

# Inverse Analysis of Experimental Scale Turbidity Currents by Deep Learning Neural Network

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## Abstract

Despite the importance of turbidity currents in environmental and resource geology, their flow conditions and mechanism are not well understood. To resolve this issue, a novel method for the inverse analysis of turbidity current using deep learning neural network (DNN) was proposed. This study aims to verify this method using artificial and flume experiment datasets. Development of inverse model by DNN involves two steps. First, artificial datasets of turbidites are produced using a forward model based on shallow water equation. To develop a inverse model, DNN then explores the functional relationship between initial flow conditions and characteristics of the turbidite deposit through the processing of artificial datasets. The developed inverse model was applied to 200 sets of artificial test data and four sets of experiment data. Results of inverse analysis of artificial test data indicated that the flow conditions can be precisely reconstructed from depositional characteristics of turbidites. For experimental turbidites, spatial distributions of grain size and thickness were accurately reconstructed. With regard to hydraulic conditions, reconstructed values of flow heights, sediment concentrations, and flow durations were close to the measured values. In contrast to the other values, there was a larger discrepancy between the measured and reconstructed values of flow velocity, which may be attributed to inaccuracies in sediment entrainment functions employed in the forward model.

# Inverse Analysis of Experimental Scale Turbidity Currents by Deep Learning Neural Network

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## Key Points:

- A new method for inverse analysis of turbidity currents using deep learning neural network was performed on turbidity current deposits in experimental scale.
- Results of inverse analysis conducted for artificial datasets proved that flow conditions can be precisely reconstructed from depositional characteristics.
- Flow conditions and deposit profiles in flume experiments were also well reconstructed except for flow velocity.

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**Abstract**

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Despite the importance of turbidity currents in environmental and resource geology, their flow conditions and mechanism are not well understood. To resolve this issue, a novel method for the inverse analysis of turbidity current using deep learning neural network (DNN) was proposed. This study aims to verify this method using artificial and flume experiment datasets. Development of inverse model by DNN involves two steps. First, artificial datasets of turbidites are produced using a forward model based on shallow water equation. To develop a inverse model, DNN then explores the functional relationship between initial flow conditions and characteristics of the turbidite deposit through the processing of artificial datasets. The developed inverse model was applied to 200 sets of artificial test data and four sets of experiment data. Results of inverse analysis of artificial test data indicated that the flow conditions can be precisely reconstructed from depositional characteristics of turbidites. For experimental turbidites, spatial distributions of grain size and thickness were accurately reconstructed. With regard to hydraulic conditions, reconstructed values of flow heights, sediment concentrations, and flow durations were close to the measured values. In contrast to the other values, there was a larger discrepancy between the measured and reconstructed values of flow velocity, which may be attributed to inaccuracies in sediment entrainment functions employed in the forward model.

## 1 Introduction

A turbidity current is a process of sediment transport into subaqueous environments such as deep lakes and ocean (Daly, 1936; Johnson, 1939). The deposits of turbidity currents are called turbidites, which are often characterized by graded bedding and sedimentary succession called the Bouma sequence (e.g., Kuenen & Migliorini, 1950; Bouma, 1962; Talling et al., 2012). Turbidite deposits have been an active area of study because of their close association with petroleum resources and their role in the destruction of sea-floor equipment, such as the submarine cables (Weimer & Slatt, 2007; Talling et al., 2015). Furthermore, turbidites are often deposited as a result of tsunami triggered turbidity currents (Arai et al., 2013), and thus are potentially useful for estimating the recurrence intervals of geohazards.

To understand the characteristics of turbidites and their implications, it is essential to study the flow behavior of turbidity currents (Talling et al., 2007). However, knowledge in this area remains limited because of difficulties in the direct observation of turbidity currents. A few in-situ measurements have been made (e.g., Xu et al., 2004; Vangriesheim et al., 2009; Arai et al., 2013; Paull et al., 2018), but hydraulic conditions measured, such as sediment concentration and flow velocity, were unclear because of the destructive nature and unpredictable occurrences of turbidity currents (Naruse & Olariu, 2008; Falcini et al., 2009; Lesshafft & Marquet, 2010; Talling et al., 2015). Therefore, inverse analysis that reconstructs the flow conditions of turbidity currents from their deposits is crucial for estimating the flow conditions in natural environments.

Prior to this research, inverse analysis of turbidity currents was conducted by Baas et al. (2000), where flow velocity was reconstructed through analyses of sedimentary structures of turbidites. The results gave an estimation of the hydraulic conditions of flow at a single location, but did not provide a reconstruction of spatial evolution of the turbidity current. By contrast, inverse analysis methods in previous studies based on numerical models provided more detailed insights to the spatial structure of flow and the evolution of flow over time (e.g., Falcini et al., 2009; Lesshafft & Marquet, 2010; Parkinson et al., 2017). However, the method proposed by Falcini et al. (2009) assumed steady flow conditions and was simplified for obtaining analytical solutions, preventing it from accurately illustrating the flow mechanism of unsteady turbidity currents that can produce normally graded bedding. Consequently, this method can not be applied to normally graded beds, which are typical characteristics of turbidites. Other studies employed the optimization method estimated the hydraulic parameters through optimizing the input parameters of numerical models, so that the resulting calculations fit the observed data from turbidites (Lesshafft & Marquet, 2010). This method can provide a relatively good reconstruction of the hydraulic conditions of turbidity currents, but has extremely heavy calculation load. Therefore, it is impossible to apply the method to natural scale turbidites, which typically run over tens to hundreds of kilometers and flow continuously for several hours (Talling et al., 2015). Optimization using the adjoint approach proposed by Parkinson et al. (2017) solved the problem of heavy calculation load, but the reconstructed values were very low in accuracy.

To resolve the aforementioned issues, Naruse and Nakao (2017) developed a new method for inverse analysis of turbidite deposits using deep learning neural network (DNN). DNN is a computational scheme that works as a universal function approximator (Liang & Srikant, 2016). Previously, it was applied to problems such as landslide susceptibility analyses (Pradhan et al., 2010) and identification of lithology from well log data (Rogers et al., 1992), where the empirical relationship between the observed data and the parameters aimed to be predicted were explored. In case of turbidity currents, however, it is impossible to obtain sufficient datasets of in-situ measurements of flow characteristics for developing a DNN inverse model. Thus, Naruse and Nakao

(2017) produced artificial datasets of turbidites using a numerical model. The produced datasets were inputted into DNN to explore the functional relationship between the deposits and the initial flow conditions. After this process, which is referred to as training, DNN becomes capable of making estimations of flow conditions from new turbidite data. Although it was proven by Naruse and Nakao (2017) that DNN is capable of reconstructing flow properties from artificial test datasets, it has yet to be tested with actual turbidite data.

In this study, we verified the ability of DNN to perform inverse analysis of turbidity currents by applying DNN to data collected from actual turbidites deposited in flume experiments. During each flume experiment, a turbidity current was generated under a controlled environment. Conditions including flow duration and initial hydraulic conditions can be set manually, and measurements of these parameters can also be conducted easily during the flow. Thus, rather than data measured in nature, data collected from flume experiment works better as a first step to verify the accuracy of DNN for inverse analysis of turbidity currents.

Here, we implemented the forward model and DNN for inverse analysis. Trained DNN was first tested with artificial data of experimental scale turbidity currents. Then, trained DNN was tested with flume experiment data. Initial flow conditions of experiments were reconstructed from deposit profiles sampled. These values of flow conditions were then fed into the forward model to reconstruct the spatio-temporal evolution of the flow. Reconstructed hydraulic conditions during the flow and the eventual grain size distribution of the deposits were compared with the measured values.

## 2 Forward Model

### 2.1 Governing Equations

The forward model implemented in this study is a layer-averaged model based on Kostic and Parker (2006). It is expanded to account for the transport and deposition of non-uniform grain size distribution discretized to multiple grain-size classes in Nakao and Naruse (2017) (Figure 1). The five governing equations are as follows:

$$\frac{\partial H}{\partial t} + U \frac{\partial H}{\partial x} = e_w U - H \frac{\partial U}{\partial x}, \quad (1)$$

$$\frac{\partial U}{\partial t} + U \frac{\partial U}{\partial x} = RC_T g (S - \frac{\partial H}{\partial x}) - \frac{1}{2} g H R \frac{\partial C_T}{\partial x} - \frac{U^2}{H} (c_f - e_w), \quad (2)$$

$$\frac{\partial C_i}{\partial t} + U \frac{\partial C_i}{\partial x} = \frac{w_i}{H} (F_i e_{si} - r_o C_i) - \frac{e_w C_i U}{H}, \quad (3)$$

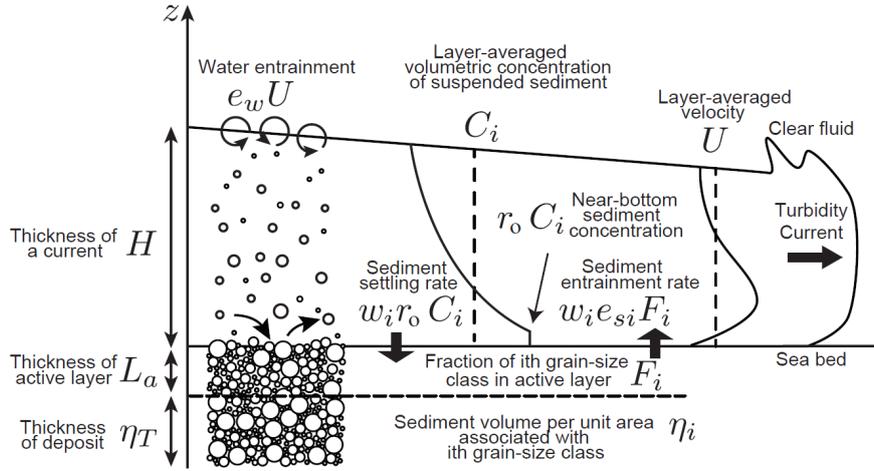
$$\frac{\partial \eta_i}{\partial t} = \frac{w_i}{1 - \lambda_p} (r_o C_i - e_{si} F_i), \quad (4)$$

$$\frac{\partial F_i}{\partial t} + \frac{F_i}{L_a} \frac{\partial \eta_i}{\partial t} = \frac{w_i}{L_a (1 - \lambda_p)} (r_o C_i - e_{si} F_i), \quad (5)$$

where the equations represent fluid mass conservation (equation 1), momentum conservation (equation 2), sediment mass conservation (equation 3), mass conservation in bed (Exner's equation) (equation 4) and sediment mass conservation in active layer (equation 5) (Nakao & Naruse, 2017).

Let  $x$  and  $t$  be the bed-attached streamwise coordinate and time, respectively. Parameters  $H$ ,  $U$  and  $C_i$  are the flow height, the layer-averaged velocity and the layer-averaged volumetric concentration of suspended sediment of the  $i$ th grain-size class, respectively. In this study, the number of grain-size classes and representative grain diameters were determined based on the grain size distribution of experiments performed (specific values noted in Section 5.1). Parameter  $C_T$  denotes the layer-averaged total

124 concentration of suspended sediment ( $C_T = \sum C_i$ ), and  $g$  is gravitational acceleration.  
 125 Parameters  $S$  and  $c_f$  are the slope gradient and the friction coefficient, respectively.  
 126 Sediment properties are described by  $R$ , the submerged specific density of sediment;  
 127  $w_i$ , the settling velocity of a sediment particle of the  $i$ th grain-size class; and  $\lambda_p$ , the  
 128 porosity of bed sediment. Parameter  $\eta_i$  is the volume per unit area of bed sediment  
 129 of the  $i$ th grain-size class and  $\eta_T$  is the sum of all  $\eta_i$  ( $\eta_T = \sum \eta_i$ ). Parameters  $L_a$ , the  
 130 active layer thickness, and  $F_i$ , the volume fraction of the  $i$ th grain-size class in active  
 131 layer, describe the structure of active layer. Parameters  $e_{si}$ ,  $e_w$  and  $r_o$  are the entrainment  
 132 rate of sediment of the  $i$ th grain-size class into suspension, the entrainment rate  
 133 of ambient water to flow, and the ratio of near-bed suspended sediment concentration  
 134 to the layer-averaged concentration of suspended sediment, respectively (Figure 1).  
 135 Closure equations of the parameters mentioned are described in Section 2.2.



**Figure 1.** Schematic diagram of processes considered in the forward model from Nakao and Naruse (2017).

## 136 2.2 Closure Equations

137 Empirical formulations from previous studies are adapted to close the govern-  
 138 ing equations. In this study, the friction coefficient  $c_f$  is assumed to be a constant  
 139 value. The particle settling velocity  $w_i$  for each grain-size class with a representative  
 140 grain diameter  $D_i$  is calculated using the relation from Dietrich (1982), which can be  
 141 expressed as follows:

$$\begin{aligned}
 w_i &= R_{fi} \sqrt{RgD_i}, & (6) \\
 R_{fi} &= \exp(-b_1 + b_2 \log(Re_{pi}) - b_3 (\log(Re_{pi}))^2 - b_4 (\log(Re_{pi}))^3 + b_5 (\log(Re_{pi}))^4), & (7) \\
 Re_{pi} &= \frac{\sqrt{RgD_i} D_i}{\nu}, & (8)
 \end{aligned}$$

142 where  $b_1$ ,  $b_2$ ,  $b_3$ ,  $b_4$  and  $b_5$  are 2.891394, 0.95296, 0.056835, 0.000245 and 0.000245,  
 143 respectively.  $e_w$  is calculated using the empirical formula from Fukushima et al. (1985)  
 144 as follows:

$$e_w = \frac{0.00153}{0.0204 + (RC_T H/U^2)}. \quad (9)$$

145 The entrainment coefficient of sediment  $e_s$  is calculated using the empirical relation  
 146 from Wright and Parker (2004):

$$e_{si} = \frac{aZ^5}{1 + (a/0.03)Z^5}, \quad (10)$$

$$Z = \alpha_1 \frac{u_*}{w_i} Re_p^{\alpha_2} S_f^{0.08}, \quad (11)$$

147 where shear velocity  $u_*$  and friction slope  $S_f$  are calculated as follows:

$$u_* = \sqrt{c_f U}, \quad (12)$$

$$S_f = \frac{c_f U^2}{RgC_T H}, \quad (13)$$

148 and the constants  $\alpha_1$  and  $\alpha_2$  are 0.586 and 1.23 respectively if  $Re_p \leq 2.36$ . If  $Re_p >$   
 149  $2.36$ , then  $\alpha_1$  and  $\alpha_2$  are 1.0 and 0.6, respectively. Kinematic viscosity of water  $\nu$  is  
 150 calculated as follows:

$$\nu = \mu/\rho, \quad (14)$$

151 where  $\rho$  and  $\mu$  denote the density of water and the dynamic viscosity of water, respec-  
 152 tively. Experimentally determined values for  $\mu$  at 20.0 °C (Rumble, 2018) were used  
 153 in the calculation of  $\nu$  in this study.

### 154 2.3 Implementation of Forward Model

155 In this study, the constrained interpolation profile (CIP) method (Yabe et al.,  
 156 2001) was used for the integration of the partial differential equations 1, 2, and 3. To  
 157 stabilize the calculation, artificial viscosity was applied with the scheme of Jameson et  
 158 al. (1981), in which the parameter  $\kappa$  was set to 0.25. The two-step Adams-Bashforth  
 159 method was used to solve ordinary differential equations 4 and 5. Interval of spatial  
 160 grids  $\Delta x$  was set to 0.05 m and time step  $\Delta t$  was set to 0.01 s. Initial values of  $S$  were  
 161 10% for all grids.

162 Dirichlet boundary condition was used for the upstream boundary, in which all  
 163 flow parameters at the upper boundary of calculation domain, including the initial  
 164 flow height  $H_0$ , the initial flow velocity  $U_0$ , the initial total volumetric concentration  
 165 of sediment  $C_{T,0}$ , and the initial volumetric concentration of each grain-size class  $C_{i,0}$ ,  
 166 were set to be constant. Parameter  $F_{i,0}$ , the initial volume fraction of the  $i$ th grain-  
 167 size class in active layer, was set to 0.25 for all grain-size classes. The downstream  
 168 boundary was the Neumann boundary condition where all parameters were set to the  
 169 same values as those of the grid adjacent to the lower boundary toward the upstream  
 170 direction. Other than the upstream boundary, all flow parameters, except the flow  
 171 height  $H$ , were initialized to zero.  $H$  was initialized to 0.0001.

172 Properties of fluids and sediment particles were set as follows. Density of fluid  $\rho$   
 173 were set to be 1000.0 kg/m<sup>3</sup>, and submerged specific density of sediment  $R = (\rho_s - \rho)/\rho$   
 174 was set differently according to the types of particles used in experiments ( $\rho_s$  is the  
 175 density of sediment particles), which are stated in Section 4.1. The porosity of bed  
 176 sediment  $\lambda_p$  was assumed to be 0.4. In this study, both the friction coefficient  $c_f$   
 177 and ratio of near-bed concentration to layer-averaged values  $r_o$  were assumed to be  
 178 constant, which were set to 0.002 and 1.5 (Kostic & Parker, 2006). Also, the thickness  
 179 of the active layer  $L_a$  was set to be constant (0.003 m). The gravitational acceleration  
 180  $g$  was 9.81 m/s<sup>2</sup>.

### 181 3 Inverse Analysis by Deep Learning Neural Network

182 In this method, initial flow conditions of turbidity currents are reconstructed  
 183 from its turbidite deposits. DNN first explores the functional relationship between the

184 initial flow conditions of turbidity currents and the resulting turbidite deposits in a  
 185 process called training. After the training process, DNN is applied to new turbidite  
 186 datasets for inverse analysis. In preparation for training, artificial training datasets  
 187 are produced using the forward model. During the training process, training datasets  
 188 are fed into DNN. DNN examines the datasets and adjusts to make good estimation  
 189 of the initial flow conditions from the deposit profile. After training, DNN, which is  
 190 now capable of predicting the initial flow conditions of new turbidites based on the  
 191 functional relationship it discovered, is tested with independent artificial datasets that  
 192 are also produced from the forward model and with flume experiment data.

### 193 3.1 Production of Training Data

194 A training dataset is a combination of randomly generated initial flow conditions  
 195 at the upstream boundary of the flow and a matching deposit profile calculated by  
 196 the forward model. A program in Python was written to generate sets of initial flow  
 197 conditions. Each set of flow conditions generated consists of an initial flow velocity  
 198  $U_0$ , an initial flow height  $H_0$ , a flow duration  $T_d$ , and the initial concentrations of each  
 199 grain-size class  $C_{i,0}$ . Other variables, such as slope, are set to constant values.

200 The forward model calculates the deposit profile of a turbidite using the randomly  
 201 generated initial flow conditions. The deposit profile is calculated as volume per unit  
 202 area for each grain-size class at a total of 80 locations within a 4 m range. Each data  
 203 point is 0.05 m away from its neighboring points. The ranges of randomly generated  
 204 initial flow conditions are shown in Table 1. In this study, 11800 training datasets were  
 205 used for training and 200 datasets were used for verification of DNN. The test artificial  
 206 datasets for verification were produced independently from the training datasets.

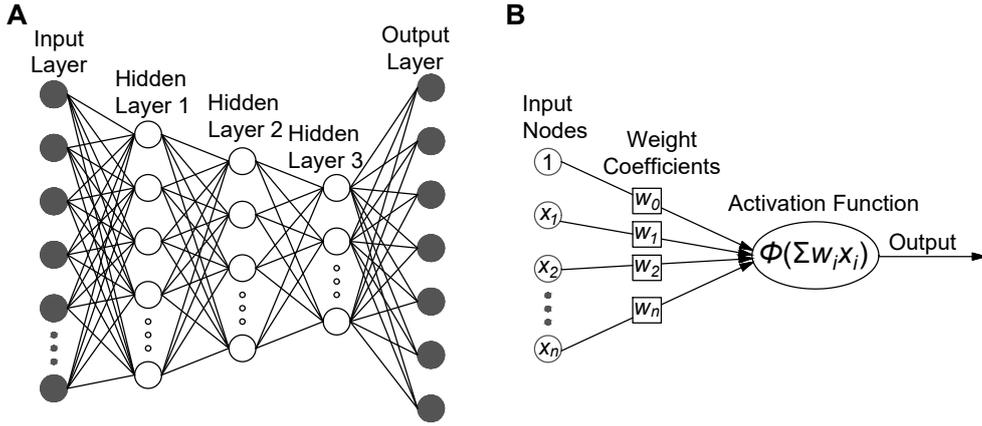
**Table 1.** Range of initial flow conditions generated for the production of training datasets.

Parameter	Minimum	Maximum
$H_0$ (m)	0.005	0.3
$U_0$ (m/s)	0.001	0.3
$C_{i,0}$	0.0001	0.02
$T_d$ (s)	1080	120

### 207 3.2 Structure of Deep Learning Neural Network

208 The type of neural network used in this study is the fully connected neural  
 209 network, which consists of an input layer, several hidden layers, and an output layer.  
 210 Each layer consists of a number of nodes. Each node connects with every node in the  
 211 adjacent layers (Figure 2A). In the input layer, the nodes hold values of the deposit  
 212 profile, i.e. the volume-per-unit-area for all grain-size classes at spatial grids. In the  
 213 output layer, the nodes hold estimates of the parameters we seek to reconstruct, which  
 214 in this case are the initial flow conditions  $U_0$ ,  $H_0$ ,  $C_{i,0}$  and the flow duration  $T_d$ . The  
 215 activation function used in this study is ReLU, which is one of the most commonly  
 216 used activation functions for DNN and is proven to perform calculations at a higher  
 217 speed than other activation functions (Krizhevsky et al., 2012).

218 Before training, the weight coefficients are set to random values. As the training  
 219 process begins, the values of deposit profile from the training datasets are fed into the  
 220 input layer. These values propagate through the hidden layers of DNN and estimates of  
 221 initial flow conditions are outputted at the output layer. At this point of training, DNN



**Figure 2.** Schematic diagrams of DNN. A. Overall structure of DNN. B. Concept of weight coefficient and activation function.

222 is yet to adapt its internal variables to the functional relationship between turbidite  
 223 deposits and initial flow conditions, thus the initial estimates are expected to be largely  
 224 off from the actual values. To explore this functional relationship, a loss function is  
 225 used to evaluate the accuracy of the estimated values. The loss function used in  
 226 this case is the mean squared error function, which is considered as one of the best  
 227 functions for regression (Wang & Bovik, 2009). The gradient of the loss function  
 228 is calculated and fed back to the hidden layers of DNN through a process called  
 229 backpropagation (Nielsen, 2015; Schmidhuber, 2015), where the internal values of  
 230 DNN are optimized toward minimizing the difference between the estimated and actual  
 231 values. This process is repeated for every epoch of calculation. An epoch is a cycle  
 232 of calculation in DNN that involves one forward pass and one backpropagation of all  
 233 training data.

234 The optimization algorithm used in this study is stochastic gradient descent  
 235 (SGD), which drastically reduces the amount of calculation involved in training with-  
 236 out compromising accuracy relative to previous algorithms of gradient descent (Bottou,  
 237 2010). In this study, Nesterov momentum is employed with SGD (Ruder, 2016). Be-  
 238 cause of difference in the order of the range of initial flow conditions, normalization  
 239 needs to be applied to training datasets before they are inputted to DNN. In this case,  
 240 all values are normalized to be between 0 and 1 for DNN to consider all parameters at  
 241 equal weights. The hyperparameters including the number of layers, number of nodes  
 242 at each layer, dropout rate, validation split, learning rate, batch size, epoch, and mo-  
 243 mentum were adjusted manually. Various combinations were attempted and the best  
 244 combination of hyperparameters was chosen based on the performance of DNN, which  
 245 is judged on the basis of the final validation loss.

246 In this study, DNN was programmed in Python using the package Keras 2.2.4.  
 247 The package Tensorflow 1.14.0 (Abadi et al., 2015) was used for backend calculation.  
 248 Calculations were conducted using GPU NVIDIA GeForce GTX 1080 Ti.

### 249 3.3 Evaluation of Trained DNN Model

250 During the verification of DNN with test artificial datasets (Section 5.1), recon-  
 251 struction result of each parameter was evaluated using bias ( $B$ ) and sample standard  
 252 deviation ( $s$ ) of residuals. The calculations were conducted using the following equa-

253 tions:

$$B = \frac{\sum x_i}{n}, \quad (15)$$

$$s = \sqrt{\frac{\sum (x_i - B)^2}{n - 1}}, \quad (16)$$

254 where  $n$  is the number of test datasets, and  $x_i$  denotes the residual of the specific recon-  
 255 structed parameter for the  $i$ th test dataset. The calculated  $s$  and  $B$  for reconstructed  
 256 parameters are listed in Table 4 and Table 5. The value of  $s$  for each reconstructed  
 257 parameter was compared with a representative value  $C_v^*$ , which is the mid-value over  
 258 the range in which the specific parameter was generated (Table 1). The confidence  
 259 interval of  $B$  was determined using bootstrap resampling method (Davison & Hink-  
 260 ley, 1997). Resampling of  $B$  was conducted for 10000 times and the 95% confidence  
 261 interval (CI) of  $B$  was determined.

262 During the verification of DNN using flume experiment data (Section 5.2), flow  
 263 parameters at the upstream end of simulation were reconstructed from the measured  
 264 properties of the deposit profile. The upstream end of simulation was set at 1.0 m  
 265 from the upstream boundary of the flume. The reconstructed parameters were in-  
 266 putted into the forward model so that the flow parameters downstream and the time  
 267 evolution of deposit profile were calculated. The calculated flow parameters down-  
 268 stream were compared with the flow conditions measured during experiments. The  
 269 deposit profile calculated from the reconstructed flow parameters was also compared  
 270 with the measured deposit profile that was used for inversion.

## 271 4 Flume Experiments

### 272 4.1 Experiment Settings

273 In this study, flume experiments were conducted using a flume located in the  
 274 basement of Building No.1 of the Faculty of Science at Yoshida Campus, Kyoto Uni-  
 275 versity. The flume was made of acrylic panels and is 4 m in length, 0.12 m in width  
 276 and 0.5 m in depth. It was submerged in a tank made of glass panels and steel sup-  
 277 porting frame. The tank was 5.5 m in length, 2.5 m in width and 1.8 m in depth.  
 278 Slope of channel floor changes at a middle point of the flume (1.0 m from the inlet),  
 279 where the upstream slope  $a$  was set to 26.8% and the downstream slope  $b$  was set  
 280 to 10% (Figure 3). Sediment was mixed with water in two mixing tanks before the  
 281 experiments.

282 During the experiments, the mixture of sediment and water was first pumped to  
 283 the constant head tank and then released into the flume from a valve at the base of  
 284 the constant head tank. Flow discharge was regulated through adjusting the degree  
 285 of opening of the valve. The amount of mixture in the constant head tank was kept  
 286 at a constant level to maintain a stable flow discharge. The damping tank at the  
 287 downstream end of the flume prevented flow from reflecting back toward the upstream  
 288 side. A pipe of freshwater supply was placed at the top of the damping tank, and  
 289 a draining pipe was placed at the bottom of the damping tank. The combination of  
 290 these two pipes kept the level of water in the tank constant and prevented reflection  
 291 of flow.

292 Four experiments were conducted in this study. Two experiments were performed  
 293 using two types of plastic particle (experiments PP1, PP2 (Section 4.3)). One of the  
 294 types of plastic particle used, polyvinyl chloride, had a density of 1.45 g/cm<sup>3</sup> and  
 295 an average grain diameter of 0.120 mm. The other type of plastic particle used,  
 296 melamine, had a density of 1.45 g/cm<sup>3</sup> and an average grain diameter of 0.220 mm.  
 297 Two experiments were performed using a mixture of siliciclastic sand and polyvinyl  
 298 chloride plastic particle (experiments BS1, BS2 (Section 4.3)). Siliciclastic sand had

299 a density of  $2.65 \text{ g/cm}^3$  and an average grain diameter of  $0.250 \text{ mm}$ . The submerged  
 300 specific gravity  $R$  of the plastic particles was  $0.45$ , whereas that of the siliciclastic sand  
 301 was  $1.65$ .

## 302 4.2 Measurements and Data Analysis

303 Before each experiment, tank water temperature was measured using a glass  
 304 alcohol thermometer. A beaker of the mixture in tank was sampled to measure the  
 305 initial concentration in tank. Flow velocity during experiment was measured with an  
 306 acoustic Doppler velocity profiler (ADVP; Nortek Vectrino Profiler). The functional  
 307 range of the ADVP used was  $4.0 - 7.0 \text{ cm}$  from the probe. To obtain the vertical  
 308 velocity profile of the flow, an actuator was used to adjust the position of the ADVP  
 309 during experiments.

310 A siphon with 10 plastic tubes was used to measure the suspended sediment  
 311 concentration of flow. The tubes were aligned vertically at  $1.0 \text{ cm}$  intervals and were  
 312 positioned such that samples were collected at  $0.0$  to  $9.0 \text{ cm}$  above bed. Aluminum  
 313 tubes with an outer diameter of  $8.0 \text{ mm}$  and an inner diameter of  $5.0 \text{ mm}$  were attached  
 314 to the outlets of the plastic tubes to stabilize its position. Sampling using siphon was  
 315 conducted when the flow reached a quasi-equilibrium state. The state of flow was  
 316 determined by observation of the development of flow. Two single-lens reflex cameras  
 317 were used to record the experiments. Flow height was determined based on the video  
 318 recorded.

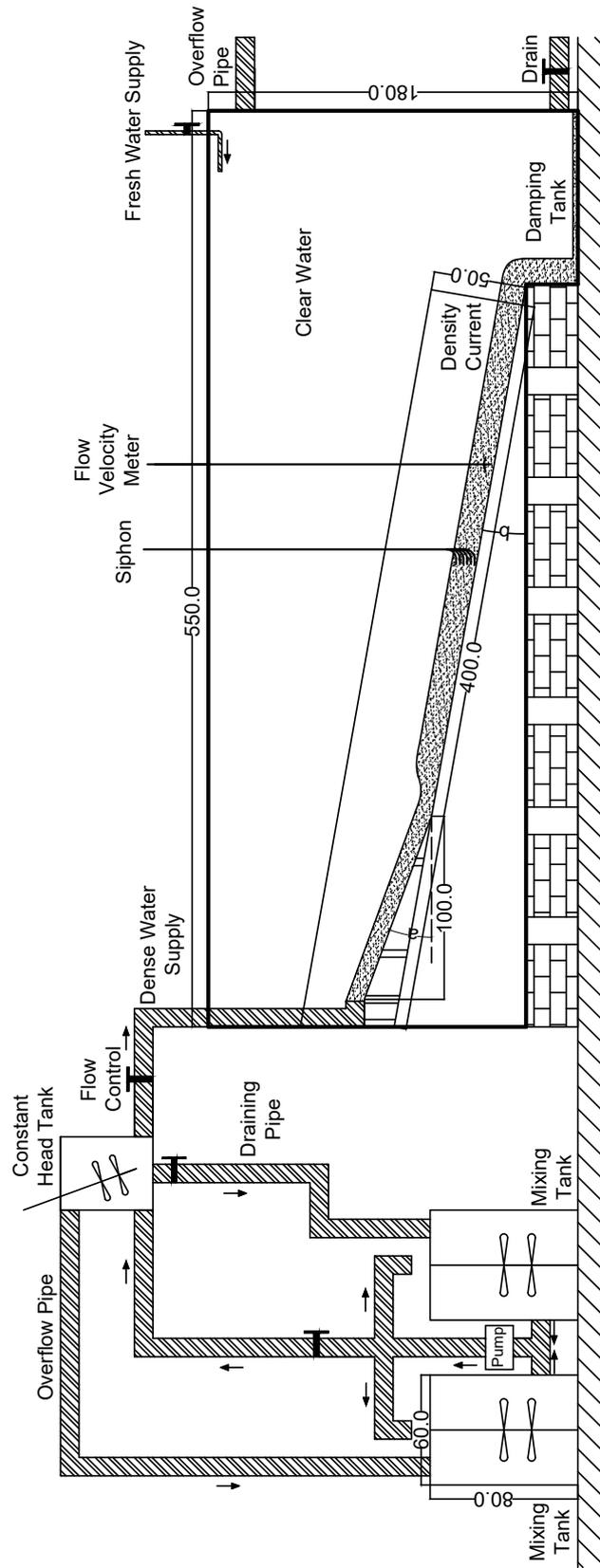
319 After the experiments, the flume was left untouched for 1 to 3 days for the  
 320 suspended sediment to settle. Afterward, photos were taken from a lateral view per-  
 321 pendicular to the flume. The lateral view of the deposited sediment could be seen  
 322 together with a ruler in the photos. The height of deposit was determined from the  
 323 photos. Water was then drained slowly from the tank with a bath pump at a rate of  
 324  $0.0002333 \text{ m}^3/\text{s}$ . After water was drained, deposited sediment was sampled at  $20 \text{ cm}$   
 325 intervals starting from the upstream boundary of the flume.

326 Samples from the siphon and tank were first weighed right after they were col-  
 327 lected. Following the weighing, samples from the deposited sediment and from siphon  
 328 and tank were dried. Samples from the siphon and tank were weighed again after dry-  
 329 ing, and the measurements were recorded for calculation of the sediment concentration  
 330 in flow and in tank. Grain size distribution analysis was conducted in a settling tube  
 331 for all sediment samples. The settling tube used was  $1.8 \text{ m}$  long. Calculation of grain  
 332 size distribution was performed using STube (Naruse, 2005). Particle settling velocity  
 333 was calculated using Gibbs (1974).

334 The measured grain-size distribution of sediment was discretized to four grain-  
 335 size classes. For experiments using plastic particle only, representative grain diameter  
 336 of grain-size classes 1, 2, 3 and 4 were set to be  $210 \mu\text{m}$ ,  $149 \mu\text{m}$ ,  $105 \mu\text{m}$  and  $74.3$   
 337  $\mu\text{m}$ , respectively. For experiments using a mixture of siliciclastic sand and plastic  
 338 particle, representative grain diameter of grain-size classes 1, 2, 3 and 4 were set to be  
 339  $297 \mu\text{m}$ ,  $210 \mu\text{m}$ ,  $149 \mu\text{m}$  and  $105 \mu\text{m}$ , respectively. Grain size distributions of sand  
 340 and plastic particle partially overlaps, thus concentrations of grain-size classes 2 and  
 341 3 were each represented by two separate parameters, i.e.  $C_{2,BS,0}$  and  $C_{2,PP,0}$ ,  $C_{3,BS,0}$   
 342 and  $C_{3,PP,0}$ , respectively. BS and PP represents the siliciclastic sand and the plastic  
 343 particle portions, respectively, of the grain-size classes 2 and 3.

344 In steady flow conditions, the relationship between the layer-averaged flow ve-  
 345 locity  $U$ , the layer-averaged sediment volumetric concentration  $C$ , and the flow height  
 346  $H$  is defined as follows (Garcia & Parker, 1993):

$$UCH = \int_a^\infty u_z c_z dz, \quad (17)$$



**Figure 3.** Schematic diagram of flume used for experiment. (All values have unit cm.)

347 where  $u_z$  and  $c_z$  are the flow velocity and sediment volumetric concentration, re-  
 348 spectively, at elevation  $z$  above bed. The relationship between the layer-averaged  
 349 flow velocity  $U$  and the velocity maximum  $U_m$  is defined with the following equation  
 350 (Altinakar et al., 1996):

$$\frac{U_m}{U} = 1.3. \quad (18)$$

351 The layer-averaged flow velocity was calculated from the velocity profile measured  
 352 by the ADVP using the relationship described by equation 18. The sediment volu-  
 353 metric concentration was calculated from siphon measurements using the relationship  
 354 described by equation 17.

### 355 4.3 Experimental Conditions

356 The experimental conditions for the four runs conducted are outlined in Table 2  
 357 and Table 3.  $C_{TT}$  is the total concentration of sediment in the mixing tank.  $C_{1T}$ ,  
 358  $C_{2T}$ ,  $C_{3T}$ ,  $C_{4T}$  are the concentrations of grain-size classes 1, 2, 3, and 4, respectively.  
 359  $C_{2T,BS}$  and  $C_{2T,PP}$ ,  $C_{3T,BS}$  and  $C_{3T,PP}$  in Table 3 are concentrations of the siliciclastic  
 360 sand portion and plastic particle portion, respectively, of grain-size classes 2 and 3.  
 361 Parameter  $x_C$  is the position of siphon downstream, whereas  $x_U$  is the position of flow  
 362 velocity meter downstream.  $x_H$  is the position where flow height was measured from  
 363 video taken. Temperature is the measured temperature of clear water in tank before  
 364 experiment.

**Table 2.** Conditions and settings of experiments conducted with two types of plastic particle.

	PP1	PP2
$C_{TT}$	0.0191	0.0276
$C_{1T}$	0.0102	0.0160
$C_{2T}$	0.00713	0.00820
$C_{3T}$	0.00146	0.00254
$C_{4T}$	0.000366	0.000817
$x_C$ (m)	1.08	2.10
$x_U$ (m)	1.46	2.48
$x_H$ (m)	1.10	1.10
Temperature ( $^{\circ}\text{C}$ )	22.5	17.0

## 365 5 Results

366 Inverse analysis was applied to deposits within a 2.6 m range downstream starting  
 367 at the end of slope  $a$  (1.0 m from the upstream boundary of the flume). Because  
 368 of the limited size of the flume, slope  $a$  was set with a steep angle (26.8%) in all  
 369 experiments. Considering the effect of instabilities from the outlet pipe and the overly  
 370 steep slope, the region with slope  $a$  was excluded from numerical simulations and  
 371 inverse analysis. As a result, for the production of artificial datasets, the upstream  
 372 boundary of simulation was set at the end of slope  $a$ . The numerical simulation was  
 373 conducted for 4.0 m downstream, but only the deposits from the upstream boundary  
 374 of simulation to 2.6 m downstream were used for training and verification. This range  
 375 corresponds to 1.0 m to 3.6 m from the upstream boundary of the flume, where samples  
 376 were collected for experiment deposits. The range of simulation was set longer than  
 377 the range used for training and verification to minimize the influence of the boundary  
 378 conditions at the downstream end of simulation.

**Table 3.** Conditions and settings of experiments conducted with a mixture of siliciclastic sand and plastic particle.

	BS1	BS2
$C_{TT}$	0.0587	0.0140
$C_{1T}$	0.0246	0.000122
$C_{2T,BS}$	0.0101	0.00385
$C_{2T,PP}$	0.00958	0.00251
$C_{3T,BS}$	0.00811	0.00503
$C_{3T,PP}$	0.00424	0.00164
$C_{4T}$	0.00202	0.000846
$x_C$ (m)	2.33	1.85
$x_U$ (m)	2.71	2.23
$x_H$ (m)	1.40	1.40
Temperature ( $^{\circ}\text{C}$ )	9.5	12.2

379 With regard to the values of hyperparameters used during the training process,  
380 the dropout rate, validation split and momentum for DNN were set to 0.5, 0.2, and  
381 0.9, respectively. Learning rate was set to 0.01, batch size to 32, number of layers to 5,  
382 number of nodes each layer to 2000, and epoch to 15000. With this setting the eventual  
383 validation loss was 0.0031 when training with plastic particle only datasets and 0.0020  
384 when training with siliciclastic sand plastic particle mixture datasets. Figures 4A and  
385 5A show that overlearning did not occur, as no deviation was observed between the  
386 resulting values of loss function for the training and validation datasets.

### 387 5.1 Verification of Inverse Model with Test Artificial Datasets

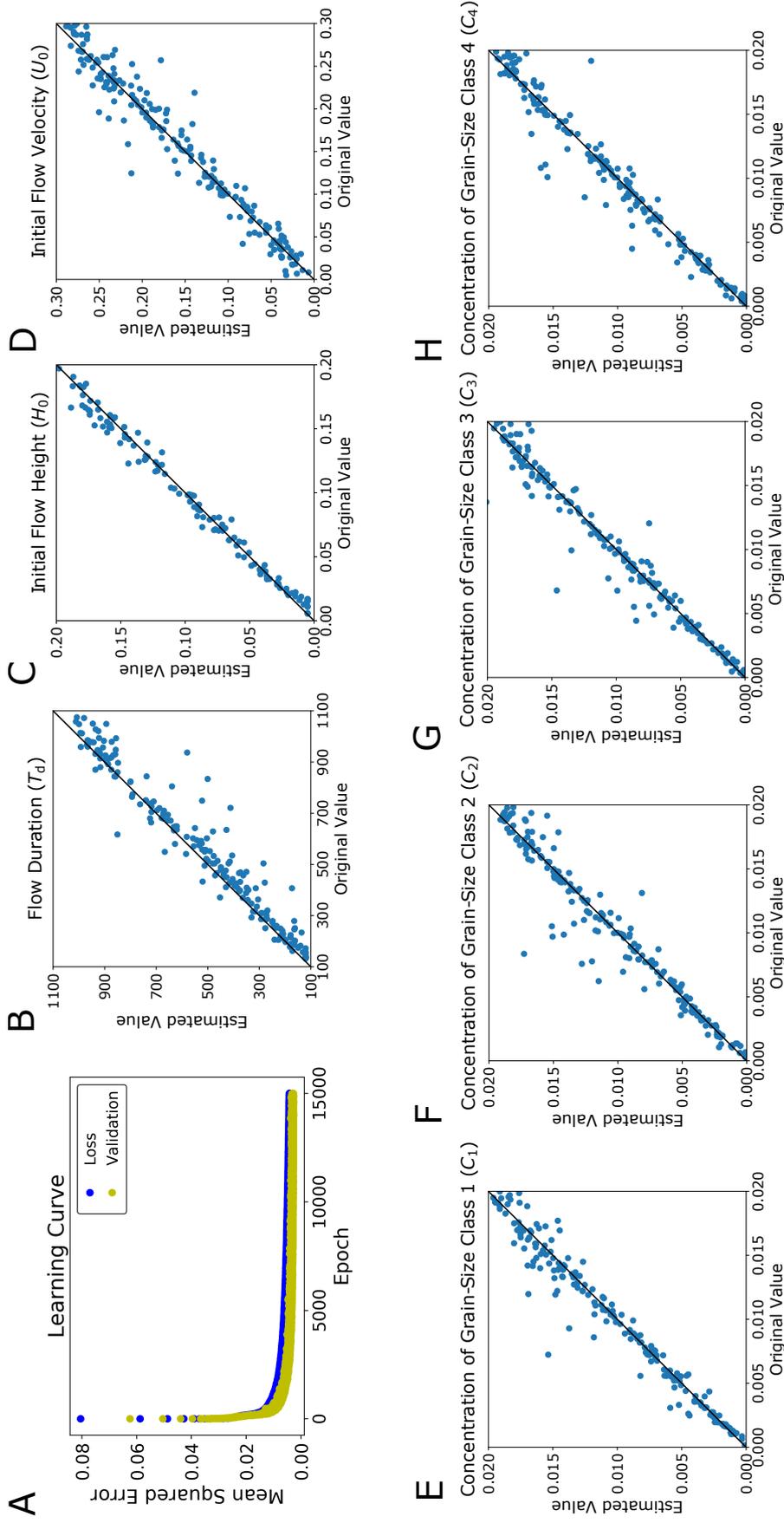
388 This section presents the results of verification with artificial test datasets. Re-  
389 sults of parameter reconstructions by DNN are shown in Figures 4 and 5. Parameters  
390 reconstructed include flow duration  $T_d$  and flow conditions at the upstream end (flow  
391 velocity  $U_0$ , flow depth  $H_0$ , sediment concentrations  $C_{i,0}$ ). Verification was performed  
392 separately for artificial datasets of experiments using plastic particle only and artificial  
393 datasets of experiments using a mixture of siliciclastic sand and plastic particle. The  
394 results are described in Sections 5.1.1 and 5.1.2.

#### 395 5.1.1 Verification with Test Artificial Datasets for Experiments Using 396 Plastic Particle Only

397 As an overall trend, the reconstructed values mostly coincided with the original  
398 values, with a few values scattering further away from the line of perfect match (the  
399 diagonal solid line) (Figure 4B, C, D, E, F, G, H). A greater degree of scattering,  
400 however, was observed for  $U_0$  in comparison to other parameters (Figure 4D). Flow  
401 duration  $T_d$  exhibited a tendency of underestimation (Figure 4B). The ranges of mis-  
402 fit ( $2s$ ) were relatively small for all parameters, which had  $2s/C_v^*$  values under 13%  
403 (Table 4). Zero lay within the 95% confidence interval (CI) of  $B$  for all parameters,  
404 except for  $T_d$  and  $C_{1,0}$ , where the CI range lay below zero for  $T_d$  and above zero for  
405  $C_{1,0}$ .

**Table 4.** Sample standard deviation and bias of the inversion result of artificial datasets for experiments using plastic particle only.

Parameters	$s$	$C_v^*$	$2s/C_v^*$	$B$	CI of $B$
$U_0$ (m/s)	0.0182	0.1505	0.121	-0.000673	(-0.00320, 0.00187)
$H_0$ (m)	0.0151	0.1525	0.0990	0.000232	(-0.00188, 0.00233)
$T_d$ (s)	64.8	600	0.108	-36.5	(-45.6, -27.8)
$C_{1,0}$	0.00114	0.01005	0.113	0.000233	(0.0000828, 0.000396)
$C_{2,0}$	0.00135	0.01005	0.134	0.000136	(-0.0000447, 0.000332)
$C_{3,0}$	0.00115	0.01005	0.114	0.000148	(-0.00000139, 0.000314)
$C_{4,0}$	0.00117	0.01005	0.116	0.000115	(-0.0000419, 0.000274)



**Figure 4.** Result of verification with independent artificial datasets for plastic particle only experiments. The black diagonal line in each graph is where values on the  $x$ -axis (the true values) equal to the values on the  $y$ -axis (the estimated values). If a point lies on this line, the reconstructed value matches the true value perfectly. A. Learning curve. B. Estimates of  $T_d$ . C. Estimates of  $H_0$ . D. Estimates of  $U_0$ . E. Estimates of  $C_{1,0}$ . F. Estimates of  $C_{2,0}$ . G. Estimates of  $C_{3,0}$ . H. Estimates of  $C_{4,0}$ .

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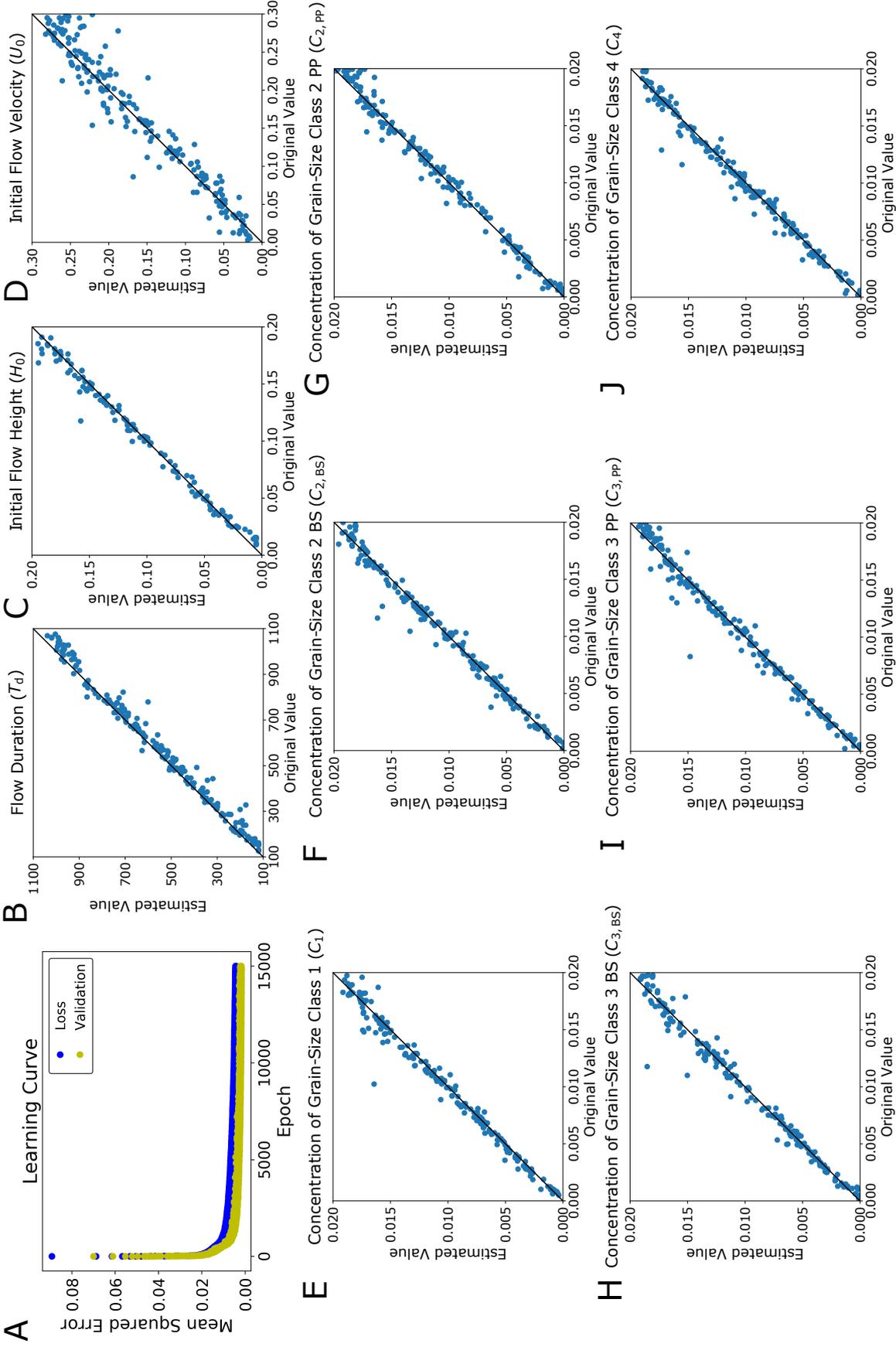
### 5.1.2 Verification with Test Artificial Datasets for Experiments Using Siliciclastic Sand and Plastic Particle

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As an overall trend, good correlations were observed for the reconstructed values and the original values of flow parameters. The reconstructed values mostly exhibited perfect matches with the original values, with a few values scattering from the line of perfect match (the diagonal black line) (Figure 5B-J). Similar to the test datasets described in Section 5.1.1,  $U_0$  exhibited a greater degree of scattering in comparison to the other parameters. A tendency of overestimation was also observed for  $U_0$  (Figure 5D).  $T_d$  showed a tendency of underestimation (Figure 5B). The range of misfit ( $2s$ ) was good for  $T_d$ ,  $H_0$ ,  $C_{1,0}$ ,  $C_{2,BS,0}$ ,  $C_{2,PP,0}$ ,  $C_{3,BS,0}$ ,  $C_{3,PP,0}$  and  $C_{4,0}$ , which had  $2s/C_v^*$  values under 8%.  $U_0$  was more scattered, having  $2s/C_v^*$  value of 14.2% (Table 5). Zero was included in the 95% confidence interval of  $B$  except for  $U_0$ ,  $H_0$ ,  $T_d$ ,  $C_{1,0}$ ,  $C_{2,BS,0}$  and  $C_{4,0}$ , where the CI ranges lay below zero for  $T_d$  and  $U_0$ , and above zero for  $H_0$ ,  $C_{1,0}$ ,  $C_{2,BS,0}$  and  $C_{4,0}$ .

**Table 5.** Sample standard deviation and bias of the inversion result of artificial datasets for experiments using a mixture of siliciclastic sand and plastic particle.

Parameters	$s$	$C_v^*$	$2s/C_v^*$	$B$	CI of $B$
$U_0$ (m/s)	0.0214	0.1505	0.142	-0.00413	(-0.00711, -0.00115)
$H_0$ (m)	0.0114	0.1525	0.0748	0.00155	(0.00000628, 0.00314)
$T_d$ (s)	30.9	600	0.0515	-23.7	(-28.1, -19.6)
$C_{1,0}$	0.000726	0.01005	0.0722	0.000215	(0.000120, 0.000320)
$C_{2,BS,0}$	0.000692	0.01005	0.0689	0.000127	(0.0000347, 0.000230)
$C_{2,PP,0}$	0.000715	0.01005	0.0711	-0.0000752	(-0.000176, 0.0000243)
$C_{3,BS,0}$	0.000818	0.01005	0.0814	0.000105	(-0.00000385, 0.000221)
$C_{3,PP,0}$	0.000816	0.01005	0.0812	0.0000754	(-0.0000336, 0.000193)
$C_{4,0}$	0.000707	0.01005	0.0703	0.000125	(0.0000303, 0.000228)



**Figure 5.** Result of verification with independent artificial datasets for siliclastic sand plastic particle mixture experiments. The black diagonal line in each graph is where values on the  $x$ -axis (the true values) equal to the values on the  $y$ -axis (the estimated values). If a point lies on this line, the reconstructed value matches the true value perfectly. A. Learning curve. B. Estimates of  $T_d$ . C. Estimates of  $H_0$ . D. Estimates of  $U_0$ . E. Estimates of  $C_{1,0}$ . F. Estimates of  $C_{2,BS,0}$ . G. Estimates of  $C_{2,PP,0}$ . H. Estimates of  $C_{3,BS,0}$ . I. Estimates of  $C_{3,PP,0}$ . J. Estimates of  $C_{4,0}$ .

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## 5.2 Inverse Analysis of Flume Experiment Data

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In this section, the deposit profiles and grain size distributions calculated are compared with the actual deposit profiles sampled from experiments (Figures 6, 8). The results of flow conditions reconstructed, including flow velocity  $U_{x_U}$ , flow height  $H_{x_H}$ , sediment concentrations  $C_{x_C}$  and flow duration  $T_d$ , are compared with the values measured during the experiments (Tables 6 and 7).  $H_{x_H}$  is the flow height  $H$  at position  $x_H$  (see the positions in Table 2) downstream when the flow reached a quasi-equilibrium state.  $U_{x_U}$  is the velocity  $U$  at position  $x_C$  downstream when the flow reached a quasi-equilibrium state.  $C_{x_C}$  is the  $C$  at position  $x_C$  downstream when the flow reached a quasi-equilibrium state. Inversion results of the plastic particle only experiments (PP1, PP2) and of the experiments using a mixture of siliciclastic sand and plastic particle (BS1, BS2) are described in Sections 5.2.1 and 5.2.2.

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### 5.2.1 Experiments using Plastic Particle Only (PP1, PP2)

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Depositional profiles in these experiments exhibited a thinning and fining downstream trend with concave-upward geometry. For both runs, the reconstructed deposit profiles of the total deposition closely matched with the sampled data (Figures 6A(1), B(1)). Grain size distributions at 1.4 m and 1.8 m downstream also demonstrated good agreement for the reconstructed and measured values for both experiments (Figures 6A(2),(3), B(2),(3)), with PP2 having a slightly better match than that of PP1.

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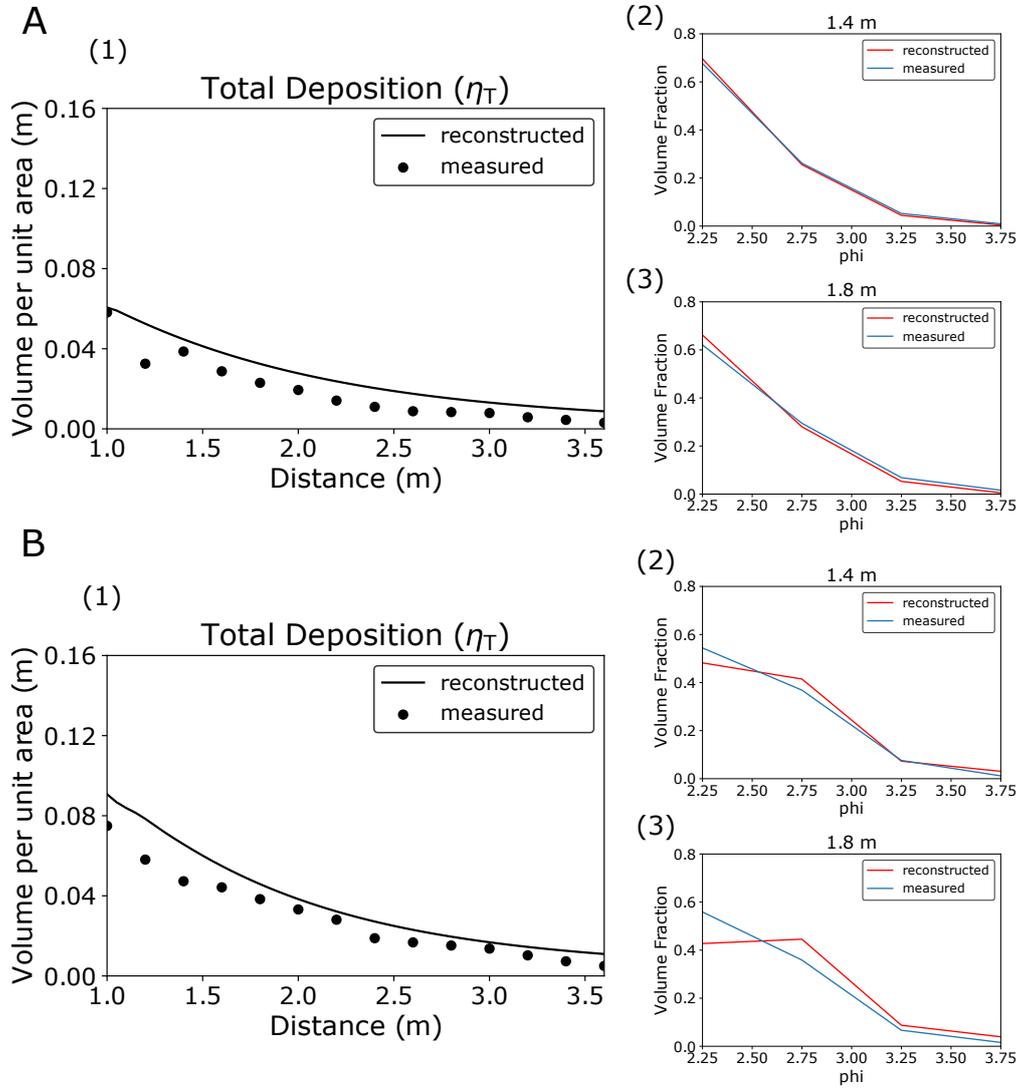
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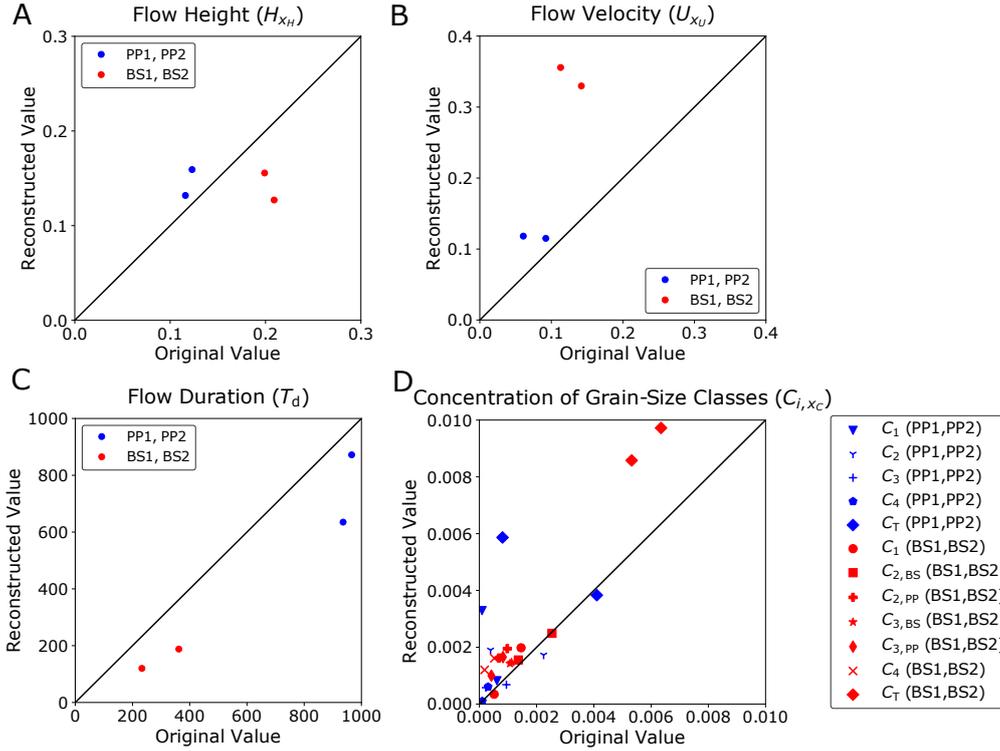
For flow height  $H_{x_H}$ , the measured and reconstructed values were in good agreement (Figure 7A). The reconstructed  $H_{x_H}$  were 0.132 m and 0.159 m for PP1 and PP2, respectively (Table 6). These values matched well with the measured values for both PP1 (0.116 m) and PP2 (0.123 m). The reconstructed  $U_{x_U}$  were 0.118 m/s (PP1) and 0.115 m/s (PP2), of which that of PP2 was close to its measured value (0.0923 m/s), whereas that of PP1 had a larger difference from its measured value (0.0608 m/s). The reconstructed values of flow duration  $T_d$  were 635 s (PP1) and 872 s (PP2), which were lower than the measured values for both PP1 (936 s) and PP2 (966 s) (Figure 7C). The reconstructed values of total concentration  $C_{T,x_C}$  were 0.00587 (PP1) and 0.00383 (PP2), of which that of PP1 did not match well with the measured value of  $C_{T,x_C}$  (0.000808), whereas that of PP2 was close to its corresponding measured value (0.00410). The reconstructed values of each grain-size class were good overall, but large deviations were observed for several reconstructions of concentrations that had very low measured values ( $< 0.1\%$ ) (Figure 7D)

**Table 6.** Flow conditions measured and reconstructed for experiments PP1 and PP2. (R.: reconstructed, M.: Measured)

Parameters	PP1 R.	PP1 M.	PP2 R.	PP2 M.
$C_{T,x_C}$	0.00587	0.000808	0.00383	0.00410
$C_{1,x_C}$	0.00329	0.0000911	0.000822	0.000612
$C_{2,x_C}$	0.00190	0.000389	0.00173	0.00224
$C_{3,x_C}$	0.000576	0.000228	0.000681	0.000944
$C_{4,x_C}$	0.000100	0.0000999	0.000607	0.000303
$H_{x_H}$ (m)	0.132	0.116	0.159	0.123
$U_{x_U}$ (m/s)	0.118	0.0608	0.115	0.0923
$T_d$ (s)	635	936	872	966



**Figure 6.** Reconstructed deposit profiles and the sampled deposit data of experiments PP1 and PP2. A. (1) Reconstructed and sampled  $\eta_T$  of PP1. (2) Grain size distribution at 1.4 m downstream. (3) Grain size distribution at 1.8 m downstream. B. (1) Reconstructed and sampled  $\eta_T$  of PP2. (2) Grain size distribution at 1.4 m downstream. (3) Grain size distribution at 1.8 m downstream.



**Figure 7.** Reconstructed vs measured flow conditions for experiments PP1, PP2, BS1 and BS2. A. Plot for  $H_{xH}$ . B. Plot for  $U_{xU}$ . C. Plot for  $T_d$ . D. Plot for  $C_{i,x_C}$ .

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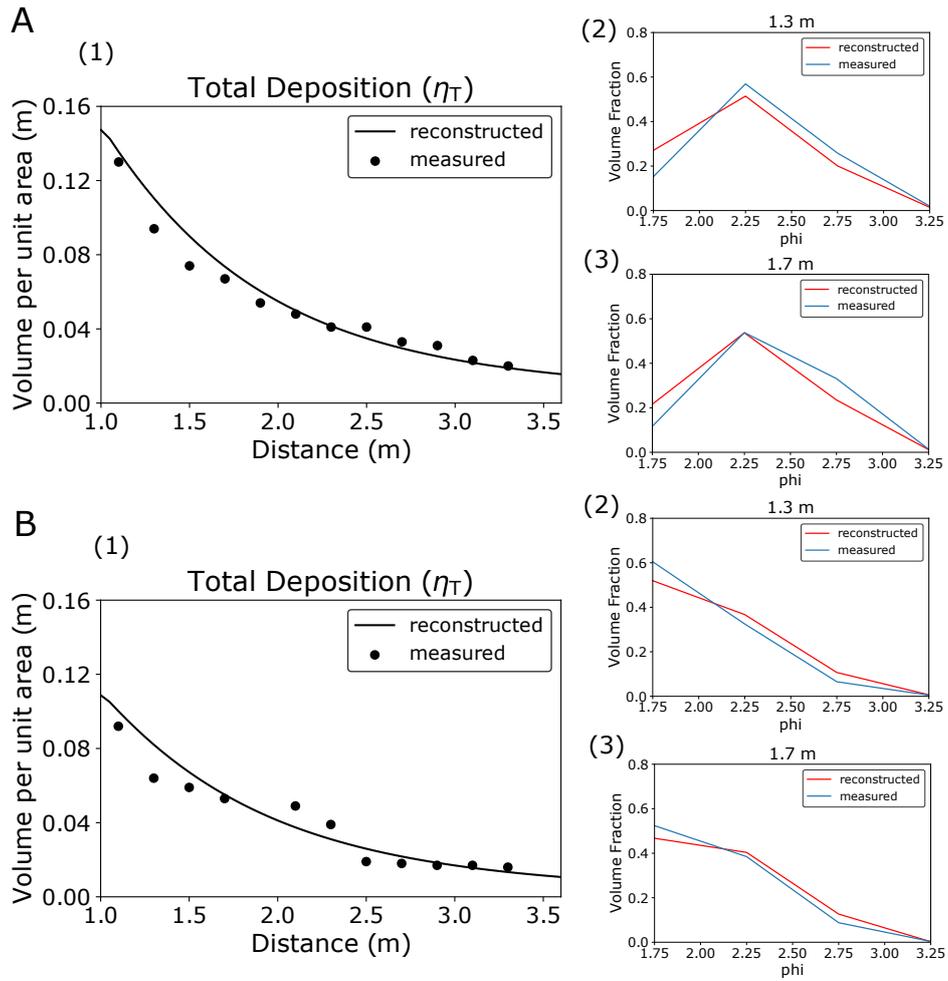
### 5.2.2 Experiments Using a Mixture of Siliciclastic Sand and Plastic Particle (BS1, BS2)

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Similar to the results for the experiments using plastic particle only, all depositional profiles in these experiments exhibited a thinning and fining downstream trend. For both experiments, the reconstructed deposit profiles of the total deposition matched closely with the sampled data (Figures 8A(1), B(1)). Grain size distributions at 1.3 m and 1.7 m downstream also demonstrated good agreement for the reconstructed and measured values for both experiments (Figures 6A(2),(3), B(2),(3)).

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The reconstructed values of flow height  $H_{xH}$  were 0.127 m and 0.156 m for BS1 and BS2, respectively (Table 7). Of these values, that for BS2 was relatively close to the measured value (0.199 m), whereas that for BS1 differs from its measured value (0.209 m). The reconstructed flow velocities  $U_{xU}$  were 0.330 m/s (BS1) and 0.356 m/s (BS2), which did not match well with the measured values 0.142 m/s (BS1) and 0.113 m/s (BS2) (Figure 7B). The reconstructed flow durations  $T_d$  were 188 s (BS1) and 120 s (BS2), which were relatively close to the measured values for both BS1 (362 s) and BS2 (233 s). The reconstructed values of total concentration  $C_{T,x_C}$  were 0.00858 (BS1) and 0.00972 (BS2), which also matched relatively well for both experiments BS1 (0.00532) and BS2 (0.00635). The reconstructed values of each grain-size class were mostly reasonable, but large deviations were observed for several reconstructions of concentrations with very low measured values ( $< 0.1\%$ ) (Figure 7D).



**Figure 8.** Reconstructed deposit profiles and the sampled deposit data of experiments BS1 and BS2. A. (1) Reconstructed and sampled  $\eta_T$  of BS1. (2) Grain size distribution at 1.4 m downstream. (3) Grain size distribution at 1.8 m downstream. B. (1) Reconstructed and sampled  $\eta_T$  of BS2. (2) Grain size distribution at 1.4 m downstream. (3) Grain size distribution at 1.8 m downstream.

**Table 7.** Flow conditions measured and reconstructed for experiments BS1 and BS2. (R.: reconstructed, M.: Measured)

Parameters	BS1	BS1	BS2	BS2
	R.	M.	R.	M.
$C_{T,x_C}$	0.00858	0.00532	0.00972	0.00635
$C_{1,x_C}$	0.000345	0.000514	0.00198	0.00146
$C_{2,BS,x_C}$	0.00155	0.00136	0.00249	0.00254
$C_{2,PP,x_C}$	0.00195	0.000974	0.00162	0.000686
$C_{3,BS,x_C}$	0.00147	0.00113	0.00143	0.00107
$C_{3,PP,x_C}$	0.00165	0.000823	0.00100	0.000418
$C_{4,x_C}$	0.00161	0.000523	0.00120	0.000182
$H_{x_H}$ (m)	0.127	0.209	0.156	0.199
$U_{x_U}$ (m/s)	0.330	0.142	0.356	0.113
$T_d$ (s)	188	362	120	233

## 6 Discussion

### 6.1 Validation of DNN as an Inversion Method for Turbidity Currents by Artificial Test Datasets

The results from verification with artificial datasets proved the ability of DNN to reasonably reconstruct the hydraulic conditions of turbidity current from its turbidite deposits. Reconstructions of initial flow conditions and the flow duration for artificial datasets (Sections 5.1.1 and 5.1.2) were good overall judging from the  $s$  and  $B$  values (Tables 4 and 5). The reconstructions of the flow duration  $T_d$ , flow depth  $H_0$ , velocity  $U_0$ , and sediment concentrations  $C_{1,0}$ ,  $C_{2,0}$  ( $C_{2,BS,0}$  and  $C_{2,PP,0}$  for mixture experiments),  $C_{3,0}$  ( $C_{3,BS,0}$  and  $C_{3,PP,0}$  for mixture experiments), and  $C_{4,0}$  were high in precision (Tables 4 and 5).

Overall, the correlations between the true and reconstructed values were obvious for all parameters reconstructed. Some scattering was observed for the reconstructed parameters, but most values stayed close within the range of perfect reconstruction. The range of misfit  $2s$  of all parameters lies below 14.3% of the matching representative value (Tables 4 and 5). A relatively greater degree of scattering is observed for  $U$  in comparison to the other parameters (Figures 4D and 5D).

With regard to the estimation bias, zero was included in the 95% confidence interval of bias for most of the parameters, proving that the reconstructed values were not significantly biased with respect to the original values. Even among the parameters where statistically significant biases were detected, their deviations were minor in comparison to the representative values of the parameters (Table 4 and 5). For example, in both artificial datasets using plastic particle only and mixture of siliciclastic sand and plastic particles, the estimation bias  $B$  for  $T_d$  had a negative value and the range of the confidence interval of  $B$  was below zero (Tables 4 and 5), indicating a tendency of underestimation for  $T_d$ . However, the bias for  $T_d$  was only 3.95% (mixture experiments) or 6.08% (plastic particle only) of the representative value of this parameter (600 s).

Thus, it can be said that this method is suitable for estimating the paleo-hydraulic conditions of actual turbidity currents. Reconstructed parameters are well correlated without any serious biases from true original values, implying that the in-

504 verse model produced in this study serves as a high precision, high accuracy estimator  
505 of flow conditions.

## 506 6.2 Verification of DNN Inversion with Flume Experiment Data

507 As a result of inversion using DNN, the overall deposit profiles were well re-  
508 constructed for all four experiments, and the reconstructed grain size distribution  
509 downstream matched with the grain size distribution sampled from experimental de-  
510 posits (Figures 6 and 8). DNN as an inverse method tries to find the combination of  
511 hydraulic conditions that best produces the deposit profiles inputted. The fact that  
512 the reconstructed hydraulic conditions accurately reproduced the deposit profiles used  
513 for inverse analysis indicates good performance by DNN with the inverse model.

514 Furthermore, the hydraulic conditions and flow duration reconstructed mostly lie  
515 within a reasonable range from the line of perfect reconstruction (Figure 7). Among  
516 the reconstructed hydraulic conditions, flow heights  $H_{x_H}$  were well reconstructed for  
517 all four experiments (Table 6 and 7). The differences between the reconstructed and  
518 measured values of  $H_{x_H}$  were less than 39.2%. The layer-averaged models with top-hat  
519 assumption (all shape factors are assumed to be unity) inevitably exhibit inaccuracies  
520 in the flow parameters due to their simplification of flow dynamics. Indeed, in previous  
521 experimental studies of the shallow water model, up to 50% deviation from the mea-  
522 sured values was observed for the calculated flow heights (Parker et al., 1987; Kostic &  
523 Parker, 2006). Considering the limitation of this model, the reconstructed flow heights  
524 in this study can be interpreted as reasonable estimates of the experimental results.

525 Flow durations  $T_d$  were also reasonably reconstructed but were underestimated  
526 for all four experiments, as observed for the artificial test datasets. The reconstructed  
527 concentration of each grain-size class  $C_{i,x_C}$  was scattered especially when the measured  
528 values were very low ( $< 0.1\%$ ), but the total concentrations of sediment exhibit rela-  
529 tively good agreement with the experimental values. In contrast to other parameters,  
530 flow velocities  $U_{x_U}$  were largely overestimated for all four experiments. Difference be-  
531 tween the reconstructed and measured  $U_{x_U}$  range from 24.6% to 215% of the measured  
532 values. Up to 25% deviation from the measured value was observed in previous re-  
533 search (Kostic & Parker, 2006), but deviation in this case exceeds that of the previous  
534 research.

535 One potential reason behind the deviations in the reconstructed values is that  
536 inverse analysis itself is essentially difficult for turbidity currents. For instance, flow  
537 parameters cannot be reconstructed when different combinations of initial conditions  
538 produce exactly the same deposit profile. However, this is unlikely because flows  
539 with a wide variety of initial conditions were tested with artificial datasets and the  
540 reconstructions were mostly reasonable, proving the ability of DNN to distinguish  
541 minor differences in the characteristics of deposits. Instead, sources of deviation could  
542 lie in the forward model employed in this study.

543 Through analysis of the result of application of DNN to flume experiment data,  
544 it became clear that there are three aspects for deviations in the reconstruction of  
545 hydraulic conditions: (1) bias inherent in the inverse model, (2) measurement errors  
546 during the experiments, and (3) inaccuracy within the forward model of turbidity  
547 currents employed in this study.

548 (1) Regarding the inherent bias in the inverse model,  $T_d$  reconstructed for the  
549 flume experiments exhibited the same tendencies of deviation during the reconstruction  
550 of artificial test datasets. Thus, deviation in the reconstruction of  $T_d$  may be considered  
551 as a systematic error originating from the internal settings of DNN.  $U_{x_U}$  exhibited  
552 similar tendencies in artificial datasets reconstruction to those in flume experiment  
553 reconstruction, but the deviation of  $U_{x_U}$  was so large for  $U_{x_U}$  in flume experiment

554 reconstruction that it seems unlikely to have originated solely from the DNN model.  
 555 The tendencies of deviation for  $C_i$  in the reconstruction of flume experiment datasets  
 556 did not match those for the artificial datasets, and thus also did not occur because of  
 557 the tendencies of the inverse model.

558 (2) The main source of deviation for sediment concentrations  $C_i$  may be inaccura-  
 559 cies in the measurements. As shown in Figure 7, some of the measured concentrations  
 560  $C_i$  were extremely small ( $< 0.1\%$ ), making them susceptible to minor disturbances  
 561 during sampling and measurements. For values that are extremely small, even minor  
 562 deviations appear to be large. Thus, for  $C_i$ , the main source of deviation may not be  
 563 the reconstructed values but the measured values.

564 (3) Inaccuracy in the forward model in describing the physical processes of tur-  
 565 bidity currents may account for deviations of the reconstructed flow velocities from  
 566 the measured values. There are several potential reasons to why the reconstruction of  
 567 flow velocity did not go as well as with the other parameters, but the most probable  
 568 reason is the inaccuracy of the entrainment function in describing the actual effect of  
 569 entrainment in flow, considering that the exponent in the calculation of dimensionless  
 570 vertical velocity in the entrainment function was determined purely by optimization  
 571 and differed greatly in previous studies (Parker et al., 1987; Garcia & Parker, 1991;  
 572 Dorrell et al., 2018). Another problem may lie in the layer averaging of flow velocity.  
 573 Even though the model uses layer averaged flow velocity for calculation, a recent study  
 574 by (Luchi et al., 2018) had pointed out that a two layer model is more suitable for  
 575 the calculation of turbidity currents. This research aims to verify DNN as a method  
 576 of inverse analysis of turbidity currents. Improvement of the forward model, including  
 577 entrainment function and velocity calculation, should be the next step in the inverse  
 578 analysis study of turbidity currents.

579 The shallow water model implemented holds certain limitations due to its sim-  
 580 plified calculation of flow dynamics. Nevertheless, the simplifications make it possible  
 581 to perform large batches of natural scale simulations. Overall, even though a certain  
 582 amount of deviation was observed for all parameters, they lie mostly within a reason-  
 583 able range for the eventual purpose of application to natural scale turbidity currents.

### 584 6.3 Comparison of DNN with Existing Methodologies

585 In comparison to previous methods of inverse analysis of turbidity currents, the  
 586 inversion method using DNN holds great advantage in terms of calculation cost and  
 587 in terms of accuracy of reconstruction. Previous inversion methods of turbidity cur-  
 588 rents seek to optimize the initial conditions to a particular set of data collected from  
 589 turbidite, which is extremely time-consuming to apply to one dataset and does not  
 590 guarantee the general applicability of the method to turbidite deposits (Lesshafft &  
 591 Marquet, 2010; Parkinson et al., 2017; Nakao & Naruse, 2017). For example, ge-  
 592 netic algorithm applied in Nakao and Naruse (2017) first initializes a population of  
 593 parameters and then optimizes the population of parameters through selection and  
 594 mutation. Eventually, the parameters left can successfully reconstruct the target tur-  
 595 bidite. However, each epoch of optimization requires the results of selection from the  
 596 previous epoch, and thus the calculation of the forward model cannot be parallelized  
 597 over epochs. In adjoint method used by Parkinson et al. (2017), control variables  
 598 within the forward model of turbidity currents are first initialized and inputted into  
 599 the numerical model. The turbidite deposit profile is calculated and compared with  
 600 the target values using a cost function. Gradients of the cost function (objective func-  
 601 tion) for control variables are calculated analytically. If the result is decided to be  
 602 less than optimal, the adjoint model will run and control variables will be adjusted by  
 603 descent method. The adjusted control variables will again be inputted into the nu-  
 604 merical model. This cycle is repeated until the reconstructed deposit profile is judged

605 to be optimal, thus the iteration of calculation cannot be performed simultaneously.  
 606 By contrast, DNN explores the general functional relationship between the turbidite  
 607 deposited and the current, making it applicable to turbidity currents in general. For-  
 608 ward model calculation to produce training datasets can be perfectly parallelized, such  
 609 that the amount of time needed for calculation reduces greatly.

610 Since parallelization of forward model calculation greatly reduced the time for  
 611 calculation, a more accurate and realistic forward model with heavier calculation load  
 612 was able to be implemented. As a result, the forward model used in this research is  
 613 much better at capturing the spatio-temporal evolution of turbidity current in com-  
 614 parison to forward model used in previous research (Falcini et al., 2009; Parkinson et  
 615 al., 2017). Falcini et al. (2009) used a steady flow forward model, whereas our forward  
 616 model is a non-steady flow model that reproduces the evolution of flow over time. The  
 617 method implemented in Parkinson et al. (2017) omitted the effect of entrainment, a  
 618 significant part of sediment transport in turbidity currents. As a result, their recon-  
 619 structed values of flow height, concentration, and grain diameter of the turbidite were  
 620 2.56 km, 0.0494%, and 103  $\mu\text{m}$ , respectively (Parkinson et al., 2017). These values  
 621 were largely off the expected range when examined based on the objective values col-  
 622 lected from the turbidite deposits. By contrast, our predictions lie relatively close to  
 623 the original values collected and the effect of sediment suspension was incorporated in  
 624 our forward model. Another improvement from previous research is that the forward  
 625 model used in this case applies to turbidite datasets of multiple grain-size classes.

626 By contrast, Lesshafft and Marquet (2010) proposed a method based on direct  
 627 numerical simulation (DNS) of the Navier-Stokes equations. However, the calculation  
 628 costs of the method were extremely high, making it unrealistic to apply the method  
 629 to natural scale turbidites. The computational cost of DNS was scaled to  $Re^3$ , thus  
 630 limiting the effectiveness of DNS to only experimental scale flows (Biegert et al., 2017).  
 631 As a result, the maximum value of Reynolds number attained in previous numerical  
 632 simulation using DNS was 15,000 (Cantero et al., 2007), which corresponds to 3.0 cm/s  
 633 for velocity and 50 cm for flow height. Thus, their methodology cannot be applied to  
 634 natural scale turbidites.

## 635 7 Conclusions

636 In this study, a new method for the inverse analysis of turbidite using DNN was  
 637 verified with actual flume experiment data. In comparison to previous methods, DNN  
 638 proved to be an efficient method for inversion of turbidity currents without compromis-  
 639 ing the accuracy of reconstruction. DNN performed well for verification with artificial  
 640 datasets, judging by the standard deviation and bias of the reconstructed parameters.  
 641 In terms of the application of DNN to experiment data, deposit profiles were well  
 642 reconstructed; however, the initial flow conditions did not match the measured values  
 643 perfectly.

644 The reconstructed flow heights  $H$  for PP1 and PP2 were 0.132 m and 0.159 m,  
 645 respectively, which were good matches with the measured values 0.116 m and 0.123  
 646 m.  $H$  reconstructed for BS1 and BS2 were 0.127 m and 0.156 m, respectively, and  $H$   
 647 measured were 0.209 m and 0.199 m, which exhibited larger differences than those  
 648 for PP1 and PP2 but were still reasonable.  $H$  was underestimated for BS1 and BS2.  
 649 On the other hand,  $T_d$  was underestimated for all experiments. For BS1 and BS2  
 650 the reconstructed values were 188 s and 120s, respectively, and measured values were  
 651 362 s and 233 s. For PP1 and PP2, the reconstructed values were 635 s and 872 s,  
 652 respectively, and the measured values were 936 s and 966 s.  $U$  was overestimated  
 653 for all experiments, with PP1 and PP2 better reconstructed than BS1 and BS2. The  
 654 reconstructed  $U$  for PP1 and PP2 were 0.118 m/s and 0.115 m/s, respectively, and  
 655 measured values were 0.0608 m/s and 0.0923 m/s. The reconstructed  $U$  for BS1 and

656 BS2 were 0.330 m/s and 0.356 m/s, respectively, and measured values were 0.142 m/s  
 657 and 0.113 m/s.

658 Overall, DNN exhibited good performance for the inversion of artificial datasets  
 659 and some parameters of the experiment data. The deposit profiles were well recon-  
 660 structed, indicating the success of DNN in the exploration of the functional relationship  
 661 between the initial conditions of flow and the resulting deposits. The results of ver-  
 662 ification with artificial datasets and flume experiments reveal that the implemented  
 663 forward model is competent in performing inverse analysis on turbidity currents, but  
 664 needs to be more robust for application to a wide range of flow conditions. Improve-  
 665 ment of forward models and parameters such as the entrainment function will be top  
 666 priority in the future. Hyperparameter settings and internal structure of DNN also  
 667 have room for improvement judging from the inversion result of artificial datasets.  
 668 Application of DNN to field datasets will be the eventual goal.

## 669 Appendix A Details of Forward Model Implemented

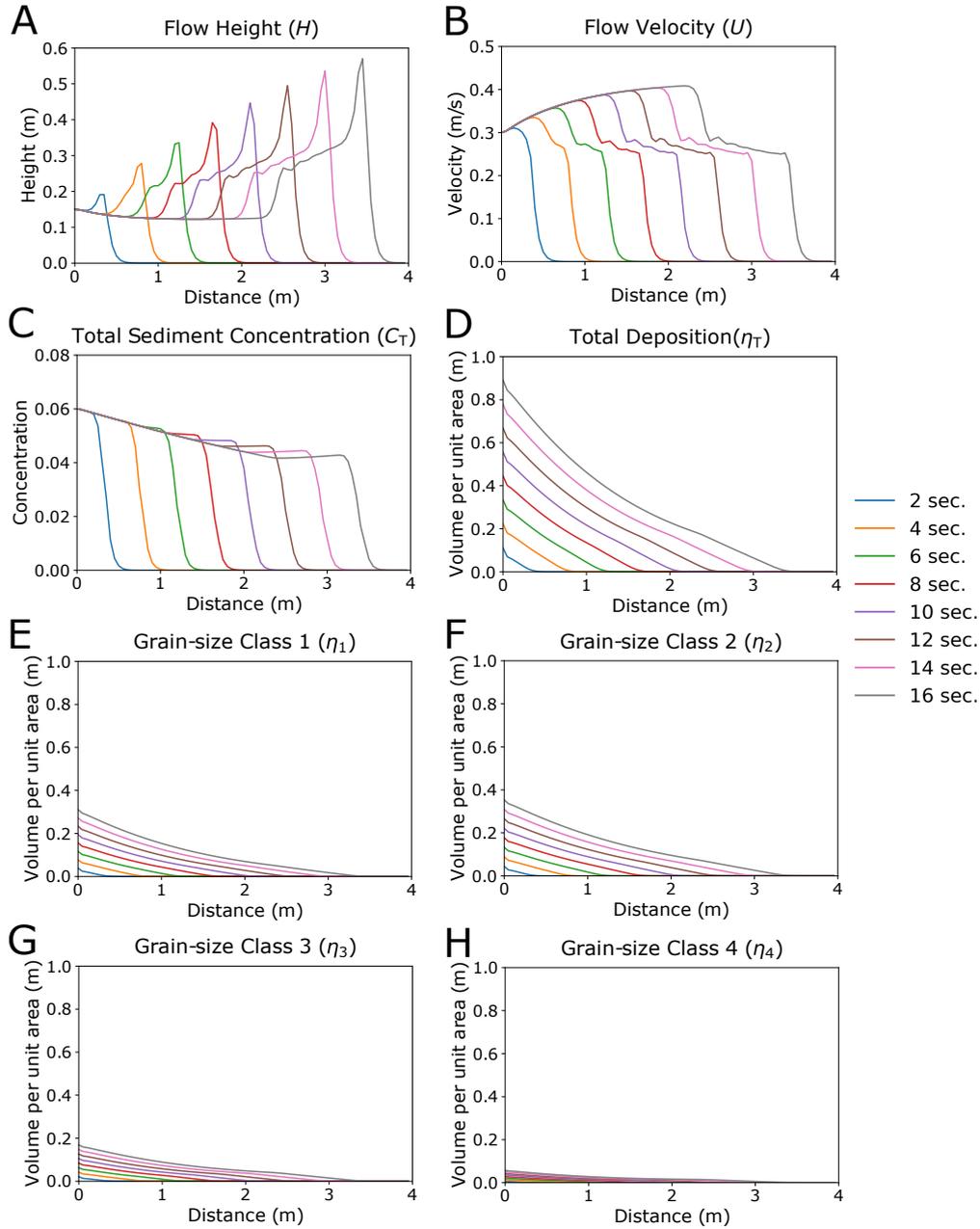
### 670 A1 Example of Forward Model Calculation

671 The forward model was tested with two sets of numerical simulations of turbidity  
 672 currents. Testing was conducted using the forward model programmed for the pro-  
 673 duction of artificial datasets for plastic particle only experiments. The settings of the  
 674 numerical simulations are listed in Table A1, whereas the time evolution of the high  
 675  $C_{T,0}$ ,  $U_0$  simulation is shown in Figure A1 and the time evolution of the low  $C_{T,0}$ ,  $U_0$   
 676 simulation is shown in Figure A2. In both cases, flow height  $H$  was greater toward the  
 677 head of the current.  $H$  at the head of the current also grew over time (Figures A1A  
 678 and A2A). Flow velocity  $U$  in the high  $C_{T,0}$ ,  $U_0$  simulation increased when the cur-  
 679 rent propagated downstream (Figure A1B), while  $U$  in the low  $C_{T,0}$ ,  $U_0$  simulation  
 680 increased initially then decreased as the current propagated (Figure A2B). The total  
 681 volumetric concentration of sediment  $C_T$  in flow decreased downstream in both cases  
 682 (Figures A1C and A2C). In the high  $C_{T,0}$ ,  $U_0$  case, a larger portion of sediment was  
 683 deposited downstream than in the low  $C_{T,0}$ ,  $U_0$  case (Figures A1D and A2D). The low  
 684  $C_{T,0}$ ,  $U_0$  case had the most sediment deposited toward the upstream end of flow.

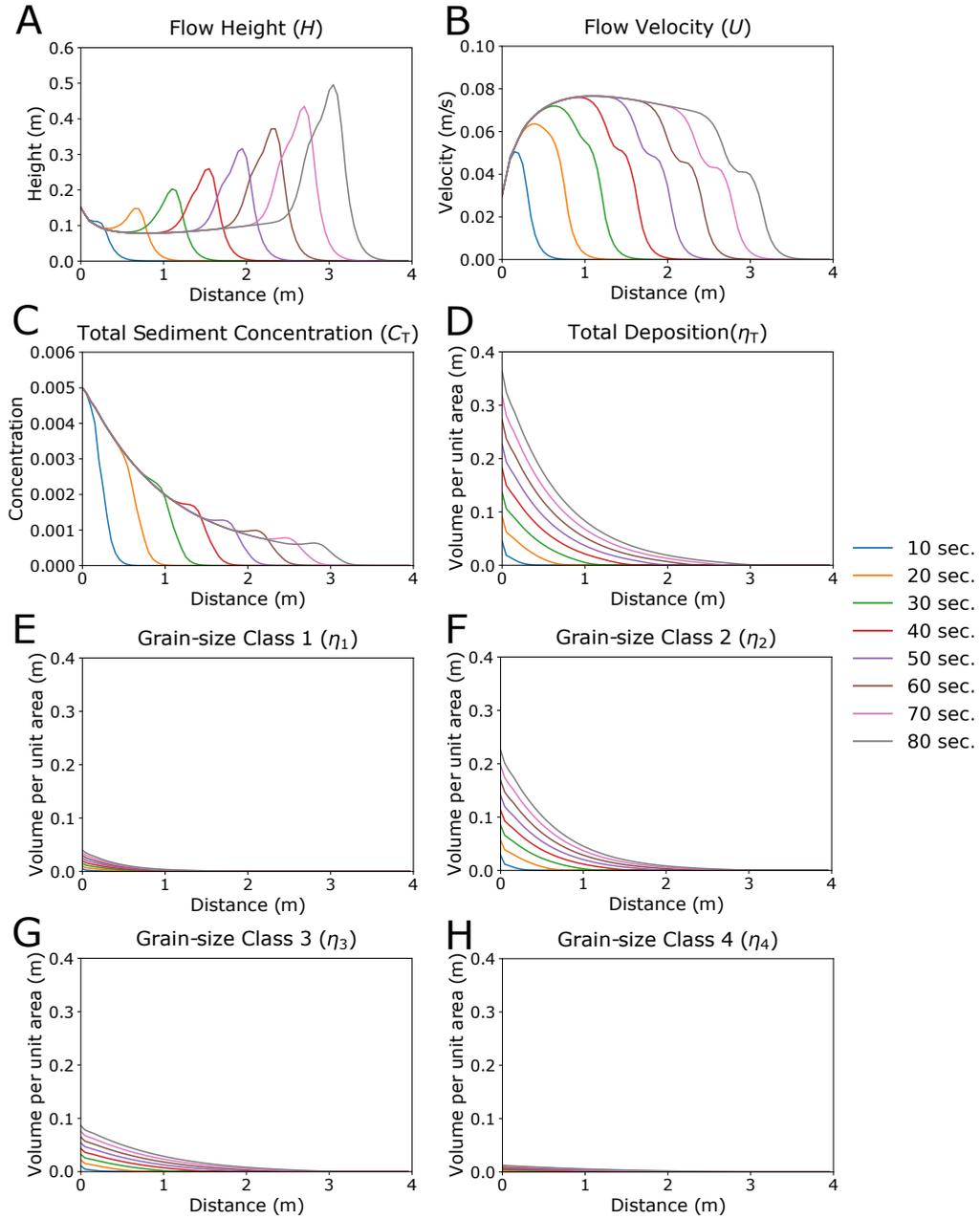
685 For the high  $C_{T,0}$ ,  $U_0$  simulation, a thicker deposit was observed for grain-size  
 686 class 1 than for grain-size classes 3 and 4 (Figures A1E,G,H). Even though the initial  
 687 concentrations of the finer grain-size classes 3 and 4  $C_{3,0}$ ,  $C_{4,0}$  were higher than that  
 688 of the coarser grain-size class 1 ( $C_{1,0}$ ), less fine sediment was deposited since it was  
 689 more likely to stay suspended and be carried beyond the lower flow boundary by the  
 690 high-velocity flow. For the low  $C_{T,0}$ ,  $U_0$  simulation, the grain-size class with a thicker  
 691 deposit has a higher initial concentration. The coarsest grain-size class, grain-size  
 692 class 1, had almost all sediment deposited near the upstream boundary, while the finest  
 693 grain-size class, grain-size class 4, had sediment spread out toward the downstream  
 694 direction (Figures A2E,H). This happened because the low-velocity flow was unable  
 695 to keep the coarse sediment suspended.

**Table A1.** Initial flow conditions of numerical simulations of turbidity currents.

	High $C_{T,0}, U_0$	Low $C_{T,0}, U_0$
$H_0$ (m)	0.15	0.15
$U_0$ (m/s)	0.3	0.03
$C_{T,0}$	0.06	0.005
$C_{1,0}$	0.01	0.00025
$C_{2,0}$	0.02	0.0025
$C_{3,0}$	0.018	0.00175
$C_{4,0}$	0.012	0.0005
$c_f$	0.002	0.002
$r_o$	1.5	1.5
Duration (s)	420	420



**Figure A1.** Example of forward model calculation with high initial flow velocity and sediment concentration (Table A1). A. Time evolution of flow height  $H$ . B. Time evolution of flow velocity  $U$ . C. Time evolution of total sediment volumetric concentration  $C_T$ . D. Time evolution of deposit profile  $\eta_T$ . E. Time evolution deposit profile of grain-size class 1  $\eta_1$ . F. Time evolution of deposit profile of grain-size class 2  $\eta_1$ . G. Time evolution of deposit profile of grain-size class 3  $\eta_1$ . H. Time evolution of deposit profile of grain-size class 4  $\eta_1$ .



**Figure A2.** Example of forward model calculation with low initial flow velocity and sediment concentration (Table A1). A. Time evolution of flow height  $H$ . B. Time evolution of flow velocity  $U$ . C. Time evolution of total sediment volumetric concentration  $C_T$ . D. Time evolution of deposit profile  $\eta_T$ . E. Time evolution deposit profile of grain-size class 1  $\eta_1$ . F. Time evolution of deposit profile of grain-size class 2  $\eta_1$ . G. Time evolution of deposit profile of grain-size class 3  $\eta_1$ . H. Time evolution of deposit profile of grain-size class 4  $\eta_1$ .

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**A2 Sensitivity Tests of Forward Model**

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The degree of sensitivity of the forward model against changes in the initial conditions of the flow and model parameters was tested (Table A2). Testing was conducted using the forward model programmed for the production of artificial datasets of plastic particle only experiments. Numerical simulations were conducted with different values of the six parameters  $H_0$ ,  $U_0$ ,  $C_{T,0}$ ,  $e_s$ ,  $r_o$ ,  $c_f$ . Other parameters were unchanged for the simulations.

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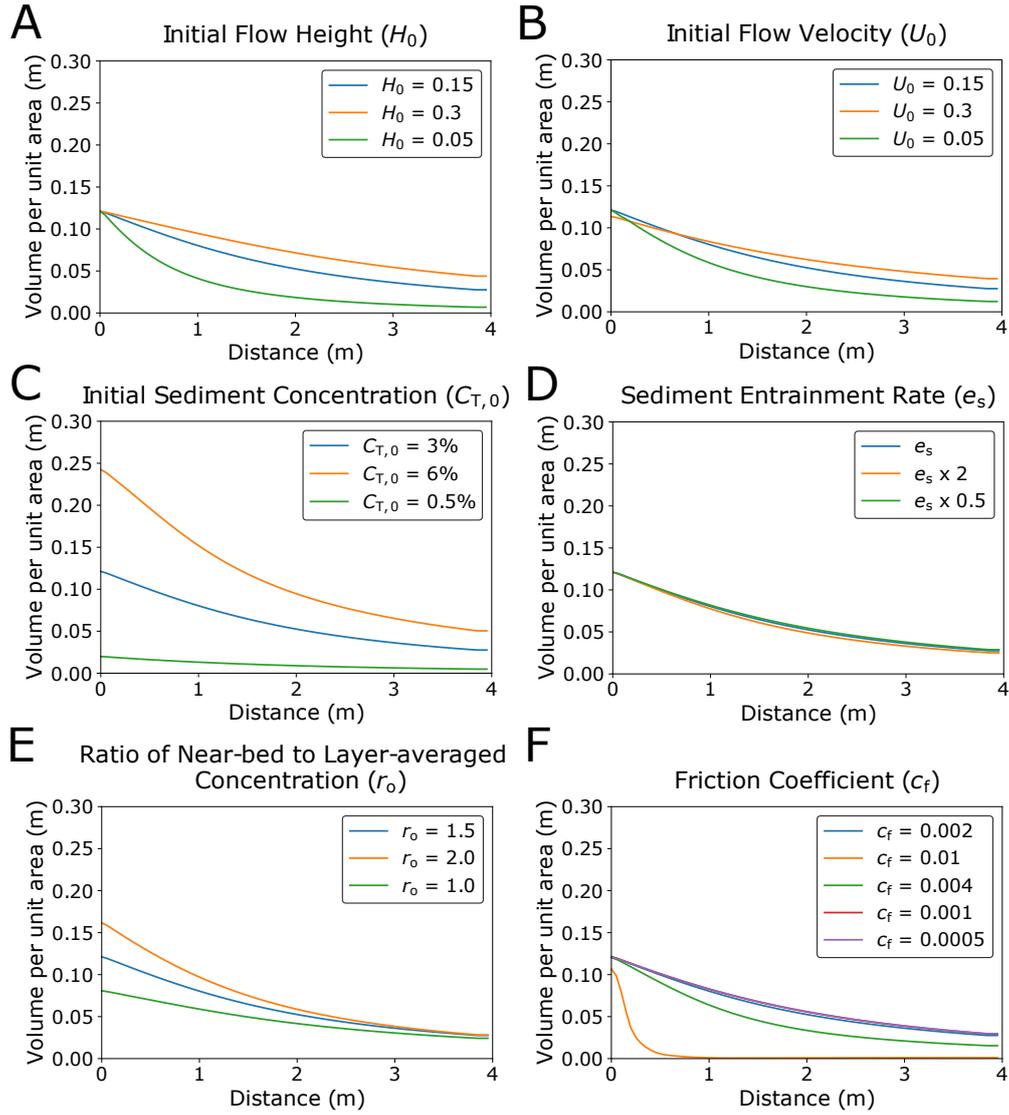
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The results of the sensitivity tests revealed that changes in deposit profile occurs when the initial flow conditions differ (Figure A3). The volume of sediment deposited increased overall as  $H_0$  increased (Figure A3A). The same trend was observed for  $U_0$ , and  $C_{T,0}$  (Figure A3B, C). Out of these three parameters, the amount of increase in deposit thickness was greatest for  $C_{T,0}$ , and smallest for  $U_0$ . With regard to model closure parameters, the resultant deposit profile exhibited nearly no change for different values of entrainment coefficient  $e_s$  (Figure A3D). Slightly lower amount of deposition was observed for greater  $e_s$ . An increase in the amount of deposition was observed as  $c_f$  decreased (Figure A3F). The thickness of deposit increased moderately when  $r_o$  increased.

**Table A2.** Settings for sensitivity tests of forward model.

Case	$H_0$ (m)	$U_0$ (m/s)	$C_{T,0}$	$e_s$	$r_o$	$c_f$
1	0.15	0.15	0.03	GP	1.5	0.002
2	0.3	0.15	0.03	GP	1.5	0.002
3	0.05	0.15	0.03	GP	1.5	0.002
4	0.15	0.3	0.03	GP	1.5	0.002
5	0.15	0.05	0.03	GP	1.5	0.002
6	0.15	0.15	0.06	GP	1.5	0.002
7	0.15	0.15	0.005	GP	1.5	0.002
8	0.15	0.15	0.03	GPx2	1.5	0.002
9	0.15	0.15	0.03	GPx0.5	1.5	0.002
10	0.15	0.15	0.03	GP	2.0	0.002
11	0.15	0.15	0.03	GP	1.0	0.002
12	0.15	0.15	0.03	GP	1.5	0.01
13	0.15	0.15	0.03	GP	1.5	0.004
14	0.15	0.15	0.03	GP	1.5	0.001
15	0.15	0.15	0.03	GP	1.5	0.0005



**Figure A3.** Sensitivity tests of deposit profile of artificial turbidites against change in initial flow conditions and closure parameters (Table A2). A. Dependency on initial flow height  $H_0$ . B. Dependency on initial flow velocity  $U_0$ . C. Dependency on initial total sediment volumetric concentration  $C_{T,0}$ . D. Dependency on sediment entrainment rate  $e_s$ . E. Dependency on ratio of near-bed to layer-averaged concentration  $r_o$ . F. Dependency on friction coefficient  $c_f$ .

713 **Notation**

714	$\alpha_1, \alpha_2$	Parameters related to sediment entrainment
715	$c_f$	Friction coefficient
716	$C_i$	Layer-averaged volumetric concentration of suspended sediment of the $i$ th grain-
717		size class
718	$C_T$	Layer-averaged total concentration of suspended sediment
719	$D_i$	Representative grain diameter of the $i$ th grain-size class
720	$e_{si}$	Entrainment rate of sediment of the $i$ th grain-size class into suspension
721	$e_w$	Entrainment rate of ambient water to flow
722	$F_i$	Volume fraction of the $i$ th grain-size class in active layer
723	$g$	Gravitational acceleration
724	$H$	Flow height
725	$L_a$	Active layer thickness
726	$R$	Submerged specific density of sediment
727	$R_{fi}$	Dimensionless particle fall velocity of the $i$ th grain-size class
728	$Re_{pi}$	Particle Reynolds number of the $i$ th grain-size class
729	$r_o$	Ratio of near-bed suspended sediment concentration to the layer-averaged concen-
730		tration of suspended sediment
731	$S$	Slope gradient
732	$S_f$	Friction slope
733	$t$	Time
734	$T_d$	Flow duration
735	$U$	Layer-averaged flow velocity
736	$u_*$	Shear velocity
737	$w_i$	Settling velocity of a sediment particle of the $i$ th grain-size class
738	$x$	Streamwise distance
739	$\eta_i$	Volume per unit area of bed sediment of the $i$ th grain-size class
740	$\eta_T$	Total volume per unit area of bed sediment
741	$\kappa$	Parameter related to artificial viscosity
742	$\lambda_p$	Porosity of bed sediment
743	$\mu$	Dynamic viscosity of water
744	$\nu$	Kinematic viscosity of water
745	$\rho$	Density of water

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750 **References**

- 751 Abadi, M., Agarwal, A., Barham, P., Brevdo, E., Chen, Z., Citro, C., . . . Zheng,  
 752 X. (2015). *TensorFlow: Large-scale machine learning on heterogeneous sys-*  
 753 *tems*. Retrieved from <http://tensorflow.org/> (Software available from  
 754 tensorflow.org)
- 755 Altinakar, M., Graf, W., & Hopfinger, E. (1996). Flow structure in turbidity cur-  
 756 rents. *Journal of Hydraulic Research*, 34(5), 713-718. Retrieved from [https://](https://doi.org/10.1080/00221689609498467)  
 757 [doi.org/10.1080/00221689609498467](https://doi.org/10.1080/00221689609498467) doi: 10.1080/00221689609498467
- 758 Arai, K., Naruse, H., Miura, R., Kawamura, K., Hino, R., Ito, Y., . . . Kasaya, T.  
 759 (2013, 11). Tsunami-generated turbidity current of the 2011 tohoku-oki earth-

- 760 quake. *Geology*, *41*(11), 1195–1198.
- 761 Baas, J. H., Van Dam, R. L., & Storms, J. E. A. (2000). Duration of deposi-  
762 tion from decelerating high-density turbidity currents. *Sedimentary Geology*,  
763 *136*(1), 71–88.
- 764 Biegert, E., Vowinkel, B., Ouillon, R., & Meiburg, E. (2017). High-resolution sim-  
765 ulations of turbidity currents. *Progress in Earth and Planetary Science*, *4*(1),  
766 33.
- 767 Bottou, L. (2010). Large-scale machine learning with stochastic gradient descent.  
768 In Y. Lechevallier & G. Saporta (Eds.), *Proceedings of compstat'2010* (pp.  
769 177–186). Heidelberg: Physica-Verlag HD.
- 770 Bouma, A. H. (1962). *Sedimentology of some flysch deposits: A graphic approach to*  
771 *facies interpretation*. Elsevier.
- 772 Cantero, M. I., Balachandar, S., & Garcia, M. H. (2007). High-resolution simulations  
773 of cylindrical density currents. *Journal of Fluid Mechanics*, *590*, 437–469.
- 774 Daly, R. A. (1936). Origin of submarine canyons. *American Journal of Science*,  
775 *31*(186), 401–420.
- 776 Davison, A. C., & Hinkley, D. V. (1997). *Bootstrap methods and their application*.  
777 Cambridge University Press. doi: 10.1017/CBO9780511802843
- 778 Dietrich, W. E. (1982). Settling velocity of natural particles. *Water Resources Re-*  
779 *search*, *18*(6), 1615–1626. doi: 10.1029/WR018i006p01615
- 780 Dorrell, R. M., Amy, L. A., Peakall, J., & McCaffrey, W. D. (2018). Particle size  
781 distribution controls the threshold between net sediment erosion and deposi-  
782 tion in suspended load dominated flows. *Geophysical Research Letters*, *45*(3),  
783 1443–1452.
- 784 Falcini, F., Marini, M., Milli, S., & Moscatelli, M. (2009). An inverse problem to  
785 infer paleoflow conditions from turbidites. *Journal of Geophysical Research:*  
786 *Oceans*, *114*.
- 787 Fukushima, Y., Parker, G., & Pantin, H. M. (1985). Prediction of ignitive turbidity  
788 currents in scripps submarine canyon. *Marine Geology*, *67*(1), 55–81.
- 789 Garcia, M., & Parker, G. (1991). Entrainment of bed sediment into suspension.  
790 *Journal of Hydraulic Engineering*, *117*(4), 414–435.
- 791 Garcia, M., & Parker, G. (1993). Experiments on the entrainment of sediment  
792 into suspension by a dense bottom current. *Journal of Geophysical Research*,  
793 *98*(C3), 4793–4807.
- 794 Gibbs, R. J. (1974). A settling tube system for sand-size analysis. *Journal of Sedi-*  
795 *mentary Petrology*, *44*(2), 583–588.
- 796 Jameson, A., Schmidt, W., & Turkel, E. (1981). Numerical solution of the euler  
797 equations by finite volume methods using runge kutta time stepping schemes.  
798 In *14th fluid and plasma dynamics conference* (p. 1259).
- 799 Johnson, D. W. (1939). *The origin of submarine canyons: A critical review of hy-*  
800 *potheses*. Columbia University Press.
- 801 Kostic, S., & Parker, G. (2006). The response of turbidity currents to a canyon–fan  
802 transition: Internal hydraulic jumps and depositional signatures. *Journal*  
803 *of Hydraulic Research*, *44*(5), 631–653. Retrieved from [https://doi.org/](https://doi.org/10.1080/00221686.2006.9521713)  
804 [10.1080/00221686.2006.9521713](https://doi.org/10.1080/00221686.2006.9521713) doi: 10.1080/00221686.2006.9521713
- 805 Krizhevsky, A., Sutskever, I., & Hinton, G. E. (2012). Imagenet classification  
806 with deep convolutional neural networks. In F. Pereira, C. J. C. Burges,  
807 L. Bottou, & K. Q. Weinberger (Eds.), *Advances in neural information pro-*  
808 *cessing systems 25* (pp. 1097–1105). Curran Associates, Inc. Retrieved from  
809 [http://papers.nips.cc/paper/4824-imagenet-classification-with-deep](http://papers.nips.cc/paper/4824-imagenet-classification-with-deep-convolutional-neural-networks.pdf)  
810 [-convolutional-neural-networks.pdf](http://papers.nips.cc/paper/4824-imagenet-classification-with-deep-convolutional-neural-networks.pdf)
- 811 Kuenen, P. H., & Migliorini, C. I. (1950). Turbidity currents as a cause of graded  
812 bedding. *The Journal of Geology*, *58*(2), 91–127.
- 813 Lesshafft, L., & Marquet, O. (2010). Optimal velocity and density profiles for the  
814 onset of absolute instability in jets. *Journal of Fluid Mechanics*, *662*, 398–408.

- 815 Liang, S., & Srikant, R. (2016). Why deep neural networks for function approxima-  
816 tion? *CoRR*, *abs/1610.04161*. Retrieved from <http://arxiv.org/abs/1610>  
817 [.04161](http://arxiv.org/abs/1610)
- 818 Luchi, R., Balachandar, S., Seminara, G., & Parker, G. (2018). Turbidity cur-  
819 rents with equilibrium basal driving layers: A mechanism for long runout.  
820 *Geophysical Research Letters*, *45*(3), 1518-1526. Retrieved from [https://](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2017GL075608)  
821 [agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2017GL075608](https://agupubs.onlinelibrary.wiley.com/doi/abs/10.1002/2017GL075608) doi:  
822 [10.1002/2017GL075608](https://doi.org/10.1002/2017GL075608)
- 823 Nakao, K., & Naruse, H. (2017, 12). Inverse analysis to reconstruct hydraulic condi-  
824 tions of non-steady turbidity currents based on multiple grain-size classes.  
825 *EarthArXiv*.
- 826 Naruse, H. (2005). Usage and advantages of an application program “stube” for set-  
827 tling tube grain-size analysis. *Journal of the Sedimentological Society of Japan*,  
828 *62*(62), 55–61.
- 829 Naruse, H., & Nakao, K. (2017, 12). Inverse analysis of turbidites by machine learn-  
830 ing. *AGU Fall Meeting Abstracts*.
- 831 Naruse, H., & Olariu, C. (2008). Hydraulic conditions of turbidity currents es-  
832 timated by inverse analysis. In *Fourth international conference on scour and*  
833 *erosion* (pp. 591–593). Japanese Geotechnical Society.
- 834 Nielsen, M. A. (2015). *Neural networks and deep learning*. Determination Press.
- 835 Parker, G., Garcia, M., Fukushima, Y., & Yu, W. (1987). Experiments on turbidity  
836 currents over an erodible bed. *Journal of Hydraulic Research*, *25*(1), 123–147.
- 837 Parkinson, S. D., Funke, S. W., Hill, J., Piggott, M. D., & Allison, P. A. (2017).  
838 Application of the adjoint approach to optimise the initial conditions of a  
839 turbidity current with the adjointturbidity 1.0 model. *Geoscientific Model*  
840 *Development*, *10*(3), 1051–1068.
- 841 Paull, C. K., Talling, P. J., Maier, K. L., Parsons, D., Xu, J., Caress, D. W., ...  
842 Cartigny, M. J. (2018). Powerful turbidity currents driven by dense basal  
843 layers. *Nature Communications*, *9*(1), 4114.
- 844 Pradhan, B., Lee, S., & Buchroithner, M. F. (2010). A gis-based back-propagation  
845 neural network model and its cross-application and validation for landslide  
846 susceptibility analyses. *Computers, Environment and Urban Systems*, *34*(3),  
847 216–235.
- 848 Rogers, S. J., Fang, J. H., Karr, C. L., & Stanley, D. A. (1992). Determination of  
849 lithology from well logs using a neural network. *AAPG Bulletin*, *76*(5), 731–  
850 739.
- 851 Ruder, S. (2016). An overview of gradient descent optimization algorithms. *CoRR*,  
852 *abs/1609.04747*. Retrieved from <http://arxiv.org/abs/1609.04747>
- 853 Rumble, J. R. (2018). *Crc handbook of chemistry and physics, 99th edition*. CRC  
854 Press.
- 855 Schmidhuber, J. (2015). Deep learning in neural networks: An overview. *Neural Net-*  
856 *works*, *61*, 85–117. Retrieved from [http://www.sciencedirect.com/science/](http://www.sciencedirect.com/science/article/pii/S0893608014002135)  
857 [article/pii/S0893608014002135](http://www.sciencedirect.com/science/article/pii/S0893608014002135) doi: <https://doi.org/10.1016/j.neunet.2014>  
858 [.09.003](https://doi.org/10.1016/j.neunet.2014)
- 859 Talling, P. J., Allin, J., Armitage, D. A., Arnott, R. W. C., Cartigny, M. J. B.,  
860 Clare, M. A., ... Xu, J. (2015). Key future directions for research on turbidity  
861 currents and their deposits. *Journal of Sedimentary Research*, *85*(2), 153–169.
- 862 Talling, P. J., Amy, L. A., & Wynn, R. B. (2007). New insight into the evolution of  
863 large-volume turbidity currents: Comparison of turbidite shape and previous  
864 modelling results. *Sedimentology*, *54*(4), 737–769.
- 865 Talling, P. J., Masson, D. G., Sumner, E. J., & Malgesini, G. (2012). Subaqueous  
866 sediment density flows: Depositional processes and deposit types. *Sedimentol-*  
867 *ogy*, *59*(7), 1937–2003.
- 868 Vangriesheim, A., Khripounoff, A., & Crassous, P. (2009). Turbidity events observed  
869 in situ along the congo submarine channel. *Deep Sea Research Part II: Topical*

- 870 *Studies in Oceanography*, 56(23), 2208–2222.
- 871 Wang, Z., & Bovik, A. C. (2009). Mean squared error: Love it or leave it? a new  
872 look at signal fidelity measures. *IEEE Signal Processing Magazine*, 26(1), 98–  
873 117. doi: 10.1109/MSP.2008.930649
- 874 Weimer, P., & Slatt, R. M. (2007). Introduction to the petroleum geology of deepwa-  
875 ter settings. *AAPG Studies in Geology*, 57, 149–227.
- 876 Wright, S., & Parker, G. (2004). Flow resistance and suspended load in sand-bed  
877 rivers: Simplified stratification model. *Journal of Hydraulic Engineering*,  
878 130(8), 796–805.
- 879 Xu, J. P., Noble, M. A., & Rosenfeld, L. K. (2004). In-situ measurements of velocity  
880 structure within turbidity currents. *Geophysical Research Letters*, 31(9).
- 881 Yabe, T., Xiao, F., & Utsumi, T. (2001). The constrained interpolation profile  
882 method for multiphase analysis. *Journal of Computational Physics*, 169(2),  
883 556–593.