# Modeling and validating a SuperDARN radar's power density profile

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November 28, 2022

#### Abstract

We have developed a model that simulates the power density profile of the Saskatoon Super Dual Auroral Radar Network (SuperDARN) radar at ionospheric altitudes. The model uses ray tracing software to project the radar system's vacuum power profile to ionospheric altitudes, taking into account the influence of the ionospheric medium on the propagation characteristics of the High Frequency radio waves. Measurements of the radar's transmissions by the Radio Receiver Instrument (RRI) in low-Earth orbit are used to validate the model during five experiments which occurred between August 4-8, 2017. Comparisons between simulated and measured RRI antenna voltages show good agreement, although there are clear instances in which the model underperforms. Nevertheless, the model demonstrates its utility as a tool for interpreting several RRI measurements of SuperDARN radars. The model also helps address a lack of knowledge of a SuperDARN radar's power profile at ionospheric altitudes. In particular, we assess the assumption that SuperDARN's scattering volume lies along the great-circle path of the transmitting beam's bearing. Comparisons between the model and RRI's measurements show that this assumption is reasonable for the five experiments investigated in this work. The model presents a new way of carrying out SuperDARN and HF radio science investigations.

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9	Key Points:
10	• SuperDARN Saskatoon's power density profile is modeled in an empirical ionosphere
11	• A simulated power density profile is validated with the e-POP RRI instrument
12 13 14	• The power density model is a useful tool for studying High Frequency radio wave propagation

#### 15 Abstract

16 We have developed a model that simulates the power density profile of the Saskatoon Super

17 Dual Auroral Radar Network (SuperDARN) radar at ionospheric altitudes. The model uses ray

- 18 tracing software to project the radar system's vacuum power profile to ionospheric altitudes,
- 19 taking into account the influence of the ionospheric medium on the propagation characteristics of
- the High Frequency radio waves. Measurements of the radar's transmissions by the Radio
- 21 Receiver Instrument (RRI) in low-Earth orbit are used to validate the model during five
- experiments which occurred between August 4-8, 2017. Comparisons between simulated and measured RRI antenna voltages show good agreement, although there are clear instances in
- 24 which the model underperforms. Nevertheless, the model demonstrates its utility as a tool for
- interpreting several RRI measurements of SuperDARN radars. The model also helps address a
- 26 lack of knowledge of a SuperDARN radar's power profile at ionospheric altitudes. In particular,
- 27 we assess the assumption that SuperDARN's scattering volume lies along the great-circle path of
- the transmitting beam's bearing. Comparisons between the model and RRI's measurements
- 29 show that this assumption is reasonable for the five experiments investigated in this work. The
- 30 model presents a new way of carrying out SuperDARN and HF radio science investigations.

# 31 Plain Language Summary

- 32 The Super Dual Auroral Radar Network (SuperDARN) radar system located in Saskatoon,
- 33 Canada is a powerful tool for studying the high-latitude ionosphere. Despite the fact that the
- 34 system has been operating nearly continuously since 1993, there is very little understanding of
- 35 what the radar's power profile looks like at ionospheric altitudes where its radar echoes
- 36 originate. This is simply due to a lack of available measurements. As a result, it is customary to
- 37 simply assume that the radar's power profile at ionospheric altitudes resembles that of its power
- profile in a vacuum projected along great-circle trajectories up to ionospheric altitudes. This
- assumption is limited because it does not adequately account for the fact that the radar's power
- 40 profile can be influenced by the ionospheric medium in which it is immersed. In this work, we
- have developed a model of SuperDARN Saskatoon's power profile at ionospheric altitudes
  which properly accounts for the influence of the ionosphere. Furthermore, we have compared
- the model's output to several measurements of the Saskatoon radar made by a radio receiver in
- low-Earth orbit to validate the model. Our results show that the model is accurate and can be
- used as a useful tool for studying SuperDARN's power profile at ionospheric altitudes.

# 46 **1 Introduction**

# 47 1.1 Background and Motivation

- 48 The Super Dual Auroral Radar Network (Chisham et al., 2007), a globally distributed network of
- 49 High Frequency (HF; 3-30 MHz) radars, monitors plasma density irregularities in the high-
- 50 latitude ionosphere. The irregularities drift at the local plasma convection velocity (Villain et al.,
- 51 1985; Ruohoniemi, 1987). Doppler velocity measurements of the irregularities from multiple
- 52 SuperDARN systems are then combined to deduce high-latitude plasma convection flows and
- the large-scale electric fields that drive them (Ruohoniemi & Baker, 1998; Bristow et al., 2016).
- 54 These fields are generated as a result of complex magnetosphere-ionosphere-thermosphere
- 55 (MIT) coupling processes. As such, SuperDARN is a vital remote sensing tool for
- 56 undertstanding complex plasma interactions in the near-Earth geospace environment.
- 57

58 The SuperDARN technique relies on coherent backscatter generated by ionospheric irregularities

- 59 to detect and measure the bulk velocity of the ionospheric plasma. The generally accepted
- 60 theory of HF radar backscatter requires that two conditions must be satisfied (Milan et al.,
- first, plasma density irregularities must be present in the scattering volume. Second,
   the incident radio wave must satisfy the Bragg condition at that location. At the decameter
- the incident radio wave must satisfy the Bragg condition at that location. At the decameter
   scale, the wavelength of SuperDARN transmissions, cross-field plasma diffusion (with respect to
- the magnetic field) is much slower than parallel diffusion (Tsunoda, 1988). F-region plasma
- 65 density irregularities are strongly field-aligned (hence the name field-aligned-irregularities;
- 66 FAIs) with the normal of wavevectors that are orthogonal to the magnetic field. Therefore, to
- 67 satisfy the Bragg condition, an incident radio wave must have an aspect angle that is close to
- orthogonal with the magnetic field. This is commonly referred to as the "aspect angle
- 69 condition".
- 70

71 Geolocating a radar echo is an important problem in HF remote sensing systems, including

- SuperDARN. A key assumption made in SuperDARN geolocation algorithms is that the
- <sup>73</sup> scattering volume where the aspect angle condition is satisfied is located along the great-circle
- path trajectory of the transmitting beam along the radar's main lobe. There is ample evidence
- that this may not always be true. For example, echoes originating from secondary maxima in a
- <sup>76</sup> SuperDARN radar's power profiles are common (Milan et al., 1997b; Burrell et al., 2015, 2018).
- 77 Furthermore, MIT coupling processes can generate large scale plasma density irregularities along
- the transmission trajectory paths which can deflect HF transmissions laterally from their nominal
- 79 great-circle path trajectories. Observations of these deflections were reported by Perry et al.,
- 80 (2016) and Warrington et al. (1997). In general, the validity and applicability of the assumption
- under various geomagnetic conditions is not well understood. The limitations of SuperDARN
   geolocation algorithms are well documented (cf. Chisham et al., 2008; Yeoman et al., 2008),
- attesting to the challenge of HF radar echo geolocation problem.
- 84
- 34
- 85 One course of action for improving SuperDARN's geolocation algorithms is to model the
- 86 system's power density profile at ionospheric altitudes with ray tracing techniques to narrow
- 87 down the location of the scattering volume. This would address to key factors in the
- 88 SuperDARN backscattering process (assuming FAIs are present): the ray paths of radio waves
- incident on the scattering volume and their power density. Accordingly, in this work, we present
- a model of the power density profile of the SuperDARN Saskatoon radar propagated into the F-
- 91 region ionosphere. We then validate the model using data collected during conjunctions between
- 92 the radar and the Radio Receiver Instrument (RRI; James et al., 2015) onboard the Cascade,
- 93 Smallsat and Ionospheric Polar Explorer (CASSIOPE; Yau et al., 2015) spacecraft.
- 94
- 95 The model has two purposes. First, it can be used to address a gap in knowledge of
- 96 SuperDARN's power density profile at ionospheric altitudes to improve our understanding of
- how the system operates, the coherent backscatter mechanism, and improve our ability to probe
- the MIT system using HF radio techniques. In particular we explore the assumption that the
- <sup>99</sup> main lobe of SuperDARN's power density profile follows a great-circle path. Second, the model
- 100 can be used as a tool for interpreting measurements of SuperDARN transmissions acquired in
- 101 low-Earth orbit, to help distinguish variations in the received signal that are a feature of the
- 102 radar's power profile from those that may be due to geophysical phenomena.

#### 1.1 Previous HF Ray Trace Modeling Work 103

Historically, SuperDARN-related modeling efforts have focused on HF ray paths and 104

determining where the rays satisfy the aspect angle condition (e.g., André, et al. 1997; Michael et 105

al., 2020). One advantage of this technique is that it does not require one to simulate plasma 106

density irregularities – it is assumed that irregularities are present everywhere, yet it is effective 107

- 108 enough to provide context to radar data and facilitate the design and deployment of new systems.
- The technique is also consistent with existing SuperDARN geolocation algorithms in that it does 109
- not consider the radar's power profile, only the geometry of the backscattering process. 110
- 111

112 In an effort to better understand SuperDARN velocity data, Ponomarenko et al. (2009)

developed a model for simulating radar echo returns generated by FAIs. Their sophisticated 113

model considered a backscatter cross-section, higher-order corrections to the aspect angle 114

condition, and geometric optics effects, but did not take the power profile of the simulated radar 115

into account. Nevertheless, their model enabled quantitative diagnoses of SuperDARN 116

backscatter characteristics. 117

118

In preparation for RRI science operations, Gillies et al. (2010) used ray tracing to predict the 119

relative O- and X-mode power profile for the SuperDARN radars located at Saskatoon, Prince 120

121 George, and Rankin Inlet. Their ray trace modeling considered the polarization state of the

transmitted signal along the ray path. It also accounted for the initial polarization state of the 122

signal at the selected radar. It did not account for the absolute power profile of the radars; the 123 authors assumed an isotropic radiator, a suitable simplification considering that the relative

124

power between the propagation modes was the quantity of interest. The work predicted that the 125 majority of SuperDARN's power should be dominated by the X-mode north of the modeled 126

radars, with a latitudinally narrow "X-mode only" feature prominent at CASSIOPE altitudes 127

during times in which the ionosphere's peak frequency is low. Neither prediction has been tested 128

- in SuperDARN-RRI experiments. 129
- 130

Seminal work by Warrington et al. (2012) on studying the influence of large-scale plasma 131

density structures on HF radio wave propagation conditions in the high-latitude ionosphere 132

included a ray trace modeling effort which used ray density as a proxy for signal strength. The 133

authors noted that transmitter radiation patterns could be specified when necessary. Their model 134

was able to accurately reproduce data collected between transmitter-receiver pairs in high-135

latitude regions that exhibited dramatic angle of arrival deviations attributed to pronounced 136

plasma density perturbations along the intervening ray paths. 137

138

An accurate representation of a transmitter's gain pattern is critical to computing its power 139

profile. Models of some SuperDARN gain patterns exist (e.g., Milan et al., 1997a; Sterne et al., 140

2011); however, they only show the far field radiation pattern in a vacuum. In this work, we are 141

interested in SuperDARN Saskatoon's power density profile at ionospheric altitudes. One could 142

conceivably model the radiation pattern of a radar at all altitudes of interest by using great-circle 143

paths to project a radar's gain pattern into the ionosphere. However, the pattern would only be 144

appropriate for operating frequencies that are well above the ionosphere's critical frequency – 145

when the index of refraction of the propagating radio waves is close to unity. Otherwise, the 146

pattern would be distorted by vertical and horizontal refraction. 147

- 149 We report on the creation and validation of a model of a SuperDARN radar's power density
- 150 profile, taking into account the influence of the ionospheric medium. This work is motivated by
- 151 Perry et al. (2016), who detected significant deviations in SuperDARN's power profile with RRI
- in the high-latitude ionosphere. They could not conclusively attribute the deviations to large
- scale plasma density irregularities along the ray path since the expected power profile of a
   SuperDARN radar at ionospheric altitudes was not well understood. This work is a first step
- towards gaining new insight into Perry et al. (2016)'s observations and using RRI to gain new
- insight into the effects of ionospheric irregularities at high-latitudes and their effects on HF radio
- 157 wave propagation.
- 158
- 159 In the next section, Section 2, we will describe how the model is created by using the provision
- 160 of high-frequency raytracing laboratory for propagation studies (PHaRLAP; hereafter referred to
- as Pharlap) ray trace toolbox (Cervera and Harris, 2014) to propagate a modeled SuperDARN far
- 162 field radiation pattern through an empirical model of the ionosphere provided by the
- 163 International Reference Ionosphere (IRI) (Bilitza et al., 2008). In Section 3, we will show
- 164 predictions of SuperDARN power density maps in the geographic region around Saskatoon and
- then validate the maps with RRI measurements. We will focus on the SuperDARN Saskatoon
- radar located due to the availability of measurements of that system's transmissions from
- 167 ionospheric altitudes. In Section 4 we will discuss the results of the validation efforts and
- discuss how the model output can help improve our understanding of HF radio wave propagation
- in the region. Finally, in Section 5, we will summarize our work and outline future
- 170 developments of the model.

# 171 **2 Methodology**

172 2.1 SuperDARN Saskatoon

173 In this work we are focused on the power profile of the SuperDARN radar at Saskatoon,

174 Saskatchewan, Canada (52.16° N, 106.53° W, geographic), which has been in operation since

175 1993. The radar (like all SuperDARN radars) uses electronic beam steering techniques to direct

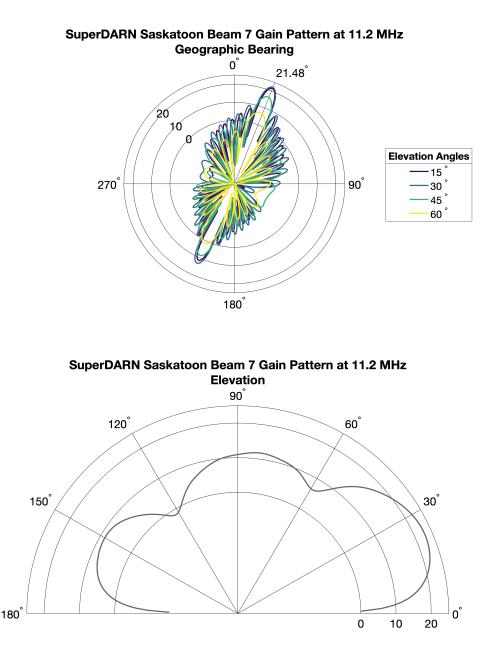
- its main lobe in 16 directions, referred to as "beams", separated by 3.24° in bearing. The radar's
- boresight, the bearing of antenna array's front main lobe, is at 23.1° (east of geographic north),
- and marks the shared border of Beams 7 and 8 (using the 0-15 beam numbering scheme in which
- the beam numbers increase in a clockwise direction). It is important to note that the radar does
- not usually transmit along the boresight, but along its numbered beams. Thus, as we will see,
  Beam 7, which is usually regarded as the "central beam", has a bearing of 21.48°, not 23.1°.
- 182

183 Figure 1 shows two polar plots of the modeled gain pattern (also commonly referred to as the

- radiation pattern) of the SuperDARN Saskatoon radar transmitting on Beam 7 at 11.2 MHz.
- 185 This gain pattern was produced using the modeling techniques and software described in Sterne
- et al. (2011). In the top panel, the pattern is shown as a function of bearing angles. The gain
- 187 pattern at several elevation angles is plotted. The Beam 7 maximum has a bearing of 21.48°, east
- 188 of geographic north. Sidelobes in the immediate vicinity of this maximum are low of the order
- 189 of 10 dBi (decibels with respect to an isotropic radiator) below the maximum. However, the
- pattern does show a significant secondary maximum directed towards the southwest. This
   secondary must be considered when analyzing radar echo returns in SuperDARN data as some
- returns may arrive along this secondary maximum from behind the radar (cf. Milan et al., 1997b;

#### 193 Burrell et al., 2015, 2018).





195

Figure 1. Polar plots of SuperDARN Saskatoon's Beam 7 gain pattern at 11.2 MHz, in dBi, as a
 function of geographic bearing at several elevation angles (top panel) and as a function of
 elevation along the bearing of 21.48° (bottom panel).

199

In the bottom panel, the gain pattern is plotted as a function of elevation angle along the bearing

of Beam 7, showing a peak of approximately 22 dBi at 22° elevation with respect to the

202 horizontal. In this work, we are only concerned with Beam 7 as it is considered the radar's

- 203 central beam, and because it simplifies validation procedures when comparing signal power
- received by RRI when the targeted radar transmits on a solitary beam. The frequency of 11.2

205 MHz was chosen since it is a common frequency used for RRI experiments with SuperDARN

Saskatoon. This frequency is also convenient since it is well enough above the expected plasma

frequency of the region such that transmissions are easily detected by RRI while still exhibiting

208 magnetoionic effects.

## 209 2.2 e-POP RRI

The Enhanced Polar Outflow Probe (e-POP) (Yau & James, 2015) is the science payload

onboard CASSIOPE, which was launched into an  $81^{\circ}$  inclination,  $325 \times 1500$  km orbit on

September 29, 2013. Among e-POP's eight instruments is a digital radio receiver, RRI, whose

scientific objectives include studying naturally and artificially generated radio emissions, and ionospheric plasma density structures. RRI is well suited for studying radio wave propagation in

the high frequency regime (HF; 3-30 MHz) of the radio spectrum and is a significant

advancement in HF radio science in the near-Earth geospace environment (Burrell, 2017). Since

the start of e-POP's science operations, RRI has conducted several conjunctive experiments with

218 SuperDARN radars, including the systems located at Saskatoon and Rankin Inlet.

219

220 SuperDARN radars are an ideal HF source for radio science experiments with RRI. The radars

are distributed throughout the mid- to high-latitudes in both hemispheres; their operating band, 8-

222 20 MHz, is almost entirely covered by RRI's operating range (10 Hz-18 MHz); and, conducting

223 RRI experiments with SuperDARN is straightforward. Indeed, coordinated RRI and

224 SuperDARN experiments have been the focus of, or featured in, several radio science

225 experiments (e.g., Burrell et al., 2015, 2018; Perry et al., 2017).

# 226 2.3 PHaRLAP

The power profile of the SuperDARN Saskatoon radar will be propagated into the terrestrial 227 ionosphere using the Pharlap ray tracing toolbox (Cervera and Harris, 2014). In this work, we 228 use Pharlap's 3D ray tracing functionality, specifically, the *raytrace 3d sp.m* method in Pharlap 229 version 4.3, which provides a variety of important radio science parameters at regular points 230 along each simulated ray including geographic location, phase path, refractive index, the angle 231 between the ray normal and the geomagnetic field, the geometric distance traveled by the ray, 232 and polarization information. Pharlap handles both the ordinary and extraordinary component of 233 the HF rays separately. The software propagates individual HF rays from the launch point 234 235 through an IRI ionosphere and international geomagnetic reference field (IGRF; Thébault et al., 2015), both of which can be easily modified or replaced by an alternative empirical model. 236

# 237 2.4 SuperDARN Saskatoon Power Density Profile Modeling Procedure

The main product of the model presented here is a power density map of the SuperDARN Saskatoon radar's radio emissions, gridded by geographic latitude, longitude, and altitude. These maps are constructed by superposing the power density contributions of each simulated ray that intersects a given bin. The power density, P(r), as a function of geometric range, r, for a given ray is:

243

244 
$$P(r) = \frac{P_T G(\theta, \phi)}{4\pi r^2} - P_{abs}(r) \qquad (1)$$

- where  $P_T$  is the power of the radar's transmitter,  $G(\theta, \phi)$ , is the gain of the transmitter as a
- function of elevation angle  $\theta$  and bearing  $\phi$  as illustrated by Figure 1, and the nominal
- transmitting power of the radar  $(P_T)$  is provided by the University of Saskatchewan SuperDARN
- engineering team. The transmitter gain pattern  $G(\theta, \phi)$  is modeled with the same software and tasking used by Sterne et al. (2011). The accurate is a Pherlan subset
- techniques used by Sterne et al., (2011). The geometric range, r, is a Pharlap output product, calculated for each simulated ray. Lastly,  $P_{abs}(r)$  accounts for the cumulative absorption
- calculated along a given ray path, described in Pederick & Cervera (2014).

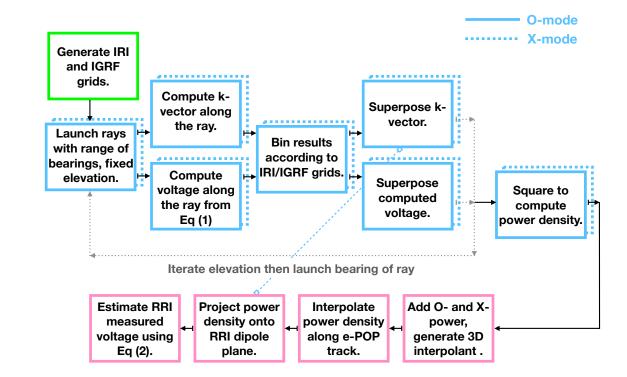


Figure 2. A flowchart describing the procedure for modeling the power density profile of the SuperDARN Saskatoon radar system, transmitting on Beam 7 at 11.2 MHz. Generating the IRI and IGRF grids is shown in green, computing the power density profile is shown in blue, and the validation procedure is shown in magenta.

Figure 2 is a flowchart describing the procedure for computing a power density profile for the 258 SuperDARN Saskatoon radar at 11.2 MHz. First, for a specified moment in time, a three-259 dimensional grid of plasma density and magnetic field values is generated with separate calls to 260 IRI and IGRF Pharlap functions (highlighted in green in Figure 2). For the results presented 261 here, the grid spans  $30^{\circ}$ - $80^{\circ}$  N in geographic latitude and  $155^{\circ}$ - $55^{\circ}$  W in longitude with a 262 resolution of 0.25°, and 50-1550 km in altitude with a 7.25 km resolution. This grid size and 263 resolution were chosen to meet a desired resolution of the power density maps and encompass a 264 significant portion of SuperDARN Saskatoon's field-of-view and CASSIOPE's altitude range. 265 In calling IRI, we used an R12 value (12-month running average sunspot number) of 15, 266 consistent with the conditions during the RRI's measurements which we will use for validation. 267 Otherwise, the default settings encoded into Pharlap were used; the IRI 2012 model using the 268 bottomside profile introduced by Bilitza et al. (2000) and the D-region model developed for IRI-269 270 1990.

- After generating the IRI and IGRF grids, 11.2 MHz rays are then launched over a range of
- bearing angles at a fixed elevation angle. As Figure 2 illustrates, the power density map
- computation procedure iterates through a range of bearing  $(0^{\circ}-360^{\circ})$  and elevation  $(10^{\circ}-90^{\circ})$
- angles, with a resolution of  $0.10^{\circ}$ . Equation 1 is calculated at points along each ray, in a step size
- no greater than 10 km. As previously discussed, the values for  $P_T$  and  $G(\theta, \phi)$  are known a priori, while the value of geometric range, *r*, and cumulative absorption  $P_{abs}(r)$  are provided as
- priori, while the value of geometric range, r, and cumulative absorption  $P_{abs}(r)$  are provided as an output of the ray tracing software. Meanwhile, the unit k-vector for the wavevector is
- computed for each step size along the ray. The power and k-vector information are then binned
- in the same grids as the ionosphere and magnetic field data.
- 281

282 Pharlap provides a phase path estimate for each ray, which is used to calculate the instantaneous

- 283 phase at any point along the rays in each bin. A wavevector is then constructed from this phase
- information, the square root of Equation 1 (a quantity proportional to an electric potential), and
- the aforementioned k-vector calculation so that multiple rays within a bin may be superposed.
- As the gray arrows in Figure 2 indicate, this process is iterated several times to account for all the desired rays.
- 287 288
- The final gridded superposed electric potential values are then squared, producing a three-
- dimensional grid containing power density values and entries. This entire process, highlighted in
- blue in Figure 2, is computed separately for the O- and X-modes. Unlike the model in Gillies et al. (2010), the model presented here does not consider the initial polarization state of the
- 292 al. (2010), the model presented here does not consider the initial polarization state of the
- transmitted rays at the radar; rather, the initial transmitting power is distributed evenly to each mode. We regard the Gillies et al. (2010) work as a higher-order correction that will be
- mode. We regard the Gillies et al. (2010) work as a higher-order correction that implemented in future versions of the model.

# 296 **3 Results and validation**

297 3.1 August 5, 2017 18:35 UT ray tracing example

Figure 3 shows Pharlap ray tracing results for August 5, 2017 at 18:35 UT. This time was specifically chosen as it coincides with an RRI experiment that will be discussed in more detail

300 shortly. The top panel shows O- and X-mode rays originating from Saskatoon, Canada,

launched along Beam 7's bearing, between from 10° and 90° in elevation angle, in 1°

302 increments. For clarity, a reduced number of rays are plotted in Figure 3 compared to the

- 303 modeling results that will be shown shortly, which consists of more rays.
- 304

A contour plot of plasma densities generated by IRI along the Saskatoon geographic meridian is

also plotted in the background. CASSIOPE's tracks during the five conjunctions which will be

used to validate the power density model are plotted in the foreground. A similar plot is shown

in the bottom panel from a different vantage point. Here, only the paths of the rays launched at  $30^{\circ}$  elevation, from 0°-360° in bearing in 2° increments are shown between 50 and 500 km

altitude. The contour plot in the background shows the IRI ionosphere at 245 km altitude.

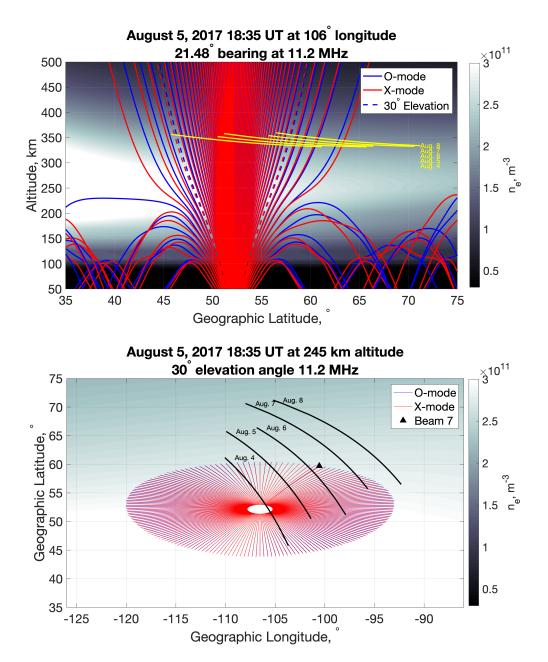


Figure 3. (top) Ray race simulation results for August 5, 2017 at 18:35 UT showing ray path for 312 11.2 MHz O- (blue) and X-mode (red) rays launched along front and rear facing maxima of the 313 Beam 7 gain pattern as a function of geographic latitude and altitude at Saskatoon's longitude. 314 Rays with a 30° elevation angle are indicated with dashed lines. (bottom) The same results but 315 for rays launched at 30° elevation for all bearings, plotted as a function of geographic longitude 316 and latitude at 245 km altitude. IRI plasma density contours are plotted in the background, with 317 the CASSIOPE tracks of five RRI conjunctions plotted in the foreground. Rays along the Beam 318 7 bearing are indicated. 319

Both panels in Figure 3 show that the ray trace predicts that rays with elevation angles above

 $\sim 25^{\circ}$  can propagate directly through to the ionosphere, while rays at lower elevation angles are internally reflected and undergo multi-hop propagation without being able to propagate vertically

- through the ionosphere downrange. A spacecraft in the region between 200 500 km altitude
- will only have access to the highest elevation rays shown in the bottom panel of Figure 3,
- according to the ray trace simulation.
- 326 3.2 August 5, 2017 18:35 UT power density map example
- 327 The main output of the SuperDARN power density model is a power density map of the radar's
- emissions in the vicinity of the radar. A map for August 5, 2017 at 18:35 UT is plotted in Figure

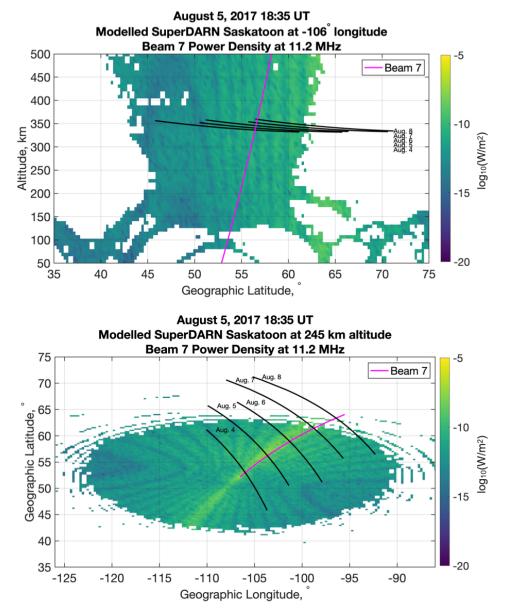


Figure 4. Simulated power density values plotted with the same axes in Figure 3. The nominal great-circle path along Beam 7's bearing is plotted in magenta in both plots. Beam 7's power density peak (bottom panel) in the front and back lobe is clear, as are the secondary peaks due to the sidelobes of the array's gain pattern (both panels).

4 with the same axes as Figure 3. This map combines contributions from the O and X

propagation modes. Two clear differences between Figures 3 and 4 are that the latter shows

binned power density values instead of ray paths, and that the elevation (top) and bearing

- (bottom) of Beam 7 are plotted in magenta.
- 338

The longitudinal slice (top panel, Figure 4) shows that, in the F region, the radar's radiated

pattern is mostly constrained between  $40 - 63^{\circ}$  latitude. Below 200 km, the latitudinal range of

the radar's radiated power is far more extensive, because of internal reflection, but it is much

more intermittent. As the bottom panel of Figure 4 shows, the rays crate a geometric shape

commonly referred to as the "iris of accessibility" (James, 2006). Within this region sub-orbital
 spacecraft such as sounding rocket payloads have access to the rays originating from the

- 345 terrestrial transmitter.
- 346

The imprint of the radar array's gain pattern is readily apparent in the bottom panel of Figure 4,

which shows the radar's power density at an altitude of 245 km. The peak forward radiated

power is centered at approximately 19° in bearing, more northerly than the great-circle path

trajectory of Beam 7 (21.48°). This is because the gain pattern (cf. Figure 1) for the higher

elevation rays (>22° elevation) is directed more northerly. Lower elevation rays are internally reflected. The peak in the radar's back lobe is also clearly defined, which, as discussed earlier,

reflected. The peak in the radar's back lobe is also clearly defined, which, as discussed earlier must be considered when interpreting SuperDARN radar echo returns. Both top and bottom

panels of Figure 4 also show modulations in power density. These are attributed to gain

- 355 pattern's sidelobes.
- 356

357 In general, the modeling results displayed in Figure 4 are consistent with a general intuition of

SuperDARN Saskatoon's power density profile. The power density peak is essentially aligned with the bearing of the transmitting beam (Beam 7 in Figure 4), constrained to a width of a few

with the bearing of the transmitting beam (Beam 7 in Figure 4), constrained to a width of degrees bearing, and it is several factors stronger than the radar's sidelobes. Below

degrees bearing, and it is several factors stronger than the radar's sidelobes. Below approximately 200 km altitude, there is an absence of power density within a few hundred

kilometers of the radar – the "skip zone". Above 200 km, the radiated power density is

contained within the iris of accessibility. Power density maps for August 4, 6, 7 and 8, 2017, are

364 provided in Supplementary Information as Figures S1 – S4, respectively.

365 3.3 Comparison to August 4 - 8, 2017 RRI conjunctions

366 We now turn our focus to validating the power density model by comparing its output to

367 measurements collected during five experiments with RRI and SuperDARN Saskatoon,

368 conducted between August 4 and August 8, 2017. In each experiment, the radar transmitted at

11.2 MHz along Beam 7. RRI was activated for 237 seconds starting at 18:57:14 UT on August

4, and approximately 23-24 minutes earlier each successive day as a result of CASSIOPE's

nodal precession. RRI was tuned to 11.205 MHz and operated in the crossed-dipole mode (Perry

et al., 2017). CASSIOPE's tracks for the five experiments are depicted with black traces in

Figure 4. Except for the August 4 experiment, when the normal of RRI's dipole plane is along

the spacecraft ram direction in the horizontal plane, CASSIOPE's attitude was slewed such that the normal of RRI's dipole plane was continuously directed at Saskatoon.

376

The procedure for validating the modeled power density profile is outlined in magenta in Figure

2. The end product of the procedure is a prediction of the total voltage measured by RRI's two

dipoles. First, a 3D interpolant of the power density profile of the radar is generated with

dimensions of geographic latitude, longitude, and altitude. MATLAB's *scatteredIinterpolant* 

- function with a linear interpolation method was used for the results presented here.
- 382

The interpolant is then sampled at points along CASSIOPE's track, producing an estimate of the power density incident on RRI along the track. This estimate is then projected onto RRI's dipole plane by computing the dot product of the wave vector associated with the modeled power density value and the unit vector normal to RRI's dipole plane. Recall that the wave vector quantity was computed and stored during the ray tracing procedure (cf. Figure 2). The normal of RRI's dipole plane is calculated from CASSIOPE's attitude information contained in its ephemeris file.

390

391 The predicted RRI voltage, V, is then calculated using (Griffiths, 1999):

- 392 393  $V = \sqrt{\frac{2IL_{eff}^2}{n^2 \epsilon_0 c}} \qquad (2)$
- 394

in which  $L_{eff}$  is the effective length of an RRI dipole antenna: 3 m (James, 2003), I is the power density projected onto RRI's dipole plane, n, is the index of refraction, provided by IRI at the spacecraft position,  $\epsilon_0$ , is the permittivity of free space, and c, is the speed of light in a vacuum.

398

The predicted voltage, V, is then compared to the measured voltage of the radar pulses received by RRI. SuperDARN pulses were extracted from RRI's data using techniques described in Perry et al., (2017). The voltage measured by both dipoles at the midpoint of the detected pulse was added, then compared to the predicted voltage. Since RRI's dipoles are orthogonal to one another, the total voltage induced by the electromagnetic energy incident on RRI's plane can be computed by applying the Pythagorean theorem to the dipole voltages.

405

Figure 5 illustrates the results of the predicted RRI voltages (black) plotted together with the
measured RRI voltages (red) as a function of CASSIOPE'e geographic latitude for the (a-e)
August 4-8, 2017 experiments. The same results are plotted as a function of geographic

longitude in Figure 6. In both plots, a 100-point moving median, equivalent to approximately 15

410 km distance along CASSIOPE's track, has been applied to the RRI data trace. Note that both

axis in the different panels of each plot have different scales and ranges in general, in order to

412 optimize the data display in each case. The magenta-colored vertical dashed line in both figures

413 marks the point at which e-POP crosses the great-circle path of Beam 7's bearing – the same

- 414 magenta line plotted in Figure 4 (both panels).
- 415

For the August 5-7 experiments, the modeled voltages values show good agreement with the

417 measured values. The difference in predicted and measured peak voltages are within

418 approximately 20%. The predicted voltage peaks are within about 2-3° in latitude (Figure 5) and

419 1-3° in longitude (Figure 6) of the measured values. The August 5 experiment shows the best

agreement; the predicted voltage profile is nearly identical to the measured RRI voltages.

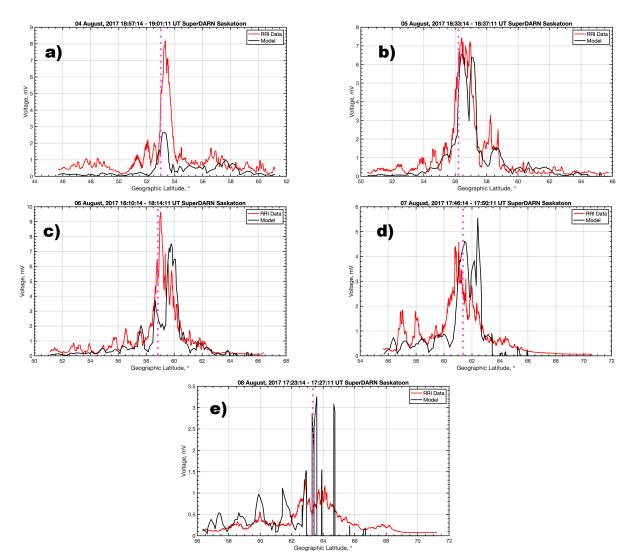


Figure 5. Predicted (black) and measured (red) RRI voltages as a function of geographic latitude
 for (a-e) August 4 – 8, 2017 RRI measurements of SuperDARN Saskatoon. Note that the
 different panels have different axis scales and/or ranges. The magenta-colored vertical dashed
 line marks the point at which e-POP crosses the great-circle path of Beam 7's bearing.

A persistent feature of note in the predicted August 4 - 7 voltage profiles is a double-hump 427 maximum. A closer inspection of the model output reveals that each is due to a separation in the 428 429 O- and X-mode power density profiles, in which the peaks located at higher latitudes are due to the O-mode. The double-hump feature is not observed in the RRI voltage profiles. The feature 430 arises in the model because we have assumed that the transmitting power of the radar is 431 distributed equally to O- and X-mode of propagation. According to Gillies et al., (2010), the 432 initial transmitting power should be weighted towards the X-mode because of the radar's 433 orientation with respect to the geomagnetic field. In reality less O-mode power should be 434 expected at RRI's altitude, which would be reflected in the model output as a reduction in the 435

436 amplitude of the higher-latitude hump. As mentioned earlier, the initial polarization at the radar

437 site will be accounted for in future versions of the model.

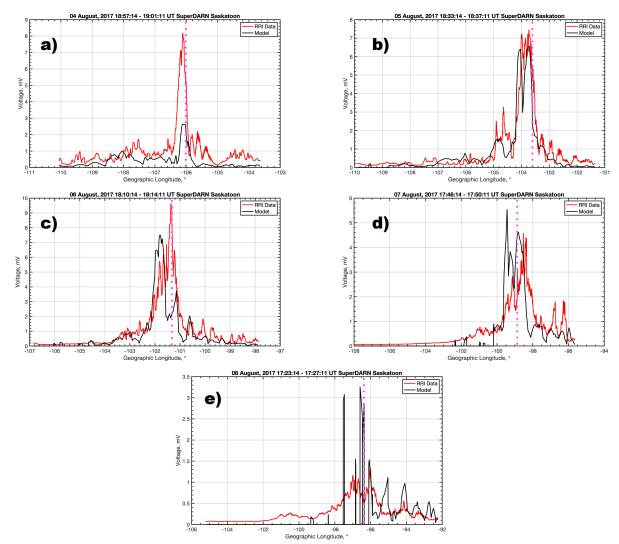


Figure 6. Predicted (black) and measured (red) RRI voltages as a function of geographic
 longitude for (a-e) August 4 – 8, 2017 RRI measurements of SuperDARN Saskatoon. Note that
 the different panels have different scales and/or ranges.

The measured voltage peaks differed from where e-POP crossed Beam 7's great-circle path by 0.5° in latitude and approximately the same amount in longitude. As Figures 5 and 6 show, the assumption that the radar's maximum power density lies along the great-circle path is reasonable for the August 4-8 experiments.

447

The August 4 and 8 experiments show considerable disagreement between the predicted and measured RRI voltages. For the August 4 experiment, the predicted voltage peak value is considerably lower than the measured peak value; however, there is good agreement on the location of the peak. Recall that for the August 4 experiment, CASSIOPE was unable to slew

452 RRI towards SuperDARN Saskatoon, unlike the August 5-8 experiments. As the separation

distance decreased, the angle between the line-of-sight to Saskatoon from CASSIOPE and the

normal to RRI's dipole plane became more oblique. The result of this was an underestimation of RRI's effective area,  $L_{eff}^2$ .

455 RRI 456

- 457 The model assumes that  $L_{eff}^2$  decreases like  $|\cos \theta|$ , where  $\theta$  is the angle between the wave
- vector of the incident radio wave and the normal to RRI's dipole plane. In reality, the
- relationship between  $L_{eff}^2$  and an incident radio wave is much more complicated, especially when
- 460  $\theta$  is in the vicinity of 90° (James, 2003). This is not accounted for in the validation procedure,
- and results in an underestimate of  $L_{eff}^2$ . This is the case for the August 4 pass when CASSIOPE
- is within a few degrees of latitude and longitude with respect to Saskatoon (52.16° N, 106.53°
  W, geographic).
- 464

There is clear disagreement between the model and RRI's measurements for the August 8 conjunction as well (cf. Figures 5 and 6). The model predicts sporadic and very low RRI voltage measurements, whereas RRI measured a subdued voltage profile. As Figure 4 shows, RRI's track during the August 8 experiment was mostly outside of the iris of accessibility, hence the predicted absence of RRI voltages. Overall, the measured RRI voltages are subdued and distorted compared to the previous days' measurements: there are two peaks rather than a

- singular peak, and the relative height of the peak is low compared to the previous conjunctions.
- 472

473 We attributed the August 8 disagreement to an overestimation of the ionospheric plasma density

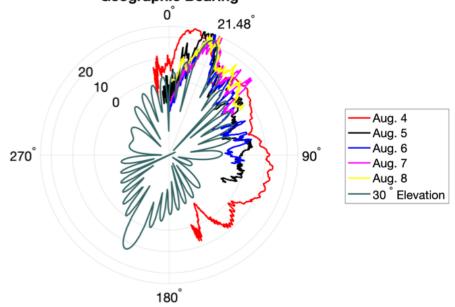
by IRI. A denser ionosphere would result in a more contracted iris of accessibility than what was

475 encountered by RRI, which is reflected by the lack of predicted RRI voltages. A more tenuous

ionosphere would have expanded the predicted iris of accessibility and would have likely
 generated a more accurate prediction of what RRI measured. We discuss this further in the

- 477 generated a more ac478 discussion section.
- 479

Figure 7 provides a summary of the August 4 – 8 measurements of SuperDARN Saskatoon in a 480 similar format to the top panel of Figure 1. Along with the Beam 7 gain pattern, the measured 481 RRI voltages (normalized to 25 dB in gain value) are plotted as a function of bearing (top panel) 482 with respect to the radar. This figure illustrates that, indeed, the relative voltage profile as a 483 function of bearing is generally consistent with the modeled gain pattern. The peak measured 484 voltages are aligned with the modeled gain pattern; the widths of the voltage peaks are also 485 consistent with the model; and the sidelobes are reduced. As discussed earlier, the largest 486 deviation between model and measurement occurred during the August 4 and 8 experiments, and 487 there is a slight northward deviation in the peak measured voltage that is especially evident 488 during the August 5 conjunction. 489



#### SuperDARN Saskatoon Beam 7 Gain Pattern vs RRI at 11.2 MHz Geographic Bearing

490

Figure 7. Summary of RRI's measurements of SuperDARN Saskatoon collected between August 4 - 8, 2017, plotted alongside the radar's modeled gain pattern, in the same format as the top panel in Figure 1.

### 494 **4 Discussion**

We have developed a model of the SuperDARN Saskatoon's radar's power density profile that
accounts for HF radio wave propagation characteristics and extends into the F-region ionosphere.
We have also compared the model with measurements of the radar's transmissions using RRI
during five conjunctions that occurred between August 4 and 8, 2017. The measurements

collected between August 5-7 show very good agreement with the model, while the August 4 and

- 500 8 measurements do not.
- 501

499

To the best of our knowledge, the model is unique in that it is the first and only to consider SuperDARN's power density profile at ionospheric altitudes. Knowledge of this quantity is critical to interpreting radar data as it indicates where the transmitter's power is directed and thus the likely location of the scattering volume. The strong agreement between the model and RRI's measurement on August 5-7 is a positive indication that the model is both accurate and a useful tool for understanding SuperDARN backscatter and interpreting RRI's measurements of the system.

- 509 4.1 Model limitations
- 510 It is important to note that the model does have its limitations. First, it only considers
- 511 transmissions at 11.2 MHz, which is at the approximate midpoint of the radar's usual
- 512 transmitting range of 10 12 MHz. Future development of the model will include a more
- 513 extensive range of transmitting frequencies used by the radar. Additionally, the power density

514 profile is exclusive to SuperDARN Saskatoon's array of log-periodic antennas. It would not be

appropriate to apply the model to another radar of different construction, for example, the

516 SuperDARN radar located at Rankin Inlet, Canada, which is made-up of twin terminated folded

- 517 dipole antennas (Sterne et al., 2011).
- 518

519 We have validated the model during daytime conditions under generally calm geomagnetic

conditions. The Kp index was 3, 3, 2, 0, and 2 during the August 4-8 conjunctions, respectively.

521 The accuracy of the model during nighttime and/or more disturbed conditions will be

522 investigated in future studies. It is expected that the model will have greater difficulty making

523 accurate predictions under more disturbed conditions since SuperDARN Saskatoon's nominal

field-of-view intersects with the auroral oval where plasma density irregularities are frequent.
 With respect to nighttime under calm geomagnetic conditions, we anticipate good agreement

526 given that the HF rays will be less affected by the depleted ionospheric density and absorption.

527

528 The model is also limited by its reliance on IRI as an empirical ionosphere. This means that the

529 power density profile model may only be accurate in a statistical sense and is likely not as

applicable during geomagnetically disturbed conditions or when transient phenomena are

present, such as auroral precipitation. The ionosphere's day-to-day variability is extremely

difficult to replicate with an empirical model. Indeed, IRI has been shown to be deficient in the

533 Canadian sector in general (Themens et al., 2017). Future developments of the model will

replace IRI with the Empirical Canadian High Arctic Ionospheric Model (E-CHAIM; Themens

et al., 2017) which has been shown to be more suitable for the region.

4.2 The August 8, 2017 conjunction

The lack of agreement between the modeled predictions and RRI measurements during the 537 August 8 conjunction is likely a consequence of IRI's constraints. As discussed earlier, we 538 hypothesize that the discrepancy is a result of IRI overestimating ionospheric density, resulting 539 in a more contracted iris of accessibility than what RRI encountered on August 8. To explore 540 this hypothesis further, we have plotted the SuperDARN backscatter data from the August 4-8 541 542 conjunctions in Figure 8, generated using the pyDARNio software package (version 1.0; The SuperDARN Data Analysis Working Group, 2021) which shows a plot of the time averaged 543 backscatter power (Signal to Noise Ratio; SNR) measured during the conjunctions. The 544 geographic latitude of the echoes was computed using the Chisham et al., (2008) virtual height 545 model. 546

547

548 The northward movement of the scatter on August 8 (yellow trace) compared to the previous days is indicative of a more tenuous ionosphere on that day, which would result in an expanded 549 iris of accessibility. This movement is especially evident north of 60°. It is challenging to 550 distinguish ionospheric and ground scatter in the radar echo data presented in Figure 8. Figure 3 551 indicates that both may be present. In either case, the scattering locations of both are expected to 552 move north as the density of the intervening ionosphere decreases since the HF rays will undergo 553 554 less refraction and will require a longer path length to satisfy the aspect angle condition or backscatter off the ground. This corroborates the hypothesis that the disparity between 555 prediction and measurements on August 8 was due to IRI's overestimation of the ionospheric 556 density, whereas IRI's prediction for August 4-7 was more accurate. 557

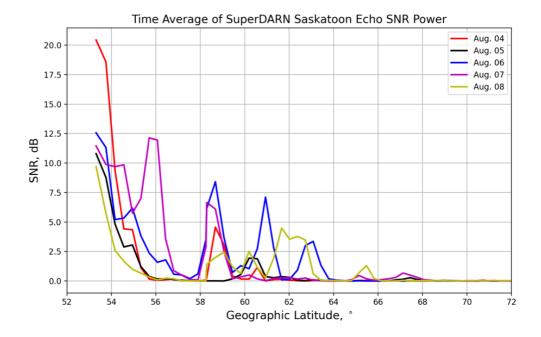


Figure 8. Time average of SuperDARN Saskatoon backscatter power (signal-to-noise ratio;
 SNR) as a function of geographic latitude, collected by Beam 7 during the RRI experiments
 August 4 – 8, 2017.

#### 562 4.3 Response to more intense solar conditions

We also investigated the model's response to a variety of solar conditions. Different IRI ionospheres were compared, with the goal of better understanding the variability of the modeled power density profile. In particular, an IRI ionosphere with driving R12 values of 60 and 100 were used (an R12 of 15 was used in modeling the August 4-8 experiments), consistent with much more intense solar activity. The only effect the higher R12 values had was to decrease the peak predicted RRI voltage (not shown) and contract the iris of accessibility. The locations of the voltage peaks were moved slightly southward, consistent with a denser ionosphere.

#### 570 4.4 Relevance to SuperDARN and HF radio science

571 The model has important implications for remote sensing the MIT system, which can be used to 572 improve the interpretation of SuperDARN data and in particular the SuperDARN geolocation 573 algorithms. Unlike determining whether radar echoes originate from in front of or behind the 574 radar, which can be accomplished with SuperDARN's interferometry data (e.g., Milan et al., 575 1997b; Burrell et al., 2015, 2018), it is unlikely that a comprehensive method of measuring the 576 effective pointing direction of the radar's transmitting beam is feasible based solely on the 577 radar's data.

578

579 Models of SuperDARN's power density profile, like the one presented and validated in this

work, are a helpful tool for enhancing SuperDARN data analysis and improving our

understanding of HF radio wave propagation characteristics in the high-latitude regions. For

example, the model can be used to simulate SuperDARN Saskatoon backscatter echoes by

considering the power density profile of the radar in addition to where the aspect angle condition

is satisfied in the radar's field-of-view. The simulated echoes can then be used as a guide for

585 interpreting backscatter data in terms of both occurrence and echo power. This new capability is

an advance on traditional SuperDARN modeling which, in general, has not accounted for the
 power density profile of the system.

588

SuperDARN data products assume that the scattering volume for a corresponding radar echo is 589 located along a great-circle path of the transmitting beam's bearing. As Figures 5-7 demonstrate, 590 this assumption is reasonably accurate under quiet geomagnetic conditions. The effective 591 pointing direction of SuperDARN Saskatoon's Beam 7, as measured by RRI, is of the order of 592 0.5° in latitude and longitude deviated from the beam's nominal great-circle path. This is a 593 negligible amount in lateral deviation – it amounts to a fraction of a beamwidth, but it is 594 approximately equivalent to a SuperDARN range gate in group delay – 45 km, a value that is 595 consistent with the expected SuperDARN echo geolocation accuracy (Yeoman et al., 2008). It 596 seems plausible, then, that the indeterminacy of the transmitting beam's main lobe at ionospheric 597 altitudes is an important, yet underappreciated, contributor to the error in SuperDARN's 598

599 geolocation methodology.

600

Little is known about the stability of the transmitting beam under more disturbed geomagnetic conditions or more regular events such as when the day/night terminator intersects with the

beam, or how such perturbations might influence the data. At polar latitudes, there is strong

evidence that the transmitting beam undergoes significant deviations (Perry et al., 2016), which

605 is consistent with other radio experiments in the polar region (Warrington et al., 1997). In both

cases, polar-cap patches were identified as the likely cause of the deviations. Even though the

607 lateral plasma density gradients that are common in the polar-cap region are rarely observed at

lower latitudes, it is still conceivable that transmitting beam deviations may occur with other
 SuperDARN radars, such as the Saskatoon radar, which may have a non-negligible effect on the

610 accuracy radar echo geolocation.

# 611 **5 Summary**

We have developed a novel model of SuperDARN Saskatoon's power density profile at

613 ionospheric altitudes. The model advances on previous SuperDARN and HF radio wave

614 propagation modeling in that it accounts for the ionospheric medium and its influence on HF

radio wave propagation and quantifies the power density profile of the radar system at

616 ionospheric altitudes. We have validated the model using RRI measurements collected between

August 4-8, 2017. The model's predictions are in good agreement with RRI's measurements,

demonstrating that the model is useful for predicting a SuperDARN radar's power density

619 profile, provided the appropriate gain pattern for the radar is used.

620

The model is a tool for interpreting SuperDARN backscatter data as it provides a method for

locating the scattering volume based on the radar's power profile, which supplements the

traditional method of determining where the aspect angle condition is satisfied in the radar's

624 field-of-view. It can also be used as a tool for interpreting RRI measurements of SuperDARN

transmissions and other ground-based transmitters; one simply needs to apply the appropriate

626 gain pattern for the targeted HF transmitter. Future developments of the model include

627 expanding the model to consider other SuperDARN transmitting frequencies and an alternate

628 empirical ionosphere.

#### 629 Acknowledgments, Samples, and Data

- This work was carried-out by KDR for an undergraduate project at the University of Calgary 630 under the supervision of GWP and AWY. The authors thank Dr. Kathryn McWilliams and the 631 SuperDARN Canada engineering team for facilitating RRI's measurements. The authors also 632 thank Dr. Glenn Hussey and Clifford Ridley for their informative discussions and advice. 633 634 Finally, the authors are indebted to the members of the e-POP Science Operations Centre, without whom RRI's measurements would not be possible. 635 636 We acknowledge funding support from the Canadian Space Agency under Grant 16SUSTSSPI, 637 Natural Sciences and Engineering Research Council of Canada under Discovery Grant 638 RGPIN/06069-2014 to AWY and the European Space Agency (ESA) Third Party Mission 639 Program. 640 641 The power density profile model code may be accessed at GWP's github page: 642 https://github.com/GWPerryNJIT. The results published in this paper were obtained using the 643 HF propagation toolbox, PHaRLAP, created by Dr. Manuel Cervera, Defence Science and 644 Technology Group, Australia (manuel.cervera@dst.defence.gov.au). This toolbox is available by 645 request from its author. 646 647 The authors acknowledge the use of data from SuperDARN, an international project made 648
- 649 possible by the national funding agencies of Australia, Canada, China, France, Italy, Japan,
- 650 South Africa, Norway, the United Kingdom and the United States of America. SuperDARN data
- 651 can be downloaded from Globus, instructions of which are provided here:
- 652 https://superdarn.ca/data-products.
- 653
- The RRI data used in this study are publicly accessible at: https://epop.phys.ucalgary.ca/data/.
- 655 The Kp-index data used in this study is publicly accessible at:
- 656 https://www.swpc.noaa.gov/products/planetary-k-index.
- 657
- The authors also acknowledge the use of the *numpy* and *Scipy* packages.

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