}), but also by the magnitude of stratospheric temperature-adjustment. Etminan et al., (2016) find this process has a slightly more negative effect on their estimate of CH_4 LW IRF, decreasing it by 2% from 0.516 W m^{-2} to 0.504 W m^{-2} . Considering the total SW effect is measured relative to adjusted LW-only SARF, this yields a larger percentage increase from their LW-only SARF (0.504 W m^{-2}) to their net SARF (0.582 W m^{-2}). In addition to the sensitivities addressed in Section 3.2, the total CH_4 SW effect also depends on state variables (and their spatial and temporal variation) that affect CH_4 LW forcing and FDH adjustment, including temperature, humidity and clouds. Text S5 discusses uncertainty associated with the calculation of SW tropopause IRF and LW SARF. Using the root-sum-square method, a total uncertainty of $\pm 12\%$ is attributed to LW SARF and a total uncertainty of $\pm 25\%$ is tentatively

288 attributed to SW tropopause IRF (see Table S2). Given the highly sensitive nature of CH₄ SW
289 tropopause IRF, we believe that Etminan et al.,’s $\pm 25\%$ estimate of uncertainty was too low, hence
290 despite the improved methodology here, we adopt the same overall uncertainty. In total, given that
291 the CH₄ SW effect is about 6% of the net SARF, this yields a total uncertainty of around $\pm 13\%$ in
292 net CH₄ SARF.

293 An alternative perspective on CH₄ SW RF can be gained by examining TOA forcings which is
294 more analogous to effective radiative forcing (ERF) estimates derived from ESMs (Shine et al.,
295 2022). For SARF, the net TOA and tropopause forcings are identical (Table 2) as is the total impact
296 of the SW forcing (i.e., 7%). However, unlike the tropopause IRF, the SW TOA IRF is not a
297 residual of changes in upward and downward irradiance and a more confident assessment is
298 possible. TOA SW IRF is 0.082 W m^{-2} (or 13% of TOA net SARF). Stratospheric heating by SW
299 absorption leads to an increase in TOA LW irradiance which drives a reduction in TOA LW SARF
300 from 0.574 to 0.530 W m^{-2} (Table 2).

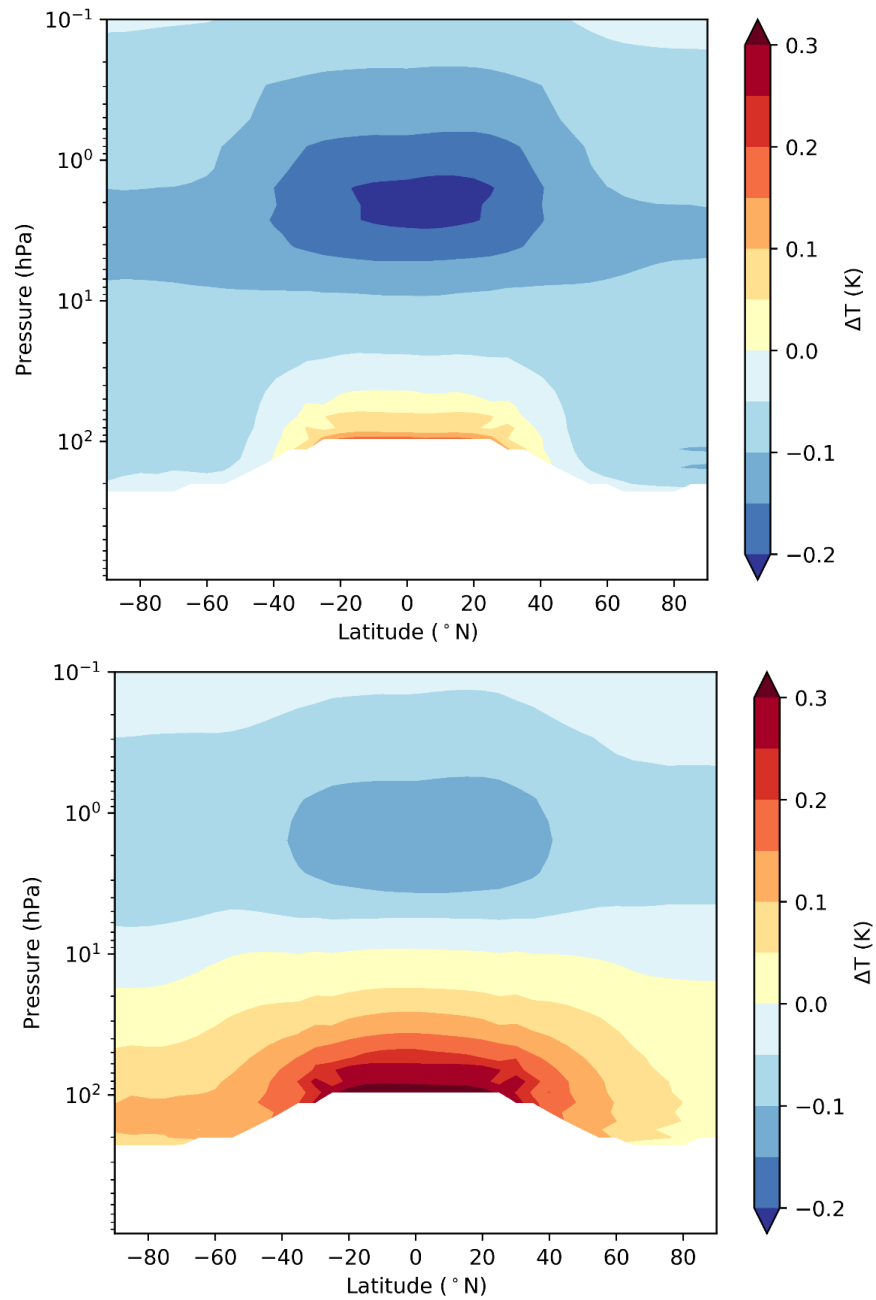


Figure 2: Annual zonal-mean ΔT (in K) following stratospheric temperature adjustment calculated using FDH for LW-only (upper panel) and LW plus SW heating rates (lower panel). Small oscillations in ΔT occur near the tropopause when SW heating rates are found in stratospheric temperature adjustment, therefore data in the lower plot have been smoothed. Derived using conditions stated in Table 1 caption.

4. Discussion

Accurate quantification of CH₄ SW RF is dependent on spatially-resolved specifications of CH₄ vertical profiles, CRE and spectrally-resolved α_s to capture their collective heterogeneous effect on SW absorption. It is also dependent on inclusion of solar mid-infrared absorption. Further studies that incorporate these factors are needed, particularly to constrain the sign and magnitude of clear-sky tropopause IRF.

A multi-model intercomparison would help quantify the accuracy of CH₄ SW absorption in ESM radiation schemes; Text S5 presents evidence that CH₄ SW IRF is significantly underestimated in some such schemes. As emphasised here, the accuracy of CH₄ SW ERF also depends on the way ESMs represent CH₄ vertical distribution and spectrally-varying α_s . This is particularly important with respect to the calculation of individual adjustments. Adjustments for a tripling of CH₄ have been calculated by Smith et al., (2018) in 11 ESMs using radiative kernels and the RFMIP (Pincus et al., 2016) fixed-SST approach. Four of these ESMs included CH₄ SW absorption bands. Although this is a small sample of models, the inclusion or omission of these bands appears to dictate the magnitude of H₂O adjustment and whether stratospheric temperature and cloud adjustments have a negative or positive effect on CH₄ forcing; this could only be confidently established if those models were run without these bands. Models that include CH₄ SW absorption simulate a negative stratospheric temperature adjustment to a CH₄ perturbation, consistent with the sign derived here from a TOA perspective. Although uncertainties are large in Smith et al., (2018), there is an indication that the total CH₄ adjustment is negative when SW absorption is included and positive when it is not.

IPCC AR6 (Forster et al., 2021) calculate CH₄ SARF using the simplified expression of Meinshausen et al., (2020) which incorporates the total CH₄ SW effect of 15% estimated by Etminan et al., (2016). The CH₄ ERF is estimated using the tropospheric adjustments of Smith et al., (2018) from the 4 ESMs that include CH₄ SW absorption. Overall, these adjustments reduce the ERF relative to TOA IRF by -14% ($\pm 15\%$; Forster et al., 2021). According to the total CH₄ SW effect of 7% presented here, IPCC AR6 CH₄ SARF and CH₄ ERF are both slightly overestimated. However, as stated above, the magnitude of tropospheric CH₄ adjustments are

uncertain; further assessment is needed by ESMs that include validated treatment of CH₄ SW absorption.

5. Conclusions

Accurate quantification of CH₄ SW RF depends on realistic spatial and temporal representations of CH₄ vertical profiles, CRE, spectrally-resolved α_s and inclusion of solar mid-infrared CH₄ absorption. At the surface, CH₄ SW IRF counteracts around 90% of CH₄ LW IRF resulting in a residual net IRF of 0.017 W m⁻². At the tropopause, CH₄ SW IRF is highly sensitive to the magnitude of the upward and downward forcing components and is significantly smaller than previous estimates at 0.002 W m⁻² ($\pm 25\%$). However, the total CH₄ SW effect can only be fully quantified when the impact of stratospheric CH₄ SW absorption is included in the calculation of CH₄ SARF. Consequently, inclusion of SW absorption enhances the LW-only SARF by 7%, from 0.574 W m⁻² to 0.613 W m⁻² for a 750 to 1800 ppb perturbation. An alternative TOA perspective demonstrates that the SW IRF is 13% of TOA net SARF; the total effect of SW absorption is still 7% (i.e., in agreement with the tropopause view) as a result of stratospheric temperature adjustment. This indicates that the perception of the role of SW RF (particularly the direct impact of SW IRF and the impact of stratospheric temperature adjustment) depends on the perspective taken. In the ERF framework, the TOA view is becoming more prevalent.

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Geophysical Research Letters

Supporting Information for

Methane's Solar Radiative Forcing

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Text S1.

SOCRATES-RF models the tropopause pressure (p_{trop}) based on the following parametric equation of Checa-Garcia et al., (2018), which describes a symmetrical tropopause between both hemispheres as:

$$p_{trop} = \delta \tanh[\beta^{-1}(|\phi| - \alpha)] + \gamma \quad (S1)$$

where ϕ is latitude, γ is the average of p_{trop} at 90° and 0° , δ defines the decrease in pressure between 90° and 0° , α defines the latitude at the centre of the transition from the 90° and 0° tropopause level (i.e. where $\tanh[\beta^{-1}(|\phi| - \alpha)] = 0$) and β defines the characteristic latitudinal width of the transition between the 90° and 0° tropopause. These parameters describe a tropopause without any seasonal dependence following the estimation of a climatological tropopause by Hansen et al., (2005). All SOCRATES-RF calculations use the combination of parameter values given in Table S1 to define p_{trop} . This yields an equatorial (0°) tropopause at 100 hPa and a polar (90°) tropopause at 230 hPa, resulting in a γ of 165 hPa (i.e. the average of 100 hPa and 230 hPa) and a δ of 65 hPa (which is defined as the difference between the tropopause at 90° and 0° divided by two i.e. $\delta = (230 \text{ hPa} - 100 \text{ hPa})/2$). At the latitude specified by α (i.e. 35°), p_{trop} equals γ ; at latitudes either side of α , the tropopause varies according to Equation S1. Figure S1 shows the variation in p_{trop} with latitude as based on the parameters defined in Table S1.

δ	β	γ	α
65 hPa	10	165 hPa	35°

Table S1: Climatological tropopause parameters used to determine tropopause pressure in all SOCRATES-RF calculations. See Checa-Garcia et al., (2018).

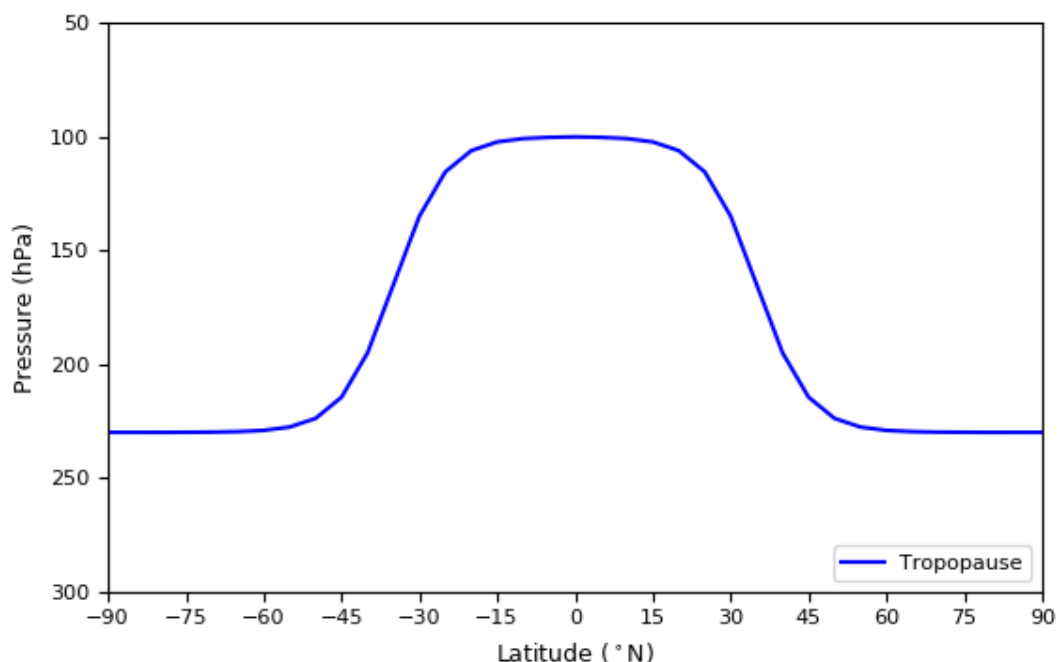


Figure S1: Variation in tropopause pressure with latitude as defined by Equation S1 and the parameters specified in Table S1. All SOCRATES-RF tropopause RFs are calculated using this definition of tropopause height.

Text S2

CH₄ mole fractions are derived from Michelson Interferometer for Passive Atmospheric Sounding (MIPAS) CH₄ climatology (Hegglin et al., 2021) created by the SPARC Data Initiative (SPARC, 2017). This dataset is comprised of monthly zonal-mean CH₄ mole fractions on a common latitude-pressure grid covering the upper troposphere to lower mesosphere (from 300 hPa to 0.1 hPa) with a latitudinal resolution of 5° from 87.5° north to 87.5° south (Figure S2). This climatology consists exclusively of measurements taken during the second MIPAS observation period (henceforth MIPAS-2) where the instrument operated with a reduced spectral resolution (but a higher vertical resolution) compared to its earlier 2002 – 2004 measurement mode (Plieninger et al., 2016). Hegglin et al., (2021) report that MIPAS-2 captures the extent and timing of high-latitude minima in CH₄ during the summer months (caused by photochemical reactions) and the semi-annual oscillation in CH₄ maxima in the tropical upper-stratosphere (caused by the upwelling of less depleted CH₄ air by the equatorial semi-annual oscillation; e.g. Choi and Holton, 1991).

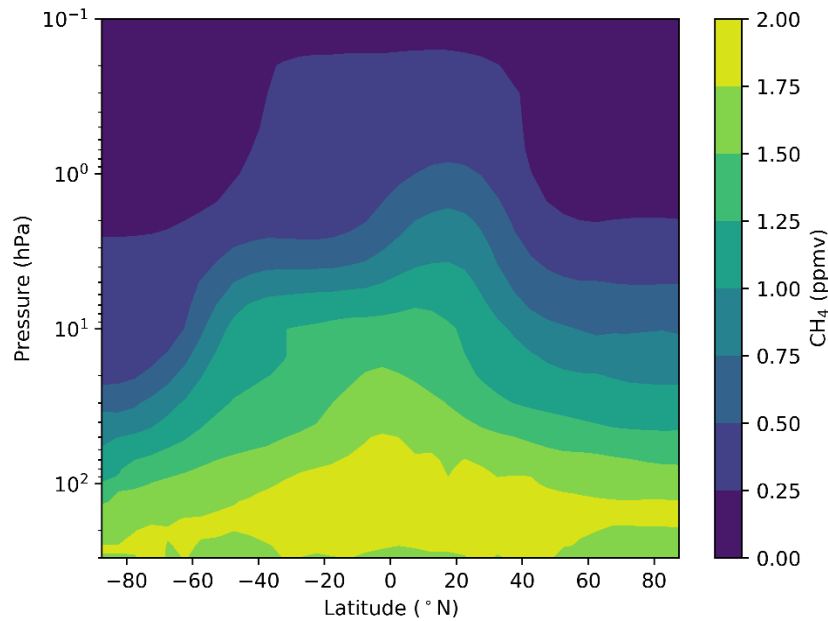


Figure S2: Zonal-mean, monthly-mean CH₄ mole fractions (in ppm) for July averaged over 2005 – 2012 showing maximum values high in the tropical upper troposphere.

Exploration of the MIPAS CH₄ climatology reveals that mole fractions in regions of the tropical upper troposphere are higher than those at higher pressure levels close to 300 hPa (where the highest mole fractions are typically expected to be found). Figure S2 displays this feature in zonal-mean, monthly-mean CH₄ for July. As shown, CH₄ maxima with values greater than 1.75 ppm (yellow filled contour) are evident across the equatorial region up to around 50 hPa. SPARC (2017) note that MIPAS-2 measurements exhibit a maximum in CH₄ in upper troposphere lower stratosphere (UTLS) region that is not present in other vertically-resolved SPARC satellite climatologies, such as the Atmospheric Chemistry Experiment-Fourier Transform Spectrometer (ACE-FTS) and the Halogen Occultation Experiment (HALOE) datasets.

This high bias is a known feature reported by several studies (e.g. von Clarmann et al., 2009; Laeng et al., 2015; Errera et al., 2016) and may be linked to instrument bias or the resulting impact of deep convection associated with the Asian monsoon region, which uplifts and traps high CH₄ mole fractions in the upper troposphere (Ricaud et al., 2014). Retrievals from the Atmospheric Infrared Sounder (AIRS) across 2003 – 2007 confirm the presence of strong CH₄ enhancement during the monsoon season in the middle to upper troposphere during July, August and September over South Asia (Xiong et al., 2009). The noted maxima in MIPAS-2 data may therefore occur as a result of such dynamical processes, which is missing from both ACE-FTS and HALOE due to sparser sampling of the

tropics by their solar occultation retrieval technique. However, it is unknown if such an enhancement in CH₄ would have been present in the pre-industrial era. Given that MIPAS-2 data are used to create monthly-mean pre-industrial fields of CH₄ in double-call RF calculations, this feature has been removed as follows:

1. At each latitude, pre-industrial CH₄ profiles are constructed using the fall-off rate of MIPAS CH₄ mole fraction, with a value of 750 ppb at 300 hPa.
2. All data points higher than 750 ppb at pressures lower than 300 hPa are masked and set to a value of 750 ppb.

This ensures that CH₄ mole fractions do not exceed 750 ppb at altitudes higher than 300 hPa. Figure S3 shows the resulting zonal-mean, monthly-mean pre-industrial CH₄ mole fractions for July.

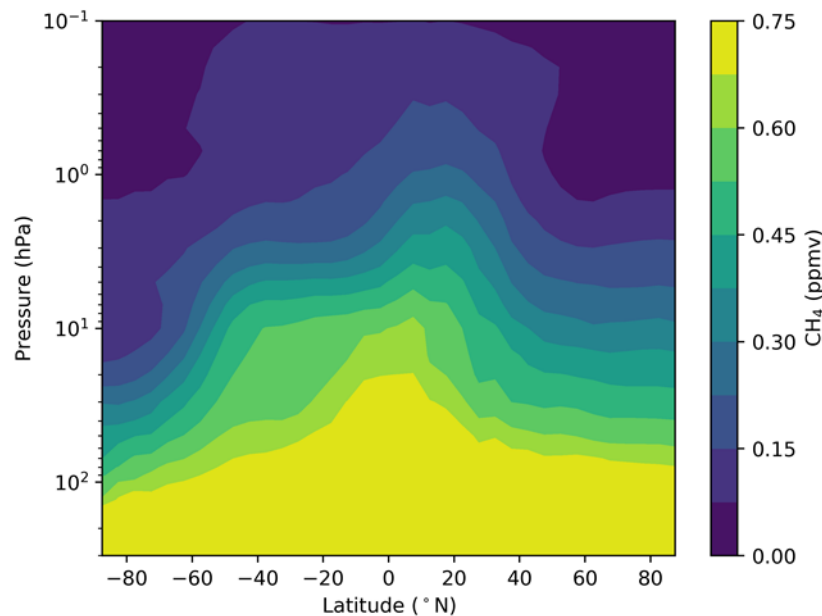


Figure S3: Pre-industrial zonal-mean, monthly-mean CH₄ for July (averaged over 2005 – 2012). Constructed using the fall-off rate of MIPAS CH₄ mole fraction with a value of 0.75 ppm at 300 hPa. The localised maxima in the tropical upper-troposphere in MIPAS CH₄ (see Text S2) has been removed by masking all data points higher than 0.75 ppm at pressures lower than 300 hPa to a value of 0.75 ppm.

Following this, MIPAS CH₄ fields are incorporated into SOCRATES-RF simulations as follows:

1. At each latitude, pre-industrial CH₄ profiles are constructed using the fall-off rate of MIPAS CH₄ mole fractions, with a value of 750 ppb at 300 hPa. This assumes that processes determining the fall off rate (and its monthly and latitudinal variation) are unchanged relative to the present day.
2. Data are regridded to a 5° x 5° spatial resolution to match the latitudinal and longitudinal coordinates of monthly-mean ERA-interim (Dee et al., 2011) climatological fields.
3. Data are interpolated to monthly-mean ERA-interim pressure levels, resulting in constant CH₄ throughout the troposphere and vertically-varying CH₄ in the stratosphere.

Text S3

Version 2.6 of the SCIAMACHY LER database (Tilstra et al., 2017) provides a climatology of Earth's surface reflectivity across 34 one-nm wavelength bands between 0.328-2.314 μm for each month of the year averaged between 2002-2012 at a 0.5° x 0.5° spatial resolution (https://www.temis.nl/surface/albedo/scia_ler.php last accessed 27/10/21). In this paper 33 bands between 0.355-2.314 μm are used to provide lower boundary conditions in SOCRATES-RF calculations. Data at 0.328 μm are excluded because surface reflectivity at this wavelength contributes a small amount to TOA reflectance (see Tilstra et al., 2017). Tilstra et al., (2017) examine the spatial, temporal and spectral behaviour of SCIAMACHY α_s and report good agreement with databases derived from the Total Ozone Mapping Spectrometer (TOMS), Global Ozone Monitoring Experiment (GOME-1 and GOME-2) and the Ozone Monitoring Instrument (OMI). Crucially, SCIAMACHY data cover a wider wavelength range than alternative datasets and thereby provide a measure of surface reflectivity at wavelengths important for CH₄ SW absorption (i.e., at 1.7 μm and 2.3 μm).

In this paper all SCIAMACHY α_s data are obtained following the MODE-LER approach - a new statistical method presented by Tilstra et al., (2017) which better determines the retrieval of surface LER from cloud-free satellite scenes. Initial handling of the database reveals α_s values above one occur multiple times in each month in several different wavebands (mostly in the VIS), predominantly across the high-latitude and polar regions. It is not known why this occurs and inspection of the accompanying metadata did not

provide an explanation. To handle this, all α_s values above one were replaced with a value of 0.95 (Huang et al., 2018).

Figure S4 presents a schematic of the method used to interpolate and extrapolate SCIAMACHY α_s to the upper wavelength limit of each band in sp_sw_260_jm2.

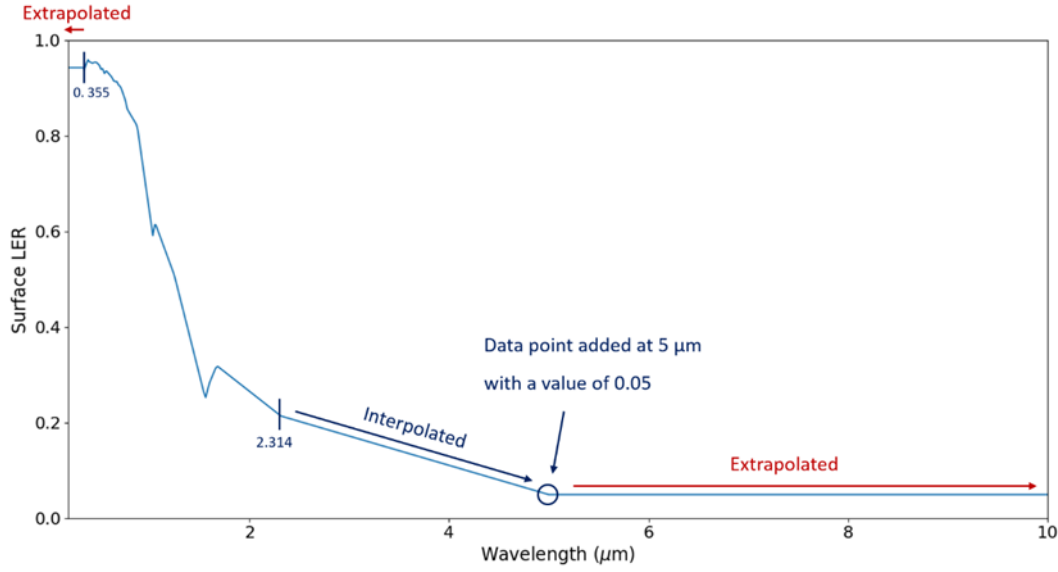


Figure S4: Schematic of the method used to assimilate SCIAMACHY spectral α_s data (Tilstra et al., 2017) into SOCRATES-RF's SW spectral file, sp_sw_260_jm2. This example demonstrates the process at a snow-covered grid-point located at 70°N, 155°W, Alaska for the month of May.

Data are linearly interpolated to the upper wavelength limit of each sp_sw_260_jm2 band between 0.355 and 2.314 μm . An additional α_s datapoint is constructed at 5 μm with a value of 0.05 to extend spectral reflectance to wavelengths greater than 2.314 μm . Huang et al., (2018) show that surface emissivity (ϵ) at wavelengths above 5 μm are around 0.95 for a range of land surface types such as water, ice, snow and desert. A value of 0.05 is therefore adopted assuming:

$$\alpha = 1 - \epsilon \quad (\text{S2})$$

To extend spectral reflectance to the upper and lower bands of sp_sw_260_jm2, data are linearly extrapolated from 5 to 10 μm and from 0.355 to 0.17 μm . Given the lack of monthly, global albedo datasets at wavelengths greater than 2.314 μm and less than 0.328 μm , this approach constructs a best estimate of spectral surface reflectance across

the 0.17 to 10 μm range. It is acknowledged that this method could be significantly improved if such global datasets became available.

Further handling of the database revealed low spectral surface albedo values (around 0.001 to 0.008) across the NIR region over the global oceans. As stated in the main text, these values are judged to inadequately represent the spectral reflectance of sea-surfaces which depend on a variety of complex factors, such as wind-induced variations in surface roughness, salinity, the effect of ocean whitecaps and marine biogeochemistry (e.g. Briegleb and Ramanathan, 1982; Roesch et al., 2002; Séférian et al., 2018). Subsequently, all sea-surface datapoints are masked and replaced with spectrally-varying α_s calculated using a SOCRATES internal subroutine. This subroutine accounts for the spectral variation of sea- α_s due to Fresnel reflection, modified to account for the impact of wind-induced variations in sea-surface roughness following Cox and Munk (1954), assuming nominal values for 10 metre wind speed of 7 m s^{-1} . Refractive indices of water are taken from Jerlov (1976) and Hale and Querry (1973), modified for nominal values of sea surface salinity of 35 g kg^{-1} . The effect of ocean whitecaps are also included from Whitlock et al., (1982); however this makes a small contribution to spectral reflectance given that whitecaps typically account for less than 2% of the sea-surface. Since sea- α_s is dependent on solar zenith angle, α_s is calculated assuming a daylight-average solar zenith angle at each grid-point and month. This scheme is of similar complexity to the representation of sea-surface albedo in Collins et al., (2018) and is more sophisticated than models such as the European Centre for Medium-range Weather Forecasts Integrated Forecasting System (ECMWF-IFS), where no spectral variation is assumed (IFS Documentation CY47R1, 2020). Much more sophisticated sea-surface parameterisations are available, for example Séférian et al., (2018). However, these require the specification of many more parameters, such as biophysical processes that govern the interaction of solar radiation with marine chlorophyll, and spatially and temporally varying surface wind speeds.

Text S4

The following description, taken from Collins et al., (2004)¹, outlines how CAM 3.0 CH₄ scale-heights are derived:

Globally, tropospheric CH₄ is assumed to be well mixed and are prescribed in units of mass mixing ratio (MMR) as follows:

$$\mu_{\text{CH}_4}^0 = 0.554 w_{\text{CH}_4} \quad (\text{S3})$$

where w_{CH_4} denotes the volume mixing ratio of CH₄. Above the tropopause CH₄ MMRs are specified as zonally-averaged quantities that vary with latitude to imitate the effect of stratospheric circulation and chemistry on CH₄. The pressure level of the tropopause (in hPa) is defined by:

$$p_{\text{trop}} = 250.0 - 150.0 \cos^2 \phi \quad (\text{S4})$$

where ϕ denotes latitude.

Following this, for $p \leq p_{\text{trop}}$, stratospheric CH₄ MMRs are defined as:

$$\mu_{\text{CH}_4} = \mu_{\text{CH}_4}^0 \left(\frac{p}{p_{\text{trop}}} \right)^{X_{\text{CH}_4}} \quad (\text{S5})$$

where X_{CH_4} denotes the mixing ratio scale heights. For $|\phi| \leq 45^\circ$,

¹ Two errors were found in this description of CH₄ mixing ratio scale-heights. Both related to Equation S7, reported originally by Collins et al., (2004) as: $X_{\text{CH}_4} = 0.2353 + 0.22549(\phi)$. See Equation S7 for corrections. We thank William D. Collins for providing clarification.

$$X_{CH_4} = 0.2353 \quad (S6)$$

and for $|\phi| > 45^\circ$,

$$X_{CH_4} = 0.2353 + 0.022549(\phi - 45) \quad (S7)$$

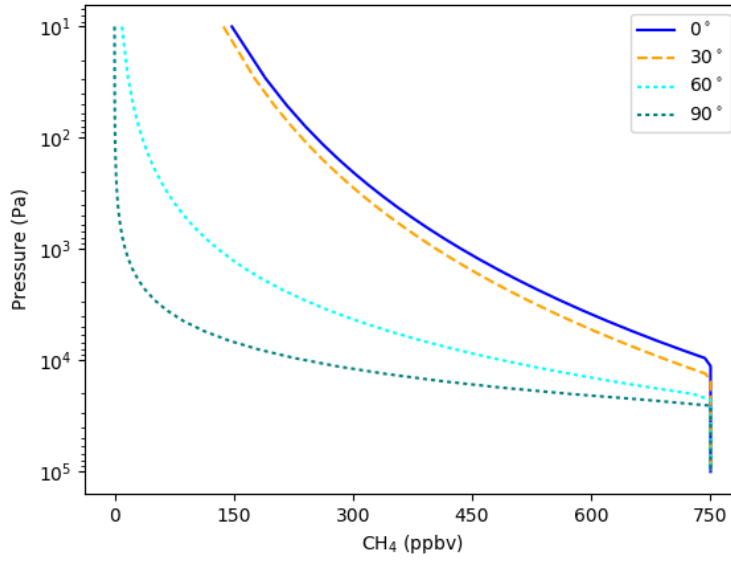


Figure S5: Zonal-mean vertically-varying CH_4 profiles at 0° , 30° , 60° and 90° latitude. CH_4 is well-mixed throughout the troposphere at 750 ppb and decreases exponentially above the tropopause as per the CAM 3.0 latitude-dependent mixing ratio scale height given by Equation S5. Note that the profile at 90° does not drop to zero completely, however it reaches a minimum of 0.04 ppb at 10 Pa. Also note that this equation does not model hemispherical and seasonal differences in mixing ratio scale heights. The profiles at the four latitudes shown here are therefore representative of both the northern and southern hemispheres for each month of the year.

Figure S6

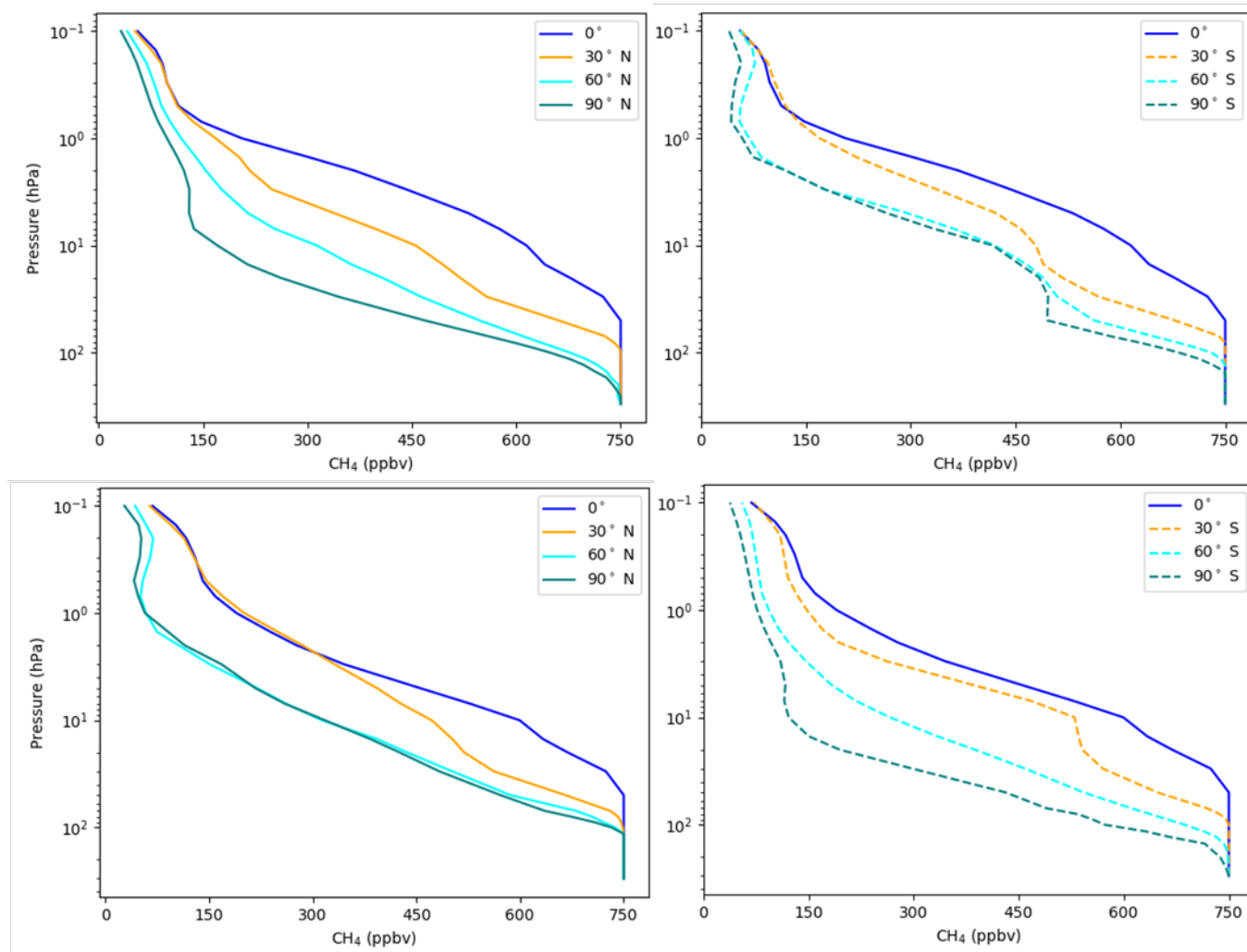


Figure S6: Pre-industrial (750 ppb) zonal seasonal-mean CH₄ profiles for DJF (upper panel) and JJA (lower panel) at 0°, 30°, 60° and 90° north (left-hand side) and south (right-hand side) constructed using MIPAS zonal-mean CH₄ fall off rate as described in Text S2.

Figure S7

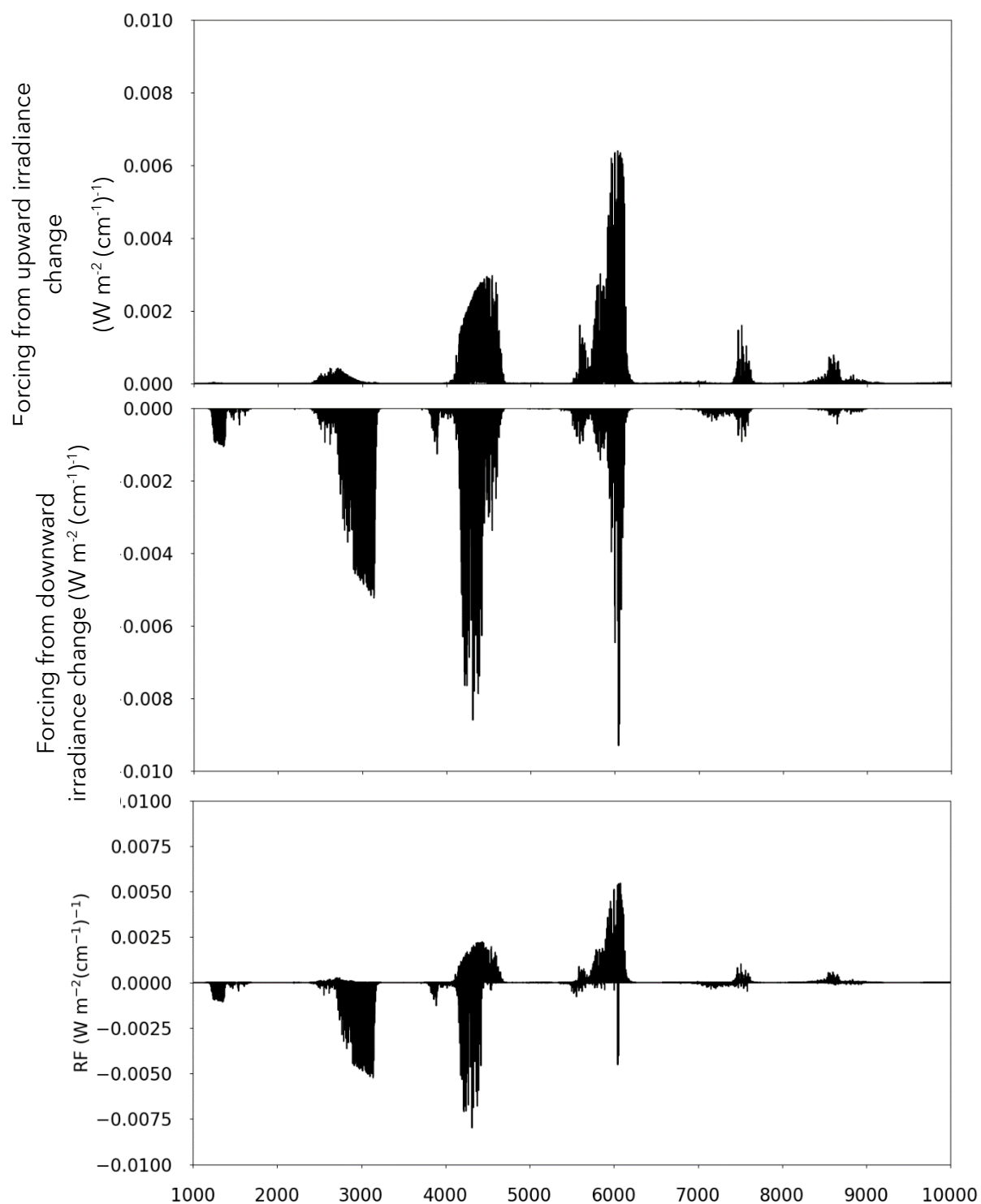


Figure S7: Spectral variation of upwelling (top) and downwelling (middle) CH_4 SW IRF components at the tropopause and the corresponding net tropopause IRF (lower) following a perturbation in CH_4 from 850 ppb to 1700 ppb using atmospheric data from Case 6 of Phase I of the Continual Intercomparison of Radiation Codes (CIRC) experiments (Oreopoulos et al., 2012).

Figure S8

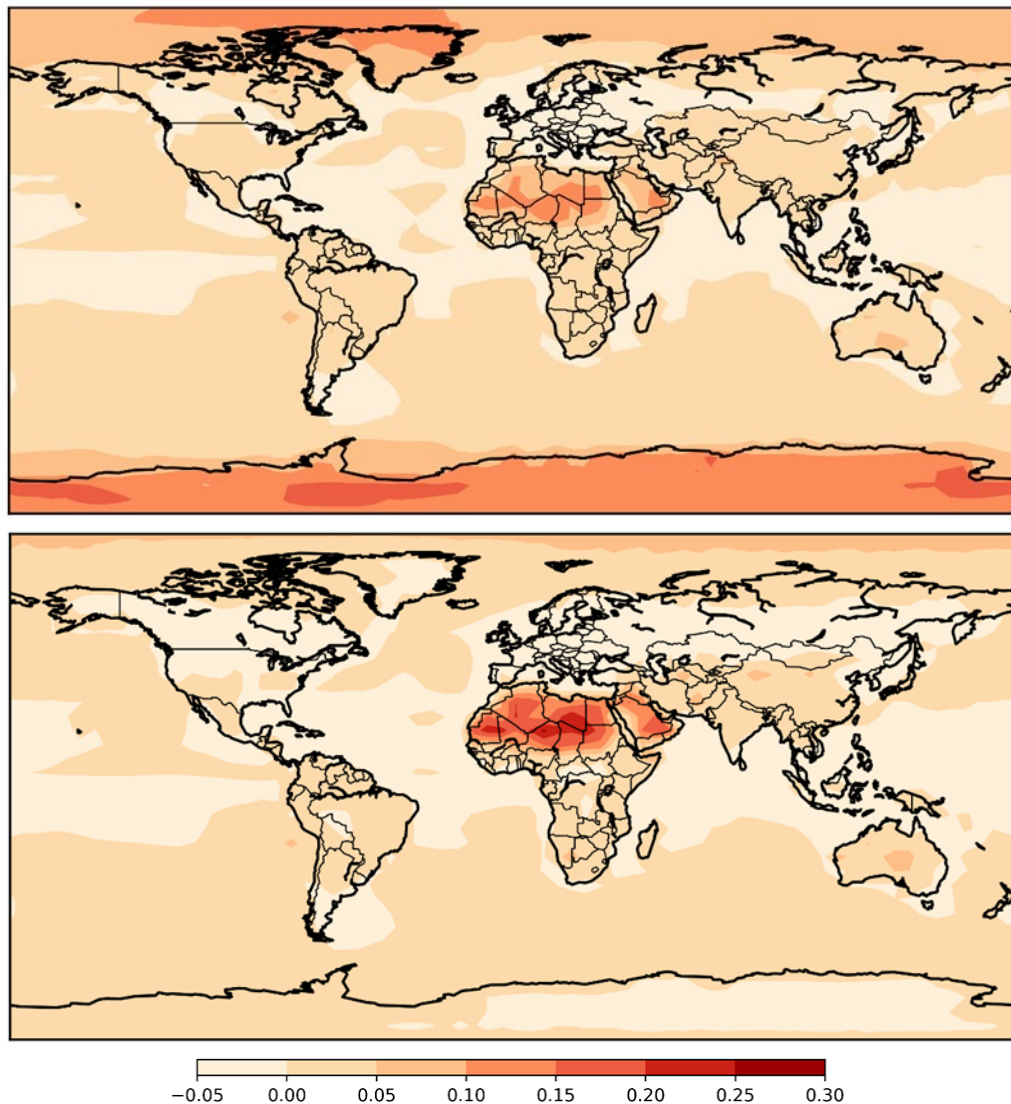


Figure S8: The impact of using spectrally-constant and spectrally-varying surface albedos on CH₄ SW IRF. Upper panel: Quasi-annual all-sky SW tropopause IRF (W m^{-2}) derived using CAM 3.0 CH₄ mole fraction scale-heights and monthly-mean fields of spectrally-constant ERA-interim α_s . Lower panel: Quasi-annual all-sky SW tropopause IRF (W m^{-2}) derived using CAM 3.0 CH₄ mole fraction scale-heights and SCIAMACHY-SSA α_s . Both are derived following a perturbation in CH₄ from 750 to 1800 ppb. Note that these estimates represent quasi-annual calculations from the four mid-season months (January, April, July and October). Analysis demonstrates that the average of these months (referred to here as the quasi-annual mean) serves as a good representation of the full 12-month annual mean for the purpose of sensitivity tests. In the main text, all reported forcings are full 12-month annual means except where indicated.

Text S5

Multi-model intercomparison studies provide a useful basis from which to assess uncertainty in radiative transfer schemes used in forcing calculations (e.g. Collins et al. 2006; *Pincus et al.*, 2015). However, there is lack of research comparing broad-band radiative transfer codes against benchmark calculations specifically for CH₄ SW RF, largely because many ESM radiation codes have historically omitted explicit treatment of CH₄ SW absorption (e.g. Collins et al. 2006). We repeated our calculations using the current 6-band SOCRATES SW spectral file configuration (sp_sw_ga7) used in the Met Office Unified Model; it significantly underestimates CH₄ SW tropopause and surface IRF by around 45% compared to the 260-band configuration. Hogan and Matricardi (2020) also report an underestimation of CH₄ SW IRF by 25 – 45% by the Rapid Radiative Transfer Model for General Circulation Models (RRTMG) (Mlawer et al., 1997) compared to benchmark calculations for perturbations across 350 to 3500 ppb. As stated in Section 2, the 260-band SOCRATES SW spectral file has been validated here against the RFM. Under a range of atmospheric conditions and CH₄ SW IRF scenarios, sp_sw_260_jm2 exhibits a mean error of 8%. Uncertainty associated with the 9-band SOCRATES LW spectral file is assessed to be 6% considering prior research (e.g. Collins et al., 2006) and validation by Walters et al., (2019).

Hodnebrog et al., (2013; 2020) discuss error in RF calculations based on a multitude of previous studies to assess uncertainty related specifically to halocarbon RF; their discussion can be used more generally to assess error in other well-mixed GHG forcings (e.g. Etminan et al., 2016) and is used here to estimate the following uncertainties: the calculation of FDH (4%), atmospheric temperature (3%) and clouds (5%) and spatial and temporal averaging (1%). Hodnebrog et al., (2013) assess uncertainty related to tropopause height to be around 5%. Here, this uncertainty is assessed to be 10% for CH₄ SW tropopause IRF and 5% for CH₄ LW SARF. A higher value is attributed to SW calculations owing to the particular sensitivity of CH₄ SW tropopause IRF to the magnitude of upward and downward forcing components.

An additional important error source relates to the specification of α_s . The method used to interpolate and extrapolate SCIAMACHY data is idealised, owing to a lack of global datasets that extend across 0.17-10 μm . An uncertainty of 20% is tentatively attributed to the use of SCIAMACHY-SSA α_s . Furthermore, idealised tests demonstrate significant sensitivity of CH₄ SW IRF to the representation of H₂O self-continuum absorption. Using the analysis of laboratory data in Anisman et al., (2022), a modified sp_sw_260_jm2 spectral file with a reduced self-continuum strength in the 2.1 μm window causes a 5% difference in January clear-sky CH₄ tropopause SW IRF (-0.039 W m^{-2}) for a 750-1800 ppb perturbation compared to the unmodified spectral file (-0.041 W m^{-2}). Similarly, the specification of H₂O fields is a potential error source given the dependence of CH₄ SW IRF on H₂O spectral overlap (Etminan et al., 2016). Collins et al., (2018) find that SW forcings differ by around 32% under clear-sky conditions and 4% under all-sky conditions, depending on which ESM background climatology is used (Table 2, Collins et al., 2018). Here, uncertainty related to H₂O fields is assessed to contribute 2% to total uncertainty in all-sky CH₄ SW tropopause IRF and LW SARF.

Finally, considering the sensitivity of CH₄ SW IRF to vertical profile specification, an uncertainty of 5% is attributed to the use of MIPAS CH₄ climatology based on the standard deviation of multi-instrument spread in CH₄ measurements (Hegglin et al., 2021). A lower uncertainty of 2% is attributed to the use of MIPAS CH₄ climatology in CH₄ LW SARF, as this calculation is not significantly dependent on CH₄ profile specification (e.g., Freckleton et al., 1998).

As summarised in the main text, we calculate uncertainty using the root-sum-square method. A total uncertainty of $\pm 12\%$ is attributed to CH₄ LW SARF and a total uncertainty of $\pm 25\%$ is tentatively attributed to CH₄ SW tropopause IRF. Overall, given that the CH₄ SW effect is about 6% of the net CH₄ SARF, this yields a total uncertainty of around $\pm 13\%$ in net CH₄ SARF.

Source of Uncertainty	Estimated Contribution to SW Uncertainty	Estimated Contribution to LW Uncertainty
Spectroscopic parameters	1%	1%
SW radiation code	8%	-
LW radiation code	-	6%
H₂O continuum	5%	Not assessed
FDH calculation	-	4%
Tropopause level	10%	5%
Atmospheric temperature	-	3%
Clouds	5%	5%
H₂O field	2%	4%
Spatial and temporal averaging	1%	1%
MIPAS CH₄ mole fractions	5%	2%
SCIAMACHY-SSA α_s	20%	-
Total uncertainty	25%	12%

Table S2: Sources of uncertainty and their estimated contribution to total uncertainty in estimates of CH₄ SW tropopause IRF and CH₄ LW SARF calculated using MIPAS CH₄ mole fractions and SCIAMACHY-SSA α_s . The total uncertainty is also shown and is calculated using the root-sum-square method. See Text S5 for references used as basis for uncertainty estimates.