Assimilation of both column- and layer-integrated dust opacity observations in the Martian atmosphere

Peter L
 Read¹, Tao Ruan², Roland M B Young³, Stephen Lewis⁴, Luca Montabone⁵, and Alexandru Valeanu²

¹Oxford University ²University of Oxford ³UAE University ⁴Open University ⁵Space Science Institute

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Abstract

A new dust data assimilation scheme has been developed for the UK version of the Laboratoire de Météorologie Dynamique (LMD) Martian General Circulation Model. The Analysis Correction scheme (adapted from the UK Met Office) is applied with active dust lifting and transport to analyze measurements of temperature, and both column-integrated dust optical depth (CIDO), τ_{ref} , (rescaled to a reference level) and layer-integrated dust opacity (LIDO). The results are shown to converge to the assimilated observations, but assimilating either of the dust observation types separately does not produce the best analysis. The most effective dust assimilation is found to require both CIDO and LIDO observations, especially for Mars Climate Sounder (MCS) data that does not access levels close to the surface. The resulting full reanalysis improves the agreement with both in-sample assimilated CIDO and LIDO data and independent observations from outside the assimilated dataset. It is thus able to capture previously elusive details of the dust vertical distribution, including elevated detached dust layers that have not been captured in previous reanalyses. Verification of this reanalysis has been carried out under both clear and dusty atmospheric conditions during Mars Years 28 and 29, using both in-sample and out of sample observations from orbital remote sensing and contemporaneous surface measurements of dust opacity from the Spirit and Opportunity landers. The reanalysis was also compared with a recent version of the Mars Climate Database (MCD v5), demonstrating generally good agreement though with some systematic differences in both time mean fields and day-to-day variability.

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Tao Ruan¹, R. M. B. Young^{1,2}, S. R. Lewis³, L. Montabone^{1,4}, A. Valeanu¹, and P. L. Read¹

5	¹ Atmospheric, Oceanic and Planetary Physics, Department of Physics, University of Oxford, Clarendon
6 7	Laboratory, Parks Road, Oxford, OX1 3PU, UK ² Department of Physics & National Space Science and Technology Center, UAE University, Al Ain,
8 9 10	³ School of Physical Sciences, The Open University, Walton Hall, Milton Keynes, MK7 6AA, UK ⁴ Space Science Institute, Boulder, CO 80301, USA

Key Points:

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12	•	Assimilation of atmospheric measurements of Mars into a global circulation model
13		is extended to include limb profiles of dust opacity.
14	•	Combining nadir and limb profiles of dust opacity enables more accurate recov-
15		ery of dust vertical structure, including elevated layers.
16	•	The climate reanalysis is significantly improved, as shown in comparisons with in-
17		dependent observations and the Mars Climate Database.

Corresponding author: Peter Read, peter.read@physics.ox.ac.uk

18 Abstract

A new dust data assimilation scheme has been developed for the UK version of the Lab-19 oratoire de Météorologie Dynamique (LMD) Martian General Circulation Model. The 20 Analysis Correction scheme (adapted from the UK Met Office) is applied with active dust 21 lifting and transport to analyze measurements of temperature, and both column-integrated 22 dust optical depth (CIDO), $\tau_{\rm ref}$ (rescaled to a reference level), and layer-integrated dust 23 opacity (LIDO). The results are shown to converge to the assimilated observations, but 24 assimilating either of the dust observation types separately does not produce the best 25 analysis. The most effective dust assimilation is found to require both CIDO and LIDO 26 observations, especially for Mars Climate Sounder (MCS) data that does not access lev-27 els close to the surface. The resulting full reanalysis improves the agreement with both 28 in-sample assimilated CIDO and LIDO data and independent observations from outside 29 the assimilated dataset. It is thus able to capture previously elusive details of the dust 30 vertical distribution, including elevated detached dust layers that have not been captured 31 in previous reanalyses. Verification of this reanalysis has been carried out under both 32 clear and dusty atmospheric conditions during Mars Years 28 and 29, using both in-sample 33 and out of sample observations from orbital remote sensing and contemporaneous sur-34 face measurements of dust opacity from the Spirit and Opportunity landers. The reanal-35 ysis was also compared with a recent version of the Mars Climate Database (MCD v5), 36 37 demonstrating generally good agreement though with some systematic differences in both time mean fields and day-to-day variability. 38

³⁹ Plain Language Summary

Data assimilation is a method of combining atmospheric observations, which are 40 inevitably uncertain and incomplete in their coverage, with a global numerical model. 41 It is commonly used for the Earth to reconstruct the best possible record of the chang-42 ing global climate. This has also been used for the Martian atmosphere in recent years, 43 using measurements of temperature, dust and ice from satellites in orbit around Mars. 44 But most previous efforts have only used measurements of the total amount of dust in 45 a vertical column from instruments that "look" vertically downwards to the Martian sur-46 face. In new work presented here, however, we also use detailed measurements of the ver-47 tical structure of the dust distribution from an instrument that "looks" towards the edge 48 of the planet. This is much more effective when atmospheric dust is not mainly concen-49 trated near the ground. Such events are reasonably common on Mars, when elevated lay-50 ers of dust are formed, which can strongly affect how the atmosphere is heated by the 51 Sun. We present examples of situations when previous methods failed to recover the cor-52 rect dust distribution, as verified against independent measurements e.g. from the Spirit 53 and Opportunity Rovers, and compare with the ESA Mars Climate Database. 54

55 1 Introduction

The dust cycle is a key component of the Martian climate, and is extremely important 56 for understanding the interannual, seasonal and synoptic evolution of the Martian en-57 vironment. (e.g. Newman et al., 2002a; Kahre et al., 2017, and references therein). In-58 tensive measurements of atmospheric temperature and dust extending over more than 59 ten Mars years (MY) now exist with unprecedented spatial coverage, thanks to various 60 orbital spacecraft. Such observations have already helped to improve our understand-61 ing of Mars' weather and climate. However, the incomplete coverage of these measure-62 ments across the planet constrains our ability to study the general circulation in full de-63 tail, particularly those aspects related to dust opacity. For instance, the Thermal Emission Imaging System (THEMIS) carried by the Mars Odyssey (MO) spacecraft can pro-65 vide multi-annual measurements of Column Integrated Dust Opacity (CIDO), but its 66 coverage in space and time is quite limited. 67

On the other hand, numerical models provide four-dimensional simulated data with 68 moderate-high temporal and spatial resolution and complete coverage in space and time, 69 but often fail to reproduce the dust cycle's full range of variability. Various authors, start-70 ing with Newman et al. (2002b) (see also Kahre et al. (2017) for a recent review) showed 71 that a global circulation model (GCM) could capture the onset and growth of regional 72 dust events, but did not realistically capture the observed interannual variability. In par-73 ticular, they could not reproduce the relatively "quiet" year of dust activity that occurs 74 immediately after a simulated global dust storm (GDS) year. Others have sought to take 75 additional factors into account, such as the finite extent of the surface dust reservoir (e.g. 76 Pankine & Ingersoll, 2004; Szwast et al., 2006) or nonlinear effects associated with the 77 "shadowing" of pockets of dust behind rocks and boulders (Mulholland et al., 2013). But 78 even the most sophisticated free-running GCMs still struggle to capture realistic inter-79 annual variability associated with dust lifting and transport. 80

To aid this task, data assimilation has become an optimal approach to provide a 81 solution that is consistent with both observations and modelled physical constraints. Data 82 assimilation corrects model-predicted variables towards observations such that the re-83 sulting solution can represent the full observed variability of the climate. This approach 84 has been widely used as an effective tool in operational weather forecasting systems or 85 climate models for analyzing meteorological variables for the Earth (e.g. Lorenc et al., 86 1991; Kalnay, 2003). This approach has already been used for a number of years to in-87 vestigate tracer/chemical evolution in the Earth's atmosphere (Collins et al., 2001; J. Wang 88 et al., 2004; Schutgens et al., 2010; Benedetti et al., 2018). Collins et al. (2001), for ex-89 ample, used an optimal interpolation approach to assimilate satellite retrievals of total 90 column aerosol optical depth (AOD) over the Indian Ocean, which reproduced the daily 91 variations of AOD at a single model grid point. J. Wang et al. (2004) used nudging to 92 assimilate AOD into a nonhydrostatic atmospheric model, which captured the observed 93 evolution of a dust event near Puerto Rico. More recently, Schutgens et al. (2010) ap-94 plied the Local Ensemble Transform Kalman filter to assimilate AOD from the AERONET 95 global surface observation network, which captured the evolution of AOD and also re-96 duced uncertainties in model estimates of the evolving aerosol distribution. At the time 97 of writing, around five major operational centres around the world use a variety of as-98 similation techniques, including optimal analysis (similar to the Analysis Correction scheme 99 presented here), variational methods or ensemble Kalman filters to analyse observations 100 of dust and aerosols from various sources (e.g. Benedetti et al., 2018, and references therein), 101 such as AOD derived from orbiting or surface-based platforms. 102

Few centres have assimilated dust profile observations, however, preferring instead 103 to focus on achieving high horizontal resolution utilizing the much more abundant AOD 104 measurements, Limited publications to date include the work of Yumimoto et al. (2008), 105 who assimilated vertical profiles of the dust extinction coefficients in a regional dust trans-106 port model. In their study, the data from a ground-based lidar network were interpo-107 lated to the vertical model levels for analyzing the model prognostic dust variables. More 108 recently, Sekiyama et al. (2010) directly assimilated the total attenuated backscatter-109 ing coefficient from the Cloud-Aerosol Lidar and Infrared Pathfinder Satellite Observa-110 tions (CALIPSO) mission into a global chemistry-transport model. The measurements 111 were averaged approximately to the model horizontal and vertical resolution before the 112 assimilation. This approach has recently been extended by Cheng et al. (2019), who also 113 assimilated CALIPSO profiles of aerosol optical depth using a 4D Ensemble Kalman Fil-114 ter approach. 115

Data assimilation has also been applied to the Martian atmosphere with success. S. R. Lewis & Read (1995) implemented the analysis correction (AC) scheme (Lorenc et al., 1991) in a simple version of a Mars General Circulation Model (MGCM) in order to assimilate temperature profiles from the Pressure Modulator Infrared Radiometer (PMIRR) instrument on-board the short-lived Mars Observer spacecraft (1993). Their

results showed that assimilation of such observations was feasible and that it improved 121 the agreement between model and observations. S. R. Lewis et al. (2007) extended this 122 approach to include dust tracer assimilation, which was combined with a full MGCM to 123 assimilate thermal profiles and CIDO rescaled to a reference level (hereafter $\tau_{\rm ref}$) using 124 retrievals from the Thermal Emission Spectrometer (TES) on-board MGS (M. D. Smith 125 et al., 2003). The performance of the data assimilation system was validated against in-126 dependent radio occultation measurements by Montabone, Lewis, Read, & Hinson (2006). 127 This showed that combined temperature and $\tau_{\rm ref}$ assimilation was able to reduce discrep-128 ancies between the model and radio occultation data below 20 km, especially when dust 129 amounts were large and changing rapidly, although some large discrepancies remained 130 due to known inconsistencies between TES temperature profiles and radio occultation 131 data. This approach was further extended by Holmes et al. (2020) to include assimila-132 tion of column dust optical depth measurements derived from Mars Reconnaissance Or-133 biter (MRO) MCS retrievals, together with measurements of water ice and ozone, into 134 a version of the LMD/UKMGCM that also advects dust and other tracers with the anal-135 vsed winds. 136

An alternative approach to assimilation of Mars observations was developed by Hoff-137 man et al. (2010) based on a complementary method using the ensemble Kalman filter 138 (EnKF, Evensen, 2003) to assimilate TES temperature retrievals. They found generally 139 improved agreement with TES temperature observations over a free-running model, and 140 in Greybush et al. (2012) the joint assimilation of TES temperatures with forcing using 141 a 2D CIDO dust field from TES was shown to improve the agreement of model and TES 142 temperatures further. This method has also been extended to include assimilation of col-143 umn dust optical depths from both MGS/TES and MRO/MCS observations (Greybush 144 et al., 2019) and a dataset is publicly available known as EMARS. 145

The approach used by S. R. Lewis et al. (2007) first assimilated TES temperature 146 profiles and $\tau_{\rm ref}$ without explicitly advecting the dust tracer field, using an empirical re-147 lation (Conrath, 1975) to prescribe the vertical distribution of dust. This system has sub-148 sequently been applied in several studies of Martian weather and climate (Montabone 149 et al., 2005; S. Lewis & Barker, 2005; Montabone, Lewis, Read, & Withers, 2006; S. R. Lewis 150 et al., 2007; Wilson et al., 2008; Rogberg et al., 2010; S. R. Lewis et al., 2016), and both 151 a three-year reanalysis covering MY 24-27 using this assimilation system and another 152 covering much of the MCS period (MY 28-32) have been published (Montabone et al... 153 2014; Holmes et al., 2020). Navarro et al. (2014) also assimilated Mars Climate Sounder 154 (MCS) temperature profiles and modified the dust vertical distribution using its corre-155 lation with temperature. However, it is essentially different from the work presented here, 156 in which dust observations are directly assimilated. The more recent study of data as-157 similation issues on Mars by Navarro et al. (2017) is more similar to the present work 158 in including some cases that assimilated MCS dust opacity profiles using the EnKF method. 159 Their study indicated some promise for this approach, although they only analysed a part 160 of MY 29 and noted some difficulties in capturing the diurnal variation in dust vertical 161 distributions. 162

The data assimilation system developed here is based on the scheme described by 163 S. R. Lewis et al. (2007). However, that scheme does not assimilate a vertically-resolved 164 dust distribution, only TES nadir retrievals of CIDO, and the model does not transport 165 the dust actively. The newly-available dataset from MCS on board MRO (Kleinböhl et 166 al., 2009) does provide vertically resolved, global measurements of the atmospheric dust 167 distribution. With this new dataset to hand, here we update the existing data assim-168 ilation system to better represent the Martian dust cycle. In later work we will use this 169 to study the formation and life cycles of regional and global dust storms in detail. 170

Section 2 describes the Mars GCM and current data assimilation scheme, and the observations of Martian dust are in Section 3. We describe how the assimilation was adapted in Section 4. Sections 5 and 6 describe verification against in-sample and out-of-sample



Figure 1: Spatial and temporal distribution of available dust opacity data from THEMIS and MCS during the study period. The colour scales show the number of measurements in 5° L_s and 3° latitude bins.

observations respectively, while Section 7 describes a systematic comparison of the Mars
 Climate Database against the reanalysis. We conclude in Section 8.

¹⁷⁶ 2 Overview of Mars GCM and data assimilation scheme

In this work the model used is based on the UK version of a three-dimensional Martian
Global Climate Model (UK-LMD MGCM, v5.1.3) (Forget et al., 1999; Mulholland et al.,
2013). The model combines a spectral dynamical solver at triangular truncation T31,
corresponding to a 96×48 longitude-latitude grid in real space, a tracer transport scheme
and dust lifting and deposition routines, along with a full range of physical parameterizations.

The equations for a hydrostatic, adiabatic and inviscid gas surrounding a rotating 183 spherical planet are cast in vorticity-divergence form. In the vertical, levels are defined 184 in terms of the terrain-following σ coordinate system using a standard finite difference 185 approach. There are 25 levels with the first three at 4, 19, and 44 m above the surface, 186 to resolve detailed surface processes represented in the model. The model top varies in 187 altitude over time but is typically at around 100 km, with a sponge layer (applying a lin-188 ear drag on eddy vorticity and divergence) in the uppermost three levels to reduce spu-189 rious reflections of vertically propagating waves. There are typically 480 dynamical and 190 96 physics timesteps per sol (where a sol is a mean solar day on Mars). 191

The radiative transfer scheme calculates atmospheric absorption and emission due 192 to carbon dioxide and airborne dust; the radiative effects of water vapour and ice are not 193 included since our focus here is on the dust cycle. We rely, therefore, on the tempera-194 ture assimilation to account for the radiative effects of clouds. The balance between in-195 coming radiative flux and thermal conduction in the soil contributes to changes in sur-196 face temperature, using a surface thermal inertia field derived from TES and Viking (For-197 get et al., 1999) and topography from the Mars Orbiter Laser Altimeter on MGS (D. E. Smith 198 et al., 2001). The surface roughness length z_0 is based on a global map compiled by Hébrard 199 et al. (2012), and implemented in the UK-MGCM by Mulholland et al. (2015). 200

The dust transport scheme includes dust lifting parameterizations, tracer advection, gravitational sedimentation and dry deposition. We assume a 1.5 µm particle size for simplicity, based on Mars Exploration Rover (MER) observations (Lemmon et al., 204 2004). The two most important distinct mechanisms responsible for the injection of dust 205 into the atmosphere are thought to be dust lifting by near-surface wind stress, and dust 206 lifting by dust devils (Newman et al., 2002a).

The data assimilation scheme is based on the analysis correction sequential esti-207 mation (AC) scheme (Lorenc et al., 1991) but with modifications specific to Mars (S. R. Lewis 208 et al., 2007). The assimilation step is computationally inexpensive compared with the 209 rest of the model, and so is performed at each dynamical timestep. S. R. Lewis et al. (2007) 210 describe the scheme in full detail. Temperature assimilations are the same as in that work, 211 except for the observational dataset used. They assimilated dust CIDO observations with-212 out advecting the dust tracer, instead setting the vertical distribution of dust opacity 213 using an empirical relation following Conrath (1975). In this work we extend the dust 214 assimilation to incorporate advective transport of radiatively active dust in the simula-215 tion model as well as to assimilate both CIDO and LIDO (Laver integrated dust opac-216 ity) observations; this is described in Sect. 4. 217

The ratio of observational error to first guess error used in the normalization factor \tilde{Q}_i (Lorenc et al., 1991, Eq. 3.20) is set to 1 for assimilation of TES temperature observations (as previously done by S. R. Lewis et al., 2007), implying that the model and observation errors are comparable. Following the study of ice opacity assimilation by Steele et al. (2014), we also set this ratio to 1 for the dust observations.

3 Observations of Martian dust

Thanks to various spacecraft in orbit around Mars since 1997, measurements of atmospheric temperature and dust exist covering the Martian atmosphere over more than ten MYs. The instruments on board these spacecraft for determining temperature and dust in the Martian atmosphere that have been used for assimilation, such as the present study, include (amongst others) TES on MGS (M. D. Smith, 2004), THEMIS on MO (M. D. Smith, 2009) and MCS on MRO (Kleinböhl et al., 2009).

TES and THEMIS dust retrievals contain CIDO data only, while MCS data con-230 tain more recent satellite observations (MCS v3 was used here for this initial proof of 231 concept) with vertically resolved, asynoptically-sampled global retrievals of atmospheric 232 profiles of temperature, and LIDO (McCleese et al., 2010). This contains information 233 on the day-to-day variability of Martian weather from the near surface to the top of the 234 middle atmosphere around 80 km altitude (Kleinböhl et al., 2009). The spacecraft have 235 different operational periods and orbits, so their retrievals have different temporal and 236 spatial coverage. Only THEMIS has overlapping observational periods with the other 237 two datasets. Further details of the THEMIS dataset, including the retrieval algorithm, 238 can be found in M. D. Smith et al. (2000, 2003) and M. D. Smith (2009). 239

The analysis in this paper focuses on part of the MCS mapping period from MY28 240 $L_s = 110^{\circ}$ (solar longitude) to MY29 $L_s = 330^{\circ}$. During this period, THEMIS CIDO 241 data and MCS LIDO data are both available. The spatial coverage of THEMIS retrievals 242 varies significantly with L_s , while MCS retrievals are more consistent and uniform ex-243 cept during the global-scale dust storm (GDS) season (roughly from MY28 $L_s = 270^{\circ}$ -244 305°; see Fig. 1(b)). Coverage was restricted to the very early stage of the storm and 245 poleward of 45°N throughout. The spatial coverage of these two datasets within the study 246 period is shown in Fig. 1. 247

Because the dust is assumed by the THEMIS dust opacity retrieval algorithm to be well mixed, the THEMIS CIDO data is commonly reported rescaled (i.e. as τ_{ref}) to a reference pressure of 610 Pa (M. D. Smith, 2009), to remove the effects of variable topography. However, this assumption may introduce uncertainty when the dust is not well



Figure 2: Sequence of operations in the new data assimilation scheme with active dust transport. Green boxes show initial conditions, blue boxes show data generated by MGCM integration, and red boxes show individual MGCM modules. Text in black applies to both CIDO and LIDO assimilation, text in blue applies to CIDO assimilation only, and text in red applies to LIDO assimilation only. Only variables related to the data assimilation scheme are included.

mixed. When an intense detached dust layer exists (Heavens et al., 2011), for example, rescaling under the well-mixed assumption could lead to an overestimate in τ_{ref} .

It is also important to note that THEMIS dust observations are provided as an in-254 frared absorption optical depth, while the modeled $\tau_{\rm ref}$ is the visible extinction optical 255 depth. Using numerical experiments, M. D. Smith (2009) determined the conversion be-256 tween IR absorption and visible extinction optical depth to be $\gamma \sim 1.3$, and Clancy et 257 al. (2003) used a scaling factor $\varepsilon = 2$ to convert from IR to visible for dust of size 1.5-258 $2.0\,\mu\text{m}$. Lemmon et al. (2004) compared visible optical depths with MER measurements 259 at 9 µm wavelength, and found agreement with Clancy et al. (2003). For simplicity, in 260 this work the dust particle size is approximated by a constant $1.5\,\mu m$, which is reason-261 ably consistent with various observational studies (Pollack et al., 1995; Clancy et al., 2003; 262 Lemmon et al., 2004). Hence the conversion factor from THEMIS IR absorption opti-263 cal depth to a model-compatible visible extinction optical depth is 2.6, and in this work 264 we imply the visible extinction optical depth when referring to CIDO and τ_{ref} , unless 265 otherwise stated. For MCS LIDO retrievals, the infrared opacities (at a wavelength cen-266 tred on 21.6 μ m) were multiplied by a factor of 7.3 to convert them to a visible equiv-267 alent (Montabone et al., 2015). 268

For a fully-independent validation of the analysis, upward-looking surface obser-269 vations provide a bottom-up view of CIDO that is independent of satellite-based datasets, 270 although only over particular locations. The MER missions, Spirit (14.57°S, 175.48°E) 271 and Opportunity (1.95 °S, 5.53 °W) provide almost continuous data coverage during MY28 272 and MY29, concurrent with the study period in this work. Both rovers carried a Pan-273 cam camera, which included solar filters at 440 nm and 880 nm wavelengths. The effec-274 tive dust particle radii based on the rovers' observations were $1.47\pm0.21\,\mu\mathrm{m}$ for Spirit 275 and $1.52\pm0.18\,\mu\text{m}$ for Opportunity (Lemmon et al., 2004). However, the 440 nm filter 276 is significantly affected by a red leak (Lemmon et al., 2015). As there should not be a 277 significant difference between the measured CIDO at 880 nm and CIDO at 700 nm (used 278 as the model visible wavelength), we therefore rescale the CIDO measurements at 880 nm 279 to the reference pressure 610 Pa, and make the comparison directly with the modelled 280 sol-averaged $\tau_{\rm ref}$ (also rescaled to 610 Pa). 281

About 10% of the THEMIS and MCS data are excluded from the assimilation and used for out-of-sample validation. The withheld data were taken as every 10th THEMIS data point and every 10th MCS vertical profile. Withholding this small fraction of the dataset should not greatly affect the assimilated results. One should not be surprised that these validation data may be correlated with the assimilated data, which does weaken their usefulness in validating the model to some extent. However, these datasets were the best available for the assimilation itself.

²⁸⁹ 4 Dust data assimilation with activated transport

Initial conditions for the prognostic variables and dust tracers are taken from a free-running spin-up run, which is run for two years prior to the start of the assimilation. Temperature assimilation using MCS data was always included, using the method described by S. R. Lewis et al. (2007), and such that temperatures were always assimilated before the dust assimilation. An identical free-running simulation without any assimilation, but with a fully active dust lifting and transport cycle tuned to reproduce plausible seasonal variations of dust loading, was run in parallel.

Previous assimilation studies using the UK-LMD MGCM excluded active dust transport, instead just correcting the temperature profiles and τ_{ref} . The dust distribution remained static in the absence of dust observations and the vertical distribution was prescribed using the Conrath (1975) distribution.

In the new scheme presented here, the data assimilation system is updated to include full dust transport, lifting and sedimentation while correcting $\tau_{\rm ref}$ and/or the model's vertical dust distribution.

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4.1 CIDO assimilation only

In this configuration only the CIDO retrievals are assimilated. The sequence of operations is shown in Fig. 2. First the dynamics timestep is integrated, and then the dust is advected to obtain the three dimensional (3D) distribution of dust mass mixing ratio $q(\mathbf{x}, m)$. The CIDO at position \mathbf{x} , predicted from this distribution ($\tau_{\rm C}$) is obtained by linearly summing up the layer-integrated dust opacity (LIDO) within each model layer ($\tau_{\rm LIDO}$). In the model the LIDO at horizontal position \mathbf{x} in model level m is given by

$$\tau_{\text{LIDO}}(\mathbf{x}, m) = q(\mathbf{x}, m) q_{\text{ext}} \frac{p_s(\mathbf{x}) \Delta \sigma(m)}{g}$$
(1)

where $q_{\rm ext} = (3Q_{\rm ext})/(4\rho r)$, $q(\mathbf{x}, m)$ is the dust mass mixing ratio, $g = 3.72 \,\mathrm{m\,s^{-2}}$ is the gravitational acceleration, $Q_{\rm ext}$ is the extinction coefficient, $\rho = 2500 \,\mathrm{kg\,m^{-3}}$ and $r = 1.5 \,\mathrm{\mu m}$ are the density and radius of dust particles, respectively. $p_s(\mathbf{x})$ is the surface pressure, and $\Delta\sigma(m)$ is the layer thickness. The reference dust opacity $\tau_{\rm ref}$ at a reference pressure $p_{\rm ref}$ is then determined from the model by

$$\tau_{\rm ref} = \left(\sum_{m} \tau_{\rm LIDO}(\mathbf{x}, m)\right) \frac{p_{\rm ref}}{p_s(\mathbf{x})} \tag{2}$$

305 306 The reference pressure p_{ref} is arbitrary; to compare the results with observations, modelled CIDO values are rescaled to 610 Pa.

The advected dust opacity and mass mixing ratio fields are then used to integrate the physical parametrizations. Finally, τ_{ref} and T are updated using data assimilated by the AC scheme, followed by increments to u and v in thermal wind balance. The dust transport scheme transports a 3D dust mass mixing ratio field, but the assimilated CIDO dust observations constrain τ_{ref} only, so the dust mass mixing ratio at each model layer must be adjusted after the assimilation. This adjustment simply consists of a multiplicative scale factor $\lambda(\mathbf{x})$, which ensures that the shape of the vertical dust profile at each horizontal grid point remains the same before and after the assimilation of τ_{ref} :

$$\lambda(\mathbf{x}) = \frac{\tau_{\text{LIDO}}(\mathbf{x}, m)'}{\tau_{\text{LIDO}}(\mathbf{x}, m)} = \frac{\tau_{\text{ref}}'(\mathbf{x})}{\tau_{\text{ref}}(\mathbf{x})}$$
(3)

where variables without and with primes are before and after assimilation, respectively. Since the extinction coefficient and layer thickness are constant within a particular time step, Eq. 1 implies that the dust mass mixing ratio $q(\mathbf{x}, m)$ is proportional to $\tau_{\text{LIDO}}(\mathbf{x}, m)$, and therefore the adjustment $\lambda(\mathbf{x})$ can be applied directly to $q(\mathbf{x}, m)$. A similar assumption was also used when assimilating AOD on Earth (e.g. Collins et al., 2001; J. Wang et al., 2004).

313 4.2 LIDO only

A more advanced method is required to make proper use of the MCS vertically-resolved dust profiles. This section describes how retrievals of dust profiles from MCS are assimilated into the model by themselves (i.e., without assimilating CIDO). Figure 2 shows the procedure for assimilation of LIDO, and it is very similar to CIDO-only, but now LIDO (τ_{LIDO}) is analyzed directly. When the model layers have a smaller vertical spacing than the MCS measurements (typical in the lower and middle atmosphere), this approach avoids the direct interpolation of observational data to the model levels.

The assimilation of the vertical dust distribution is not yet widely used in Earth 321 aerosol modelling or forecasting. Work to date includes Yumimoto et al. (2008), who as-322 similated vertical profiles of dust extinction coefficients into a regional dust transport 323 model. Data from a ground-based LIDAR network were interpolated to the appropri-324 ate vertical model levels. Sekiyama et al. (2010) directly assimilated total attenuated backscat-325 tering coefficients measured by the CALIPSO spacecraft into a global chemistry-transport 326 model. In that case they averaged the observations over the model's horizontal and ver-327 tical resolution before assimilation. 328

Since MCS does not (in general) take data at the same levels as those used in the 329 model, we need to pre-process the observed dust distribution. Our approach differs in 330 approach from both Yumimoto et al. (2008) and Sekiyama et al. (2010), who both used 331 the CALIPSO satellite LIDAR meaurements of dust opacity in the Earth's atmosphere. 332 Here, MCS dust retrievals are reported as dust opacities at atmospheric pressures typ-333 ically 1-1.5 km apart, but their intrinsic vertical resolution is about 5 km (Kleinböhl et 334 al., 2009), so the dataset oversamples the actual MCS measurements. First we integrate 335 the MCS dust retrievals vertically with a 5 km grid spacing in order to recover the ob-336 served LIDO. This ensures that the assimilated data has the same vertical resolution as 337 the actual measurements. This preserves smaller-scale vertical variability in the mod-338 elled dust profile that would be unresolved in the MCS observations. 339

The approach used here resembles the assimilation of thermal profiles into the UK-LMD MGCM (S. R. Lewis et al., 2007). First, we use the modelled dust opacities $\tau_{\text{LIDO}}(m)$ to predict the dust opacity $\tau_{\text{LIDO}}^{(back)}(i)$ within observation layer *i*. Model layers that overlap more than one observed layer are split linearly in $\ln p$ among the observed layers. Figure 3 shows an example where three model layers overlap one observed layer. In this instance the modelled dust opacity of the observed layer is

$$\tau_{\text{LIDO}}^{(\text{back})}(i) = \tau_{\text{LIDO}}(m-1) \frac{\ln[p(i)/p(m-1/2)]}{\ln[p(m-3/2)/p(m-1/2)]} + \tau_{\text{LIDO}}(m) + \tau_{\text{LIDO}}(m+1) \frac{\ln[p(m+1/2)/p(i+1)]}{\ln[p(m+1/2)/p(m+3/2)]}$$
(4)

The LIDO increment within this observation layer is then

$$\Delta \tau_{\rm LIDO}(i) = \tau_{\rm LIDO}(i) - \tau_{\rm LIDO}^{(back)}(i) \tag{5}$$



Figure 3: Schematic showing the calculation of the dust opacity increment $\Delta \tau_{\text{LIDO}}(i)$ within observation layer *i*, given the observed LIDO within that layer $\tau_{\text{LIDO}}(i)$ (right) and the modelled dust opacities $\tau_{\text{LIDO}}(m)$ at the overlapping model levels (left).

From which the LIDO increment at each model layer due to observation layer i is

$$\Delta \tau_{\rm LIDO}(m-1) = \frac{\ln[p(i)/p(m-1/2)]}{\ln[p(i)/p(i+1)]} \Delta \tau_{\rm LIDO}(i)$$
(6)

$$\Delta \tau_{\text{LIDO}}(m) = \frac{\ln[p(m-1/2)/p(m+1/2)]}{\ln[p(i)/p(i+1)]} \Delta \tau_{\text{LIDO}}(i)$$
(7)

$$\Delta \tau_{\rm LIDO}(m+1) = \frac{\ln[p(m+1/2)/p(i+1)]}{\ln[p(i)/p(i+1)]} \Delta \tau_{\rm LIDO}(i)$$
(8)

The horizontal assimilation then uses these increments to update the modelled dust field, following the standard procedure for $\tau_{\rm ref}$.

³⁴⁹ Dust is transported in terms of dust mass mixing ratio, so the assimilation needs ³⁵⁰ to correct this quantity. As dust mass mixing ratio is proportional to LIDO, it is mul-³⁵¹ tiplied by a factor $\eta = \tau'_{\text{LIDO}}(\mathbf{x}, m)/\tau_{\text{LIDO}}(\mathbf{x}, m)$, where the primed and unprimed quan-³⁵² titles are the corrected and uncorrected LIDO values.

4.3 Joint CIDO and LIDO

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To take advantage of both datasets simultaneously, we can assimilate both CIDO and 354 LIDO together. In principle, one could use the measured dust profiles to correct the dust 355 in model layers where there are observations, then use the CIDO data to correct the rest 356 of the column. This avoids unnecessarily adjusting the vertical distribution using CIDO 357 when part of the distribution has already been corrected using LIDO data. However, THEMIS 358 and MCS measurements are not normally taken at the same time and place, so it is dif-359 ficult to use both simultaneously at one location. Therefore we instead assimilate the 360 LIDO and CIDO datasets independently. Testing suggested marginally better results when 361 CIDO was assimilated first, so in the joint assimilation temperature is assimilated first, 362 followed by CIDO, and finally LIDO. 363



Figure 4: Global mean $\tau_{\rm ref}$ (rescaled to 610 Pa) over the study period for the free-running simulation (green), the CIDO-only reanalysis (red), LIDO-only reanalysis (cyan), and joint CIDO/LIDO reanalysis (magenta). THEMIS CIDO observations used in the assimilation are shown as triangles. Each point is an average over five sols.

³⁶⁴ 5 Verification I.: in-sample observations and free-running model

The methods described above were used to analyse various combinations of THEMIS and MCS observations obtained during Mars Years 28 and 29, representing a typical pair of years that include dusty seasons both with and without a planet encircling event. In this section we present results that compare assimilated analyses with a free-running model simulation with full dust transport and seasonal variability and evaluate the convergence of the assimilation towards the input data. Further results and figures can be found in Section S.1 of the Supplementary Material.

The model-predicted and assimilated values of $\tau_{\rm ref}$ (rescaled to 610 Pa) from each 372 variant of the scheme were interpolated to the positions of THEMIS CIDO measurements. 373 Figure 4 (red line) shows the global mean $\tau_{\rm ref}$ of these interpolated data over the course 374 of the study period. This shows that all three of the reanalyses converge to the assim-375 ilated THEMIS data outside the GDS period in MY28, but in contrast to the other vari-376 ants, the LIDO-only assimilation overestimates the peak $\tau_{\rm ref}$ during the GDS and mis-377 represents the timing of its onset. The free-running simulation, on the other hand, cap-378 tures some of the variability, but completely misses the development of the GDS around 379 MY28 $L_s = 300^{\circ}$. 380

5.1 Assimilating CIDO only vs. MCS dust observations

It is also useful to compare the dust reanalysis with the observed time-zonal mean dust distribution. A set of vertical dust distributions (MCS-binned observations hereafter) were produced by sampling MCS dust profiles in 5° horizontal grids during daytime (local time 06:00–18:00) and nighttime (local time 18:00–06:00), after binning the data in $L_s = 5^\circ$ intervals. The model results were interpolated to the same grid and averaged over the same L_s time windows, and restricted to altitudes where MCS-binned observations were available before taking the zonal mean.

A comparison is shown in Fig. 5. With CIDO assimilation, the top of the dust layer 389 when detached dust layers are absent is broadly similar to observations, as is the free-390 running model (Fig. 5a middle frame). However, elevated detached dust layers (Fig. 5b) 391 cannot be reproduced in either the CIDO-only reanalysis or the free-running model. Such 392 detached dust layers were observed in MCS night-time retrievals (Heavens et al., 2011) 393 and later confirmed by other instruments (M. D. Smith et al., 2013; Guzewich et al., 2013). 394 In the version of the model used in this study, dust tends to be lifted to a lower height 395 than the observed detached dust layer, and it is then well mixed all the way to the ground. 396 Hence a successful reanalysis is likely to require assimilating vertically resolved dust mea-397 surements (i.e. LIDO) to reproduce the detached dust layers in a reanalysis. 398

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5.2 Assimilating LIDO only vs. THEMIS dust observations

In the LIDO-only assimilation, the period with a GDS is captured (Fig. 4, cyan line) despite the limited MCS coverage during that period (Fig. 1b). These observations are sufficient for the reanalysis to capture the initial condition and northern boundary condition of the GDS, but the inferred peak in $\tau_{\rm ref}$ is higher and later than in the observations.

The GDS onset occurs in the southern hemisphere, so the available observations 405 during the dust storm period capture this in the assimilation at a later time than if these 406 observations had been in the southern hemisphere. Once MCS data becomes available 407 in the southern hemisphere again during the "cleanup" of the storm, $\tau_{\rm ref}$ returns towards 408 the observed values, but again later than observed in the THEMIS data. The larger peak 409 in the LIDO-only assimilation suggests sedimentation in the model is not efficient enough 410 to remove dust transported southward from the northern boundary of the GDS into the 411 unobserved regions during the peak of the storm. 412

413 414

5.3 Joint CIDO/LIDO assimilation vs. THEMIS and MCS dust observations

Assimilating CIDO improves the dust horizontal spatial distribution, while assimilating LIDO improves the vertical distribution. By assimilating both we capture features such as the inter-annual variability of the global mean τ_{ref} in the THEMIS observations between GDS and non-GDS years (Fig. 4, magenta line). The global mean τ_{ref} is reproduced during the MY28 GDS, as in the CIDO-only assimilation (red line), and unlike the LIDO-only assimilation (cyan line).

⁴²¹ During the "quiet" dust season ($L_s = 0^\circ - 180^\circ$), jointly assimilating CIDO and ⁴²² LIDO gives a reasonable agreement with THEMIS observations (Fig. 4, magenta line). ⁴²³ Overestimates in the LIDO-only assimilation during MY29 $L_s = 0^\circ - 90^\circ$ were reduced ⁴²⁴ by assimilating CIDO as well. $\tau_{\rm ref}$ is slightly overestimated by the joint assimilation where ⁴²⁵ $\tau_{\rm ref} < 0.3$ and slightly underestimated where $\tau_{\rm ref} > 2$.

In the zonal-time mean dust opacity profile compared with MCS binned observations centered at $L_s = 122.5^{\circ}$ (Fig. 5b, bottom panel), the joint assimilation of CIDO



(a) $L_s = 352.5^{\circ}$, without detached dust layers. From top: MCS observations, free-running model, CIDOonly reanalysis.



(b) $L_s = 122.5^{\circ}$, with detached dust layers. From top: MCS observations, free-running model, CIDO-only reanalysis, LIDO-only reanalysis, joint CIDO/LIDO reanalysis.

Figure 5: Night-time (18:00-06:00 local time) zonal-time mean dust opacity (km⁻¹) during MY28 (a) without and (b) with detached dust layers. Time averages are over $5^{\circ}L_s$.



Figure 6: Seasonal evolution of the zonal mean $\tau_{\rm ref}$ (CIDO rescaled to 610 Pa), (a) in the free-running simulation, (b) in the joint CIDO/LIDO reanalysis. Because dust is generally well mixed in the lower atmosphere, and hence varies strongly with surface pressure, CIDO is rescaled to the 610 Pa pressure surface to account for Mars' topography.

and LIDO produces very similar results to the assimilation of LIDO-only (Fig. 5b, 4th panel).

430 5.4 Comparison between free-running model and joint CIDO/LIDO re-431 analysis

Figure 6 shows zonal mean column dust opacity $\tau_{\rm ref}$ in the free-running simulation 432 and the joint CIDO/LIDO reanalysis over MY28–29. Within each Martian year, the sea-433 sonal variability of the dust opacity in both free-running simulation and reanalysis ex-434 hibits at least some features that are generally consistent with spacecraft observations 435 (M. D. Smith, 2008). Global dust opacity is higher during the second half of the year, 436 and relatively quiet during the first half of the year. In the free-running simulation (Fig. 6a), 437 however, the active dust period in each MY lasts longer than in observations (M. D. Smith, 438 2009), while the dust opacity in the reanalysis (Fig. 6b) shows more realistically inter-439 mittent seasonal variability within each dusty season. The peak in dust opacity is also 440 sharper in L_s in the reanalysis, and tends to shut down prior to the decline in solar forc-441 ing that occurs towards the end of northern winter. 442

The interannual variability in the reanalysis (Fig. 6b) is essentially the same as in observations (e.g. M. D. Smith (2008, Fig. 8a), M. D. Smith (2009, Fig. 6) and Montabone et al. (2015)). Global dust storms (GDS) do not happen every MY, but the reanalysis successfully reproduces the observed GDS around MY28 $L_s \approx 265^\circ - 310^\circ$. The initi-



Figure 7: As Fig. 1, but for out-of-sample dust retrievals.

ation and duration of the GDS in the reanalysis are also consistent with THEMIS dust
retrievals (M. D. Smith, 2009, Fig. 6 upper panel).

The mild dusty season in MY29 following the MY28 GDS also suggests a more realistic interannual variability in the renanalysis. The free-running simulation displays some variability, with a slightly stronger dusty season in MY28 than in MY29 (Fig. 6a), but remains some considerable way away from the observations.

⁴⁵³ 6 Verification II.: independent, non-assimilated observations

In this section, the reanalysis from the new dust assimilation system is validated 454 against non-assimilated data, including the independent upward-looking measurements 455 from the Mars Exploration Rovers (MER) "Spirit" and "Opportunity" (Bell III et al., 456 2003; Lemmon et al., 2004). In order to have a more comprehensive validation, about 10% of the THEMIS and MCS data were withheld from the assimilation, and they are 458 also used as an out-of-sample validation. Those withheld data were selected from 1 in 459 every 10 of the data (for THEMIS) and of the profiles (for MCS). It would not be sur-460 prising to see that these selected THEMIS and MCS data for validation may have cor-461 relation with the data assimilated into the model, and this, to some degree, compromises 462 their application to validate the reanalysis. The completely independent datasets from 463 the MER landers, however, provide a complementary way of validation that does not suf-464 fer from these correlations. Hereafter, the reanalysis/assimilation used refers to the joint 465 assimilation of CIDO and LIDO, as described in section 4.3. 466

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6.1 Out-of-sample THEMIS dust observations

The distribution of non-assimilated (out-of-sample) THEMIS CIDO retrievals is shown in Fig. 7a. Coverage is similar to the full THEMIS dataset (Fig. 1a), and while it has only 10% of the data points, its distribution in latitude and L_s reflects the spread in the full THEMIS dataset. To compare with the out-of-sample THEMIS data, both the observations and model results were rescaled to the 610 Pa pressure level to account for Mars' topography, and the model results were interpolated both horizontally and in time to the out-of-sample data points.



Figure 8: 5-sol global mean $\tau_{\rm ref}$ over MY28–29 showing the free-running simulation (green), joint CIDO/LIDO reanalysis (magenta), and out-of-sample THEMIS observations (triangles). Compare Fig. 4 for the in-sample observations.

Figure 8 shows the comparison with the global mean $\tau_{\rm ref}$ using out-of-sample observations. The free-running simulation tracks the observations up to $L_s = 240^{\circ}$ of MY28, but fails to capture the subsequent GDS. It does predict a marginally milder dust season in MY29 compared to MY28, as observed, but does not reproduce the observed dust in either case.

The reanalysis performs significantly better, capturing the MY28 GDS as well as the precursor initiation events and subsequent decay, and the interannual variability during MY29's dusty season. The magnitudes in the reanalysis are also more consistent with observations than the free-running model, although the maximum τ_{ref} during the MY28 GDS is still lower than observations. Measurement uncertainties in the THEMIS data may be 20% or higher, however (M. D. Smith, 2004), so the reanalysis could still be broadly consistent with the observations at the peak of the GDS.

⁴⁸⁷ During the "quiet" season, the free-running model predictions generally fit the THEMIS
⁴⁸⁸ data well, especially during MY29, at least in a global average sense. During this period
⁴⁸⁹ both free-running model and observations fall within the minimum observational uncer⁴⁹⁰ tainty, which is 0.104 for the visible extinction opacity (M. D. Smith, 2009).

⁴⁹¹ Correlations between τ_{ref} in the out-of-sample THEMIS observations and the free-⁴⁹² running simulation and reanalysis are shown in Fig. 9. As with the in-sample observa-⁴⁹³ tions, the free-running simulation generally underestimates dust loading, mainly in the ⁴⁹⁴ dusty season. The reanalysis produces significantly better correlations with out-of-sample



Figure 9: Scatter plots showing individual $\tau_{\rm ref}$ points comparing the out of sample THEMIS observations with the free-running model and various reanalyses over the period shown in Fig. 4. Colours show the data density as the number of points per square of side $\tau_{\rm ref} = 0.05$. Red lines show the linear least square fit, with *m* the fitting coefficient, r^2 the coefficient of determination, and *err* the standard error in *m*. Black lines show m = 1. (a) Free-running simulation, (b) joint CIDO/LIDO reanalysis.

⁴⁹⁵ THEMIS observations. $\tau_{\rm ref}$ in the reanalysis only slightly overestimates observations where ⁴⁹⁶ $\tau_{\rm ref} < 0.5$, and slightly underestimates them where $\tau_{\rm ref} > 2$.

6.2 Out-of-sample MCS dust observations

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Figure 7b shows the distribution of out-of-sample MCS dust profiles. As with THEMIS there is a similar distribution pattern to the full dataset (Fig. 1b). The model results were averaged over several pseudo-height ranges (0–10 km, 10–20 km, 20–30 km, 30–40 km, and 40–80 km), assuming a 10 km scale height and a 610 Pasurface pressure. Within each pseudo-height range, the mean difference in dust opacity between the out-of-sample data and the joint CIDO/LIDO reanalysis was calculated for several latitude bands and is shown in Fig. 10 for the "quiet" and "dusty" seasons.

During the "quiet" season (Fig. 10a), in southern high latitudes $(90 \,^\circ\text{S} - 50 \,^\circ\text{S})$ 505 the free-running model significantly underestimates the dust opacity below 30 km, while 506 the reanalysis reproduces the observations significantly better. In the midlatitudes of both 507 hemispheres (50 $^{\circ}$ S - 15 $^{\circ}$ S and 15 $^{\circ}$ N - 50 $^{\circ}$ N) the free-running simulation again un-508 derestimates the observations, although this underestimate decreases at higher altitude. 509 The reanalysis generally falls within or close to the observational uncertainty, with the 510 largest differences at 30–40 km. In the tropics $(15 \,^{\circ}\text{S} - 15 \,^{\circ}\text{N})$, the free-running simu-511 lation generally underestimates the dust opacity. The reanalysis errors are generally larger 512 than the MCS observational uncertainties except for 0-10 km. It is an improvement over 513 the free-running simulation for 10-20 km, 20-30 km, and 40-80 km, but overestimates the 514 dust opacity for 0–10 km and 30–40 km, with absolute differences larger than those in 515 the free-running simulation. In northern high latitudes $(50 \text{ }^{\circ}\text{N} - 90 \text{ }^{\circ}\text{N})$ the uncertain-516 ties in the MCS observations are about 50% larger than elsewhere. The free-running model 517 falls within the observational uncertainties at all altitudes. The renanalysis also falls within 518 observational uncertainty except below 10 km. 519



(b) "Dusty" season $(L_s = 180^\circ - 360^\circ)$.

Figure 10: Mean difference in dust opacity between the out-of-sample MCS observations and the free-running simulation (green) and joint CIDO/LIDO reanalysis (magenta). In each of the two seasons, the globe is split into latitude bands: (a) $90 \,^{\circ}\text{S} - 50 \,^{\circ}\text{S}$, (b) $50 \,^{\circ}\text{S} - 15 \,^{\circ}\text{S}$, (c) $15 \,^{\circ}\text{S} - 15 \,^{\circ}\text{N}$, (d) $15 \,^{\circ}\text{N} - 50 \,^{\circ}\text{N}$, and (e) $50 \,^{\circ}\text{N} - 90 \,^{\circ}\text{N}$. Grey dashed lines are the average uncertainties in the MCS observations.

Figure 10b shows the same for the dusty season $(L_s = 180^\circ - 360^\circ)$. In general, the reanalysis agrees better with the MCS observations than the free-running simulation. The maximum error in the free-running simulation increases from south to north, which may be due to the difficulty in predicting frontal dust storms in northern high latitudes during dusty season.

In southern high latitudes, the free-running simulation tends to underestimate the 525 dust opacity above 10 km, and the reanalysis tends to overestimate the dust opacity be-526 low 30 km, but fall within the observational uncertainty above 30 km. In southern mid-527 latitudes, the reanalysis is closer to the observations than in the tropics and northern 528 midlatitudes, and at 10-20 km, 20-30 km, and 40-80 km is within or close to observa-529 tional uncertainty. In the tropics, the free-running simulation is similar to the northern 530 middle latitudes, except below 10 km where it slightly overestimates the dust opacity. 531 The reanalysis tends to underestimate the dust opacity between 10 and 30 km, and over-532 estimate the dust opacity above 30 km. In the northern midlatitudes, the reanalysis un-533 derestimates the dust opacity below 30 km, but the differences are still smaller than the 534 free-running simulation. Above 30 km, the reanalysis overestimates the dust opacity with 535 differences larger than the free-running simulation. In northern high latitudes, the re-536 analysis falls either within or close to the MCS observational uncertainties, with lower 537 differences at higher altitudes. 538

6.3 Independent Pancam observations from Spirit

The reanalysis and free-running simulation were compared with Spirit and Opportunity observations by interpolating τ_{ref} horizontally and in pressure to the rover locations at Gusev Crater and Meridiani Planum respectively. Figure 11 shows these values at the locations of the two rovers during MY28–29.

Figure 11a shows τ_{ref} at the Spirit rover site. During the relatively "quiet" dust season Spirit Pancam observations are normally below $\tau_{ref} = 0.3$. Although the freerunning simulation agreed well with THEMIS τ_{ref} observations globally during this season (see Fig. 8), at the Spirit landing site it generally underestimates the dust loading. τ_{ref} only reaches ~0.1 during the "quiet" season at this location. During the dusty season the free-running simulation suggests an increase in τ_{ref} at Spirit's location, but its increase does not match the increase in the observations.

Conversely, the reanalysis agreed better with the Spirit Pancam data. It captured 551 the annual and interannual variability in the data well. During the "quiet" season, the 552 reanalysis reproduces the magnitude and variation in τ_{ref} at the Spirit landing site. Un-553 derestimates are mainly during MY29 $L_s = 60^{\circ} - 120^{\circ}$. The reanalysis captures the 554 increase of dust loadings during MY29 $L_s = 140^{\circ} - 160^{\circ}$, but not the peak $\tau_{\rm ref}$. The 555 MY28 GDS is reflected by an increase in $\tau_{\rm ref}$ to 3.5 at the Spirit landing site. The re-556 analysis reproduces the initiation and decay of the MY28 GDS, and also the variabil-557 ity of dust loading during MY29. Nevertheless, the reanalysis dust loading during the 558 first peak in MY29 ($L_s \approx 160^\circ$) still does not reach the maximum observed by Spirit. 559 It is worth noting that, although the free-running simulation fails to produce the observed 560 amount of dust in both dusty seasons, it does exhibit some interannual variability. 561

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6.4 Independent Pancam observations from Opportunity

Figure 11b shows τ_{ref} at the Opportunity rover site compared with the free-running model and reanalysis. The observed dust loading at Meridiani Planum has a similar evolution to that seen at Gusev Crater, but with slightly higher values in general. Both the freerunning simulation and the reanalysis underestimate the peak of the MY28 GDS, though τ_{ref} in the reanalysis is much closer to the measurements. The reanalysis also better reproduces the observed variability of dust loading during both dusty seasons, but both the free-running simulation and the reanalysis underestimate the dust loading during the "quiet" season.

A similar discrepancy was also noticed in earlier studies when comparing datasets 571 from different instruments. Montabone et al. (2015) found a systematic underestimate 572 of dust opacity over the Opportunity landing site in Meridiani Planum in both TES and 573 THEMIS datasets starting at the spring equinox, up to a factor ~ 2 during northern sum-574 mer, which may have been linked to the likely presence of clouds. Since the THEMIS 575 data assimilated in this study falls within the same period, it is not surprising to find 576 577 a similar discrepancy in our reanalysis compared with the Opportunity data. Lemmon et al. (2015) also raised problems with Opportunity 880 nm data around $L_s = 30^{\circ} -$ 578 130°. The source of this discrepancy remains an open problem. 579

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7 Validation of the Mars Climate Database against the reanalysis

Reanalyses produced by data assimilation are important for verifying models against "reality" or "ground truth". If the reanalysis is of sufficiently high quality, then it can be used as a surrogate for the real atmosphere when validating and verifying model output (e.g. ERA-40, Uppala et al., 2005, for the Earth's atmosphere). We used the reanalysis described in previous sections to verify the Mars Climate Database v5.2 (MCD, S. Lewis et al., 1999; Millour et al., 2015) against our "real" atmosphere.

The quantities of interest produced in both the reanalysis and the MCD are sur-587 face pressure, surface temperature, air temperature, density, and zonal and meridional 588 velocities. We did not compare dust diagnostics, because the reanalysis dust distribu-589 tion is likely to have more high frequency variability than the MCD, due to the 15-minute 590 timescale over which observations are assimilated; these short time periods are not rep-591 resented in the MCD, which is forced by dust fields changing over the timescale of a day. 592 There are also significant differences between the way the dust is treated in the MCD 593 and in the GCM used to construct the reanalysis, such as the GCM used here assumes 594 a fixed dust particle radius of 1.5 μ m, while the MCD uses a two-moment scheme which 595 retains information about the full dust particle size distribution (Madeleine et al., 2011). 596 This deficiency of our model will be improved in the future. 597

For each 30° L_s period in the reanalysis (MY28 $L_s = 120^{\circ}$ to the end of MY29) we computed monthly means and day-to-day variability in the same way as the MCD. First we interpolated horizontally from the MCD grid ($5.625^{\circ} \times 3.75^{\circ}$) to the reanalysis grid ($5^{\circ} \times 5^{\circ}$). Then we interpolated atmospheric quantities linearly in log p to 30 fixed pressure levels spaced by 2.5 km up to 40 km pseudo-altitude above a reference pressure of 610 Pa (assuming a scale height of 10 km), and spaced by 5 km above that.

The monthly mean for variable X at (longitude, latitude, pressure) position (i, j, k)

is

$$\overline{X_{ijk}} = \frac{1}{N} \sum_{t=1}^{N} X_{ijk,t}$$
(9)

where t = 1...N includes all times within a 30° L_s period. Since the orbit is elliptical, N varies with season. The day-to-day variability is (Forget et al., 2015, Eq. 5):

$$\operatorname{Var}(X_{ijk}) = \sqrt{\frac{1}{N} \sum_{t=1}^{N} \left(\overline{X_{ijk,t}^{1\,\mathrm{sol}}} - \overline{X_{ijk,t}^{10\,\mathrm{sols}}}\right)^2} \tag{10}$$

where $\overline{X_{ijk,t}^{1 \text{ sol}}}$ and $\overline{X_{ijk,t}^{10 \text{ sols}}}$ are running means over 1 sol $(t \pm 0.5 \text{ sols})$ and 10 sols $(t \pm 5 \text{ sols})$ respectively. The day-by-day variability removes the diurnal cycle, along with any long-term trend, leaving the variability associated with the day-to-day "weather". For atmospheric quantities we calculated zonal-monthly means at (latitude, pressure) points



(a) Spirit, at 14.57 $^{\circ}\mathrm{S},$ 175.48 $^{\circ}\mathrm{E}$ (Gusev Crater).



(b) Opportunity, at 1.95 °S, 5.53 °W (Meridiani Planum).

Figure 11: τ_{ref} from the reanalysis (magenta) and free-running model (green) compared with (a) Spirit and (b) Opportunity Pancam observations. Each point is averaged over one sol. -21-



Figure 12: Differences between reanalysis means and MCD means for the period $L_s = 180 - 360^{\circ}$ in MY29, separated into 30 L_s segments. Positive means the reanalysis value is larger than the MCD value. Surface quantities are monthly means and atmospheric quantities are zonal-monthly means. Grey shows points with missing data (where all points along a latitude circle are below the surface). Black lines show orography at intervals of 4 km between -4 km and +20 km above the geoid (negative values are dashed).

(j,k):

$$\overline{X_{jk}} = \frac{1}{I} \sum_{i=1}^{I} \overline{X_{ijk}}$$
(11)

where I is the number of longitude points above the surface. To compute zonal-monthly means for the day-to-day variability, we computed the root-mean-square of the day-today variability along each latitude circle. Note this gives us a measure of the day-to-day variability at points along each latitude circle, rather than the spread of values along the latitude circle.

The MCD contains day-to-day variability as a function of position for each month and dust scenario, and monthly means as a function of local time of day at zero longitude. To obtain equivalent monthly means for comparison with the reanalysis, we averaged over all local times of day after interpolating each column to the required pressure levels.

Figures 12 and 13 summarise the differences between our reanalysis and the MCD 614 during the second half of the Martian year. Because this paper focuses on dust, we con-615 centrate on the dusty season from $L_s = 180 - 360^{\circ}$. The patterns of mean and vari-616 ability differences between the reanalysis and MCD were generally similar in all seasons 617 (with one exception, discussed below) and for all quantities when comparing MY28 with 618 MY29, so these summary figures show only MY29. Full sets of figures for all compar-619 isons between the reanalysis and MCD for all quantities over all months analysed are in-620 cluded as Supplementary Material, Section S.2 and Figs S4–S15. 621

The surface temperature reanalysis was generally cooler than the MCD at most places and times, particularly between 45–60° latitude in each hemisphere. Exceptions were the polar regions and Tharsis, Arabia, and Elysium, which were persistently warmer in the



Figure 13: As Fig. 12, but for day-to-day variability. Where the ratio is greater than 1, the reanalysis variability is larger than the MCD variability over that period.

reanalysis. The largest differences were near the edge of the polar icecap. The latitude 625 at the edge of the polar icecap has a large day-to-day variability during a given month. 626 as the edge of the polar cap moves over time, and so the day-to-day variability ampli-627 tude will transition from small (with ice) to large (without ice). This is present in both 628 the reanalysis and MCD day-to-day variability (Fig. S4). This also means the monthly 629 mean difference between the MCD and the reanalysis is sensitive to the position of the 630 edge of the polar cap, and so we see large differences between the reanalysis and the MCD 631 in that region. There is a strong warm bias at southern polar latitudes between $L_s =$ 632 $240^{\circ}-330^{\circ}$, due to the permanent CO₂ polar ice cap, which is present in the MCD but 633 is not simulated in the version of the model used in the reanalysis. 634

Differences between the MCD and reanalysis surface pressure vary considerably with 635 season. The surface pressure is persistently higher in the reanalysis in the summer hemi-636 sphere, and lower in the reanalysis in the winter hemisphere. There are also rings in the 637 difference maps around regions with the most extreme elevation changes on the planet, 638 particularly Olympus Mons, Elysium Mons, and Hellas (Fig. 12). The most likely rea-639 son is that the UK version of the Mars GCM (reanalysis) uses a spectral dynamical core, 640 while the LMD Mars GCM (MCD) uses a grid point dynamical core. Quantities sensi-641 tive to surface elevation, such as surface pressure, will have large differences purely as 642 a result of the topography being represented differently in the two models. The rings them-643 selves are characteristic of the Gibbs phenomena that occur when a step function is spec-644 trally decomposed, and are therefore likely to be spurious. 645

In the equatorial region the atmospheric temperature reanalysis is typically cooler than the MCD close to the surface, warmer around 100 Pa, cooler between 1 and 10 Pa, and warmer above 1 Pa. This pattern is repeated in most months. Both poles are typically warmer in the reanalysis than in the MCD, at least in the lower atmosphere, with a warm "tongue" in the difference maps extending into the stratospheric polar region in the winter hemisphere. There is significantly more day-to-day variability in the reanalysis near the surface at the winter pole.

Density differences are generally small (up to 0.2 in $\log_{10} \rho$), but with some pat-653 terns. At low latitudes the density is generally lower in the reanalysis below 10 Pa, and 654 higher in the reanalysis above this level. In the polar regions the density is nearly always 655 lower in the reanalysis in the winter hemisphere. The day-to-day variability is lower in 656 the reanalysis where the "warm tongue" appears in the air temperature maps. The hor-657 izontal striations in Fig. 12 are likely to be artifacts: there are different vertical grids in 658 the reanalysis and MCD, which are interpolated to a common pressure grid for compar-659 ison, and density covers several orders of magnitude, so differences in interpolation will 660 be magnified. 661

Zonal velocities in the reanalysis are, in general, more westward than in the MCD. 662 This means that the eastward mid-latitude jets, the most prominent features of the monthly 663 means (Fig. S12), are weaker in the reanalysis. These differences can be quite large – up 664 to 50 m s^{-1} in magnitude. At low altitudes near the equator, the zonal flow is more east-665 ward in the reanalysis than in the MCD. The day-to-day variability is generally lower 666 in the reanalysis in the winter hemisphere, and larger in the reanalysis in the summer 667 hemisphere (Fig. 13), with the exception of the winter pole near the surface, which has 668 a high day-to-day variability in the reanalysis. 669

In general the meridional velocity reanalysis has a stronger upper-level equatorial and midlatitude meridional circulation than does the MCD: flow away from the sub-solar point is strengthened in the reanalysis between 0.1–1 Pa during the dusty season, compared with the MCD. Like the zonal velocity, typically the day-to-day variability is generally lower in the reanalysis in the winter hemisphere, and larger in the reanalysis in the summer hemisphere.

The only exception to the similar results for MY28 and MY29 was $L_s = 270 -$ 676 300° , which contains the build up to and peak of the MY28 global dust storm. Figure 14 677 shows the differences between the reanalysis and MCD during this period for both years. 678 The main difference between the two years is that during MY28 the day-to-day variabil-679 ity is significantly larger in the reanalysis than in the MCD, when compared with the 680 corresponding period during MY29. This applies to the surface temperature (particu-681 larly near the equator), and atmospheric temperatures, density, and both zonal and merid-682 ional velocities. Throughout the reanalysis sequence the day-to-day variability in the re-683 analysis is typically 1–2 times that in the MCD. This is likely due to the greater impor-691 tance of shorter timescales in the reanalysis than in the simulations used to generate the MCD, such as a 15 minute timescale in the temperature field, as that is the interval be-686 tween successive calls to the data assimilation procedure, and this may be expected to 687 stimulate variability on timescales shorter than a day. However, during the MY28 global 688 dust storm period this difference is amplified. During MY28 there is a clear warm anomaly 689 (and corresponding low-density anomaly) between 1 and 100 Pa in the reanalysis com-690 pared with the MCD, compared with the corresponding period during MY29. A second 691 major differences is that the monthly mean meridional velocity (i.e. the cross-equatorial 692 flow) between 0.1–1.0 Pa is stronger by almost 10–20 m s⁻¹ in the reanalysis compared 693 with the MCD during the MY28 global dust storm, while during the corresponding pe-694 riod in MY29 the difference is $\pm 5 \text{ m s}^{-1}$. 695

696 8 Conclusions

The data assimilation system integral to the UK-LMD Mars GCM described by S. R. Lewis et al. (2007) has been updated. That work assimilated temperature and CIDO, prescribing a vertical dust distribution using an empirical function of height (Conrath, 1975). The new scheme adds activated dust lifting, transport, and deposition schemes, assuming a single dust particle size. It also assimilates vertically resolved dust profiles via LIDO, either instead of CIDO or in addition to it. This update has been prompted by the acquisition of vertically-resolved dust profiles by MCS on board MRO (McCleese et al., 2010).



Figure 14: Reanalysis vs. MCD for the MY28 global dust storm period $L_s = 270 - 300^{\circ}$. Both the difference in monthly means and the ratio between day-to-day variability are shown. The corresponding period in MY29 is shown (in the same column format) in Figs 12 and 13. The units only apply to the mean differences.

When CIDO is assimilated by itself, the assimilation can reproduce the observed 704 interannual variability of the dust horizontal spatial distribution, including the gener-705 ation and dissipation of the MY28 GDS (Fig. 4, red line). This inter-annual variability 706 cannot be reproduced by our free-running model, although some degree of tuning can 707 be done to reproduce the "quiet" $(L_s = 0^\circ - 180^\circ)$ and "dusty" $(L_s = 180^\circ - 360^\circ)$ 708 periods during a single year. Even without assimilating any vertical information, the CIDO 709 assimilation reduces systematic errors in the model's estimate of the dust vertical dis-710 tribution. However, it misses detached dust layers that form during northern spring and 711 summer (Fig. 5b), which have been a challenge for Mars GCMs to reproduce. 712

Conversely, when LIDO is assimilated by itself the model can reproduce some fea-713 tures of the detached dust layers (Fig. 5b, 4th panel), with reasonable interannual vari-714 ability, which the free-standing model is unable to reproduce. Rafkin (2012) discussed 715 the difficulty of producing this detached dust layer in model simulations, especially in 716 a relatively coarse resolution GCM. Similar detached dust layers were reproduced in mesoscale 717 model simulations of "rocket dust storms" by Spiga et al. (2013) and have since been pa-718 rameterised in a coarse resolution GCM (C. Wang et al., 2018) with some success. Nev-719 ertheless, being able to reproduce them in a reanalysis provides a valuable alternative 720 means of investigating their observed characteristics and impact on Mars' atmospheric 721 circulation. However, the limb-viewing MCS does not continuously observe the lowest 722 part of the atmosphere, where the dust concentration is generally highest except where 723 there are detached layers. Consequently assimilating LIDO only does not reproduce the 724 observed global average $\tau_{\rm ref}$ as well as when CIDO is assimilated by itself (see Fig. 4). 725

Once combined together, the joint assimilation of CIDO and LIDO benefits from
 information about the total column dust opacity and horizontal distribution from the
 THEMIS dataset, and the vertical distribution from the MCS dataset. The evolution of
 the MY28 GDS is tracked well, and some features of the detached dust layers are also
 reproduced well.

The joint assimilation of CIDO and LIDO is a powerful tool that helps us to reconstruct the Martian climate as well as individual dust events. For example, it is difficult to retrieve a complete vertical dust distribution from MCS measurements during the MY28 GDS (Fig. 1b). Using the joint assimilation, CIDO provides information to constrain the model where vertical profiles are sparse, so assimilation can map the fourdimensional dust distribution during the MY28 GDS in the absence of complete observations.

The reanalysis was validated against out-of-sample THEMIS and MCS dust ob-738 servations, as well as upward-looking MER Pancam observations at 880 nm. The reanal-739 ysis successfully reproduced the observed interannual and intraseasonal variability in the 740 original THEMIS data, and generally improved the representation of the dust vertical 741 distribution compared to free-running simulations. In general, the free-running model 742 tends to underestimate $\tau_{\rm ref}$, particularly during the dusty season, which can be signifi-743 cant during major dust storm events. Hence assimilating dust observations serves to greatly 744 reduce the model uncertainty during such events. Although free-running simulations were 745 able to simulate the pattern of dust loading during the quiet season, they failed to sim-746 ulate the vertical dust distribution. This is consistent with the general observation that 747 dust accumulates close to the ground, below the base of typical MCS dust profiles. 748

At the Spirit rover location the reanalysis captured the dust variability and intensity during both quiet and dusty seasons (Fig. 11a). At the Opportunity rover location, the reanalysis captured the dust variability (Fig. 11b) and improved the pattern of dust loading over the free-running simulation during the dusty season. However, there were persistent underestimates of $\tau_{\rm ref}$ by the reanalysis, particularly during northern spring and summer. This is likely to be due to a systematic disagreement between THEMIS and Opportunity data (Montabone et al., 2015).

A systematic comparison between the reanalysis and the MCD during MY28 and 756 MY29 exhibited a number of trends, although agreement between the reanalysis and the 757 MCD was generally good. Day-to-day variabilities were typically larger in the reanal-758 ysis by a small factor (1-2 times), likely associated with the stimulation of fluctuations 759 on timescales comparable with the interval between successive calls to the data assim-760 ilation procedure (15 minutes). We also found some uncertainty in the position of the 761 edge of the polar caps in the MCD, since the main differences in surface temperature arise 762 along those latitudes. The reanalysis also showed significantly warmer poles than in the 763 MCD at most times of year, along with weaker mid-latitude jets, and a stronger cross-764 equatorial flow during the MY28 global dust storm. 765

The combined CIDO-LIDO dust reanalysis significantly improves the estimation 766 of Martian horizontal and vertical dust distributions over a free-running model and over 767 CIDO or LIDO alone. The reanalysis provides a solution generally consistent with the 768 available Martian dust observations. Assimilation has considerable potential as a tool 769 for studying individual dust lifting events and for mapping Mars' three-dimensional dust 770 distribution over time. Elsewhere, we will report on two case studies using the scheme, 771 investigating a southward-moving regional dust storm during MY29, and the global dust 772 storm during MY28. 773

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Supporting Information for "Assimilation of both column- and layer-integrated dust opacity observations in the Martian atmosphere"

Tao Ruan¹, R. M. B. Young^{1,2}, S. R. Lewis³, L. Montabone^{1,4},

A. Valeanu¹ and P. L. $Read^1$

¹Atmospheric, Oceanic and Planetary Physics, Department of Physics, University of Oxford, Clarendon Laboratory, Parks Road,

Oxford, OX1 3PU, UK

²Department of Physics & National Space Science and Technology Center, UAE University, Al Ain, United Arab Emirates

 $^3{\rm School}$ of Physical Sciences, The Open University, Walton Hall, Milton Keynes, MK7 6AA, UK

 $^4\mathrm{Space}$ Science Institute, Boulder, CO 80301, USA

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- 1. S1. Additional figures for Verification against in-sample observations
- 2. S2. Additional figures for Reanalysis vs. MCD comparisons
- 3. Figures S1 to S15

Introduction In the following two sections we include additional figures that supplement the information provided in the main text. Section S1 presents additional figures showing the performance of the reanalyses against in-sample observations. Section S2 presents additional figures that provide in-depth comparisons between various climatological fields from the reanalysis and diagnostics obtained from the Mars Climate Database version 5.2.

S1. Verification against in-sample observations

The methods described in the main text were used to analyse various combinations of THEMIS and MCS observations obtained during Mars Years 28 and 29, representing a typical pair of years that include dusty seasons both with and without a planet encircling event. In this section we present results that evaluate the convergence of the assimilation towards the input data.

The model temperatures (either from the free-running model or the assimilated model state) were interpolated to the temperature retrieval pressure levels, before being converted to a global mean difference between model and observations. Figure S1 shows this difference averaged over several pseudo-height ranges (0–10 km, 10–20 km, 20–30 km, 30–40 km, and 40–80 km), assuming a 10 km scale height and a 610 Pasurface pressure. MCS retrievals are reported at pressure levels separated by around 1–1.5 km, while the true instrumental resolution is about 5 km (Kleinböhl et al., 2009), so this grouping of vertical levels smoothes the oversampled MCS retrievals over a distance larger than the true observational vertical resolution. Figure S2(a,b) shows the correlation between the assimilated THEMIS observations and the CIDO-only reanalysis and free-running model. Figure S3 shows the global-time mean difference in dust opacity between the in-sample observations and both the reanalysis and free-running model (see main text).

S2. Reanalysis vs. MCD comparisons Figures S6–S15 contain the complete set of comparisons made between the reanalysis and the Mars Climate Database. The following

quantities are included: surface temperature (Figs S4-S5), surface pressure (Figs S6-S7), atmospheric temperature (Figs S8-S9), density (Figs S10-S11), zonal velocity (Figs S12-S13), and meridional velocity (Figs S14-S15).

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Surface quantities are monthly means (i.e. over 30 L_s) and atmospheric quantities are zonal-monthly means. Comparisons are made each month between MY28 $L_s = 120 -$ 150° and MY29 $L_s = 330 - 360°$ (20 months in total). Each plot shows the reanalysis monthly mean, MCD monthly mean, monthly mean difference (reanalysis minus MCD), reanalysis day-to-day variability, MCD day-to-day variability, and day-to-day variability ratio (reanalysis : MCD). The layout and format of each figure is the same.

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Figure S1. Global-time mean temperature difference between MCS observations and the freerunning model (green) and CIDO-only reanalysis (red), during MY29 $L_s = 110^{\circ} - 330^{\circ}$. Grey dashed lines show the average uncertainty in the MCS observations (Kleinböhl et al., 2009).





Figure S2. Scatter plots showing individual τ_{ref} points comparing the assimilated THEMIS observations with the free-running model and various reanalyses over the period shown in Fig. 4 of the main text. Colours show the data density as the number of points per square of side $\tau_{ref} = 0.05$. Red lines show the linear least square fit, with m the fitting coefficient, r^2 the coefficient of determination, and *err* the standard error in m. Black lines show m = 1.



Figure S3. Global-time mean dust opacity difference between observed MCS dust opacities and the free-running model (green), CIDO-only reanalysis (red) LIDO-only reanalysis (cyan), and joint CIDO/LIDO reanalysis (magenta), during MY29 $L_s = 110^{\circ} - 330^{\circ}$. Grey dashed lines show the average error in the MCS observations.



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Figure S4. Surface temperature during MY28.



Figure S5. Surface temperature during MY29.



Figure S6. Surface pressure during MY28.



Figure S7. Surface pressure during MY29.



Figure S8. Zonal mean temperature during MY28.



Figure S9. Zonal mean temperature during MY29.



Figure S10. Zonal mean density during MY28.

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Figure S11. Zonal mean density during MY29.



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Figure S12. Zonal mean zonal velocity during MY28.



Figure S13. Zonal mean zonal velocity during MY29.



Figure S14. Zonal mean meridional velocity during MY28.

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Figure S15. Zonal mean meridional velocity during MY29.