

# Attenuation of ocean surface waves in pancake and frazil sea ice along the coast of the Chukchi Sea

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## Abstract

Alaskan Arctic coastlines are protected seasonally from ocean waves by presence of coastal and shorefast sea ice. This study presents field observations collected during the autumn freeze up of 2019 near Icy Cape, a coastal headland in the Chukchi Sea of the Western Arctic. The evolution of the coupled air-ice-ocean-wave system during a four-day wave event was monitored using drifting wave buoys, a cross-shore mooring array, and ship-based measurements. The incident wave field was attenuated by coastal pancake and frazil sea ice, reducing significant wave height by 1 m over less than 5 km of cross-shelf distance spanning water depths from 13 to 30 m. Spectral attenuation coefficients are evaluated with respect to wave and ice conditions and the proximity to the ice edge. Attenuation rates are found to be three times higher within 500 m of the ice edge, relative to values farther in the ice cover. Attenuation rates follow a power-law dependence on frequency, with an exponent in the range of  $(2.3, 2.7) \text{ m}^{-1}$ .

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3 **Attenuation of ocean surface waves in pancake and**  
4 **frazil sea ice along the coast of the Chukchi Sea**

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12 **Key Points:**

- 13 • Buoy observations are used to calculate spectral attenuation rates of surface waves  
14 in pancake and frazil sea ice near the coast of Alaska.
- 15 • Consistently higher attenuation is observed near the ice edge than further in the  
16 ice cover.
- 17 • Attenuation rates follow a power-law dependence in frequency and are applica-  
18 ble to parametrization schemes in wave forecast models.

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## Plain Language Summary

Changes in the Arctic sea ice cover have consequences for coastal Alaskan regions. Nearshore sea ice melts earlier and forms later in the year, exposing the coastlines to increased ocean wave energy and storm surges. Recent reports show that erosion along the Arctic coasts is on the rise and poses a threat to local habitats and human communities. This study aims to improve our understanding of the protective role of sea ice by measuring wave energy across the nearshore ice cover. Using drifting buoys deployed inside and outside fragmented sea ice, we monitored ocean waves during a storm event typical for coastal regions in the Chukchi Sea. We found that the wave heights were reduced by 1 m over 5 km distance and the effects of this type of ice on waves were consistent with previous studies. Thanks to high resolution of our measurements, we were able to determine that the dampening effect was stronger immediately next to the ice edge. Our measurements may be applied to improve present and future operational and climate models used to forecast and understand wave activity near the Arctic coasts.

## 1 Introduction

The Arctic region is a rapidly changing environment, characterized by increasing rates of summer sea ice decline, rising temperatures and lengthening open-water seasons. Arctic coastlines are considered particularly vulnerable to these changing conditions, which

pose a unique threat to biological systems, human communities, and infrastructure (Forbes, 2011). Indeed the erosion rates along the Arctic coast have been accelerating (Lantuit et al., 2012; A. E. Gibbs & Richmond, 2015; Gibbs et al., 2019), and the length of ice free season appears directly related to the higher erosion rates (Barnhart et al., 2014). While warmer temperatures are an obvious driver of erosion, in particular in areas with permafrost bluffs, mechanical processes associated with wave activity are likewise considered a leading contribution (Overeem et al., 2011). The effect of coastal sea ice in dissipating large waves and decreasing the magnitude of storm surges is reduced as the open water season lengthens and extends further into the autumn period of increased storminess in the Alaskan Arctic (Atkinson, 2005; Fang et al., 2018). At the same time, a rise in surface wave activity has been observed in the Chukchi and Beaufort Sea (X. L. Wang et al., 2015; Thomson et al., 2016), linked to increased fetch distance due to larger open water extent during ice-free season (Thomson & Rogers, 2014). Together, these changes in ice and wave conditions are expected to accelerate Arctic coastal erosion.

Quantifying the role of sea ice presence in coastline protection requires an accurate representation of wave and sea ice interactions, which span a wide range of conditions typical for the coastal Arctic. The complex and potentially nonlinear processes that govern these interactions pose a challenge to both observations and numerical models. In recent years, however, considerable progress has been achieved on both fronts (Squire, 2018). The present study focuses on a set of conditions that are becoming increasingly commonplace. The Alaskan Arctic in autumn is becoming increasingly defined by wave activity, especially conditions of frazil and pancake ice forming in wave fields of 2-3 m significant wave height (Thomson et al., 2018; Roach et al., 2018). Pancake ice is specifically associated with new ice formation in dynamic wave conditions (Doble et al., 2015). Wave attenuation in this type of ice is dominated by dissipative processes (as opposed to scattering, see Squire et al. (1995); Kohout and Meylan (2008)) and is typically formulated as an exponential decay of spectral wave energy  $E(x, f)$  with distance  $x$  travelled in ice, i.e.  $E(x, f) = E(0, f)e^{-\alpha(f)x}$ . Here,  $\alpha(f)$  is the spectral attenuation coefficient and its quantity is determined by a number of physical mechanisms with varying levels of contribution based on wave and ice conditions. Quantifying these small scale effects using process-based models presents a significant challenge (Shen & Squire, 1998) and are often difficult to reconcile with in situ observations. In operational wave models (WAVEWATCH III, SWAN), forecasters sometimes use empirical parametrizations

83 where the dissipative attenuation rate  $\alpha$  follows a power law in frequency  $\alpha(f) \propto f^n$ ,  
84 a relation demonstrated in four independent field studies by Meylan et al. (2018). Co-  
85 efficients of this function are obtained using the best fit to the  $\alpha$  values determined from  
86 in situ measurements in the Marginal Ice Zone (MIZ), where waves propagate from open  
87 water into ice cover. Most prior MIZ measurements, and most operational wave mod-  
88 els, are predominantly for deep-sea applications and large domains. The applicability of  
89 these parametrizations to coastal sea ice in nearshore conditions has not been investi-  
90 gated.

91 In addition to in situ observations, laboratory experiments using wave tanks pro-  
92 vide insight into wave-ice interactions under controlled and repeatable conditions (Ta-  
93 ble 1 in Parra et al. (2020) provides a useful overview). Tank measurements of wave dis-  
94 sipation are confined to smaller spacial scales than in situ observations, and typically they  
95 report attenuation rates of  $10^{-1} - 10^{-2} \text{ m}^{-1}$  (Herman et al., 2019; Shen, 2019), which  
96 are two or three orders of magnitude larger than those found in field experiments (Doble  
97 et al., 2015; Rogers et al., 2016). The additional damping is often attributed to mech-  
98 anisms associated with the inherent physical constraints of the laboratory setup such as  
99 overwash (Meylan et al., 2015), sidewall effects and properties of the materials simulat-  
100 ing the ice cover. However, other wave dissipation field measurements conducted on small  
101 spacial scales in the proximity to the ice edge have also reported reported higher dissi-  
102 pation rates (Rabault et al., 2017; Asplin et al., 2018) that are similar to those obtained  
103 in wave tanks.

104 Here we study a nearshore region of the Chukchi Sea, with water depths ranging  
105 from 13 to 30 m. We present an observational dataset capturing a four-day long wave  
106 event, with an aim to determine the magnitude of cross-shore wave attenuation in coastal  
107 pancake and frazil sea ice, and further constrain empirical models for use in coastal ap-  
108 plications. Sampling took place at the end of November 2019, when the Chukchi and Beau-  
109 fort region of Alaska experienced an unusually late onset of winter ice accompanied by  
110 increased wave activity. The measurements provide a unique record of the increasingly  
111 frequent open water conditions at a time of year when the coastline historically would  
112 have been protected by ice. The resulting dataset offers a relatively high spatial reso-  
113 lution (150 m over a 3000 m transect within sea ice) in close proximity of the ice edge.

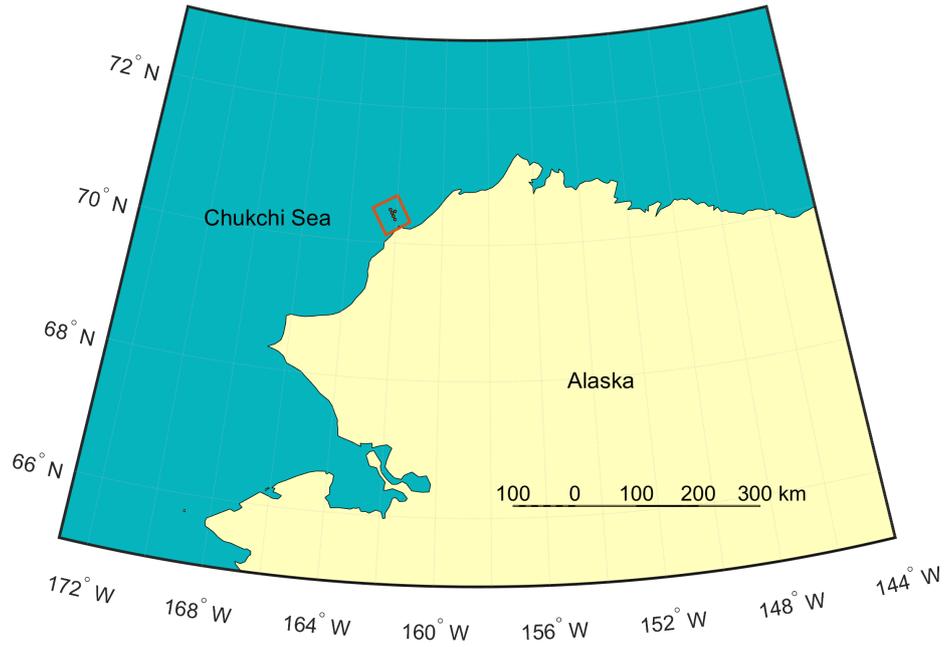
114 In Sections 2.1 and 2.2, we describe our experiment setup and the conditions at the  
115 site with an emphasis on wave and sea ice observations. In Section 2.3, we consider the  
116 effects of intermediate and shallow depth on our measurements and discuss our approach  
117 to evaluating spectral dissipation throughout the event. Section 3 presents apparent spec-  
118 tral attenuation rate with respect to ice type and proximity to ice edge and a compar-  
119 ison with past observations in the MIZ. Section 4 provides a discussion of the uncertainty  
120 in our ice edge estimate and a brief analysis of the evolution of sea ice and temperature  
121 throughout the event.

## 122 2 Methods

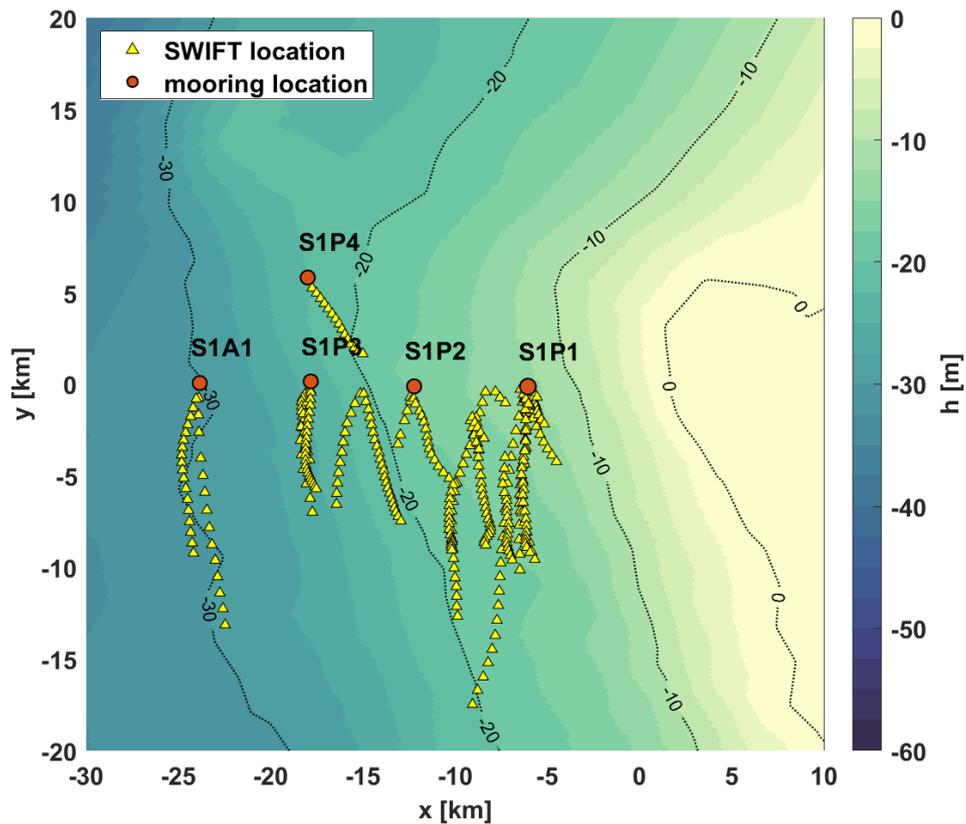
### 123 2.1 Description of field experiment

124 Data presented in this study were collected over a four-day period during the Coastal  
125 Ocean Dynamics in the Arctic (CODA) research cruise on the *R/V Sikuliaq* in Novem-  
126 ber 2019 near a barrier island system west of Icy Cape headland in the Chukchi Sea (Fig-  
127 ure 1a). The coastline is shaped by sand and gravel islands and barrier spit extensions  
128 from land. The larger Icy Cape region is considered erosional, with average shoreline change  
129 of -0.4 m/yr (A. E. Gibbs & Richmond, 2015). Starting on 21 November 2019, a low pres-  
130 sure system passed over the Icy Cape area and created an energetic wind sea, referred  
131 here as wave event. Drifting coastal and pancake ice was present near the coast, atten-  
132 uating the incoming wave field. The cruise objective was to study interactions of ocean  
133 surface waves and sea ice in nearshore conditions by means of mooring arrays and op-  
134 portunistic sampling of surface waves. Shoreward wave propagation was sampled over  
135 a transect of 20 km using five moorings and six drifting wave buoys, allowing us to ob-  
136 serve dissipative effects of ice as a function of proximity to the ice edge.

137 The array consisted of five moorings positioned in the cross-shore direction at depths  
138 increasing from 13 m to 30 m and locations  $x = 5$  to 25 km (Figure 1b).  $x$  here refers  
139 to the cross-shore distance from the coast increasing in the positive direction (reverse  
140  $x$ -axis in Figure 1b). The furthest offshore mooring (denoted as S1A1) was equipped with  
141 an Acoustic Doppler profiler (Nortek Signature1000) on a seafloor tripod, sampling waves  
142 and currents at 2 Hz. The remainder of the array (labeled S1P1 - S1P4) comprised seafloor  
143 pressure and temperature loggers (RBR Duet) each with additional temperature loggers  
144 (Onset HOBOS) strung along the sub-surface moorings. One of the moorings (S1P4) had



(a)



(b)

Figure 1: (a) Location of the Icy Cape study site within state of Alaska. (b) Detail of the study site, including a local coordinate system and ocean bathymetry (in meters) of the site obtained from the ETOPO1 dataset (Amante & Eakins, 2020). Circles represent mooring locations and triangles show trajectories of all drifting wave buoys deployed during the wave event.

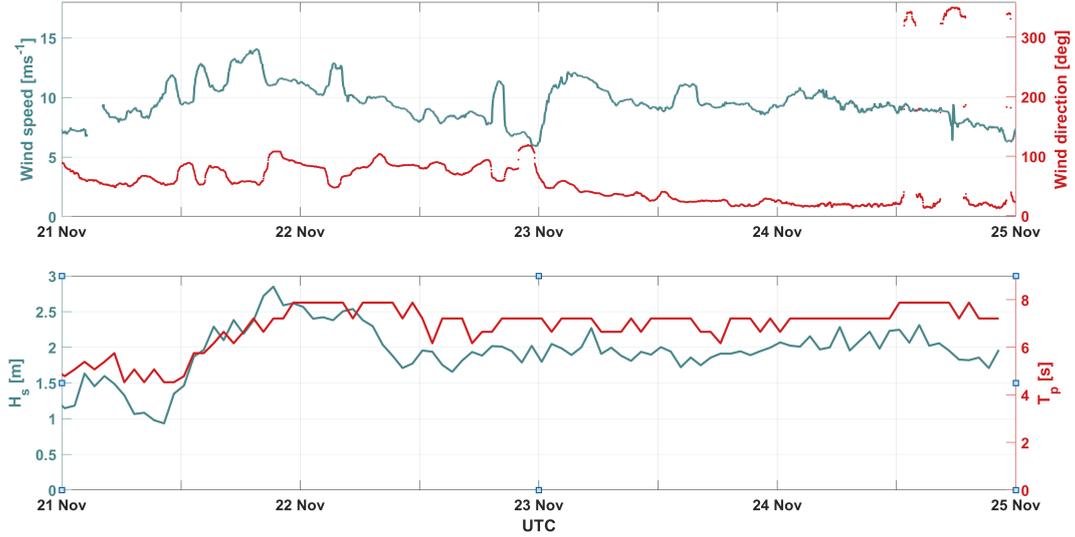


Figure 2: Top: Wind speed (blue) and direction (red) throughout the event wave event at Icy Cape obtained using ship-based anemometers on board *R/V Sikuliaq*. Bottom: Significant wave height (blue) and wave period (red) recorded at the location of the S1A1 mooring using the Nortek Signature1000.

145 an additional turbidity sensor and a SWIFT (Surface Wave Instrument Float with Track-  
 146 ing) (Thomson, 2012) buoy attached at the surface.

147 In addition to the continuous mooring observations, several freely drifting SWIFT  
 148 buoys were deployed from *R/V Sikuliaq* (Figure 1b). The goal was to obtain complemen-  
 149 tary data in the Lagrangian reference frame and cover a range of ice types and condi-  
 150 tions relative to the evolving ice edge, as well as to sample the wave activity in the along-  
 151 shore direction. Six SWIFT buoys were used throughout the experiment, deployed along  
 152 the cross-shore transect defined by the moorings S1P1 - S1A1 and recovered when in need  
 153 of maintenance or upon drifting too far away from the study site. Each buoy was equipped  
 154 with an inertial measurement unit (IMU), GPS, and radio and satellite transmitters. Some  
 155 carried additional instruments including anemometers, cameras, water and air temper-  
 156 ature loggers, providing further insight into the evolution of the wave event.

## 157 2.2 Wave event

158 The observations were collected over the course of a wave event at Icy Cape (Fig-  
 159 ure 2). The ship arrived at the site as the waves were building, and they peaked at the

160 end of the first day with 3 m recorded significant wave height  $H_s$  (integrated over fre-  
 161 quency domain  $0.0098 < f < 0.4902$  Hz). Wave heights remained at approximately 2  
 162 m for the remainder of data collection. Both wind and wave directions were from the north-  
 163 east and later from the north, with incident wave angle ranging between  $40^\circ$  and  $0^\circ$  with  
 164 respect to the cross-shore direction. Wind speed recorded by the ship-based anemome-  
 165 ters varied between 6-14 m/s, peaking on 21 November 2019. The peak wave period in-  
 166 creased from 5 s to 8 s over the course of the event.

### 167 *2.2.1 Sea ice measurements*

168 The type and extent of sea ice during the event is reconstructed from three inde-  
 169 pendent sources: images recorded every 5 s by cameras mounted on the SWIFT buoys,  
 170 hourly visual observations from *R/V Sikuliaq*, and Synthetic-aperture radar (SAR) im-  
 171 ages provided by RADARSAT-2. In addition, a small number of physical samples of pan-  
 172 cake and grease ice were collected using dip nets to determine thickness and quality. Dur-  
 173 ing the peak of the event, when pancake ice was most consolidated, the measured thick-  
 174 ness of samples ranged between 7-10 cm.

175 SAR images (Figure 3) are used to evaluate the extent and evolution of sea ice. Backscat-  
 176 ter characteristics are a measure of surface roughness and depend on the acquisition mode,  
 177 incidence angle, weather conditions, etc. In this case, ice appears as a bright area in the  
 178 early (Figure 3a) and late (Figure 3f) Sentinel-1 images, while it shows up as a low sig-  
 179 nal in the four RADARSAT-2 images (Figure 3b,c,d,e), probably because of the higher  
 180 sea state in open water. Images obtained on November 20, 22, 23 show that the ice edge  
 181 was located between the S1P1 and S1P2 moorings for the duration of our observations,  
 182 while the images taken after the event on November 25, 26 suggest that the sea ice re-  
 183 treated and became more patchy.

184 Cameras mounted on the masts of the SWIFT buoys record a low resolution im-  
 185 age every 5 seconds. The images were reviewed manually, and any unusable (i.e., blurry,  
 186 obscured by icing or darkness) images were discarded. The remaining images were sub-  
 187 jectively analysed and assigned an integer code on a scale of 0-12 to characterize the ice  
 188 type. This categorization was introduced in Rogers et al. (2018) and previously applied  
 189 to observations from the Arctic Sea State Experiment in 2015 (Thomson et al., 2019)  
 190 (see their Table 1). The ice codes range from less to more solid ice, and can be broadly

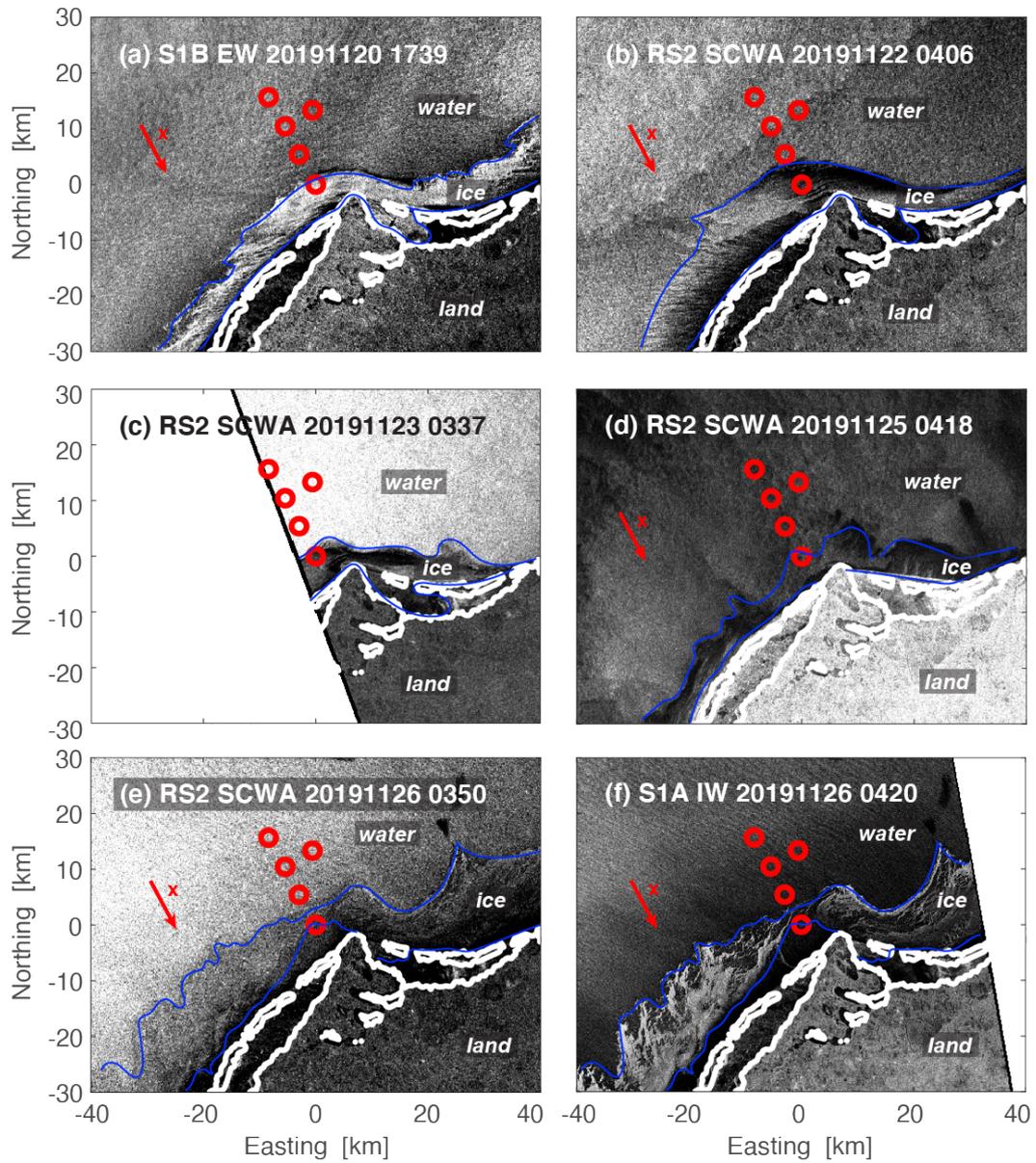


Figure 3: SAR images acquired by Sentinel-1 and RADARSAT-2 in the Icy Cape area before (a), during (b,c), and after (d,e,f) the wave event. Satellite, acquisition mode, date, and time are noted. White contours indicate the coast line, and the region interpreted as being ice-covered is contained by the blue contours. Red circles indicate mooring locations.

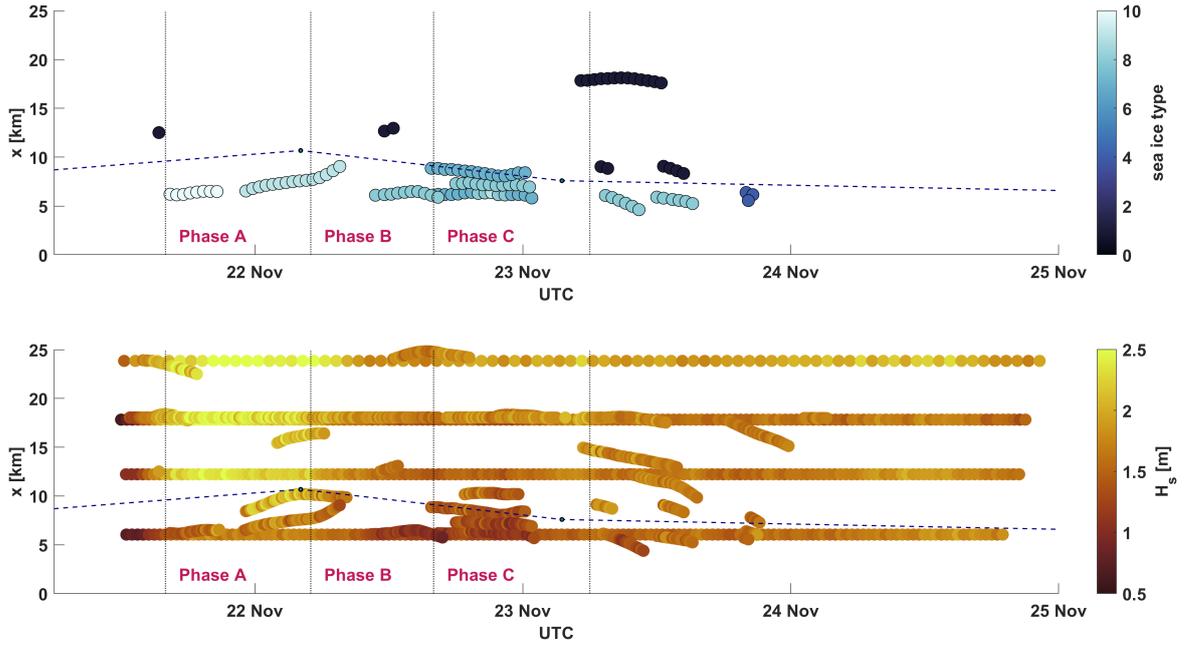


Figure 4: Top: Evolution of sea ice type during the wave event at Icy Cape, as a function of time and cross-shore distance. The plot combines information about ice type from SWIFT mounted cameras and visual observation upon deployment and recovery (colored circles), as well as ice edge location from RADARSAT-2 imagery (markers connected by dashed line). Vertical lines delineate event phases considered in further analysis. Bottom: Evolution of significant wave height during the event as a function of time and cross-shore distance, combining data reported by SWIFT buoys and mooring instruments.

191 grouped as 0-1 for open water and possible grease ice, 2-4 for frazil ice, 5-8 for brash ice  
 192 and small to medium-sized pancakes, and 9-12 for substantial pancake ice. Visual ob-  
 193 servations (from the ship) of sea ice evolution agree well with the dotted line represent-  
 194 ing interpolation between RADARSAT-2 ice edge estimates. Collectively these data al-  
 195 low us to quantify the location and type of sea ice throughout the event and to analyse  
 196 its impact on the incident wave field in Section 3.

197 Complementary to drifting cameras, hourly ship-based observations of sea ice were  
 198 performed according to the Arctic Ship-based Sea Ice Standardization Tool (ASSIST)  
 199 observation protocol (<http://www.iarc.uaf.edu/icewatch>). Additionally, information about  
 200 ice type was logged each time the ship stopped to take measurements. All the above records  
 201 have been combined into a single dataset presented in Figure 4 (top).

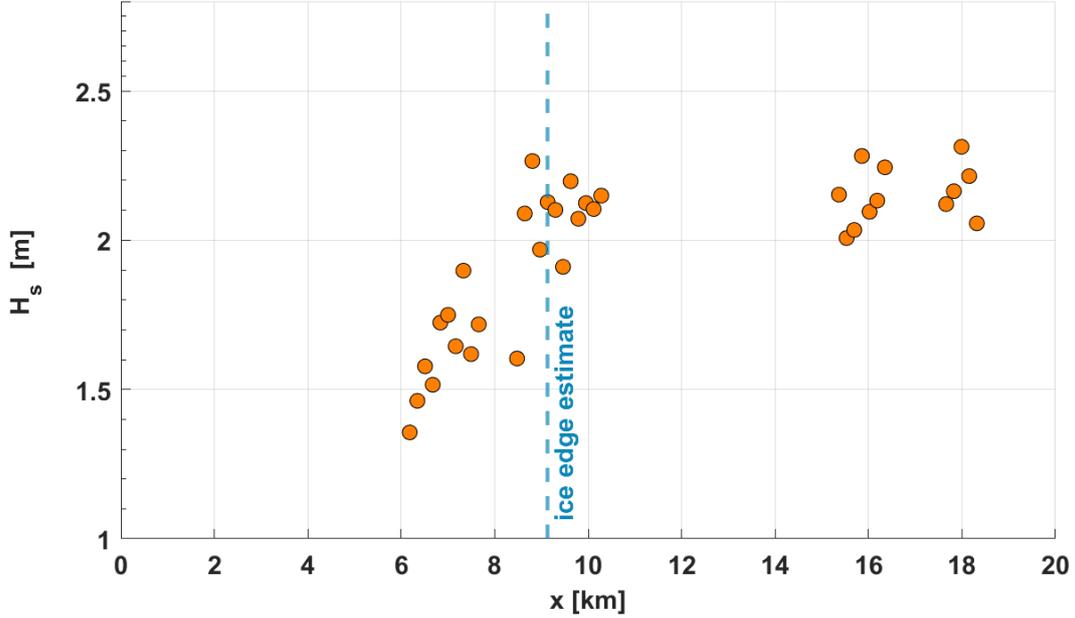


Figure 5: Significant wave height as a function of cross-shore distance at the peak of the event (Phase A). Vertical dotted line represents the ice edge estimate derived from RADARSAT-2 imagery. Note that the vertical axis starts at 1 m.

### 2.2.2 Wave measurements

The measurements of wave activity were collected in both Eulerian (cross-shore mooring array, a SWIFT buoy moored to the sea floor) and Lagrangian reference frame (drifting SWIFT buoys). This setup offers a good overview of the spatial evolution of the event. In particular, data from drifting SWIFT buoys reveal that the direction of surface currents was dominantly alongshore, and none of the quantities considered in our analysis evolved considerably in the drifting reference frame. This provides a good justification for neglecting alongshore coordinate  $y$  in the reference frame from Figure 1b and confining our analysis to cross-shore direction  $x$ .

Figure 4 (bottom) shows spatial and temporal evolution of the significant wave height calculated using data from all mooring arrays and SWIFT buoys. The event peaked on 21 Nov 2019 with incident waves reaching 3 m, and slowly died out over the period of the next three days. Figure 5 further illustrates the dissipative effect of sea ice at the peak of the event (Phase A), reducing the wave height by 0.7 m over less than 3 km.

216 The cross-shore timelines of sea ice and wave evolution offer a useful overview of  
 217 the event progression (Figure 4). *R/V Sikuliaq* arrived at the site as the waves were build-  
 218 ing up on 21 November. Peak of the wave activity started at 16:00 UTC and lasted for  
 219 the next 12 hours, during which substantial pancake ice was observed. Both wave heights  
 220 and size of ice floes started to decline on 22 November, as the pancake floes became mushy  
 221 and unconsolidated. On 23 November, only patches of frazil ice were observed and ac-  
 222 tive deployments of wave buoys were concluded. On the last day of the event we recorded  
 223 very sparse sea ice presence as the wave activity continued to decrease.

224 In the following analysis of spectral wave dissipation, we consider only SWIFT buoy  
 225 measurements collected during Phase A, B and C in Figure 4 to ensure statistical ro-  
 226 bustness of wave data within the ice cover. Figure 6 shows mean energy density spec-  
 227 tra measured in each phase and binned by the distance from the ice edge estimate de-  
 228 rived from SAR imagery.

229 A hard spectral cutoff  $E_n(f)/E(f, x) < 1/10$  has been applied to the data to pro-  
 230 duce Figure 6 and all further analysis in order to avoid spurious negative biases in at-  
 231 tenuation that originate in instrument noise.  $E_n(f)$  here is the spectral energy of the noise  
 232 and has been empirically determined to follow  $f^{-4}$  with an equivalent height  $H_n = 0.10$   
 233 m that is specific to the instrumentation and post processing method. Thomson et al.  
 234 (2020) show that noise can manifest as a flattening, or ‘rollover,’ in the high frequency  
 235 tail of spectral attenuation rates as waves dissipate in the ice and energy density approaches  
 236 the noise floor. The above cutoff substantially reduces this effect at a cost of discard-  
 237 ing a large portion of observations at the high frequency tail, as evidenced in Figure 6  
 238 where the lines furthest from the ice edge do not extend across the full frequency range.

239

## 240 **2.3 Analysis of spectral energy dissipation**

### 241 *2.3.1 Nearshore effects*

242 Effects of wave shoaling and refraction between locations S1P4 and S1P1 (Figure  
 243 1b) were estimated to be less than 5% and 10%, respectively over the considered time  
 244 period. The shallow water effect of nonlinearity, relative to dispersion, can be estimated  
 245 using the Ursell number  $Ur = \frac{a}{\kappa^2 h^3}$ , where  $a$  is the wave amplitude,  $\kappa$  is the wavenum-  
 246 ber magnitude and  $h$  is the water depth (Ursell, 1953). The magnitude of  $Ur$  at the shal-

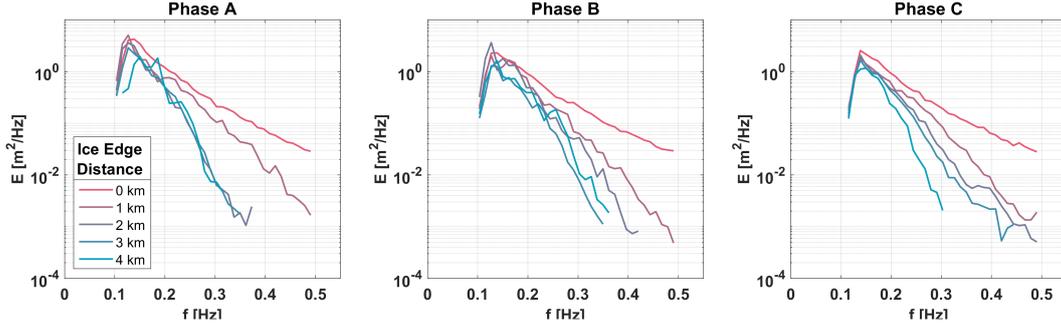


Figure 6: Mean ocean wave spectra binned by distance from the ice edge (line labeled 0 km represents mean incident spectrum) for phases A,B and C. Each bin comprises between 1 and 109 spectral estimates, each with 32 underlying degrees of freedom.

247 lowest observation (water depth 13 m) is only 0.2 during the most active phase of the  
 248 event, indicating that nonlinear triad interactions are weak. We conclude that the ef-  
 249 fects of intermediate and shallow depths on waves likely played a negligible role compared  
 250 to the dissipating effects of sea ice. Figure 5 corroborates this by showing a large change  
 251 in significant wave heights at the ice edge, and no trend outside of the ice.

### 252 *2.3.2 Binning method*

253 We choose phases A, B and C for analysis to categorize conditions that are qual-  
 254 itatively similar with respect to wave activity and sea ice type to ensure statistical sta-  
 255 tionarity. All quantities observed in a given phase are considered to be representative  
 256 and their variance in time is neglected.

257 To further simplify our analysis and increase sample size, we divide the ice covered  
 258 area into equidistant intervals along the  $x$ -axis, and bin average available wave data to  
 259 obtain  $E_i(f) = \bar{E}(f, x \in \langle x_i, x_{i+i} \rangle)$ , where  $i$  denotes the bin number and  $x_i$  refers to  
 260 delimiters of the intervals. Throughout this section, the bin size is set to  $\delta x = 150$  m,  
 261 chosen as the best compromise between spatial resolution and robustness of the data.

262 The location of the ice edge  $x_0$  is set to the mean distance interpolated from the  
 263 SAR imagery for each phase of the event, combined with available in situ observations  
 264 from the ship log and cameras on the SWIFT buoys. The incident wave field is obtained  
 265 as the mean spectral energy density of all measurements where  $x > x_0$ ,  $E_0(f) = \bar{E}(f)|_{\{x\} > x_0}$ .

266 **2.3.3 Attenuation coefficient**

267 In one dimension, we can express attenuation of the directional spectral energy den-  
 268 sity  $E_0$  across above defined bins as

$$269 \quad E_i(f) = E_0(f)e^{-\alpha_i(f)\Delta x_i} \quad (1)$$

270 where  $\Delta x_i = x_0 - x_i + \delta x/2$  corresponds to the mid-distance of the bin from the ice  
 271 edge  $x_0$ ,  $\alpha$  is the attenuation coefficient and  $E_i$  is the mean spectral density within bin  
 272  $i$ .

273 Correcting for the mean incident angle  $\theta$  of the wave direction with respect to the  
 274  $x$  axis, the attenuation coefficient for bin  $i$  is

$$275 \quad \alpha_i(f) = \frac{\cos \theta}{\Delta x_i} \ln \frac{E_0(f)}{E_i(f)} \quad (2)$$

276 **3 Results**

277 **3.1 Spectral attenuation rates**

278 In all three phases, we see a strong dependence of attenuation on the distance from  
 279 the ice edge. In particular, the values of  $\alpha(f)$  in the first 3 bins (corresponding to ap-  
 280 proximately 500 m distance from the ice edge) are substantially higher than in the far-  
 281 ther bins (Figure 7). The left plot (Figure 7) shows values of  $\alpha$  within 500 meters from  
 282 the ice edge ( $\alpha_{<500}$ ), while the middle plot shows attenuation coefficients farther in the  
 283 ice ( $\alpha_{>500}$ ). While the uncertainties of  $\alpha_{<500}$  are large due to the comparative scarcity  
 284 of wave data, there is a clear indication that the proximity to the ice edge plays a sig-  
 285 nificant role in the dissipation rate across all three phases. Meanwhile, the differences  
 286 between phases and associated ice types ranging from solid pancake ice in Phase A to  
 287 frazil ice in Phase C all lie within the standard error of the mean, obscuring any indi-  
 288 cations regarding their relative effects on the wave field. Combining data from all three  
 289 event phases while distinguishing only by proximity to ice edge (Figure 7, right) reduces  
 290 the uncertainty and shows the attenuation rate  $\alpha_{<500}$  to be approximately three times  
 291 higher than  $\alpha_{>500}$ . The dotted lines in Figure 7 show inferred attenuation rates from other  
 292 studies and are discussed in more detail below.

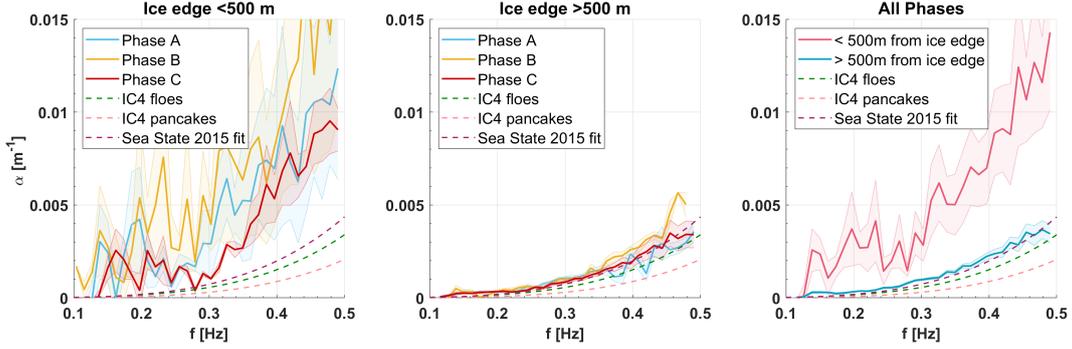


Figure 7: Attenuation coefficients at the Icy Cape wave event. Dashed lines show SWAN IC4 parametrizations (Rogers, 2019) and polynomial fit of Sea State 2015 dataset (Cheng et al., 2017). Shaded areas represent standard error of the mean taken over considered bins. Left: Attenuation coefficient during phases A, B and C using bins within 500 m of the ice edge. Center: Attenuation coefficient during phases A, B and C using bins farther than 500 m of the ice edge. Right: Attenuation coefficient averaged over phases A, B and C.

293 **3.2 Comparison with existing datasets**

294 A number of existing field observations of wave dissipation in pancake and frazil  
 295 ice, both in the Southern Ocean and more recently in western Arctic, allow us to dis-  
 296 cuss the above results in a wider context. While the magnitudes of the attenuation re-  
 297 ported by comparable experiments span several orders of magnitude ( $10^{-3} < \alpha(f) <$   
 298  $10^{-5}$ ), the spectral behaviour for loose, non-compact sea ice generally follows a power  
 299 law fit  $\alpha \propto f^n$  with  $n$ -values falling between 2 and 4, or a two-term polynomial fit  $\alpha =$   
 300  $af^2 + bf^4$  proposed by Meylan et al. (2014). In Figure 8, we have explored these op-  
 301 tions, along with a power law with an offset ( $\alpha = af^b + c$ ) while distinguishing between  
 302 data collected near and further from the ice edge. In both cases, all three formulas pro-  
 303 duce nearly identical fits using the nonlinear least squares method.

304 Wave dissipation data collected during the Sea State campaign in the Beaufort Sea  
 305 during autumn 2015 (Cheng et al., 2017) provide a convenient comparison. The wave  
 306 and ice conditions as well as instruments bear many similarities to the Icy Cape obser-  
 307 vations, even though the Sea State cruise sampled farther offshore in the deep-water MIZ.  
 308 The Sea State 2015 results, using the formula  $\alpha = af^3 + bf^4$ , are shown as a purple

309 dashed line in Figure 8. The new CODA 2019 results described herein are the solid lines,  
310 and they agree well in the region  $> 500$  m from the ice edge. However, in the region  $<$   
311  $500$  m from the ice edge, our new results exceed the those previous dissipation estimates  
312 by a factor of three.

313 The other two dashed lines in Figure 8 show the shapes of two empirical parametriza-  
314 tions used in WAVEWATCH III (Rogers et al., 2018; The WAVEWATCH III<sup>®</sup> Devel-  
315 opment Group, 2016) and SWAN wave models (Rogers, 2019) to represent dissipation  
316 effects of sea ice on waves, distinguishing broadly between ice floes (defined as between  
317 10 and 25 m in diameter) and pancake ice. These parametrizations (denoted IC4M2 in  
318 both models) follow similar frequency dependence as our best fits albeit smaller in mag-  
319 nitude, in particular the ‘pancake’ option which would be considered representative of  
320 the conditions at Icy Cape. We conclude that the data obtained is consistent with the  
321 IC4 parametrizations, and might be used to further constrain the wave models. Table  
322 1 summarizes fitting parameters used in generating these curves, along with their con-  
323 fidence intervals.

324 Figures 7 and 8 suggest that the proximity to the ice edge plays a dominant role  
325 in the magnitude of wave dissipation. While the variability in  $\alpha$  could be partially ex-  
326 plained by inhomogeneity of the ice cover and uncertainty in our ice edge estimate, this  
327 signal remains consistent throughout all analysed wave conditions and ice types. This  
328 suggests that the attenuation of the incident wave field is not constant throughout the  
329 ice cover, and disproportionately larger energy loss occurs in the vicinity of the ice edge.  
330 While our data offer no indication of what physical processes might cause this effect, this  
331 result is consistent with visual observations where waves undergo almost instantaneous  
332 damping in high frequencies as they travel past the ice edge. However, this visual ob-  
333 servation may merely reflect the high dissipation rate of higher frequencies in general,  
334 which is ubiquitous in this and similar studies. The more novel feature of our results is  
335 that they indicate faster dissipation near the ice edge at *all* frequencies. Most previous  
336 field studies have far less spatial resolution (e.g., buoy spacing of 10 km in Sea State 2015)  
337 compared with the present study, and thus it is possible that similar differences near the  
338 ice edge were obscured.

339 The higher dissipation rates near the ice edge are qualitatively consistent with lab-  
340 oratory studies which are inherently measuring dissipation near the ice edge (e.g. R. Wang

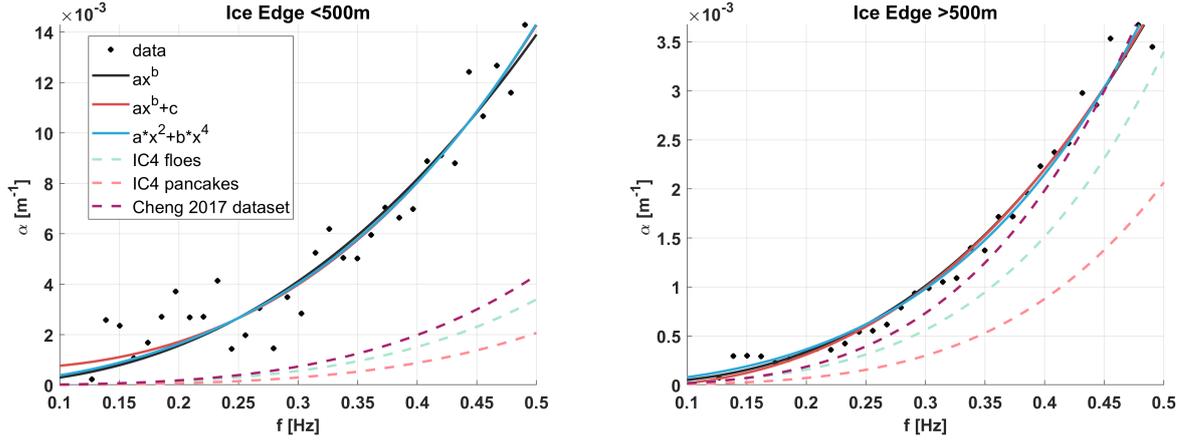


Figure 8: Selected polynomial fits to attenuation coefficients measured within and past 500m distance from the ice edge, compared with IC4 parametrization (Rogers, 2019) and polynomial fit of Sea State 2015 dataset published in (Cheng et al., 2017).

341 & Shen, 2010; Cheng et al., 2019; Parra et al., 2020), and prior field measurements of  
 342 dissipation over small spatial scales (e.g. Rabault et al., 2017; Asplin et al., 2018).

Fitting formula	Ice edge < 500 m	Ice edge > 500 m
$af^b$	$a = 0.072(0.048, 0.097)$ $b = 2.3(2.0, 2.8)$	$a = 0.026(0.022, 0.030)$ $b = 2.7(2.5, 2.9)$
	$R^2 = 0.91$	$R^2 = 0.99$
$af^b + c$	$a = 0.091(0.038, 0.14)$ $b = 2.7(1.9, 3.5)$ $c = 6 \times 10^{-4}(-5.3 \times 10^{-4}, 1.7 \times 10^{-3})$	$a = 0.026(0.02, 0.032)$ $b = 2.7(2.4, 3)$ $c = -3.4 \times 10^{-5}(-1.6 \times 10^{-4}, 1 \times 10^{-4})$
	$R^2 = 0.92$	$R^2 = 0.99$
$af^2 + bf^4$	$a = 0.038(0.026, 0.049)$ $b = 0.078(0.015, 0.14)$	$a = 0.0076(0.0062, 0.0091)$ $b = 0.036(0.028, 0.044)$
	$R^2 = 0.92$	$R^2 = 0.98$

Table 1: Parameter estimates of the polynomial fits for  $\alpha(f)$  in Figure 8 evaluated using the nonlinear least squares method, along with 95% confidence interval.  $R^2$  represents a measure of goodness of fit.

343 Differences in attenuation rates in Figure 7 challenge our assumption that  $\alpha(f)$  in  
 344 (1) is homogeneous across the ice cover. Squire (2018) suggests that formula (1) can be  
 345 replaced with a more generalized (nonlinear) form  $d_x E = -\alpha E^n$  to address existing  
 346 issues of fitting exponential function to observational data. In particular, several exper-  
 347 iments (Kohout et al., 2014; Montiel et al., 2018) suggest that wave heights exceeding  
 348 3 m reduce linearly rather than exponentially. This implies a reduction of linear expo-  
 349 nential growth rate for larger wave heights. This is qualitatively reversed from the sit-  
 350 uation at Icy Cape, in which the dissipation is higher nearer the ice edge, where wave  
 351 heights are largest. However, there are two caveats. Firstly, wave heights in the present  
 352 case are below 3 m, where Kohout et al. (2014); Montiel et al. (2018) predict linear ex-  
 353 ponential decay. Secondly, our analysis indicates reduction of attenuation coefficients with  
 354  $x$ ; we do not explicitly compute dependence on wave height.

## 355 4 Discussion

356 The attenuation of waves approaching the Arctic coasts has broad implications for  
 357 a range of coastal processes and practical applications. The ability of wave forecast mod-  
 358 els to predict this attenuation is dependent on both a skilled understanding of what hap-  
 359 pens right at the ice edge (and ability to determine where the ice edge is located), as well  
 360 as an understanding of the coupled processes by which the waves and ice evolve. Here  
 361 we discuss both issues.

### 362 4.1 Ice edge uncertainty

363 We have investigated the possibility that the difference between  $\alpha_{>500}$  and  $\alpha_{<500}$   
 364 is a spurious result of our analysis. In particular, our estimate of the ice edge location  
 365 predominantly relies on our interpretation of the first three SAR images in Figure 3 (pan-  
 366 els (a), (b) and (c)), in addition to in situ observations obtained from ship logs and SWIFT  
 367 cameras. While this behaviour is found in all three phases considered, the uncertainty  
 368 of our ice edge estimate and the temporal averaging might prevent us from fully resolv-  
 369 ing effects on the scale of 500 m. We attempted to reduce this uncertainty by focusing  
 370 on a four hour time window around the satellite image on 22 November (Figure 3) which  
 371 provides our best estimate of the farthest cross-shore location of the ice edge during the  
 372 peak of the event. While the dataset confined in this smaller time window is not suffi-  
 373 ciently robust to allow full analysis similar that in Section 3 (only a small number of cross-

374 shore bins are populated with data), the same effect is observed, even though it is largely  
375 confined to the bin nearest to the ice edge, i.e., within the first 150 m. This is consis-  
376 tent with the supposition that the bulk of the damping occurs immediately after the waves  
377 enter the ice cover. The range of 500 m reported in our full analysis could be a result  
378 of a ‘smearing’ effect originating in the temporal averaging and uncertainty in ice edge  
379 location. To fully explore whether this effect is real and the physical processes involved  
380 would require a more persistent observation at the ice edge with even higher spatial res-  
381 olution.

## 382 4.2 Evolution of sea ice and sea surface temperature

383 Figure 4 indicates that both sea ice coverage and type changed rapidly during the  
384 event and that this ultimately had the dominant impact on the wave energy in the nearshore.  
385 Ice type transitioned from consolidated pancakes 10 cm thick at the beginning of the wave  
386 event to patches of grease ice towards the end, while the ice edge retreated shoreward  
387 past our instrument range. Sea surface temperature measurements show a strong cross-  
388 shore gradient, with freezing temperatures coinciding with the ice edge (Figure 9) and  
389 remaining above freezing in the open water during the initial stages. In later phases the  
390 sea ice retreated while the temperature difference decreased, although the air temper-  
391 ature remained well below freezing. We estimated that the net surface heat flux remained  
392 negative despite the sea ice retreat, albeit with an increase from approximately  $-200 \text{ W/m}^2$   
393 to  $-100 \text{ W/m}^2$  in the nearshore in later stages. Remote sensing imagery in Figure 3 sug-  
394 gests that the ice coverage in the broader Icy Cape region was patchy during and after  
395 the event. Combining the patterns of sea surface temperature and satellite imagery, it  
396 seems likely that the ice retreat at Icy Cape was caused by a combination of advection  
397 and local melting. Intrusion of warmer water was detected at 20 m depth on the third  
398 day of the event. We speculate that such temporary ice retreats might be a common episodic  
399 phenomenon during the autumn freeze up, with associated effects on wave attenuation  
400 (or lack thereof).

401 The presence of sea ice in the vicinity of Arctic coast has a leading order effect on  
402 shoreline erosion (Barnhart et al., 2014). As the duration of the open water season in  
403 the area increases, so does the sensitivity of the coastlines to storm surges in the autumn  
404 months. In 2019, the onset of coastal sea ice in the observed area was uncharacteristi-  
405 cally late, leaving the shoreline exposed to wave events, such as the one documented in

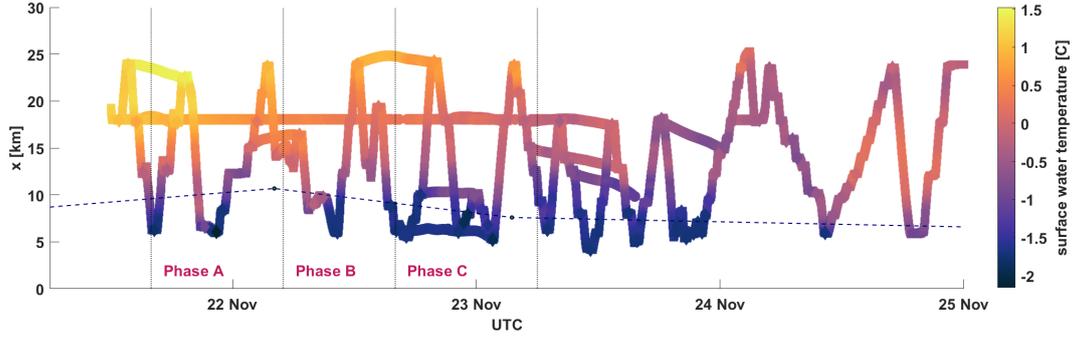


Figure 9: Top: Evolution of surface water temperature during the wave event at Icy Cape as a function of cross-shore distance and time. The plot combines information from SWIFT buoys, flow-through temperature sensor on the *R/V Sikuliaq* and temperature recorded in the top 1m of ship-based CTD casts. Dashed line representing sea ice extent estimate is added for reference.

406 this study, all through November. The above discussion suggests that oceanographic pro-  
 407 cesses may be as important as the local surface heat fluxes to determining the presence  
 408 and fate of ice. This implies that coupled ocean–sea ice models might be necessary to  
 409 reliably predict the cross-shore propagation of wave energy flux (which is determined by  
 410  $\alpha$ ) in this region.

## 411 5 Conclusions

412 The following conclusions are made from our analysis:

- 413 • Spectral energy dissipation in pancake and frazil ice measured  $> 500$  m from the  
 414 ice edge is consistent with published observations of similar conditions.
- 415 • Higher attenuation rates are observed near the ice edge, suggesting that a linear  
 416 exponential attenuation formula may not be valid universally across the ice cover.  
 417 Further measurements capable of resolving wave activity in the immediate prox-  
 418 imity of the ice edge are needed to understand this effect and underlying phys-  
 419 ical processes.
- 420 • Power dependence on frequency is found to be consistent across the ice cover, even  
 421 though coefficients of proportionality are not.

- 422 • Spectral attenuation rates observed during the event are compatible with the IC4  
 423 parametrization scheme used in WAVEWATCH III and SWAN and can be applied  
 424 to constrain these wave models.
- 425 • Coupled ocean-wave-sea ice models might be necessary to represent the evolution  
 426 of nearshore ice and wave conditions in the autumn season due to complexity of  
 427 the interplay between thermodynamic and oceanographic drivers.

428 These results may be applied on synoptic, seasonal, and decadal time scales to under-  
 429 stand the diminishing protection of Arctic coasts by sea ice and the increasing poten-  
 430 tial for wave-driven coastal erosion.

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 445 Repository (R2R) archive for the Coastal Ocean Dynamics in the Arctic (CODA) cruise  
 446 (doi:10.7284/908599). Data also will be available from the US Arctic Data Center.

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