Observing Upper Ocean Stratification during Strong Diurnal SST Variation Events in the Suppressed Phase of the MJO

Je-Yuan Hsu¹, Ming Feng², and Susan Wijffels³

¹National Taiwan University ²CSIRO Oceans and Atmosphere ³Woods Hole Oceanographic Institution, Physical Oceanography

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Abstract

Six ALAMO floats are deployed within the tropical warm pool of the eastern Indian Ocean, to study the thermal stratification in the diurnal warm layer (DWL) during strong diurnal SST variation (DV SST) prior to the onset of Madden-Julian Oscillations (MJO). Strong DV SST of > 2 °C is measured by four floats before the passage of a MJO event (i.e., during the suppressed phase), when the peak insolation > 1000 W m⁻² and the wind speed < 3 m s⁻¹. Even after the occurrence of daytime peak SST, the temperature gradient in the DWL can still extend to > 10 m until the midnight, which may be driven by the turbulent mixing at the base of DWL. Interestingly, the foundation SST (SST_{fnd}) at three floats increases rapidly from 26.4 °C to > 27.6 °C over two days, coincident with the shoaling of surface mixed layer depth (MLD) by more than 20 m. The strongly stratified near surface layer may sustain higher SSTs and enhance air-sea heat fluxes until the onset of stronger winds. The KPP mixing scheme used in a 1-D model can simulate the observed DV SST magnitude reliably, but fail to predict the rapid increase of SST_{fnd}. The magnitude of DV SST is affected by the near surface stratification, but the SST_{fnd} is modulated by the evolution of stratification above the MLD. Future field measurements in the upper ocean during diurnal warming are proposed to help improve air-sea flux simulations and the forecast of MJOs.

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4	Je-Yuan Hsu ¹ , Ming Feng ^{2, 3} and Susan Wijffels ⁴
5	¹ Institute of Oceanography, National Taiwan University, Taipei, Taiwan
6	² CSIRO Oceans and Atmosphere, Perth, Australia
7	³ Centre for Southern Hemisphere Oceans Research (CSHOR), Hobart, Australia
8	⁴ Woods Hole Oceanographic Institution, Woods Hole, USA
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11	
12	
13	
14	Corresponding Author: Je-Yuan Hsu
15	Email: jyahsu@ntu.edu.tw
16	
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22 Abstract

23 Six ALAMO floats are deployed within the tropical warm pool of the eastern Indian Ocean, to 24 study the thermal stratification in the diurnal warm layer (DWL) during strong diurnal SST 25 variation (DV SST) prior to the onset of Madden-Julian Oscillations (MJO). Strong DV SST of > 2 °C is measured by four floats before the passage of a MJO event (i.e., during the suppressed 26 phase), when the peak insolation > 1000 W m⁻² and the wind speed < 3 m s⁻¹. Even after the 27 28 occurrence of daytime peak SST, the temperature gradient in the DWL can still extend to > 10 m 29 until the midnight, which may be driven by the turbulent mixing at the base of DWL. 30 Interestingly, the foundation SST (SST_{fnd}) at three floats increases rapidly from 26.4 $^{\circ}$ C to > 27.6 31 °C over two days, coincident with the shoaling of surface mixed layer depth (MLD) by more 32 than 20 m. The strongly stratified near surface layer may sustain higher SSTs and enhance air-33 sea heat fluxes until the onset of stronger winds. The KPP mixing scheme used in a 1-D model 34 can simulate the observed DV SST magnitude reliably, but fail to predict the rapid increase of 35 SST_{fnd} . The magnitude of DV SST is affected by the near surface stratification, but the SST_{fnd} is 36 modulated by the evolution of stratification above the MLD. Future field measurements in the 37 upper ocean during diurnal warming are proposed to help improve air-sea flux simulations and 38 the forecast of MJOs.

40 **1. Introduction**

41 Air-sea heat fluxes over high sea surface temperatures (SST) of the tropical warm pools 42 (TWPs) have a critical influence on the atmospheric general circulation. These TWPs frequently 43 feature low wind conditions and thus experience ubiquitous and strong diurnal variations of SST 44 (DV SST). Latent and sensible heat fluxes, modulated by the SST, can affect the onset and 45 timing of intra-seasonal weather systems such as the MJOs (Zhang 2005; Maloney 2009; Seo et 46 al. 2014; Sobel et al. 2014). While the intrinsic time scale of MJOs is longer than a week, the 47 coupled model forecasts of MJOs can be influenced significantly by DV SSTs through impacts 48 on daily-mean SST (Bernie et al. 2008; Rupert and Johnson 2015; Demott et al. 2015). Modeling 49 the diurnal variations of the upper ocean remains challenging (Kawai and Wada 2007). To 50 forecast the DV SST accurately, turbulent mixing in the upper ocean must also be simulated 51 accurately, which involves the forecast on both the shear and density stratification for inducing 52 shear instability mixing. Therefore, exploring the diurnal SST variations and associated 53 evolution of stratified layers near the ocean surface may be crucial for improving the MJO 54 forecast.

55 The evolution of stratified layers (Fig. 1) above the seasonal thermocline during the DV 56 SST has been discussed in several previous studies such as Brainerd and Gregg (1993a) and 57 Sutherland et al. (2016). Under low wind conditions, the absorption of insolation during the daytime forms a sharp vertical temperature gradient in the surface mixed layer, termed the 58 59 diurnal thermocline (Kudryavtsev and Soloviev 1990; Caldwell et al. 1997). The diurnal warm layer (DWL, Sui et al. 1997; Matthews et al. 2014; Sutherland et al. 2016; Moulin et al. 2018; 60 61 Hughes et al. 2020), spanning the ocean surface to the base of diurnal thermocline, can modify 62 the magnitude of DV SST through the inhibition of turbulent mixing across its base (Fairall et al.

63 1996b; Bellenger and Duvel 2009; Moulin et al. 2018). Below, the layer between the DWL and 64 the top of the seasonal thermocline is termed the remnant layer (RL, Brainerd and Gregg 1993a; 65 Caldwell et al. 1997), which can be regarded as a fossil mixed layer formed during the previous 66 night due to convective cooling. The penetrative solar radiation can not only form the DWL, but 67 also restratify the RL (Brainerd and Gregg 1993a). After a late afternoon peak, the SST drops 68 around sunset when outgoing turbulent heat fluxes and longwave radiation exceed the insolation 69 (Moulin et al. 2018), and thus induces nighttime convective mixing. The induced convective 70 mixing will then mix the surface mixed layer by destratifying the DWL and RL (Fig. 1c). It will 71 set the foundation SST (SST_{fnd}, which is closed to the nighttime minimum SST) similar with that 72 of the previous night.

73 Consecutive days of strong insolation and low wind speeds is one of the most important 74 features of the suppressed phase of MJO. This can result in the stratification of surface mixed 75 layer (e.g., Bernie et al. 2005; Moum et al. 2014) and generate higher daily-mean SST (Shinoda 76 and Hendon 1998; Bernie et al. 2005). The insolation captured in the DWL will be redistributed 77 through the deeper surface mixed layer by nighttime convective mixing. This heat can gradually 78 increase the SST_{fnd} and the stratification across the base of the DWL (Sui et al. 1997). In turn it 79 may suppress cold entrainment, allow SSTs to increase and through associated air-sea fluxes, 80 drive up the accumulation of atmospheric boundary layer heat and moisture for developing the 81 deep convection in MJOs (Zhang and Ling 2017)

The concept of an ocean barrier layer (BL) is first proposed (Lukas and Lindstrom 1991) to identify the discrepancy between the surface mixed layer depth (MLD) and the top of seasonal thermocline, mostly due to opposing salinity stratification. The BL is later defined as the layer between the isothermal layer depth (ILD) and MLD (Sprintall and Tomczak 1992; McPhaden

86	and Foltz 2013; Chi et al. 2014). Strong precipitation may form a low-salinity layer near the sea
87	surface (as illustrated in Fig. 1d and h). The density stratification due to salinity will result in a
88	shallow MLD, while the temperature remains nearly homogenous (i.e., the definition of
89	isothermal layer IL) or even can be cooler at the surface compared to below (McPhaden and
90	Foltz 2013; Chi et al. 2014). The presence of salinity-based BL increases the required vertical
91	mixing to access the seasonal thermocline (Chi et al. 2014). On the other hand, the temporal
92	change of MLD is an useful indicator for identifying the strength of wind-driven mixing
93	(Balaguru et al. 2015), because the nighttime convective mixing cannot destratify the
94	temperature gradient in the seasonal thermocline abruptly. Thus, clearly salinity and temperature
95	gradients in the upper ocean are important to the dynamics of stratification and turbulent mixing.
96	Though the importance of DWL for inhibiting the turbulent mixing has been explored by
97	many previous studies (e.g., Moulin et al. 2018; Hughes et al. 2020), the thickness of DWL is
98	often estimated by finding the depth of an isotherm with high temperature (Matthews et al.
99	2014). However, because the density stratification modulated by the temperature structure can
100	affect the efficiency of turbulent mixing, exploring the factors to the extension of temperature
101	"gradient" in the DWL should be more crucial than focusing on the thickness of a high-
102	temperature layer. This study will discuss the evolution of temperature gradient above the MLD,
103	by using both the observations and model simulations during a strong DV SST event.
104	In November 2018, a collaborative field campaign between CSHOR and China's First
105	Institution of Oceanography was conducted to explore the air-sea interaction in the Indonesian-
106	Australian Basin of the TWP (Feng et al. 2020). A shelf version of the Bailong buoy system
107	(Appendix A) was deployed off the northwest coast of Australia, along with six rapidly profiling
108	ALAMO floats from MRV Systems and two Teledyne Web EM-APEX floats. The observations

off Australia's Northwest Shelf collected during the suppressed phase of one MJO event around
the end of 2018 (Feng et al. 2020) will be reviewed in section 2. Definitions and the estimation
of different upper ocean stratified layers during diurnal warming will be described in section 3.
Results and discussions on SST variations and evolution of upper ocean layers will be presented
in sections 4 and 5, respectively. Section 6 will compare the observations with the model results
using the K-profile parameterization (KPP) of vertical mixing.

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116 2. ALAMO Float Measurements in the Field Experiment

The field array was deployed at 115.3 °E and 16.8 °S on Nov 22nd 2018. Two ALAMO 117 118 floats (9206 and 9208) failed within two days of deployment. Floats 9205 and 9209 were 119 deployed at the southwest of floats 9207 and 9210. The distance between the floats was always less than 50 km before Dec 5th. The remaining four ALAMO floats initially drifted 120 northwestward (Fig. 2). The floats except 9205 turned east since Dec 2nd, but still remained in a 121 122 cold filament whose SST ~ 26.5 °C. The drifting velocity of the floats is similar with the 123 satellite-measured geostrophic current (not shown in this study). That is, the float trajectories 124 should be mainly affected by the geostrophic currents of a strong cold-core cyclonic eddy located 125 just east of the float array, not the wind-driven current. No rains fell until the passage of an MJO 126 around the middle of December 2018 (Feng et al. 2020). 127 Two types of CTD sensors, Seabird SBE-41 (9207) and RBR (9205, 9206, 9208, 9209 128 and 9210), were mounted on the ALAMO floats. The floats continuously profiled the temperature and salinity in the upper 600 m from Dec 1st to 5th 2018. Only the ascending profiles 129 130 are used for the analysis in this study, to avoid the CTD measurements being contaminated by

132 resolution of Seabird profiles was 1 m in the upper 50 m. The RBR sensors returned a vertical resolution of 0.1 m in the upper 5 m, and 1 m from 5 to 50-m depth. We estimated the SST on 133 134 the RBR profiles by finding the peak temperature in the upper 1 m, often at ~ 0.2 -m depth, after 135 excluding those with salinity measurements < 32 psu (which indicates air in the samples). 136 ALAMO floats also recorded the temperature during surface drift, when the pressure was less 137 than 0.2 dbar. Though the Seabird CTD sensors did not have the measurements in the upper 1 m, 138 the first surface value of temperature at about 0.1 or 0.2 dbar was always several degrees higher 139 than those reported at pressures < 0 dbar, similar with the typical difference between the air

the wake effect. The time interval between ascending profiles was 3 to 4 h. The vertical

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temperature and SST in the region. Thus, we used the first surface value of temperature on theSeabird surface reports as the SST at float 9207.

On Nov 28th 2018, the temperature in the upper 40 m at all four floats was about 26.3 °C 142 and is largely vertically mixed until Dec 1st (Fig. 3). The float-measured SST was ~ 0.2 °C 143 higher than the skin SST measured by the Himawari-8 satellites on Nov 30th (Fig. 2), consistent 144 145 with previous studies (Beggs et al. 2013). The SST measurements taken by Himawari-8 satellites were described in appendix A. A strong DV SST event occurred between Dec 2nd and 4th (section 146 147 4), ~ 10 days before the passage of the deep convection in one MJO (often termed active phase; 148 Feng et al. 2020) arrived in the region. Except for float 9205, the temperature in the upper 40 m increased up to 28 °C after Dec 3rd. Temperatures greater than 27 °C in the upper 40 m were then 149 sustained until the onset of the MJO's active phase around Dec 13th. Float 9205 measured 150 temperatures of > 26 °C down to 60-m depth on Dec 4th, deeper than that at the other three floats 151 152 at 40-m depth, while the salinity in the upper 40 m at float 9205 was higher than that at the other floats by over 0.3 psu after Dec 2^{nd} . 153

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- 155 3. Definitions of Upper Ocean Stratified Layers and DV SST 156 Various criteria for estimating the depth of the MLD have been proposed (Sprintall and 157 Roemmich 1999; de Boyer Montégut et al. 2004; Suga et al. 2004), such as a temperature 158 difference that near the ocean surface (Wyrtki 1964) or density gradient criteria (Lukas and 159 Lindstrom 1991). Some studies estimate the MLD by finding the difference of potential density $\Delta \rho$ between $\rho(MLD)$ and $\rho(z_0)$ exceeding some arbitrary constant (e.g., $\Delta \rho = 0.1$ kg m⁻³ in Chi et 160 al. 2014), where z_0 is the reference depth closed to the ocean surface. The definition of z_0 is 161 162 required to exclude the unknown spikes of density gradient due to turbulence near the sea 163 surface. 164 A key question is, which z_0 is shallow enough to represent the "surface" value of $\rho(z_0)$ 165 (Brainerd and Gregg 1995)? Different studies choose z_0 to avoid the impact of diurnal near 166 surface temperature stratification, e.g., $z_0 = 5$ m in McPhaden and Foltz (2013) or $z_0 = 10$ in de 167 Boyer Montégut et al. (2004), due to either an interest in longer timescales or limited vertical 168 observations. The daily development of temperature stratification in a DWL can complicate the 169 estimates of MLD (as the example in Fig. 1f), conflicting with the simple concept of a persistent 170 MLD at the top of seasonal thermocline. A more careful definition of z_0 needs to be used in the
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173 **3.1 Depth of stratified layers in a diurnal cycle**

presence of a DWL.

The absorption of insolation forms a DWL with strong temperature stratification in the surface mixed layer, developing within a few hours of sunrise (Moulin et al. 2018). Here, the diurnal warm layer depth is defined as the shallowest depth greater than 3 m (to avoid very near

surface turbulence) where vertical temperature gradients weaken to less than 0.02 °C m⁻¹ over a 177 178 5-m span. That means temperature gradient is small below the DWL. This criterion focuses on 179 illuminating the roles of strong near surface temperature gradients (including the diurnal 180 thermocline) in stabilizing the upper ocean (Moulin et al. 2018), instead of the thickness of a 181 high-temperature layer (e.g., Matthews et al. 2014; section 5). In other word, in this study, a 182 "thicker" DWL occurs when a strong temperature gradient induced by the diurnal warming (> 0.02 °C m^{-1}) extends more deeply. The uncertainties of estimating DWL depth due to the choices 183 184 of the values of temperature gradient and spanning depth are studied in supporting information 185 A.

The MLD and ILD are estimated by using the estimated DWL as the reference depth z_0 . We estimate the MLD by fulfilling two criteria at the same time: potential density difference $\Delta \rho$ $= \rho(MLD) - \rho(z_0) > 0.15 \text{ kg m}^{-3}$ (McPhaden and Foltz 2013) and potential density gradient $\partial \rho / \partial z$ $< -0.01 \text{ kg m}^{-4}$ (Brainerd and Gregg 1995), where the axis z is positive upward. The ILD is estimated in the same way as the MLD (Chi et al. 2014), but using the temperature-dependent profiles of ρ , which are computed assuming a constant salinity – here, the average of salinity in the upper 5 m.

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194 3.2 Temperature versus salinity-driven BL

The barrier layer (BL) described by Lukas and Lindstrom (1991) exists due to salinity stratification for driving the discrepancy between the MLD and ILD (Fig. 1d and h. McPhaden and Foltz 2013; Chi et al. 2014). They also report a DWL near the ocean surface during the daytime. Because the strong temperature stratification in the DWL stabilizes the upper ocean, the effect of the DWL on vertical mixing (Moulin et al. 2018) can be similar to that of near surface

200 freshening (Smyth et al. 1997), and thus form a BL. Here, we will define a temperature-driven 201 BL (TBL) and salinity-driven BL (SBL). The TBL is identified if DWL > 10 m. The SBL is 202 identified if ILD - MLD > 15 m (see supporting information A for more details on the criteria). 203 Example profiles of salinity and temperature are used to illustrate the difference between 204 TBL and SBL (Fig. 4). A sharp diurnal thermocline with a temperature gradient extending to > 20-m depth, i.e., a thick DWL, is captured by float 9210 in the afternoon on Dec 2nd. Strong 205 206 density stratification in the TBL is due to the absorption of insolation. In comparison, at float 9209, the salinity stratification between 20 and 40-m depth around the 7am of Dec 3rd leads to 207 the separation of the ILD and MLD, due to a shallower SBL. The TBL and SBL (Chi et al. 2014) 208 209 driven by temperature and salinity stratification, respectively, may both increase the upper ocean 210 stratification, and thereby the threshold for the shear needed to induce vertical mixing.

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212 3.3 Magnitude of DV SST and foundation SST

213 The Group for High Resolution Sea Surface Temperature (GHRSST) defines the 214 foundation SST (SST_{fnd}) as the SST not affected by the diurnal variability, or the SST before 215 solar heat gain begins early in the day. Estimating the SST_{fnd} is difficult without using reliable 216 measurements of air-sea heat fluxes or models for predicting diurnal variation (Zhang et al. 217 2016). Several studies find the SST_{fnd} by averaging the nighttime SSTs before sunrise - from 12 218 to 5:30 am (Karagali and Høyer 2014; Zhang et al. 2016). Following this, we define SST_{fnd} as the 219 mean SST from 1 to 5 am. The magnitude of DV SST is then the difference between the SST_{fnd} 220 and following peak SST (SST_{max}).

222 4. SST Warming and Air-sea Heat Fluxes

223	We compute the SST_{fnd} and the magnitude of DV SST from Dec 1 st to 4 th 2018 (section
224	3), by using the float measurements of SST. On Dec 1^{st} , the observed SST_{fnd} at all floats is
225	similar, ~ 26.4 °C, and reaches the peak (~ 26.9 °C) at around 3 - 4 pm. Compared with the
226	magnitude of DV SST reported by the previous studies (e.g., from 0.5 to 1.3 °C in Moulin et al.
227	2018), the magnitude of DV SST ~ 0.5 °C on Dec 1^{st} is not extremely high (Fig. 5b-e). From Dec
228	2^{nd} to 3^{rd} , the significant DV SST > 2 °C occurs at all floats, termed a strong DV SST event in
229	this study. The highest SST of 29 °C occurs at the float 9209 on Dec 3 rd .
230	Air-sea heat fluxes are computed (Fig. 6) using the atmospheric measurements from the
231	surface buoy and float-measured SST (supporting information B), based on the COARE 3.0
232	algorithm (Fairall et al. 1996a; Fairall et al. 2003). The trend of daily-mean air-sea net heat flux
233	is consistent with that of upper ocean heat storage rate. The latent (LH) plus sensible heat flux
234	(SH) drops from 220 to 80 W m ⁻² between Dec 1^{st} and 3^{rd} , because the wind speed decreases
235	down to 2 m s ⁻¹ (Fig. 5a). The peak of downward shortwave radiation before Dec 2^{nd} is already >
236	900 W m^{-2} (appendix A), and low vertical shear of horizontal current is observed at the
237	surrounding EM-APEX float (not shown in this study), presumably due to the low wind speed.
238	That is, this strong DV SST event in the beginning of December mainly results from the
239	decreasing wind speed since Dec 1 st , not increasing insolation.
240	The warming of SST_{fnd} varies significantly between each float from Dec 2 nd to 4 th (Fig.
241	5), even when their separation is less than 50 km. The SST_{fnd} at floats 9207 and 9210 increases
242	from 26.6 to 27.7 °C from Dec 2^{nd} to 3^{rd} in one day, right after the significant DV SST > 2 °C on
243	Dec 2^{nd} . The SST _{fnd} at float 9209 increases from 26.4 to 27.6 °C between Dec 2^{nd} and 4^{th} . The

244 SST_{fnd} at float 9205 is ~ 27 °C, smaller than the other three floats by 0.8 °C on Dec 4th. The

equivalent mean net heat flux at float 9205 is therefore ~ 8.4 W m⁻² lower than that at the floats 9207 and 9210 from Dec 2^{nd} to 4^{th} (Fig. 6).

On the other hand, though the LH is suppressed by the low wind speed from Dec 2nd to 4th, the SST_{fnd} warming will eventually favor the LH once the wind speed increases (Hsu et al. 2019), e.g., the LH is up to -220 W m⁻² at wind speed > 6 m s⁻¹ on Dec 6th (not shown in this study). The relative humidity rises to more than 80 % after Dec 7th as well (appendix A). In other words, the warming of SST_{fnd} in a short time period may largely enhance the "efficiency" for accumulating of air-sea heat fluxes and moisture during the suppressed phase of MJOs (Maloney 2009).

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255 5. Upper Ocean Stratified Layers during Strong DV SST

256 5.1 Evolution of DWL

257 The strength of the ocean stratification is tracked through computing the Brunt-Väisälä frequency N² (=-(g/ ρ)($\partial \rho/\partial z$), where g is the gravity constant). The DWL is thin (< 10 m) on Dec 258 1^{st} , associated with a small DV SST, ~ 0.5 °C. Starting from Dec 2^{nd} , the solar insolation forms a 259 260 "thick" DWL as a TBL that extends its temperature gradient to more than 20-m depth, resulting in strong density stratification ($N^2 > 1.0 \times 10^{-4} \text{ s}^{-2}$) near the ocean surface (Fig. 5). Note that the 261 262 criteria for estimating the DWL depend on only the temperature gradient. The extension of N^2 263 results from the vertical structure of temperature instead of salinity. We further compare the estimates of DWL with the other definition of diurnal warm laver 264

265 DWL* (with the superscript *) in the aspect of a high-temperature layer. The DWL* can be found

266	at the depth of an isotherm in each day, whose temperature T^* equals to $\alpha^*SST_{max} + (1-\alpha)SST_{fnd}$
267	or is at least 0.1 °C higher than the SST _{fnd} , assuming the $\alpha = 0.3$ (Matthews et al. 2014).
268	Considering the variations of SST_{max} and SST_{fnd} between different days, the DWL^* can only be
269	used for comparing with the DWL in the individual day. During the strong DV SST event (Fig.
270	7), both DWL and DWL^* reach the peak in the evening or at the midnight (after 7 pm each day),
271	within several hours after the occurrence of SST_{max} in the afternoon (~4 pm each day; section 4).
272	The turbulent diffusivity after the diurnal peak transports the warm water and extents the
273	temperature gradient to the deeper layer until the midnight. The DWL can be a reliable indicator
274	for identifying the strong DV SST in the consecutive days.
275	When the measured DV SST at all floats is > 2 °C (Fig. 5), not only a thin layer with high
276	temperature is formed near the sea surface, a TBL with the significant extension of strong
277	temperature gradient also appears in the upper ocean. Interestingly, even the thin and high-
278	temperature (~ SST_{max}) layer within the upper 5 m disappears before 6 pm on Dec 3 rd (Fig. 7),
279	the extension of temperature gradient in the DWL is not shoaled to < 5 m until the midnight of
280	Dec 4 th . The deepening of DWL before the nighttime convective mixing may be induced by the
281	shear at the base of DWL (Matthews et al. 2014; Hughes et al. 2020). Because the density
282	stratification is affected by the temperature structure, the extension of temperature gradient
283	driven by the vertical shear (Hughes et al. 2020) may nonlinearly affect the turbulent diffusivity
284	above the MLD, and thereby the cooling of SST from the daytime peak.
285	
286	5.2 Stratification above the MLD
207	The MID is estimated by using the DWI as the reference depth to evoid the temperature

287 The MLD is estimated by using the DWL as the reference depth to avoid the temperature
288 gradient in the DWL (section 3). Because of the large N² below the MLD, the estimated MLD

captures the top of the seasonal thermocline reliably, ~ 50 m before Dec 2^{nd} . Strong nighttime convective mixing occurs above the MLD (N² < 0 shaded by the white color in Fig. 5), mainly driven by latent heat flux and longwave cooling. During the strong DV SST event (section 4), the MLD at the floats except 9205 is shoaled by 20 m, consistent with the change of SST_{fnd} from 26.6 to > 27.7 °C. At float 9205, the MLD and SST_{fnd} are nearly constant. Because the salinity at float 9205 is higher than that at the other floats, we suspect that different vertical structure of salinity between the floats may be associated with the variation of MLD shoaling.

Because the trajectories of the floats are slightly different, the measured vertical structure of the salinity between the floats is not the same during the strong diurnal SST warming. On Dec 3rd, float 9209 measure a fresh-water layer with salinity ~34.4 psu near the sea surface. It results in a SBL at 30-m depth during the diurnal warming. Because the simulated magnitude of DV SST at float 9205 is still similar with the observation, the presence of SBL in the subsurface layer may not affect the DV SST significantly.

Except at float 9205, the average of N^2 between 20 and 40-m depth (part of the surface 302 mixed layer on Dec 1st) increases from 5.0×10^{-5} to 1.0×10^{-4} s⁻² from Dec 2nd to 4th. It shoals 303 the MLD by > 20 m. The restratification rate $\partial N^2/\partial t$ is ~3.5 × 10⁻¹⁰ s⁻³, much faster than that 304 reported by Brainerd and Gregg (1993a) during the daytime ($< 4 \text{ m s}^{-1}$ and peak insolation ~ 700 305 W m⁻²), ~ 1.6×10^{-10} s⁻³. The upper ocean becomes stably stratified in a few days. More 306 importantly, though the wind speed increases to > 6 m s⁻¹ after Dec 5th, the MLD at floats 9207 307 and 9209 is still about 30 m, shallower than that at 50-m depth before Dec 2^{nd} (Fig. 3). The 308 decrease of MLD agrees with the increase of SST_{fnd} from 26.5 °C to > 27 °C before and after the 309 310 strong DV SST event.

311	Clearly, the increase of SST_{fnd} is inversely proportional to the shoaling of MLD,
312	consistent with the model results reported by Bernie et al. (2005). The SST_{fnd} can be higher if the
313	same amount of heat content is accumulated in a shallower MLD. The shoaling of MLD also
314	coincidentally occurs after the strong DV SST > 2 °C except float 9205. Because the extension
315	of temperature gradient in the DWL can inhibit the vertical mixing efficiently, we speculate that
316	the thick DWL as a TBL may reduce the nighttime convective mixing for eroding the
317	stratification above the MLD. It prolongs the period for the penetrative solar radiation to
318	restratify the RL and shoal the MLD in the following day.
319	Despite of it, there may be some other factors for causing the restratification in the RL.
320	The role of penetrative solar radiation is studied in a one-dimensional model. The change of heat
321	absorption at different layers does not affect the density stratification below 30-m depth
322	significantly (section 6.2). The effect of horizontal advection is also studied. According to the
323	satellite measurements, the temperature advection at the sea surface may be insignificant near the
324	floats (supporting information C). However, without sufficient float measurements as direct
325	evidences, it is hard to quantify the temperature advection driven by the warm patch at 115 $^{\circ}\text{E}$
326	and 15.3 °S (Fig. 2), and thereby the restratification of RL (Brainerd and Gregg 1993b). That is,
327	the cause of the rapid restratification of RL is still in doubt in this study. Understanding the
328	mechanism for changing the MLD is crucial for predicting the SST_{fnd} variations in the future.
329	

330 6. Simulations of SST Variations using the KPP mixing scheme

The strong DV SST > 2 $^{\circ}$ C is observed at all ALAMO floats, associated with the 331 332 extension of temperature gradient in the DWL to the deeper layers (section 5). Can a numerical 333 model simulate the near surface temperature stratification we observed during these strong DV 334 SST events accurately? More importantly, which factors may be crucial for simulating the upper 335 ocean stratification during the diurnal warming? Compared with the multi-layer models such as 336 PWP3D (Price et al. 1986), the K-profile parameterization (KPP) can better simulate the DV 337 SST (Kawai and Wada 2007), and has been used in several ocean models (e.g., Shinoda and 338 Hendon 1998; Bernie et al. 2005). We will use the KPP in a one-dimensional Regional Oceanic 339 Modeling System (ROMS; Shchepetkin and McWilliams 2005) to simulate the evolution of 340 upper ocean stratification at the ALAMO array. Details of model settings and parameters in KPP 341 are described in appendix B.

342

343 6.1 Simulated SST and density stratification N^2

We compare the model, with a fine vertical resolution near the ocean surface (section 344 6.2), with the observations (Fig. 8a-d). On the first day of model simulations (Dec 1st), the KPP 345 simulated the SST reliably, including the DV SST of 0.5 °C. The simulated N^2 near the sea 346 surface is similar to that observed, as found in Bernie et al. (2005). After Dec 1st, the model still 347 predicts the SST at float 9205 well, including the SST_{max} of 29 °C on Dec 3rd. At float 9209, the 348 349 simulated SST agrees with the observed SST well until the midnight of Dec 4th, i.e., before the 350 shoaling of MLD from 40 to 20-m depth. The model results of SST at floats 9207 and 9010 differ from the observations significantly since Dec 2nd, consistent with the timing of rapid 351 352 restratification in the RL.

The simulated temperature and salinity are used to compute the N^2 for discussing the evolution of upper ocean stratified layers (Fig. 8). For the floats 9205 and 9209, which have the similar MLD with the observations before Dec 4th, the simulated magnitude of DV SST agrees 356 with the observed DV SST. The occurrence of strong DV SST mainly results from the air-sea heat fluxes in the one-dimensional process. Though the simulated SST_{max} at floats 9207 and 357 9210 can still be > 28 °C on Dec 3rd, different N² above the MLD thereby SST_{fnd} results in the 358 359 discrepancy of SST between the model and observations. Even the KPP mixing scheme 360 simulates the DV SST magnitude and a highly stratified layer near the sea surface reliably, the 361 failure on predicting the stratification above MLD may affect the simulated SST_{fnd} thereby the 362 SST_{max}.

363 On the other hand, compared with the observed DWL, the thickness of simulated DWL is all less than 10 m, thinner than the observations after Dec 1^{st} (Fig. 8e-1). The simulated N² in the 364 DWL (> 1.0×10^{-3} s⁻²) is two times larger than the float measurements from 12 to 4 pm between 365 Dec 2nd and 4th. Most heat with high temperature gradient is accumulated near the sea surface in 366 367 the model, unlike the observed temperature gradient extending to the deeper layer. In other 368 words, though the magnitude of DV SST is similar, the structure of DWL between the model 369 results and observations can still differ significantly.

370

371

6.2 Effect of penetrative solar radiation on the RL's restratification

372 Considering the importance of penetrative solar radiation for inducing diurnal warming 373 and restratifying the RL (Brainerd and Gregg 1993b), different coefficients based on five water 374 types are used in the parameterization of penetrative solar radiation (Paulson and Simpson 1978; 375 appendix B) during the model simulations (Fig. 9). The model results at float 9209 will be 376 discussed, because the SST difference at float 9209 between the model results and observations is not significant until the warming of SST_{fnd} on Dec 4th. Before Dec 3rd, the DWL in the upper 5 377 378 m can be simulated by all model runs with different water types. The simulated magnitude of DV

SST is similar with the observation. On Dec 3^{rd} , the simulation using the water type I forms a 379 shallower DWL with higher N^2 than the other water types. Because only the water type I has the 380 significant DV SST > 2 °C, the value of N^2 in the DWL may be the most dominant factor for 381 382 simulating the magnitude of DV SST. Interestingly, the model results using the water type II or III have faster SST_{fnd} warming 383 from Dec 2^{nd} to 4^{th} (~ 0.7 °C) than those using other water types. Though their simulated N² in 384 the DWL is smaller than that in water type I, the N^2 in the DWL is not completely destratified by 385 the nighttime convective mixing since the midnight of Dec 3^{rd} . The remaining N² in the water 386 387 types II and III implies that the simulated nighttime convective mixing may entrain less cold 388 water from the seasonal thermocline to the ocean surface than that in the water type I. That is, the 389 extension of temperature gradient to the deeper layer may more efficiently inhibit the nighttime

390 convective mixing in the model simulations. The evolution of upper ocean stratification,

including both DWL and MLD, is important to the forecast on the SST variations.

392

393 6.3 Effects of vertical resolution in the upper ocean

Several previous studies discuss the importance for using the vertical resolution $\Delta z \le 1$ m in the simulation of DV SST (e.g., Bernie et al. 2005; Hughes et al. 2020). It may be sometimes impractical to use the vertical resolution of 1 m in the entire ocean model for a climate forecast. The effect of vertical resolution in different depth ranges of the upper ocean is thus studied by using fine and coarse grids in the SST simulations, respectively, as detailed in Fig. 10a. The difference of simulated SST_{max} (Fig. 10) between the fine and coarse grids is

400 negligible during the weak DV SST (e.g., Dec 1^{st} and 2^{nd}), but significant during the strong DV

401 SST (e.g., Dec 3rd). High vertical resolution in the upper 20 m may directly affect the

accumulated heat near the ocean surface for the SST warming, by simulating the detailed 402 structure of N^2 in the DWL, especially during the strong DV SST. The simulated SST_{fnd} in the 403 coarse grids is slightly warmer than that in the fine grids on Dec 4th, after the strong diurnal 404 warming on Dec 3^{rd} . The Δz from 20 and 60-m depth may affect the simulated nighttime 405 406 convective mixing, and thereby the SST cooling from the daytime peak to the nighttime 407 minimum. Therefore, the Δz in the upper 20 m and from 20 and 60-m depth has a different impact on the SST variation. The Δz in the upper 20 m may affect the simulated N² in the DWL 408 and SST_{max} in the afternoon. The Δz from 20 and 60-m depth may affect the simulations of the 409 410 nighttime convective mixing and SST_{fnd}.

411

412 6.4 Parameters in the KPP mixing scheme

Because the KPP run with a high vertical resolution near the ocean surface fails to predict the structure of DWL and rapid SST_{fnd} warming at three ALAMO floats (section 6.1), the values of the mixing parameters in the K_p parameterizations are explored to seek improvements on the simulations of SST_{fnd} at float 9209 (Fig. 11). We will discuss the parameters Ri_c and Ri₀ (appendix B), which directly affect the vertical diffusivity K_p within and below the OBL, respectively.

419 Compared to the simulation using the default setting of mixing parameters ($Ri_c = 0.3$ and 420 $Ri_0 = 0.7$), decreasing the Ri_c from 0.3 to 0.1 (i.e., assuming the OBL is thinner) has negligible 421 effects to the simulated SST. Changing the thickness of the OBL may not affect the simulated 422 SST_{max} significantly, presumably due to the similar K_{ρ} below the OBL. On the other hand, 423 decreasing Ri_0 from 0.7 to 0.3 (i.e., more difficult for inducing shear instability) increases the 424 SST_{max} significantly, but has negligible effect on the SST_{fnd} . Inhibiting the vertical mixing by 425 restricting the depth range of K_{ρ} may affect the prediction of SST_{max}. If the Ri₀ alternatively

426 increases from 0.7 to 1 (i.e., larger depth ranges of K_p to transport heat at the base of OBL), the

427 KPP still fails to simulate the rapid SST_{fnd} warming. That is, the Ri₀ which affects the turbulent

428 diffusivity below the DWL is the most important parameter for the simulated DV SST magnitude 429 in the KPP.

430

445

431 7. Summary and Conclusion

432 Six ALAMO floats with high vertical resolution ≤ 1 m in the upper 50 m are deployed off the northwest Australia on Nov 22nd 2018. Four floats measure strong DV SST of up to 2 °C in 433 the beginning of December, under low wind speed (~ 2 m s^{-1}) and sunny conditions. A rapid 434 SST_{fnd} warming is observed at three floats (9207, 9209 and 9210), rising from 26.4 to more than 435 27.6 °C between Dec 2nd and 4th. The increase of SST_{fnd} at float 9205 is ~ 0.5 °C, lower than that 436 437 at the other floats. The warming rate of SST_{fnd} varies significantly, even though the distance between floats is less than 50 km. Because of the rapid SST_{fnd} warming, the latent plus sensible 438 heat flux at floats 9207 and 9210 is ~8.4 W m⁻² higher than that at float 9205 from Dec 2^{nd} to 4^{th} . 439 440 To emphasize the presence of a strong temperature gradient above the surface mixed 441 layer during the diurnal warming, a diurnal warm layer depth (DWL) is defined here by finding the mean temperature gradient $\partial T/\partial z$ from z_0 to (z_0+5) -m depth > 0.02 °C m⁻¹. In other word, a 442 443 thick DWL will extend its strong temperature gradient to the deeper layer. Under the strong DV SST ~ 2 °C, the averaged N² within the thick DWL (~ 20 m) is more than 1.0×10^{-4} s⁻². Below 444 the DWL, N² in the surface mixed layer increases from 5.0×10^{-5} to 1.0×10^{-4} from Dec 2nd to

4th. This restratification rate is faster than that reported by previous studies (e.g., Brainerd and 446

Gregg 1993a). This fast restratification below the DWL may prolong the period of high SST inthe tropical warm pool.

449 The KPP mixing scheme in a 1-D ROMS model is used to simulate the SST and upper 450 ocean stratified at the float positions. The simulated SST agrees well with the observed SST at 451 float 9205, including the strong DV SST > 2 $^{\circ}$ C. However, the model fails to simulate the rapid 452 SST_{fnd} warming at the other three floats, presumably due to the other factors for simulating the 453 restratification above the MLD. Factors impacting the simulation of SST variations in the KPP 454 are discussed. High vertical resolution in the upper 20 m of < 1 m is required for reliably 455 simulating the magnitude of DV SST. Decreasing the mixing parameter Ri₀ for inhibiting the 456 turbulent diffusivity in the diurnal thermocline can directly increase the peak SST in the model 457 simulations. Changing the mixing parameter Ri_c has negligible effects to the SST simulations. 458 In summary, a stable sunny atmosphere with low wind speed is favorable for the 459 formation of a thick DWL during the suppressed phase of the MJOs. The extension of 460 temperature gradient in the DWL is studied using the high vertical resolution measurements in 461 the first time, and can be more than 20 m during the strong DV SST > 2 °C. The shoaling of 462 MLD can also increase the SST_{fnd}. Though the KPP mixing scheme can simulate the DV SST 463 magnitude reliably by using high vertical resolution near the sea surface, it fails to predict the 464 increase SST_{fnd}, mainly due to different upper ocean density structure between the model and 465 observations. Questions still remain regarding to the factors for shoaling the observed MLD in a 466 short time period. Future field measurements on turbulent diffusivity within the DWL in TWPs 467 are important for improving the ocean mixing approaches in the global coupled models for the 468 MJO forecast.

469

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482 Appendix A. Buoy and Satellite Measurements

483	The Bailong buoy system from the First Institution of Oceanography (Cole et al. 2011)
484	includes the atmospheric measurements near the ocean surface (Fig. 12) and subsurface ocean
485	measurements in the upper 500 m (Feng et al. 2020). The atmosphere data is sampled in every 10
486	minutes. During the strong DV SST event from Dec 2 nd to 4 th 2018, the peak insolation is more
487	than 1000 W m ⁻² , and the wind speed at 4-m height above the sea surface is about 2 m s ⁻¹ . Strong
488	diurnal variation of air temperature is found, ~ 1 °C on Dec 2^{nd} . The relative humidity (RH) is
489	about 60%, and then increases until the onset of the MJO. The temporal variation of downward
490	longwave radiation is small, from 400 to 450 W m ⁻² .
491	The infrared sensor mounted on the Japanese geostationary Himawari-8 satellite
492	measures the skin SST in four spectral bands (8.59, 10.40, 11.24 and 12.38 μ m) for every 10
493	minutes with the horizontal resolution < 2 km (Kramar et al., 2016). The product of Himawari-8
494	SST reprocessed by Dr. Christopher Griffin
495	(http://opendap.bom.gov.au:8080/thredds/catalog/abom_imos_ghrsst_archive-
496	1/v02.0fv03test/Continental/L3C-01hour/ABOM-L3C_GHRSST-SSTskin-AxIH08/2018/) in
497	Bureau of Meteorology Australia as an hourly dataset is available online. We also use the near-
498	real time sea surface height anomalies data processed by Integrated Marine Observing System
499	(IMOS) data portal (Baird and Ridgway 2012; Deng et al. 2010), to understand the distribution
500	of eddies around floats. The geostrophic current is computed, mostly northward and $< 0.2 \text{ m s}^{-1}$
501	along the trajectories of floats.
502	

503 Appendix B. KPP Mixing Scheme in the 1-D ROMS Model Simulation

504 **B.1. ROMS model description**

505 The one-dimensional ROMS model with the KPP mixing is used to simulate the SST 506 warming at each ALAMO floats. The atmosphere measurements taken by the FIO buoys, 507 including insolation, downward longwave radiation, air temperature, air pressure, atmosphere 508 wind and relative humidity, are used as the forcing at the floats, by assuming the spatial variation 509 of atmosphere condition negligible within the distance of 80 km. The temporal resolution is 10 min. The profiles at each float from 9 to 11 pm on Nov 30th are averaged as the initial conditions. 510 511 We use the parameterization of penetrative solar radiation Q proposed by Paulson and 512 Simpson (1977) for simulating the change of upper ocean stratification, which can be expressed 513 as $\begin{pmatrix} & & & \\ & & & & \\ & & & & & \end{pmatrix}$

514
$$Q = Q_0 \left(r \exp\left(\frac{z}{\mu_1}\right) + (1 - r) \exp\left(\frac{z}{\mu_2}\right) \right)$$
(S2)

where Q_0 is the insolation, r, μ_1 and μ_2 are the coefficients based on the data of five different water types in Jerlov (1976) (**Error! Reference source not found.**;

<u>https://www.myroms.org/wiki/Jwtype</u>). According to the description in the ROMS model,
water type I is used for open Pacific Ocean; water type IA is used for open Indian Ocean; water
type IB is used for open Atlantic Ocean; water type II is used for Azores; water type III is used
for North Sea. Most model results are simulated using the coefficients of water type I, except
those in section 6.2.

B.2. KPP mixing scheme

The K-profile parameterization (KPP. Large et al. 1994) used in many ocean models, such as HYbrid Coordinate Ocean Model (HYCOM; Chassignet et al. 2007), computes the turbulent diffusivity by assuming a shape function in the OBL. It differs to the turbulence kinetic energy (TKE) closure scheme (e.g., Mellor and Yamada 1982), which uses a prognostic TKE energy equation and length scale of mixing. In the KPP mixing scheme, after prescribing the surface forcing, the depth h of ocean boundary layer (OBL) will be first determined by using a critical bulk Richardson number Ri_c (default = 0.3). The diffusivity in the OBL, K_{Ric} , is computed based on the surface flux and h, assuming a nondimensional vertical shape function. For the diffusivity below the OBL, K_{Ri0}, the shear instability mixing occurs only when the local gradient Richardson number Ri is smaller than a critical gradient Richardson number Ri_0 (default = 0.7). The total diffusivity K_{ρ} is constituted by the K_{Ric} , K_{Ri0} and background diffusivity $K_{\rho 0}$ (default = 1.0×10^{-6}).

542 **References**

- 543 Baird, M. E., and K. R., Ridgway, 2012. The southward transport of sub-mesoscale lenses of
- 544 Bass Strait Water in the centre of anti-cyclonic mesoscale eddies, Geophysical Research
- 545 Letter, 39, L02603, doi:10.1029/2011GL050643.
- 546 Balaguru, K., G. R., Foltz, L. R., Leung, E., D'Asaro, K. A., Emanuel, H., Liu, and S. E.,
- 547 Zedler, 2015. Dynamic Potential Intensity: An improved representation of the ocean's
- 548 impact on tropical cyclones, Geophysical Research Letter, 42, 6739–6746,
- 549 doi:10.1002/2015GL064822.
- 550 Beggs, H., L., Majewski, G., Kruger, R., Verein, P., Oke, P., Sakov, X., Huang, L., Garde and C.
- 551 Tingwell, 2013. Report to GHRSST14 from Australia Bluelink and IMOS.
- 552 <u>http://imos.org.au/fileadmin/user_upload/shared/SRS/Beggs_Australian_Report_GHR</u>
 553 SST14_10Dec2013.pdf
- Bellenger, H. and J. Duvel, 2009. An Analysis of Tropical Ocean Diurnal Warm Layers. J.
- 555 Climate. 22, 3629–3646. https://doi.org/10.1175/2008JCLI2598.1
- 556 Bernie, D. J., S. J. Woolnough, J. M. Slingo, and E. Guilyardi, 2005. Modeling Diurnal and
- 557 Intraseasonal Variability of the Ocean Mixed Layer, Journal of Climate, 18, 1190-1202.
- 558 Bernie, D. J., E. Guilyardi, G. Madec, J. M. Slingo, S. J. Woolnough and J. Cole, 2008. Impact
- of resolving the diurnal cycle in an ocean–atmosphere GCM. Part 2: A diurnally coupled
- 560 CGCM. Climate Dynamics, 31, 909-925, doi: 10.1007/s00382-008-0429-z
- 561 Brainerd, K. E. and M. C. Gregg, 1993a. Diurnal restratification and turbulence in the oceanic
- surface mixed layer: 1. Observations. Journal of Geophysical Research: Oceans, 98, 22645-
- 563 22656, doi:10.1029/93JC02297.

- Brainerd, K. E., and M. C. Gregg, 1993b. Diurnal restratification and turbulence in the oceanic
 surface mixed layer: 2. Modeling, Journal of Geophysical Research, 98, 22657–22664,
 doi:10.1029/93JC02298.
- 567 Brainerd, K. E., and M. C. Gregg, 1995. Surface mixed and mixing layer depths, Deep Sea
- 568 Research Part I: Oceanographic Research Papers, 42, 1521-1543.
- 569 Caldwell, D. R., R-C. Lien, J. N. Moum, and M. C. Gregg, 1997. Turbulence decay and
- restratification in the equatorial ocean surface layer following nighttime convection. Journal
 of Physical Oceanography, 27, 1120–1132.
- 572 Chassignet, E. P., H. E., Hurlburt, O. M., Smedstad, G. R., Halliwell, P. J., Hogan, A. J.,
- 573 Wallcraft, R., Baraille and R., Bleck, 2007. The HYCOM (HYbrid Coordinate Ocean
- 574 Model) data assimilative system. Journal of Marine Systems 65, pp. 60-83. doi:

575 https://doi.org/10.1016/j.jmarsys.2005.09.016

- 576 Chi, N.-H., R.-C. Lien, E. A. D'Asaro and B. B. Ma, 2014. The surface mixed layer heat budget
- 577 from mooring observations in the central Indian Ocean during Madden–Julian Oscillation
- 578 events. Journal of Geophysical Research: Oceans. 119, 4638–4652,
- 579 doi:10.1002/2014JC010192.
- 580 Cole, R., J. Kinder, Chun Lin Ning, W. Yu and Yang Chao, ""Bai-Long": A TAO-hybrid on
- 581 RAMA," OCEANS'11 MTS/IEEE KONA, Waikoloa, HI, 2011, pp. 1-10. doi:
- 582 10.23919/OCEANS.2011.6106952
- 583 Cronin, M. F., and McPhaden, M. J., 1998. Upper ocean salinity balance in the western
- equatorial Pacific, Journal of Geophysical Research, 103, 27567–27587,
- 585 doi:10.1029/98JC02605.

586	de Boyer Montégut, C., G., Madec, A. S., Fischer, A., Lazar, and D. Iudicone, 2004. Mixed layer
587	depth over the global ocean: An examination of profile data and a profile-based
588	climatology, Journal of Geophysical Research, 109, C12003, doi:10.1029/2004JC002378.
589	DeMott, C. A., N. P. Klingaman, and S. J. Woolnough, 2015. Atmosphere-ocean coupled
590	processes in the Madden-Julian oscillation, Reviews of Geophysics, 53, 1099-1154.
591	Deng, X, D. A. Griffin, K. Ridgway, J.A. Church, W.E. Featherstone, N. White and M. Cahill,
592	2010. Satellite altimetry for geodetic, oceanographic and climate studies in the Australian
593	region, in: Vignudelli S., A. Kostianoy, and P. Cipollini and J. Benveniste (eds.), Coastal
594	Altimetry, Springer-Verlag, Berlin. ISBN: 978-3-642-12795-3. e-ISBN: 978-3-642-12796-
595	0. doi: 10.1007/978-3-642-12796-0_18
596	Drushka, K., S. T. Gille and J. Sprintall, 2014. The diurnal salinity cycle in the tropics. Journal of
597	Geophysical Research: Oceans, 119, 5874-5890, doi: 10.1002/2014jc009924
598	Drushka, K., W., Asher, A., Jessup, E., Thompson, S., Iyer, D., Clark, 2019. Capturing Fresh
599	Layers with the Surface Salinity Profiler. Oceanography. 32. 76-85.
600	10.5670/oceanog.2019.215.
601	Fairall, C. W., E. F., Bradley, D. P., Rogers, J. B., Edson, and G. S., Young, 1996a. Bulk
602	parameterization of air-sea fluxes for Tropical Ocean-Global Atmosphere Coupled-Ocean
603	Atmosphere Response Experiment, Journal of Geophysical Research, 101, 3747-3764,

- 604 doi:10.1029/95JC03205.
- 605 Fairall, C. W., E. F., Bradley, J. S., Godfrey, G. A., Wick, J. B., Edson, and G. S. Young, 1996b.
- 606 Cool-skin and warm-layer effects on sea surface temperature, J. Geophys. Res., 101, 1295–
- 607 1308, doi:10.1029/95JC03190.

608	Fairall, C.W., E. F. Bradley, J. E. Hare, A. A. Grachev, and J. B. Edson, 2003. Bulk
609	Parameterization of Air-Sea Fluxes: Updates and Verification for the COARE Algorithm. J.
610	Climate, 16, pp 571-591.
611	Feng, M., Y., Duan, S., Wijffels, JY., Hsu, C., Li, H., Wang, Y., Yang, H., Shen, J., Liu, C.,

- 612 Ning, and W., Yu, 2020. Tracking air-sea exchange and upper ocean variability in the
- 613 Southeast Indian Ocean during the onset of the 2018-19 Australian summer monsoon.
- 614 Bulletin of the American Meteorological Society. https://doi.org/10.1175/BAMS-D-19-
- 615 <u>0278.1</u>.
- 616 Gargett, A. E., and G. Holloway, 1984. Dissipation and diffusion by internal wave breaking,
- 517 Journal of Marine Research, 42, 15-27
- Guilyardi, E., 2006. El niño-mean state-seasonal cycle interactions in a multi-model ensemble.
 Climate Dynamics. 26. 329–348
- 620 Hsu, J.-Y., H., Hendon, M., Feng, and X., Zhou, 2019. Magnitude and phase of diurnal SST
- 621 variations in the ACCESS-S1 model during the suppressed phase of the MJOs. Journal of
- 622 Geophysical Research: Oceans, 124, 9553–9571. https://doi.org/10.1029/2019JC015458
- 623 Hughes, K. G., J. N. Moum and E. L. Shroyer, 2020. Evolution of the Velocity Structure in the
- 624Diurnal Warm Layer. Journal of Physical Oceanography, 50, 615-631, doi: 10.1175/jpo-d-
- 625 19-0207.1
- 626 Jerlov, N. G., 1976. Marine Optics, 14, Elsevier Oceanography Series
- 627 Karagali, I., and Høyer, J. L., 2014. Characterisation and quantification of regional diurnal SST
- 628 cycles from SEVIRI. Ocean Science, 10, 745–758. <u>http://dx.doi.org/10.5194/os-10-745-</u>
- 629 <u>2014</u>.

- 630 Kawai, Y. and Wada, A., 2007. Diurnal sea surface temperature variation and its impact on the
- 631 atmosphere and ocean: A review. Journal of Oceanography, 63, 721-744, doi:
- 632 10.1007/s10872-007-0063-0
- 633 Kramar, M., A. Ignatov, B. Petrenko, Y. Kihai, P. Dash, 2016. Near real time SST retrievals
- from Himawari-8 at NOAA using ACSPO system. Proc. SPIE 9827, Ocean Sensing andMonitoring.
- 636 Kudryavtsev, V.N. and A.V. Soloviev, 1990. Slippery Near-Surface Layer of the Ocean Arising
- 637 Due to Daytime Solar Heating. Journal of Physical Oceanography, 20, 617–628,
- 638 https://doi.org/10.1175/1520-0485(1990)020<0617:SNSLOT>2.0.CO;2
- Kunze, E., 2003. A review of oceanic salt-fingering theory. Progress in Oceanography, 56, 399417. doi: https://doi.org/10.1016/S0079-6611(03)00027-2
- 641 Large, W. G., J. C. McWilliams, and S. C. Doney, 1994. Oceanic vertical mixing: A review and
- a model with a nonlocal boundary layer parameterization, Reviews of Geophysics, 32, 363-403.
- Lee, C., K.-I., Chang, J. H., Lee, and K. J., Richards, 2014. Vertical mixing due to double
- diffusion in the tropical western Pacific, Geophysical Research Letter, 41, 7964–7970,
- 646 doi:10.1002/2014GL061698.
- Lukas, R., and E. Lindstrom, 1991. The mixed layer of the western equatorial Pacific Ocean, J.
- 648 Geophys. Res., 96(S01), 3343–3357, doi:10.1029/90JC01951.
- 649 Maloney, E. D., 2009. The Moist Static Energy Budget of a Composite Tropical Intraseasonal
- 650 Oscillation in a Climate Model, Journal of Climate, 22, 711-729.

- 651 Matthews, A.J., D.B. Baranowski, K.J. Heywood, P.J. Flatau, and S. Schmidtko, 2014. The
- 652 Surface Diurnal Warm Layer in the Indian Ocean during CINDY/DYNAMO. Journal of

653 Climate, 27, 9101–9122, <u>https://doi.org/10.1175/JCLI-D-14-00222.1</u>

- McPhaden, M. J. and G. R. Foltz, 2013. Intraseasonal variations in the surface layer heat balance
- of the central equatorial Indian Ocean: The importance of zonal advection and vertical
- 656 mixing. Geophysical Research Letters. 40, 2737–2741, doi:10.1002/grl.50536.
- Mellor, G. L., and T. Yamada, 1982. Development of a turbulence closure model for geophysical
 fluid problems, Rev. Geophys., 20, 851–875, doi:10.1029/RG020i004p00851.
- Moulin, A. J., J. N. Moum and E. L. Shroyer, 2018. Evolution of Turbulence in the Diurnal
- 660 Warm Layer. Journal of Physical Oceanography, 48, 383-396, doi: 10.1175/jpo-d-17-0170.1
- Moum, J. N., S. P. de Szoeke, W. D. Smyth, J. B. Edson, H. L. DeWitt, A. J. Moulin, E. J.
- Thompson, C.J. Zappa, S. A. Rutledge, R. H. Johnson, and C. W. Fairall, 2014. Air–Sea
- 663 Interactions from Westerly Wind Bursts During the November 2011 MJO in the Indian
- 664 Ocean. Bull. Amer. Meteor. Soc., 95, 1185–1199, https://doi.org/10.1175/BAMS-D-12-
- 665 00225.1
- Paulson, C. A., and J. J. Simpson, 1977. Irradiance measurements in the upper ocean, Journal of
 Physical Oceanography, 7, 952-956.
- Price, J. F., R. A. Weller and R. Pinkel, 1986. Diurnal cycling: Observations and models of the
- upper ocean response to diurnal heating, cooling, and wind mixing. Journal of Geophysical
- 670 Research: Oceans, 91, 8411-8427, doi: 10.1029/JC091iC07p08411
- 671 Ruppert, J. H., and R. H. Johnson, 2015. Diurnally modulated cumulus moistening in the
- 672 preonset stage of the Madden-Julian oscillation during DYNAMO, Journal of the
- 673 Atmospheric Sciences, 72, 1622–1647.

- 674 Seo, H., A. C. Subramanian, A. J. Miller, and N. R. Cavanaugh, 2014. Coupled Impacts of the
- 675 Diurnal Cycle of Sea Surface Temperature on the Madden–Julian Oscillation, Journal of
- 676 Climate, 27, 8422-8443, https://doi.org/10.1175/JCLI-D-14-00141.1
- 677 Shchepetkin, A. F., and J. C. McWilliams, 2005. The regional oceanic modeling system (ROMS)
- a split-explicit, free-surface, topography-following-coordinate oceanic model. Ocean
- 679 Modelling, 9, pp. 347-404.
- 680 Shinoda, T. and H. H. Hendon, 1998. Mixed Layer Modeling of Intraseasonal Variability in the
- Tropical Western Pacific and Indian Oceans. Journal of Climate, 11, 2668–2685,

682 https://doi.org/10.1175/1520-0442(1998)011<2668:MLMOIV>2.0.CO;2

- 683 Shinoda, T., T. G. Jensen, M. Flatau, and S. Chen, 2013. Surface Wind and Upper-Ocean
- 684 Variability Associated with the Madden–Julian Oscillation Simulated by the Coupled
- 685 Ocean–Atmosphere Mesoscale Prediction System (COAMPS), Monthly Weather Review,
- 686 141, 2290-2307.
- 687 Sobel, A., S. Wang, and D. Kim, 2014. Moist Static Energy Budget of the MJO during
- 688 DYNAMO, Journal of the Atmospheric Sciences, 71, 4276-4291.
- 689 Soloviev, A., and R., Lukas, 2006. The near-surface layer of the Ocean. Springer.
- 690 Sprintall, J., and M. Tomczak, 1992. Evidence of the barrier layer in the surface layer of the
- 691 tropics, J. Geophys. Res., 97, 7305–7316, doi:10.1029/92JC00407.
- Sprintall, J., and D. Roemmich, 1999. Characterizing the structure of the surface layer in the
 Pacific Ocean, J. Geophys. Res., 104, 23,297–23,311.
- 694 Suga, T., K. Motoki, Y. Aoki, and A. M. Macdonald, 2004. The North Pacific climatology of
- 695 winter mixed layer and mode waters, Journal of Physical Oceanography, 34, 3–22.

- 696 Sui, C., X. Li, K. Lau, and D. Adamec, 1997. Multiscale Air–Sea Interactions during TOGA
- 697 COARE. Monthly Weather Review, 125, 448–462, <u>https://doi.org/10.1175/1520-</u>
- 698 <u>0493(1997)125<0448:MASIDT>2.0.CO;2</u>
- 699 Sutherland, G., L. Marié, G. Reverdin, K.H. Christensen, G. Broström, and B. Ward, 2016.
- 700 Enhanced Turbulence Associated with the Diurnal Jet in the Ocean Surface Boundary
- 701Layer. Journal of Physical Oceanography, 46, 3051–3067, https://doi.org/10.1175/JPO-D-
- 702 15-0172.1
- 703 Vijith, V., P. N. Vinayachandran, B. G. M. Webber, A. J. Matthews, J. V. George, V. K.
- Kannaujia, A. A. Lotliker and P. Amol, 2020. Closing the sea surface mixed layer
- temperature budget from in situ observations alone: Operation Advection during BoBBLE.
- 706 Scientific Reports, 10, pp. 7062, doi: https://doi.org/10.1038/s41598-020-63320-0
- Wyrtki, K., 1964. The thermal structure of the eastern Pacific Ocean, Deutsche hydrographische
 Zeitschrift, 8, 6-84
- 709 Zhang, C., 2005. Madden-Julian Oscillation, Reviews of Geophysics, 43, RG2003,
- 710 doi:10.1029/2004RG000158.
- 711 Zhang, H., H. Beggs, L. Majewski, X. H. Wang, and A. Kiss, 2016. Investigating sea surface
- temperature diurnal variation over the Tropical Warm Pool using MTSAT-1R data, Remote
- 713 Sensing of Environment, 183, 1-12.
- 714 Zhang, C. and J. Ling, 2017. Barrier Effect of the Indo-Pacific Maritime Continent on the MJO:
- 715 Perspectives from Tracking MJO Precipitation. Journal of Climate, 30, 3439–3459,
- 716 https://doi.org/10.1175/JCLI-D-16-0614.1
- 717 Zhao, Z., 2018. The global mode-2 M2 internal tide. Journal of Geophysical Research: Oceans,
- 718 123, 7725–7746. https://doi.org/10.1029/2018JC014475

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Table

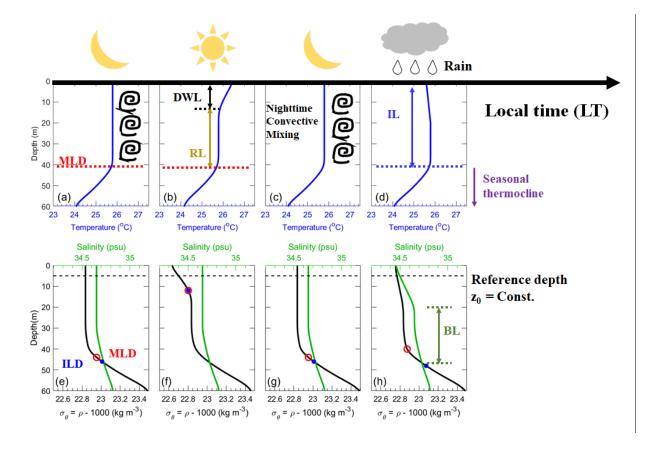
	Water Types				
	Ι	IA	IB	II	III
r	0.58	0.62	0.67	0.77	0.78
μ_1	0.35	0.60	1.00	1.50	1.40
μ ₂	23.00	20.00	17.00	14.00	7.90

Table. 1. The coefficients used in the parameterization of penetrative solar radiation (Eq. S2)

based on the data of five different water types reported by Jerlov (1976).

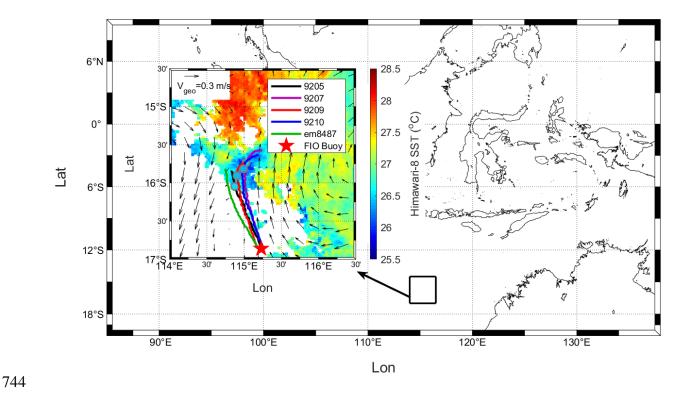
729 **Figure**





732 Fig. 1. Illustration on the surface mixed layer depth (MLD), diurnal warm layer (DWL), remnant 733 layer (RL), isothermal layer (IL) and barrier layer (BL), using assumed temperature (blue lines in 734 a-d) and salinity profiles (green lines in e-h). The RL exists between the DWL and MLD. 735 Examples of estimating MLD and ILD based on the criteria used in the previous studies are 736 presented in (e)-(h). The MLD (red circles) is estimated when the potential density difference $\rho(MLD) - \rho(z_0) > 0.1 \text{ kg m}^{-3}$ (Chi et al. 2014), assuming the constant reference depth z_0 at 5-m 737 738 depth (black dashed lines in e-h). The isothermal layer depth (ILD. Blue dots) is estimated in the 739 same way as the MLD but using temperature-dependent profiles of p, which assumes the salinity 740 as the mean in the upper 5-m ocean. The formation of BL (difference between MLD and ILD > 741 20 m) is due to the precipitation freshening the salinity in the upper 20 m, and is regarded as the

- salinity-driven barrier layer (SBL) in section 3. The temperature near the sea surface in (d) is
- 743 slightly cooled by rains (Druksha et al. 2019).



745 Fig. 2. Trajectories of four ALAMO floats (color lines) and one EM-APEX float (green line) in 746 the map of SST (color shading) measured by Himawari-8 satellite (appendix A) within the region of TWP (black box in the big map) around 12 am on Dec 1st 2018. The ALAMO floats and EM-747 748 APEX floats were deployed near the FIO buoy (red pentagram) at 115.3 °E and 16.8 °S in the northwest Australia on Nov 22nd, and drifting northwestward (color lines are trajectories from 749 Nov 22nd to Dec 6th 2018) due to the geostrophic current (black arrows in the small map, and 750 751 their magnitude can be referenced to that on the upper-left corner). The color dots connected to the ALAMO float trajectories are the float-measured SST around the midnight of Dec 1st 2018. 752 The geostrophic current V_{geo} is estimated using the IMOS sea surface height anomalies on Dec 753 754 1st 2018 (appendix A).

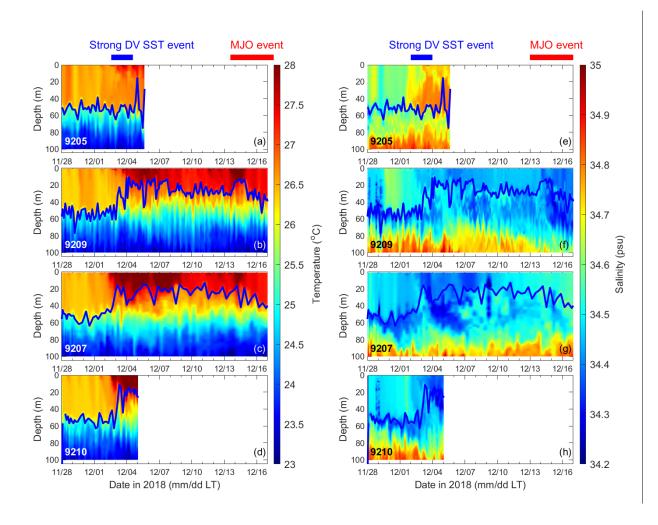


Fig. 3. Measurements of temperature (a-d) and salinity (e-h) in the upper 100 m taken by four

ALAMO floats (9205, 9209, 9207 and 9210). The period of strong DV SST and MJO events are

761 described in Feng et al. (2020). Blue lines are the estimated MLD (section 3).

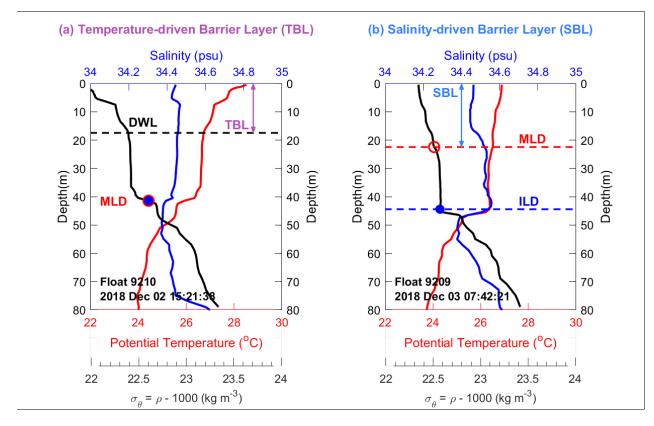


Fig. 4. Examples of temperature-driven barrier layer TBL (depth range indicated by the purple double arrow in a) and salinity-driven barrier layer SBL (depth range indicated by the light blue double arrow in b) based on the profiles of potential density anomaly σ_{θ} (= ρ -1000; black lines), potential temperature (blue lines) and salinity (red lines) at floats 9210 and 9209. The surface mixed layer depth MLD (red circles and red dashed line), isothermal layer depth ILD (blue dots and blue dashed line) and diurnal warm layer depth DWL (black dashed line) are estimated. The TBL is identified if DWL > 10 m. The SBL is identified if ILD - MLD > 10 m. More discussions on the criteria for identifying TBL and SBL can be found in supporting information A.

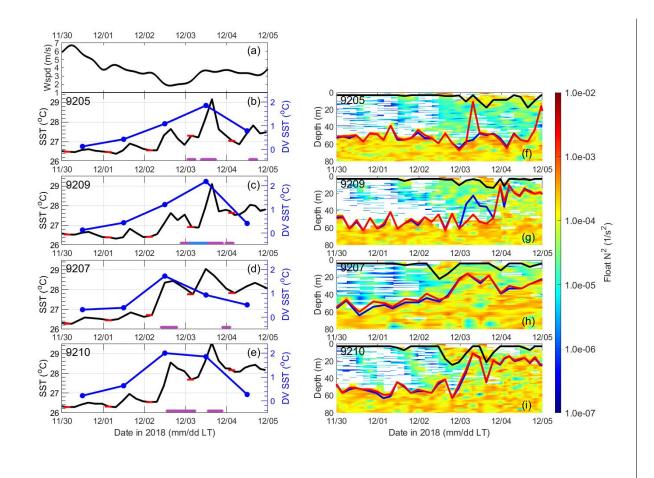




Fig. 5. Wind speed measurements on the buoy (a), and measurements of SST (black lines in b-e; referenced to the left axis), DV SST (blue dots connected with lines in b-e; referenced to the right axis) and N² (f-i; N² < 0 is shaded in white color) on four ALAMO floats (9205, 9207, 9209 and 9210). In (b)-(e), the red lines mark the foundation SST (SST_{fnd}; referenced to the left axis) in each day, and the purple and light blue bars in (b)-(e) mark the period where TBL and SBL exist, respectively. The TBL and SBL are identified using the estimates of DWL (black lines), MLD (blue lines) and ILD (red lines) in (f)-(i).

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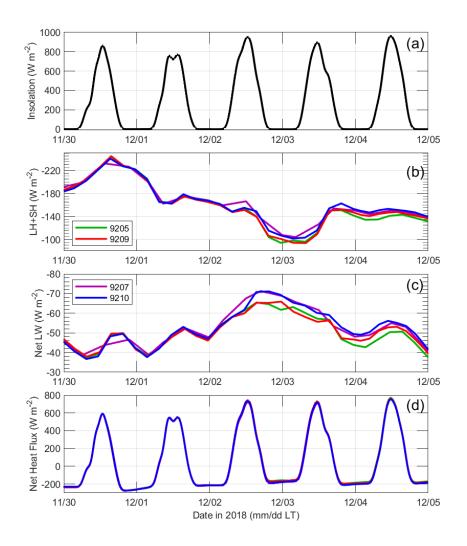
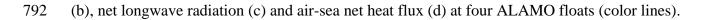


Fig. 6. Measurements of insolation on the buoy (a), and estimates of latent plus sensible heat flux



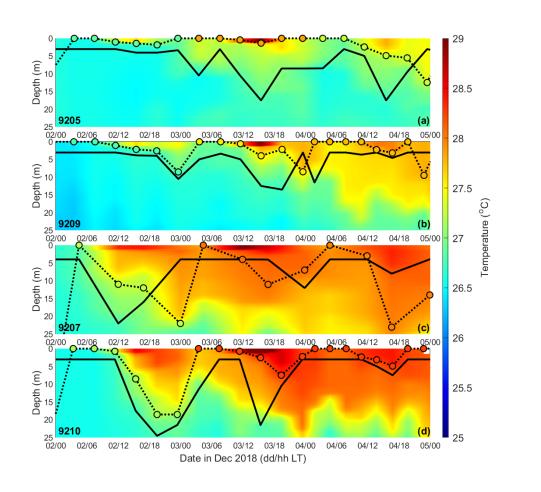


Fig. 7. Measurements of temperature (color shading) in the upper 25 m, estimated DWL (black
lines) and estimated DWL* (black dashed line connected with color dots) from Dec 2nd to 4th
2018 at four ALAMO floats (a-d). The color dots are the temperature T* of the isotherm, used for
finding the DWL*.

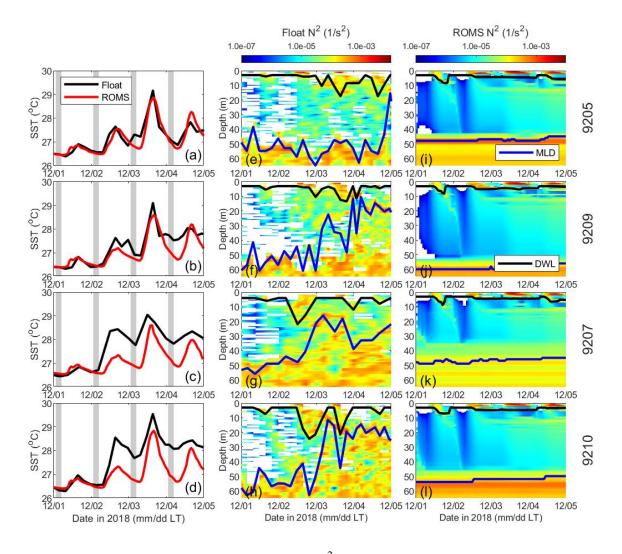


Fig. 8. Simulations of SST (red lines in a-d) and N^2 (color shading in i-l) using the KPP in the ROMS, compared with the observations of SST (black lines in a-d) and N^2 (color shading in e-h) at all ALAMO floats (each row). The period for computing the SST_{fnd} in each day is marked by the grey area in (a)-(d). The DWL (black lines) and MLD (blue lines) are estimated in (e)-(l), respectively.

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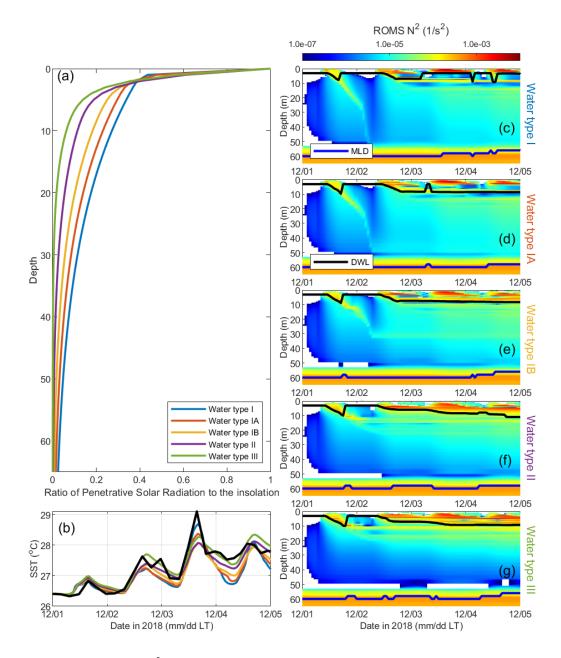
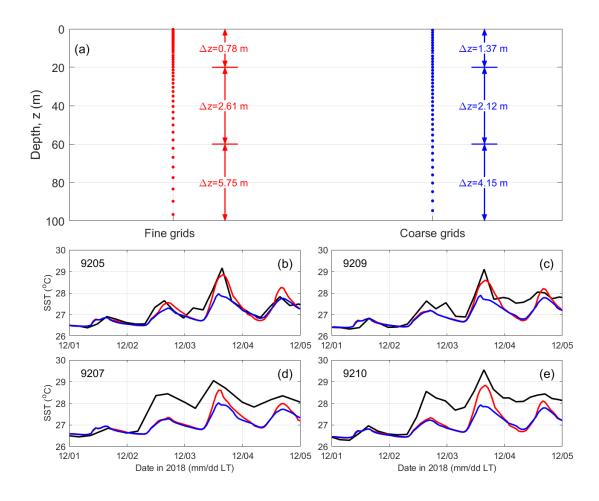


Fig. 9. Model results of N² (color shading in c-g) and SST (b) at float 9209 by varying the
penetrative solar radiation based on five water types (colored lines in a and b). The DWL (black
lines) and MLD (blue lines) are estimated in (c)-(g), respectively. The ratio of penetrative solar

- 814 radiation to the insolation at different depth during the model simulations is shown in (a). The
- 815 black line in (b) is the float-measured SST.



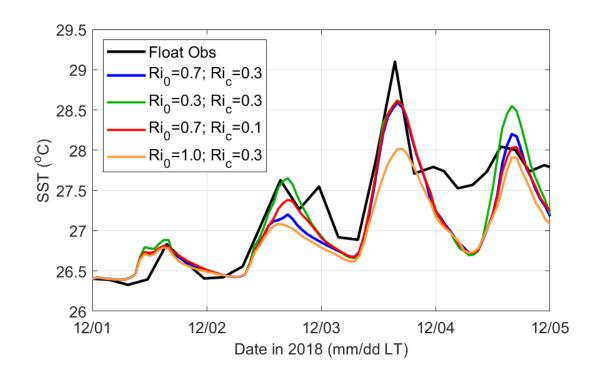
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817 Fig. 10. Model results of SST at four ALAMO floats (b-e) simulated using two types of vertical

grids (a) in the KPP: fine (red lines) and coarse grids (blue lines), with the comparison to the

819 float measurements (black lines), where Δz in (a) is the average of vertical resolution in different

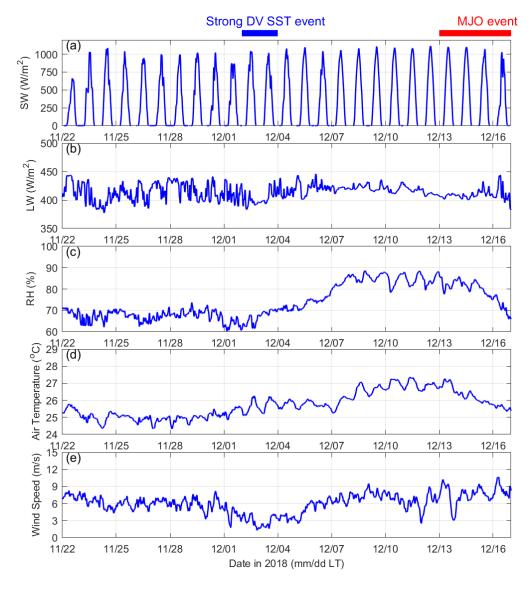
- 820 range of depth.
- 821
- 822
- 823



826 Fig. 11. Simulations of SST at float 9209 by different set of the mixing parameters Ri₀ and Ri_c

827 (blue: $Ri_0 = 0.7$ and $Ri_c = 0.3$; green: $Ri_0 = 0.3$ and $Ri_c = 0.3$; red: $Ri_0 = 0.7$ and $Ri_c = 0.1$; orange:

 $Ri_0 = 1.0$ and $Ri_c = 0.3$), with the comparison to the float observation (black).





831 Fig. 12. Measurements of downward shortwave radiation SW (a), downward longwave radiation

832 LW (b), relative humidity RH at 3-m height above the sea surface (c), air temperature at 3-m

height above the sea surface (d), and wind speed at 4-m height above the sea surface (e) on the

FIO buoy. The period of the strong DV SST and MJO events are described in Feng et al. (2020).

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Figure1.

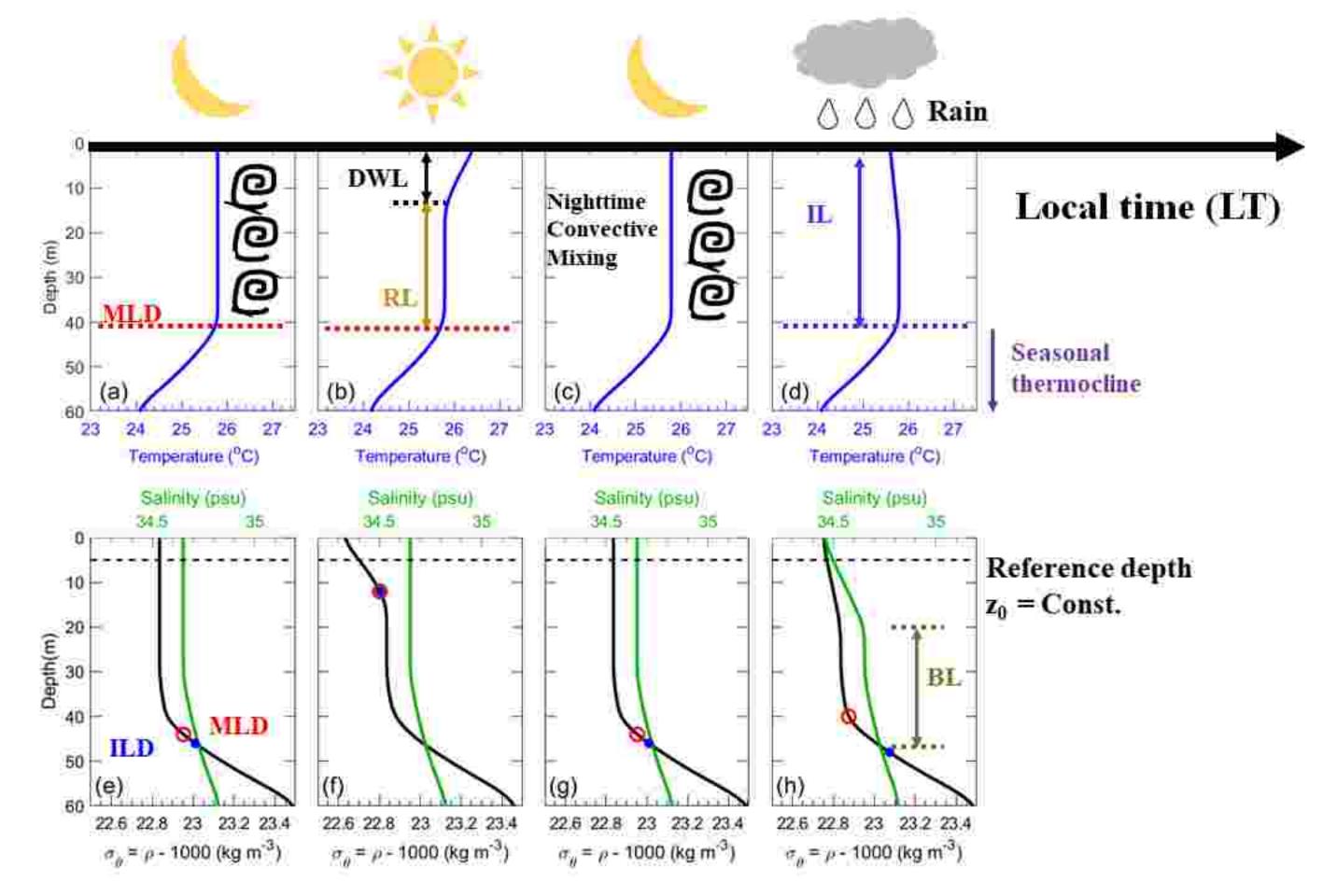
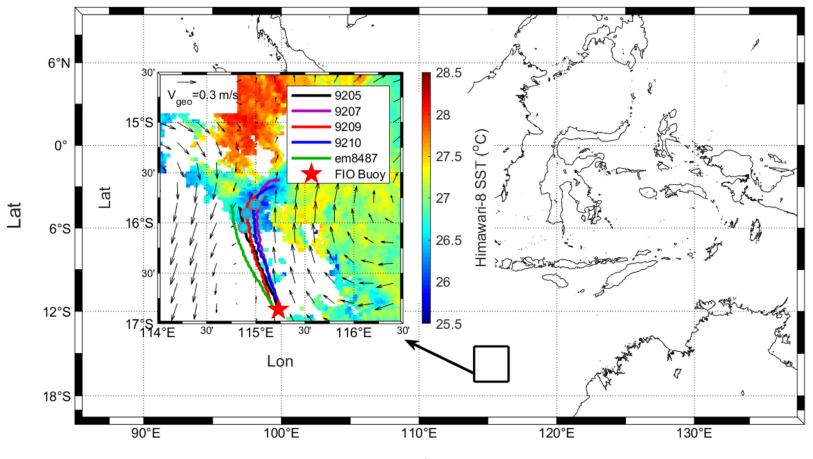
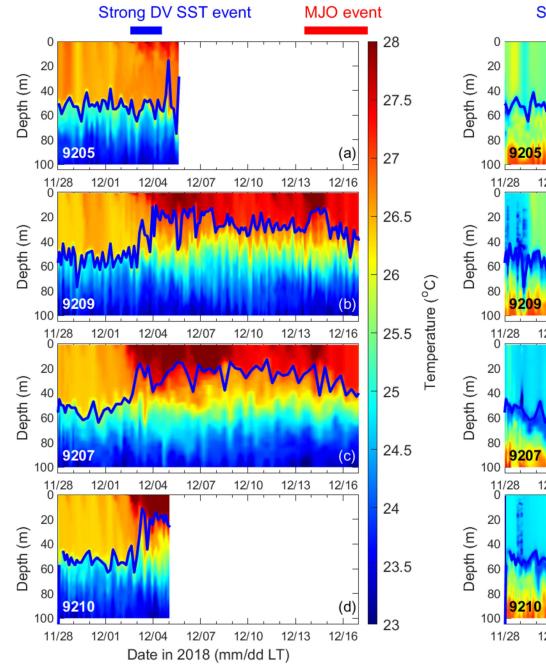


Figure2.



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Figure3.



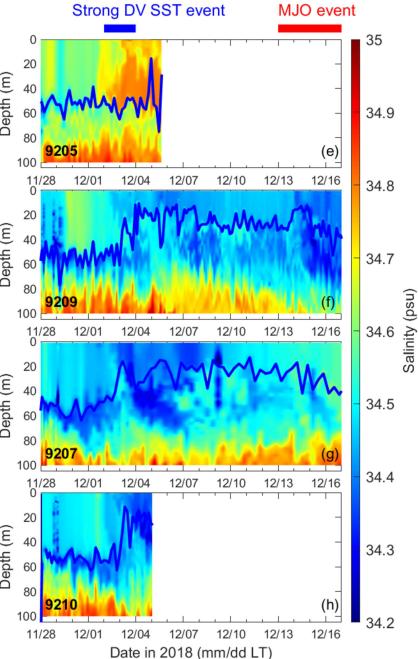
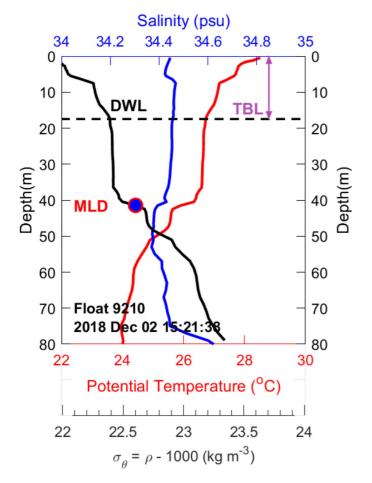


Figure4.





(b) Salinity-driven Barrier Layer (SBL)

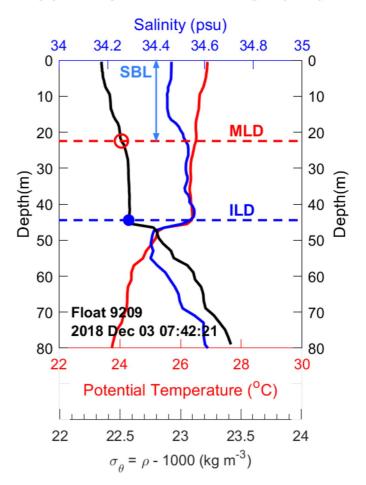


Figure5.

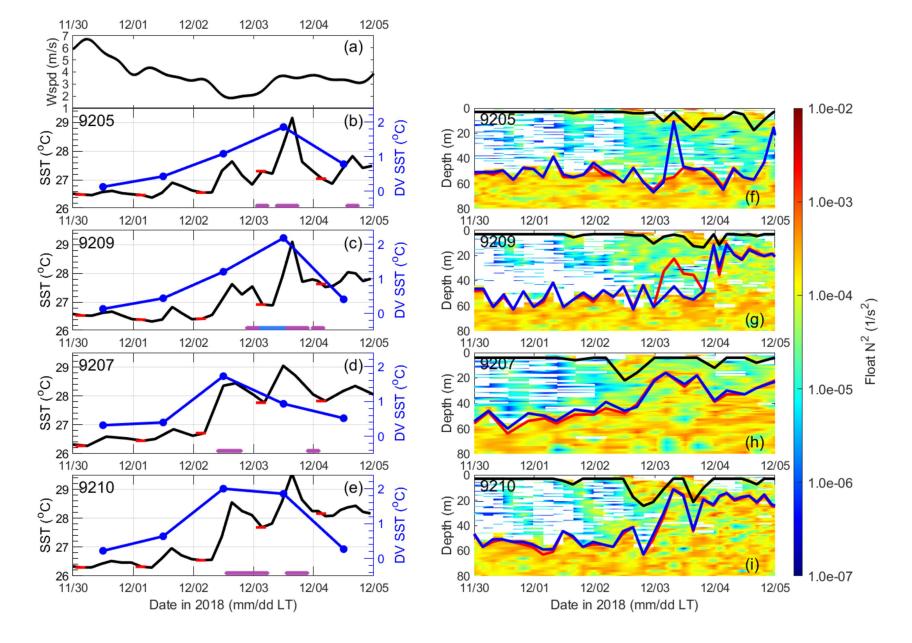


Figure6.

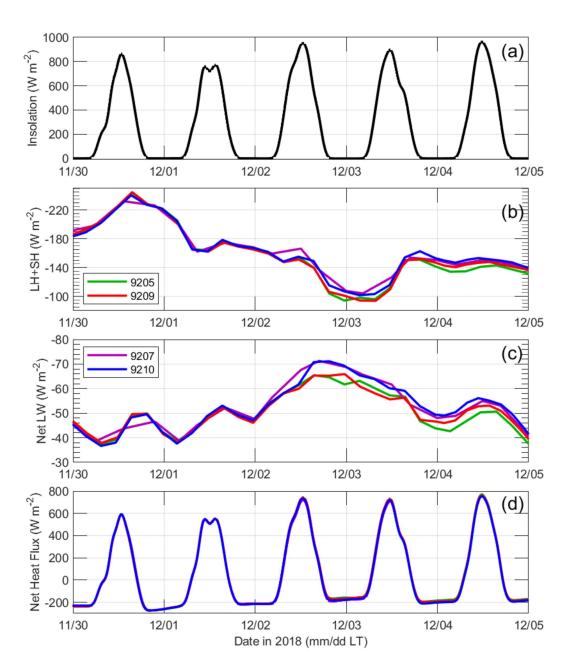


Figure7.

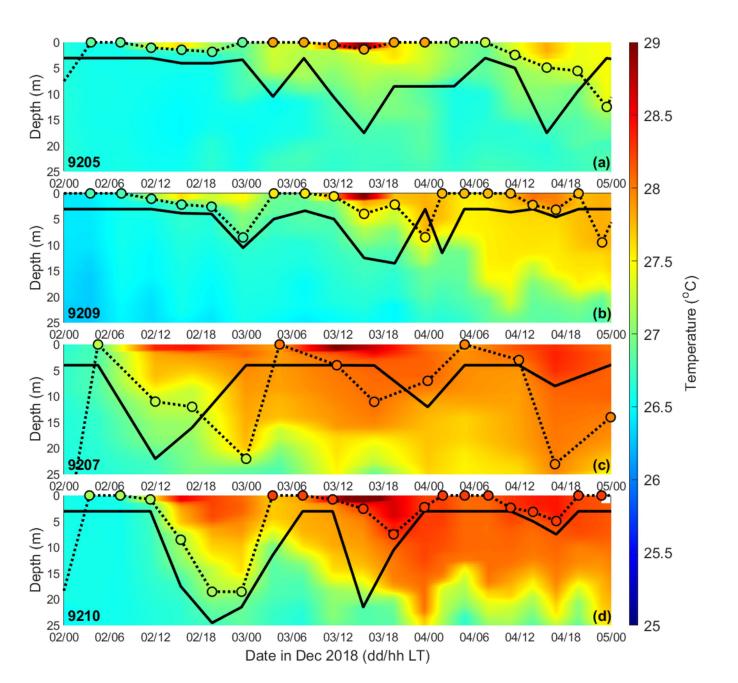


Figure8.

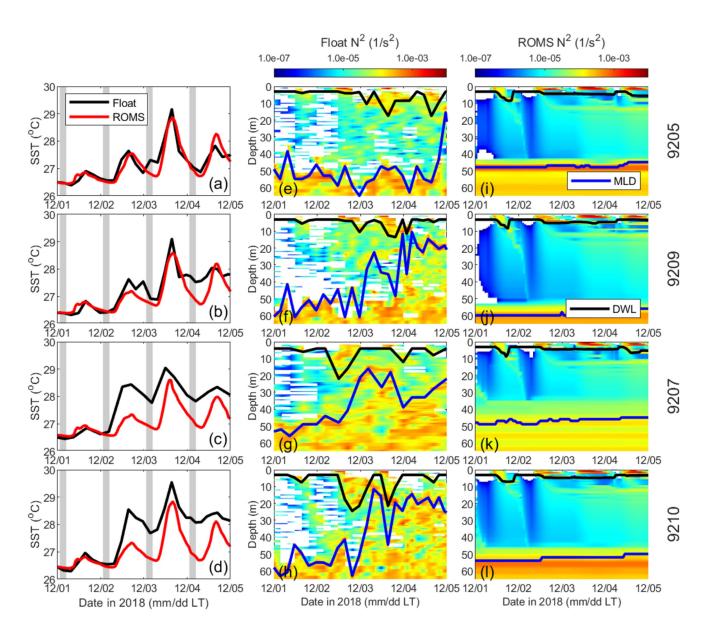


Figure9.

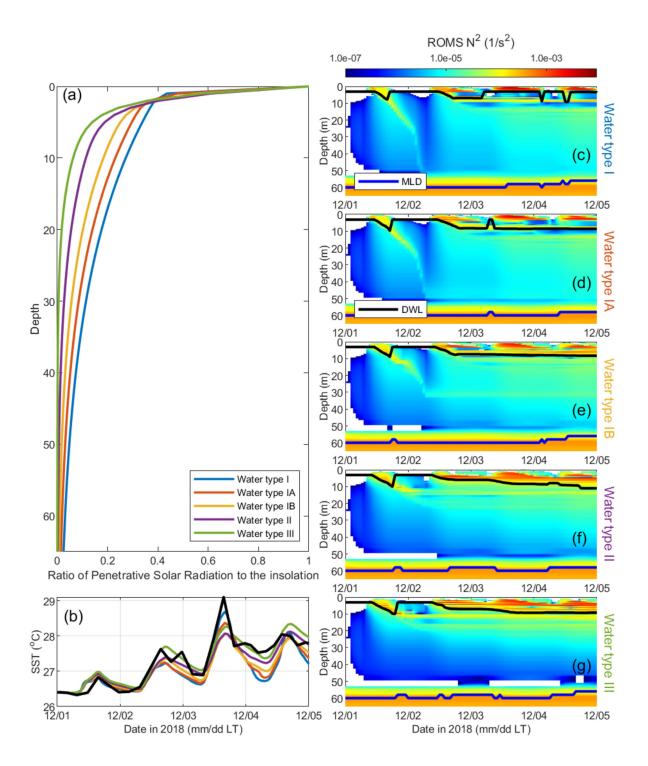


Figure10.

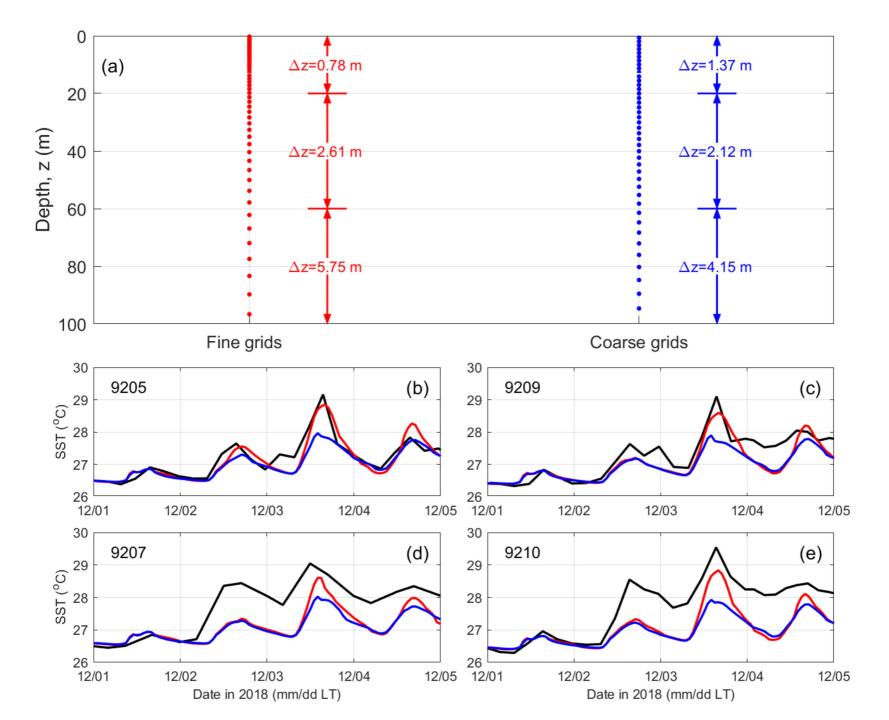


Figure11.

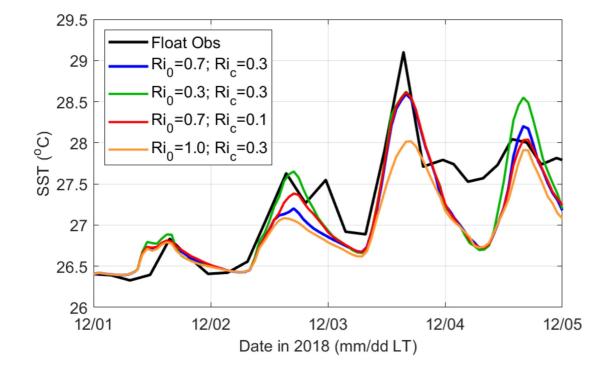


Figure12.

