

# Relative tsunami hazard from segments of Cascadia subduction zone for Mw 7.5-9.2 earthquakes

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## Abstract

Tsunamis from earthquakes of various magnitudes have affected Cascadia in the past. Simulations of Mw>7.5–9.2 earthquakes constrained by earthquake rupture physics and geodetic locking models show that Mw>8.5 events initiating in the middle segments of the subduction zone can create coastal tsunami amplitudes comparable to those from the largest expected event. The simulations reveal that the concave coastline geometry of the Pacific Northwest coastline focuses tsunami energy between latitudes 44°-45° in Oregon. The possible coastal tsunami amplitudes are largely insensitive to the choice of slip model for a given magnitude. These results are useful for identifying the most hazardous segments of the subduction zone and demonstrate that a worst-case rupture scenario does not uniquely yield the worst-case tsunami scenario at a given location.

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1 Relative tsunami hazard from segments of  
2 Cascadia subduction zone for  $M_w$  7.5-9.2  
3 earthquakes

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## Abstract

26           Tsunamis from earthquakes of various magnitudes have affected Cascadia in the past. Simulations  
27 of  $M_w$  7.5–9.2 earthquakes constrained by earthquake rupture physics and geodetic locking models show  
28 that  $M_w \geq 8.5$  events initiating in the middle segments of the subduction zone can create coastal tsunami  
29 amplitudes comparable to those from the largest expected event. The simulations reveal that the concave  
30 coastline geometry of the Pacific Northwest coastline focuses tsunami energy between latitudes  $44^\circ$ - $45^\circ$  in  
31 Oregon. The possible coastal tsunami amplitudes are largely insensitive to the choice of slip model for a  
32 given magnitude. These results are useful for identifying the most hazardous segments of the subduction  
33 zone and demonstrate that a worst-case rupture scenario does not uniquely yield the worst-case tsunami  
34 scenario at a given location.

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## Plain Language Summary

37           Offshore earthquakes along the Pacific Northwest coast of the U.S. and Canada (Cascadia) can  
38 have magnitudes as high as 9.2, as was probably the case for an earthquake in the year 1700 CE that  
39 resulted in a large tsunami in Cascadia and across the Pacific Ocean. To learn more about the future  
40 tsunami hazard in the region, we design computer models of tsunamis from a wide range of earthquake  
41 scenarios. We find that almost regardless of the earthquake source details, events larger than magnitude 8.5  
42 near the coast of Oregon can create large and widespread tsunamis along the US west coast. These are  
43 consequences of the geometry of offshore earthquake faulting and the concave shape of coastline in the  
44 region.

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### 47           Key Points:

- 48           • A  $M_w=8.5$  event in central Cascadia (Oregon) can create coastal tsunami amplitudes  
49           comparable to those from the largest possible event.
- 50           • The concave coastline contributes to larger coastal tsunami amplitudes in central Cascadia.
- 51           • The choice of slip model does not significantly affect the distribution of coastal tsunami  
52           amplitudes in Cascadia.

53

## 54 1. Introduction

55 Coupling between the subducting Juan de Fuca plate and the overriding North America at  
56 the shallow-dipping Cascadia subduction zone [Crosson & Ownes, 1987] is expected to cause  
57 future large earthquakes and tsunamis. Previous studies [Atwater, 1987; Satake *et al*, 2003;  
58 Goldfinger *et al*, 2017] show that the ~1100-km plate boundary extending from British Columbia  
59 to northern California has generated large tsunamis in the past. The tsunami from the most recent  
60 great Cascadia earthquake or sequence of events, with an inferred total magnitude of 9 [Melgar,  
61 2021], devastated the American and Japanese coasts on 26 January 1700 [Satake *et al*, 1996] as  
62 shown by Native American oral traditions [Heaton & Snavely, 1985] and detailed Japanese  
63 written accounts [Atwater *et al*, 2015]. Geological and paleoseismic evidence also indicates  
64 earlier prehistoric tsunamigenic events [Atwater *et al*, 1991; Darienzo & Peterson, 1995; Peters *et*  
65 *al*, 2003; Goldfinger *et al*, 2012].

66 A major challenge to modeling future Cascadia tsunami hazards is the large uncertainty in  
67 recurrence interval. Very large Cascadia events seem to occur on average every 400-500 years,  
68 and smaller events ( $M_w \lesssim 8.7$ ) are thought to occur every ~200 years. However, the uncertainties  
69 in these measurements are sometimes on the same order of magnitude as the measurements  
70 themselves [e.g., Kelsey *et al*, 2005; Goldfinger *et al*, 2012]. Hence, it is challenging to assess  
71 tsunami hazards from different rupture segments along the coast [e.g., González *et al*, 2009].  
72 Along-strike segmentation in the subduction zone is usually associated with variabilities in outer  
73 wedge morphology and structure [Watt & Brothers, 2021]. Another source of uncertainty is the  
74 limited constraints on the location and size of slipping segments of the subduction zone, and thus

75 the magnitude of future earthquakes, due to the absence of well-constrained information on the  
76 lateral extent of past ruptures [e.g., Witter *et al*, 2011; Goldfinger *et al*, 2017].

77         Tsunami hazard in Cascadia has been modeled as a function of various recurrence  
78 intervals and earthquake magnitudes, yielding results expressed as hazard curves and inundation  
79 maps for sites along the coast [e.g., Thio & Somerville, 2009; Thio *et al*, 2010; Priest *et al*, 2013;  
80 Park *et al*, 2017]. Such studies, which are also conducted for other, better-documented subduction  
81 zones [e.g., Satake, 2015] are useful in planning response to tsunamis [Lindell & Prater, 2010].  
82 However, they often do not distinguish between hazards due to rupture from various sections of  
83 the subduction zone.

84         Our study aims to identify the most hazardous segments of the Cascadia subduction zone  
85 by considering a range of earthquake scenarios and resulting tsunamis. We use a set of rupture  
86 scenarios derived by scaling the slip distribution prescribed by locking models as initial  
87 conditions to simulate Cascadia tsunamis. We then compare our tsunami simulation results to  
88 those based on dynamic ruptures as well as perturbed versions of these rupture scenarios to  
89 identify the contribution of various portions of the subduction zone. Finally, we use simple  
90 numerical experiments with synthetic bathymetry to investigate the contribution of Cascadia  
91 coastal morphology to the distribution of tsunami amplitudes. The result helps us better  
92 understand the contribution of source and coastal components to Cascadia tsunamis.

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## 94 **2. Methods**

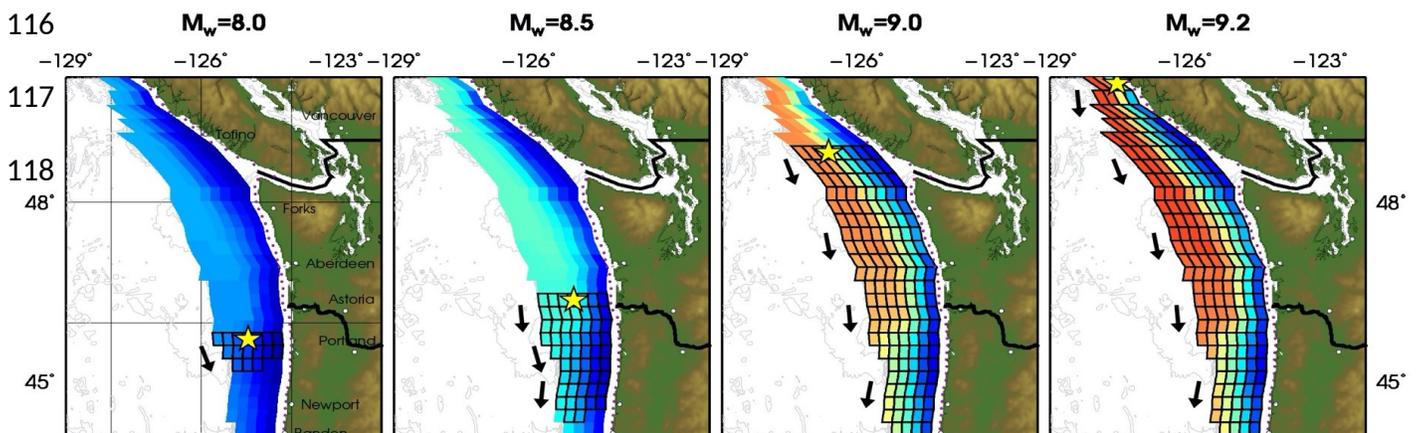
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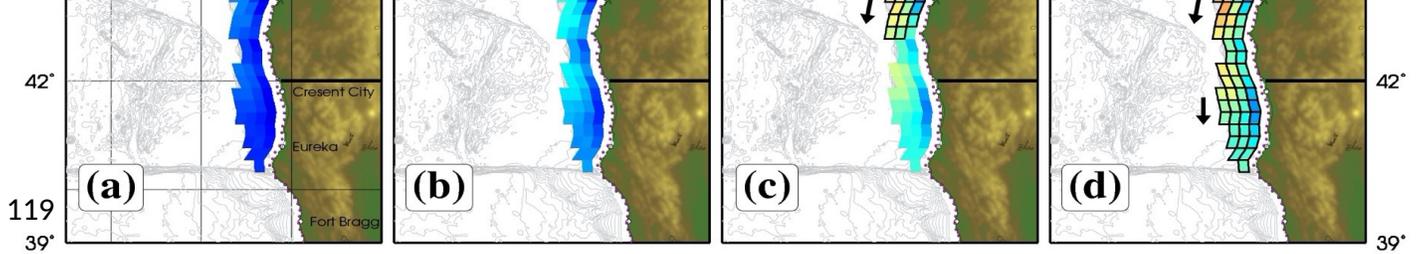
## 96 2.1 Rupture Model & Scaling of Slip

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98 Our earthquake simulations are based on locking models that estimate the slip deficit on  
99 the plate interface needed to match geodetic observations [Li *et al*, 2018]. If all the deficit were  
100 released in a single earthquake, the model would yield the maximum possible slip and the largest  
101 earthquake. However the next event may not release all the accumulated stress. We use the  
102 Gamma locking model [Schmalzle *et al*, 2014] to represent high levels of slip-deficit extending to  
103 the trench (Fig. 1d). In contrast with models with more uniform slip distributions, releasing the  
104 accumulated slip mostly confined near the trench would result in pulse-like, relatively short-  
105 period tsunami fronts. Due to very large seafloor uplift in the immediate vicinity of the trench,  
106 such a model results in relatively larger coastal tsunami waves, especially in the near-field. By  
107 assuming an average recurrence interval  $T_r$  for stress release, as inferred from offshore turbidite  
108 deposits [Goldfinger *et al*, 2012, 2017], we convert the slip rate deficit in the locking model to  
109 slip. An average value of  $T_r$ , i.e., 320 years [Goldfinger *et al*, 2017] results in a maximum slip  
110 amplitude of  $\sim 20$  m along the trench in northern Cascadia.

111 We discretize the rupture into  $25 \times 25$  km blocks (Fig. 1), each of which is considered as a  
112 pure double-couple source with the dip of the slab, the azimuth of the trench, and a slip angle of  
113  $90^\circ$ . We then calculate a surface deformation field [Mansinha & Smylie, 1972] using the average  
114 slip value within the block. Because smaller earthquakes require less accumulated stress over  
115





120 **Figure 1.** Scaled dislocation fields at the ocean floor calculated from the locking model for (a)  $M_w=8.0$ , (b)  $M_w=8.5$ ,  
 121 (c)  $M_w=9.0$ , and (d)  $M_w=9.2$ . Meshes show sample north-south rupture scenarios. Black arrows denote the  
 122 direction of rupture propagation from the hypocenters marked by yellow stars. <sup>n</sup>er,  
 123 the main trends of coastal amplitude distribution are not affected by the choice of scaling  
 124 equations (see supplementary material).

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133 We construct a field of ocean floor dislocations for six magnitudes:  $M_w=7.5, 8.0, 8.5, 9.0,$   
 134  $9.1$  and  $9.2$  (Fig. 1). The two largest magnitudes are selected to obtain better resolution of the  
 135 tsunami hazard for the potentially largest events. The hypocenter is not positioned at the trench to  
 136 guarantee that the rupture nucleates between the surface and the base of the seismogenic zone  
 137 near  $\sim 30$  km depth [Wang & Tréhu, 2016]. For each magnitude we start rupture scenarios at the

138 northernmost block for the chosen geometry and propagate the rupture southward until it is large  
139 enough to yield the desired magnitude. Rupture propagates along strike and dip with speeds  
140 between 2 - 3 km/s before reaching the bottom of the seismogenic zone, mimicking an elliptical  
141 rupture. This process is then repeated by moving the hypocenter one block south, resulting in a  
142 new scenario. This approach leads to 118 rupture scenarios. In our model, ruptures of  $M > 8.5$   
143 earthquakes primarily propagate along strike because the down-dip rupture extent saturates,  
144 resulting an elongated rupture area (Fig. 1). This process yields a frequency-magnitude  
145 distribution similar to the Gutenberg-Richter distribution due to the constraints resulting from  
146 rupture areas on the overall seismic moment [Stein & Wyssession, 2003; Fig. S2]. Our arbitrary  
147 choice of rupture directivity (north to south) has little effect because of the proximity of the  
148 coastlines to the trench ( $\sim 150$  km) and the large dominant wavelength of the tsunami  
149 [Rabinovich, 1997] prevent the rupture duration from significantly affecting the coastal tsunami  
150 amplitudes. Numerical experiments with various modes of rupture in flat oceans and near  
151 coastlines of different morphologies reveal that the contribution of rupture directivity is not  
152 significant in the near field (Fig. S3), as also shown by Williamson *et al* [2019].

153

## 154 **2.2 Tsunami Simulation Method**

155 We simulate tsunamis from each scenario using the MOST algorithm [Titov *et al*, 2016]  
156 that solves the fully nonlinear version of the shallow water approximation of the Navier-Stokes  
157 equation. MOST has been extensively validated through comparisons with laboratory and field  
158 data using standard international protocols [Synolakis *et al*, 2008]. Simulations are performed for  
159 4.5 hr time windows, allowing the tsunami to propagate along the entire coast. These simulations

160 use 0.5 s time steps to satisfy the stability-resolution requirement [Courant *et al*, 1928]. Wave  
161 height calculations are truncated at a depth of 30 m along the coastlines (typically at a distance of  
162 1-4 km) to avoid nonlinear shoaling effects, especially in the presence of large offshore  
163 deformation values. Therefore, no run-up values are calculated. Although run-up typically  
164 increases the tsunami hazard for generally linear coastal bathymetry and in the absence of bays,  
165 the distribution of coastal amplitudes (at shallow depth) will almost match that of run-up [Plafker,  
166 1997]. We calculate time histories of tsunami amplitudes at 100 virtual gauges along the coastline  
167 at a depth of ~35 m and one gauge at the entrance of the Strait of Juan de Fuca, north of Fork. We  
168 use GEBCO bathymetry interpolated to a spatial resolution of 18 arc-seconds to ensure enough  
169 (~20) grid points per wavelength [Shuto *et al*, 1986].

170 MOST was developed to model static sources and cannot be directly applied to kinematic  
171 ruptures [Titov & Synolakis, 1998]. Hence, we apply it to the discretized rupture blocks each of  
172 which happens at time  $t_i$  after the origin time. For each block, calculations are terminated after a  
173 duration  $\Delta t$  and the outputs are fed into MOST as initial conditions for the next block. The  
174 process continues until rupture ends after which the problem turns into a regular tsunami  
175 propagation. Although this approach introduces discontinuities in both water surface elevation  
176 and velocity, it produces results comparable to those from fully kinematic algorithms such as  
177 GeoClaw [Berger *et al*, 2011, González *et al*, 2011] that have been verified for Cascadia [Melgar  
178 *et al*, 2016]. The discrepancy between our results and those from the previous studies is largely  
179 due to the scaling equations. As discussed earlier (Fig. S1) Geller's [1976] scaling equations  
180 predict larger slip for a given rupture. While kinematic rupture properties do not significantly  
181 affect near-field tsunami propagation, we consider kinematic ruptures for a more comprehensive  
182 view of tsunami behaviors.

### 183 3. Cascadia Earthquake and Tsunami Scenarios

#### 184 3.1 Tsunami Simulation Results

185 We analyze the tsunami hazard for our rupture scenarios (Fig. 1) by simulating the  
186 resulting tsunamis. Our magnitude range of  $M_w=7.5-9.2$  accommodates both the largest expected  
187 rupture and the smallest rupture with noticeable ocean floor deformation ( $\sim 1\text{m}$ ). Note that the  
188 smallest rupture does not necessarily reach the ocean floor. Tsunami simulations for the rupture  
189 scenarios are shown in Fig. 2 (also see supplementary material). We find that various earthquake  
190 magnitudes can create similar tsunami amplitudes at a given location depending on the position of  
191 the hypocenter. For example, as shown in Figs. 2b and 2c, a  $M_w=8.5$  earthquake with a  
192 hypocenter in the south ( $\sim 43^\circ\text{N}$ ) and a  $M_w=9.0$  earthquake starting further north ( $\sim 48^\circ\text{N}$ ) both  
193 produce tsunami amplitudes of  $\sim 8\text{ m}$  at Newport ( $\sim 44.5^\circ\text{N}$ ). This partly reflects geometrical  
194 spreading wherein the energy flux in the propagating tsunami decreases with increasing distance.  
195 However, as seen in Fig. 2, this trend is not monotonic.

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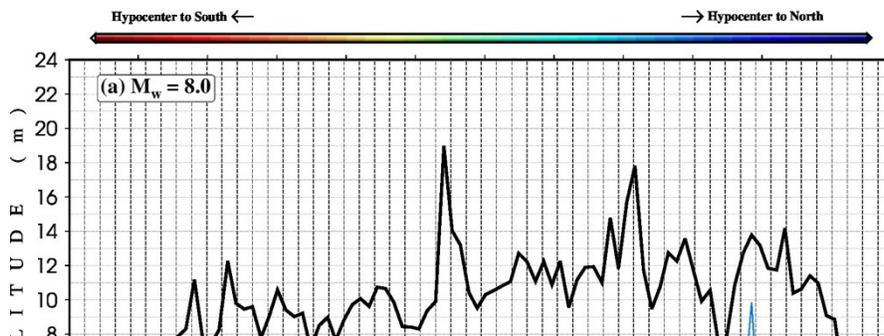
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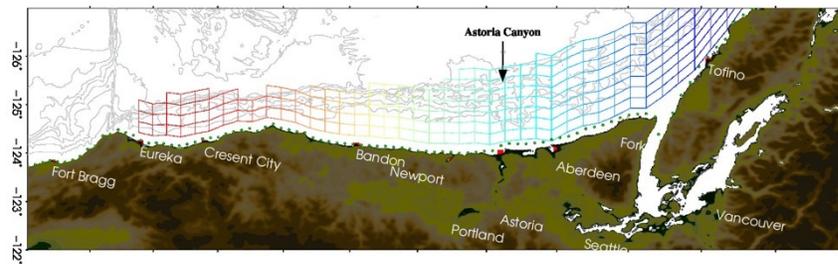
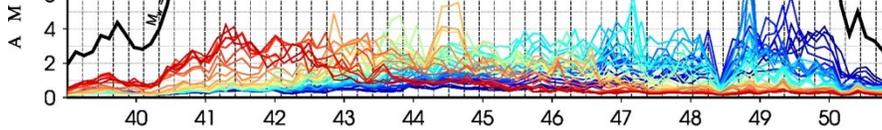
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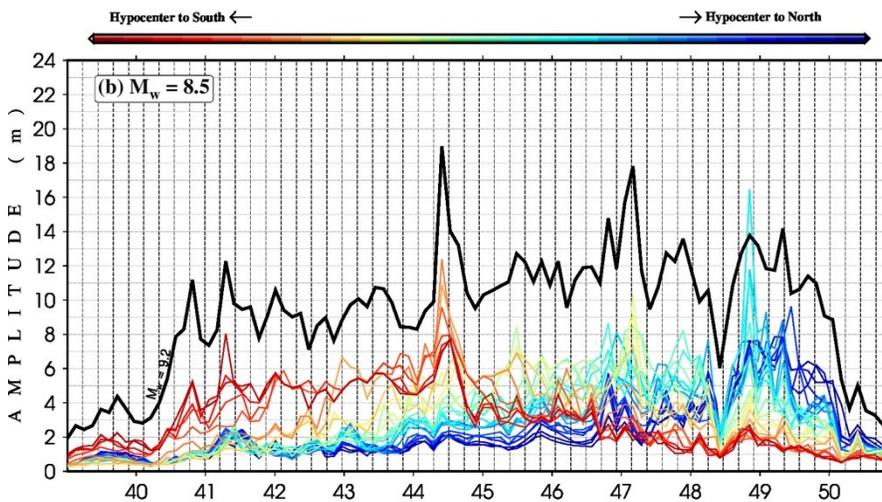




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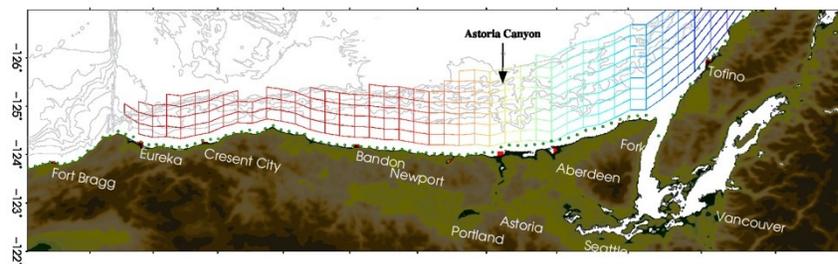


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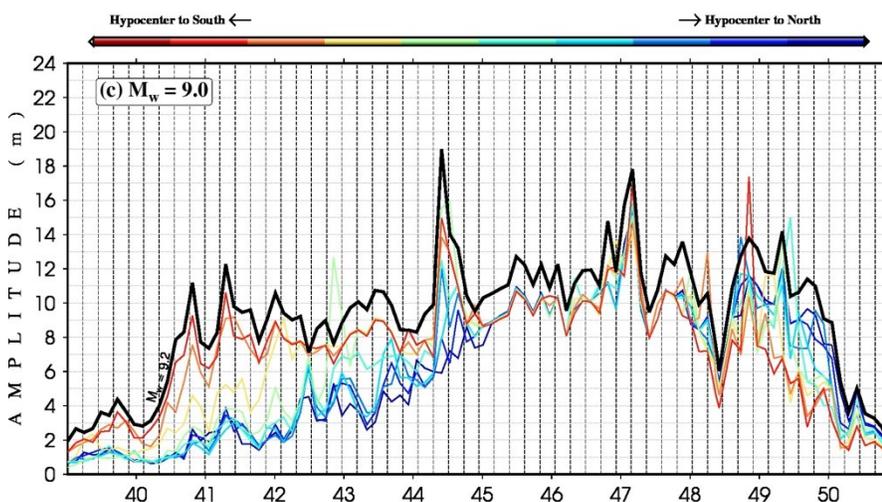
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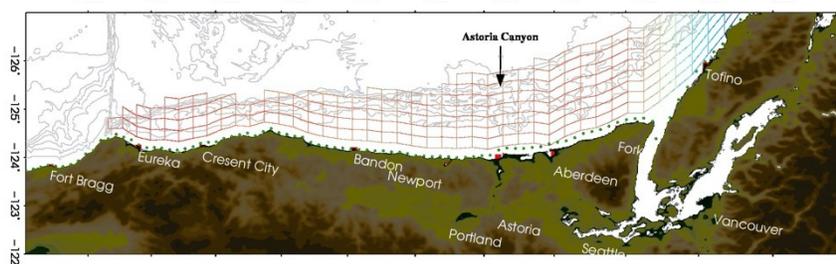
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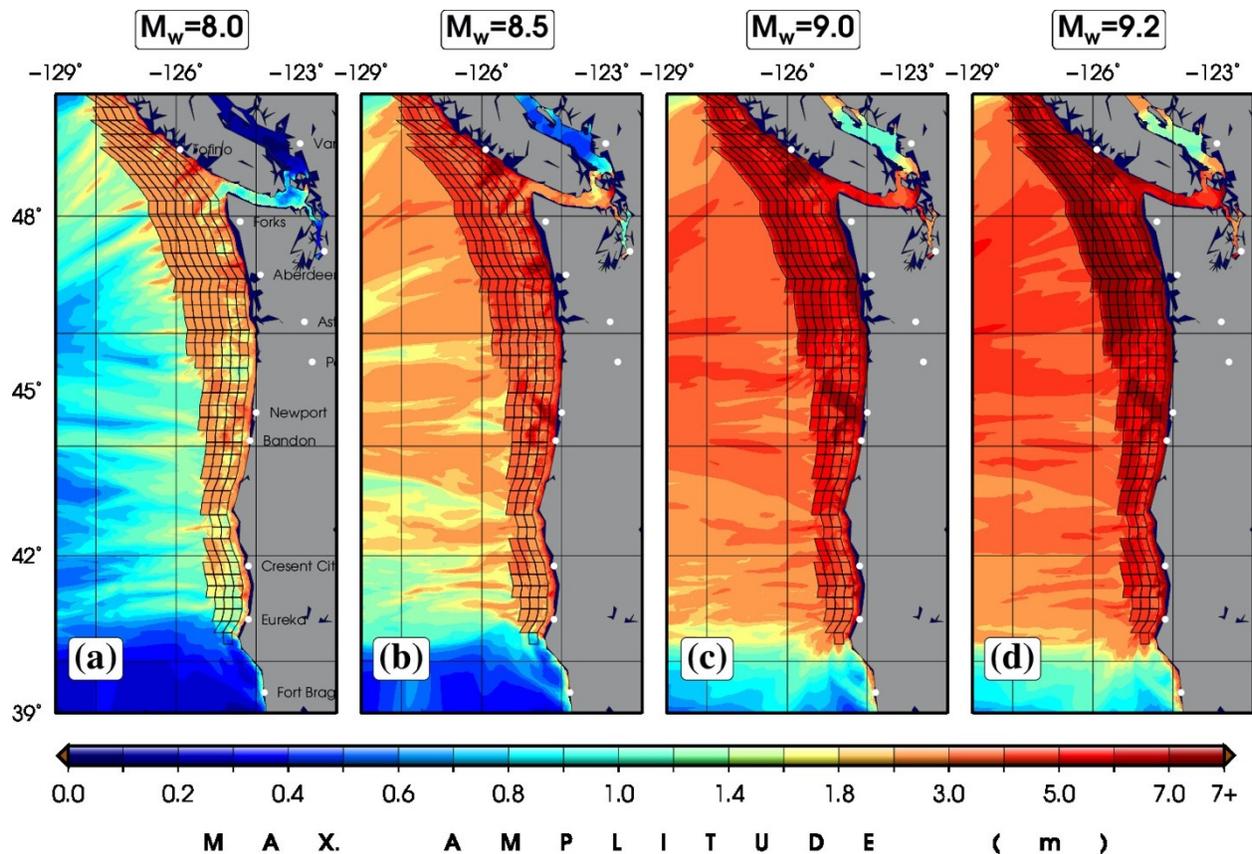


The curves are color-coded according to hypocenter latitude with hot colors in the south and cold colors in the north. Black curve shows coastal tsunami amplitudes for the  $M_w=9.2$  event; [Bottom] Cascadia coastline and the rupture blocks are colored to provide a sense of hypocenter latitude. Gray contours indicate bathymetry. Large population centers are shown for reference.

### 213 3.2 Effects of Coastal Geometry

214 Spatial variation of near-shore tsunami amplitudes from large sources mostly reflects the  
215 influence of piecewise coastal slopes [Kânoğlu & Synolakis, 1998], because the near-shore  
216 bathymetry of Cascadia varies little with latitude (with the exception of the Astoria Canyon  
217 [Griggs & Kulm, 1970]). In fact, bathymetric profiles across all latitudes within 400 km from  
218 shoreline have a correlation coefficient of  $\approx 0.75$ . Hence, in the absence of major bathymetric  
219 features the largest amplitudes occur at the latitudes with the largest earthquake slip, due to  
220 geometrical spreading and directivity [Ben-Menahem & Rosenman, 1972; Aki & Richards,  
221 2002]. The latter causes the waves to interfere constructively in a direction perpendicular to the  
222 rupture, focusing tsunami energy onto the closest shorelines (Fig. 3).

223



12 **Figure 3.** Cumulative maximum tsunami amplitudes for (a)  $M_w=8.0$ , (b)  $M_w=8.5$ , (c)  $M_w=9.0$ , and (d)  $M_w=9.2$  source scenarios. Each panel shows the largest tsunami amplitudes across all rupture scenarios for a given magnitude.

224

225           The simulations reveal that in the absence of significant local bathymetric features, the  
226 concave geometry of the coast between 43° - 48°N concentrates amplitudes in central Cascadia  
227 (between 44° - 45°N; around Newport, Oregon) especially from ruptures in central Cascadia, in  
228 agreement with edge wave theory [Munk *et al*, 1956]. We carried out numerical tsunami  
229 simulations in flat oceans along a narrow, shallow continental shelves to study the effects of a  
230 coastline curvature (Figs. S3 and S4) on tsunami amplitudes. These experiments show that  
231 coastline concavity increases the tsunami energy in the nadir (here, mid-latitudes) by focusing the  
232 energy of edge wave modes along the coast, on the continental shelf. Another amplitude peak  
233 offshore, which approaches the shoreline by increasing curvature, results from the concentration  
234 of tsunami reflection at the focal point of the curved shoreline. The cluster appears at half the  
235 radius of coastline curvature (i.e., focal point) of coastline analogous to that predicted by  
236 geometric optics for concave mirrors (see Fig. S4 in supplementary material, and the  
237 supplementary video SV1).

238           Given the shoreline's large radius of curvature (~1000 km), the former effect is more  
239 pronounced and can increase coastal amplitudes in central Cascadia by more than ~10%. We  
240 attribute the relatively larger amplitudes near Oregon in all the scenarios (Fig. 2) to this  
241 phenomenon. Although this effect makes Oregon coast almost as hazardous as northern regions  
242 (near Washington), it does not violate the generalizations that smaller earthquakes create smaller  
243 tsunamis, and that shorelines closer to large fault slip experience larger tsunami waves, which are  
244 generally true for linear coastlines and in the absence of bays [Davies *et al*, 2018] (Figs. S3 to  
245 S5).

246 Another interesting simulation result is the apparent relative immunity of northern  
247 California coastlines to Cascadia tsunamis, especially from  $M_w < 9.0$  earthquakes with hypocenters  
248 in the north. At first glance, this is surprising given to the region's proximity to a slip cluster near  
249 the southern tip of the rupture. However, it results from both the end of rupture and the convex  
250 promontory near Eureka (Fig. 2), south of where coastal amplitudes drop. Simulations suggest  
251 that the coastal morphology creates and scatters free edge waves (supplementary material),  
252 making the California shorelines virtually sheltered.

### 253 3.3 Choice of Slip Model

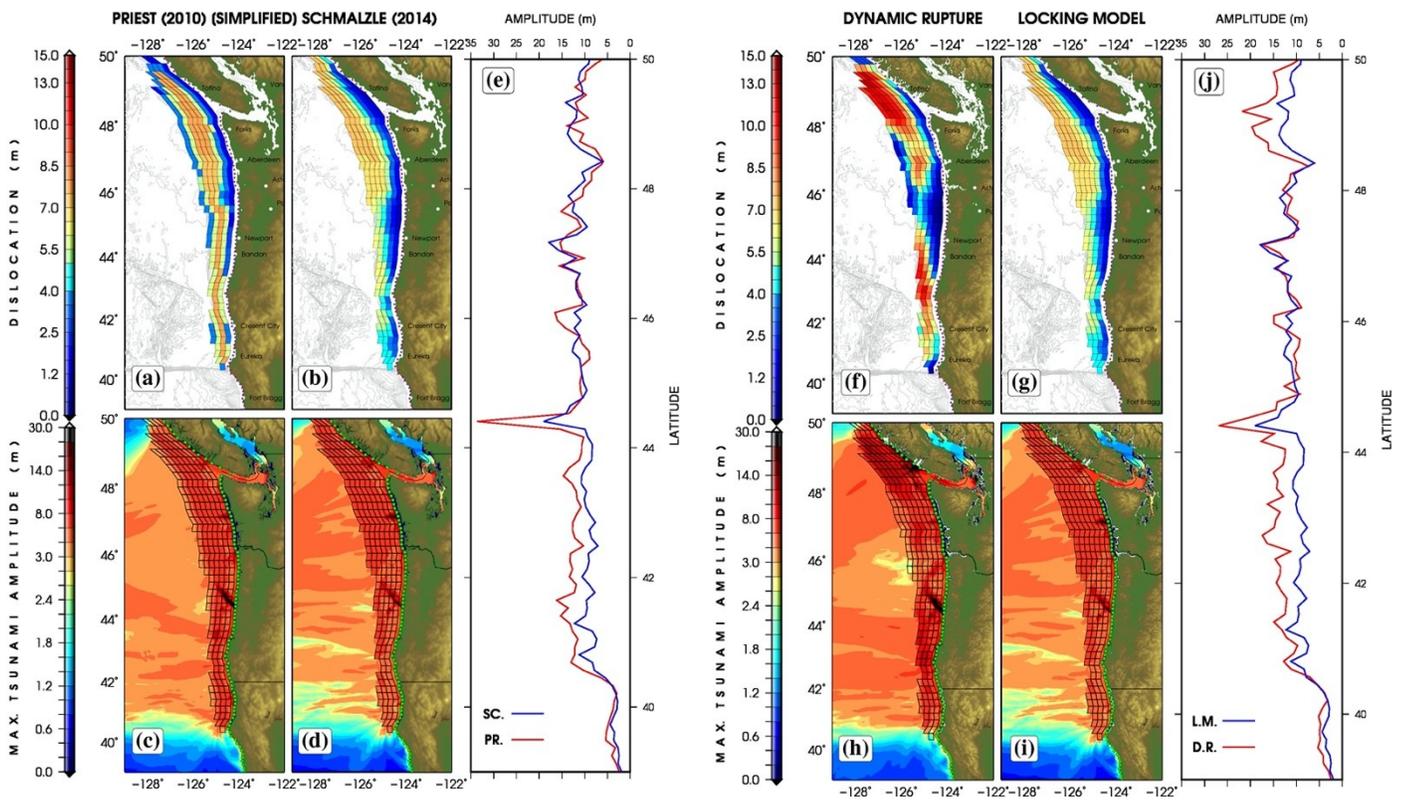
254 The distribution of potential slip from the locking model shows two main clusters (Fig.  
255 1d). The larger cluster (both in area and slip magnitude) is in the north, close to British Columbia,  
256 and the smaller one is around  $44^\circ\text{N}$ . However, the tsunami simulation results are not significantly  
257 affected by the choice of slip models on a regional scale. We simulate tsunamis using a simple  
258 slip model [modified from Priest *et al*, 2010] and found similar results to those from our more  
259 physical model. As shown in Fig. 4a-e, the coastal amplitudes show a correlation coefficient of  
260 0.8. We also used a dynamic rupture model with identical recurrence intervals, derived from a  
261 Gaussian locking model [Schmalzle *et al*, 2014; Ramos *et al*, 2021]. Such models consider the  
262 dynamic interaction of fault stresses and frictional strengths, and near-trench slip can be amplified  
263 due to constructively interfering free-surface reflections within the accretionary wedge. The  
264 resulting tsunami simulations yield a similar (correlation coefficient  $CC=0.8$ ) distribution of  
265 coastal amplitudes (Fig. 4j).

266 The absolute values of coastal tsunami amplitudes from these simpler models can locally  
267 vary from our modeling results by up to 30%. However, the general trend of tsunami amplitudes

268 remains similar. In the absence of conclusive geodetic and seismic constraints on fault locking,  
 269 we think our model adequately represents potential future ruptures and consider the 30%  
 270 discrepancy as illustrating the uncertainty.

271

272



**Figure 4.** (a) A simplified locking model [Priest *et al.*, 2010] produces similar tsunami amplitudes to those from our model shown in (b), both in the Pacific (c,d) and along the coastline (e). Also, a dynamic  $M_w=9.2$  rupture derived from the Gaussian locking model (f) and our choice of locking model (g) result in similar tsunami amplitudes both in the Pacific (h,i) and along the coastline (j).

274 Simulation of the tsunami from a perturbed version of our choice of locking model  
275 (created by introducing white noise equal to 50% of the maximum to the deformation field of  
276 parent model) yields a significantly different distribution ( $CC=0.3$ ) of coastal tsunami amplitudes  
277 (Fig. S6). We attribute this discrepancy to the disruption of large-scale slip clusters which  
278 changes the dominant period of the tsunami. Such smaller wavelengths significantly alter the  
279 interaction of the tsunami with the shoreline, thus resulting in a different pattern of coastal  
280 amplitudes. Otherwise, given the similar bathymetry along strike, different rupture models with  
281 comparable dimensions of slip clusters (barring an absence of large slip deficit in central Cascadia;  
282 Li *et al.*, 2018) would result in similar tsunamis.

283

## 284 **4 Discussion and Conclusions**

285 Simulations of tsunamis from physics-based  $M_w=7.5-9.2$  earthquake rupture scenarios  
286 show that largest and most widespread coastal tsunami amplitudes result from ruptures at or  
287 starting from mid-latitudes in central Cascadia. This result is almost independent of the choice of  
288 slip model as long as the dimensions of major slip clusters are preserved. Such ruptures,  
289 especially with  $M_w>8.5$ , can create tsunami amplitudes exceeding 50% of those from the largest  
290 expected  $M_w=9.2$  rupture (Fig. S7a). Statistical analysis using the metric  $MT$  [Salaree & Okal,  
291 2020] suggests that the near-field propagation patterns of tsunamis from  $M_w>8.5$  events are very  
292 similar (supplementary material). This effect is important because realistic estimates of the  
293 expected loss are valuable in designing mitigation policy [Stein & Stein, 2014]. Although smaller  
294 earthquakes generate smaller tsunamis, their expected amplitudes (up to 12 m using our choice of

295 scaling law) are significant. Thus, smaller earthquakes that are more likely to occur in the near  
296 future may create comparable – though more localized – damage than the less frequent worst-case  
297 scenario [e.g., Priest *et al*, 2010; Thio *et al*, 2010].

298         We also find that the large-scale morphology of Cascadia's coastline focuses and  
299 defocuses tsunami energy. Simulations (Fig. S4 and video SV1) show that coastline curvature can  
300 increase the coastal tsunami amplitude by more than 10%. Comparison of Cascadia with other  
301 subduction zones where coastal curvature is insignificant (i.e., Chile with curvature of  $\sim 0.017$ )  
302 shows why such heightened tsunami amplitudes are not observed in these regions. Our  
303 simulations show that the southernmost sites (Fort Bragg, Eureka and Crescent City, Figs. 2 and  
304 3) show almost no change in the tsunami amplitudes for events with increasing magnitude above  
305  $M_w=9.0$  (Fig. 2), due to the large promontory at  $\sim 42.5^\circ\text{N}$  separating the concave coastline from  
306 that to the south (Fig. S5a). Examination of the along-strike tsunami amplitudes (Fig. 3) reveals  
307 that the relative differences throughout the simulation area are small (sometimes  $<1\text{m}$ ) for  
308 earthquakes larger than  $M_w=9.0$ .

309         Our findings have implications for similar tectonic settings such as the Chile and Alaska  
310 subduction zones that have experienced large and heterogeneous megathrust ruptures. The  
311 bathymetry in these regions is also similar to Cascadia, i.e., almost uniform bathymetric slopes  
312 and large-scale geometric coastal morphology, in contrast with regions of more complex  
313 bathymetry such as Japan. Similarly, by identifying the most hazardous segments of the  
314 subduction zone, our results can be used to assist in selecting sites for DART tsunami sensors or  
315 novel technologies such as SMART cables [Howe *et al*, 2019]. This is important for near-field

316 tsunami warning because these instruments are mostly deployed on the up-dip side of trenches,  
317 where the identification of the main areas contributing to the tsunami hazard is crucial.

## 318 **Acknowledgments**

319           The manuscript significantly benefited from invaluable discussions with Jean-Paul  
320 Ampuero, Amanda Thomas and Kelin Wang. Some figures were drafted using the Generic  
321 Mapping Tools [Wessel & Smith, 1998]. This study was supported by a National Science  
322 Foundation grant (PREEVENTS geosciences directorate No. 1663769).

## 323 **Data and materials availability**

324           All bathymetry data used in the main text or the supplementary materials are available via  
325 the General Bathymetric Chart of the Oceans (<https://www.gebco.net>). The tsunami simulation  
326 code is maintained and distributed by the US National Oceanic and Atmospheric Administration  
327 (<https://nctr.pmel.noaa.gov/nthmp/>). Rupture data and tsunami simulation results are available via  
328 Deep Blue Data at <https://doi.org/10.7302/xe96-3z26>

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